## Anisotropic clustering in the Baryon Oscillation Spectroscopic Survey



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## Motivation for studying Redshift Space Distortions

Growth function G(a): δ(k, a) = aG(a)δ<sub>i</sub>(k)
In General Relativity G(a) is determined once H(a) is specified/measured; generically this relation is different in modified gravity models

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## SDSSIII WiggleZAnistropic Clustering: $P(k_{\perp}, k_{\parallel})$



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Outline

Our basic model for galaxy clustering

- Anisotropic galaxy clustering
  - Alcock-Paczynski effect
  - Redshift space distortions
- First results from BOSS
  - Error budget and future prospects

# Galaxy clustering lightning theory review

 Theory I: underlying matter power spectrum (determined at z >~ z<sub>CMB</sub>, neglecting V)

Theory II: Expansion history H(0 < z < z<sub>GAL</sub>)

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## Matter Power Spectrum

- Entire P(k) (not just BAO) acts as standard ruler determined by CMB
- We marginalize over the (negligible) uncertainty



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**Z**2

## Theory II: geometry

We measure  $\theta$ ,  $\phi$ , and z for each galaxy, and use a cosmological model to convert to comoving coordinates  $z_1$ 

 $\Theta$ 

#### $\chi(z)$ (or $D_A(z)$ )

I/H(z)

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## SDSSIII Theory II: Alcock-Paczynski

 $\xi(r_P, \pi)$  appears anisotropic if you assume the wrong cosmological model (constrain  $\eta_{AP} = D_A * H$ )

#### $\chi(z) =_0 \int^z c \, dz' / H(z')$

BAO in  $\xi_0(s)$  determines "geometric mean"  $D_V \propto (D_A^2 H^{-1})^{1/3}$ 



 $X(z)^*\Delta\theta$ 

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## **Redshift Space Distortions**

#### $\theta, \phi, redshift$

depends on the geometry of the universe

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 $\chi(z) = \chi_{true} + v_p/aH(a)$ 

 $\chi(z) =_0 \int^z c dz' / H(z')$ 

comoving coordinates: x, y, z

## SDSSIII Redshift Space Distortions (RSD)

#### real to redshift space separations

 $\nabla \cdot \mathbf{v_p} = -aHf \, \delta_m$ 

#### $|v_P| \sim d \sigma_8/d \ln a = \sigma_8 * f$

isotropic squashed along line of sight

X

 $f=d\,\ln\,\sigma_8\,/d\,\ln\,a\,\approx\,\Omega_m{}^\gamma$ 

## SDSSIII RSD: linear theory (Kaiser 1987)

 $\delta_g^s(k) = (b + f\mu_k^2)\delta_m^r(k)$ 

 $\mu_k^2 = k_z^2 / k^2$ 

## SDSSIII Legendre Polynomial moments: P(k)

#### General Expansion

 $P(k,\mu_k) = \sum_{k} P_{\ell}(k)L_{\ell}(\mu_k)$ 

#### Linear theory prediction

 $\begin{pmatrix} P_0(k) \\ P_2(k) \\ P_4(k) \end{pmatrix} = P_m^r(k) \begin{pmatrix} b^2 + \frac{2}{3}bf + \frac{1}{5}f^2 \\ \frac{4}{3}bf + \frac{4}{7}f^2 \\ \frac{8}{35}f^2 \end{pmatrix}$ 

## SDSSIII Legendre Polynomial moments: $\xi(r)$

#### General Expansion

$$\xi(s,\mu_s) = \sum_{\ell} \xi_{\ell}(s) L_{\ell}(\mu_s)$$

#### Relation to $P_{\ell}(k)$

$$\xi_{\ell}(s) = i^{\ell} \int \frac{k^2 dk}{2\pi^2} P_{\ell}(k) j_{\ell}(ks)$$

#### SDSSIII Modeling RSD: Reid and White 2011 (arXiv: 1105.4165)

 $\xi_0,\,\xi_2$  sufficient to constrain  $b\sigma_8,\,f\sigma_8$  ; MOST of 2d clustering information retained



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## SDSSII

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#### Fitting to 2d clustering

- Use full model of  $\xi_{0,2}$  (s  $\geq 25 \text{ h}^{-1} \text{ Mpc}$ ) to constrain:
- growth of structure (f $\sigma_8$ )
- $D_V \propto (D_A^2/H)^{1/3}$ 
  - Alcock-Paczynski ( $\eta_{AP} \propto D_A(z_{eff}) * H(z_{eff})$ )
    - marginalizing over shape of underlying linear P(k),  $b\sigma_8$ ,  $\sigma_{FOG}^2$

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#### Alcock-Paczynski in multipoles



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# DR9 spectroscopic results: preliminary!

- DR9 data final (public July 2012), clustering/ covariances ~final, cosmological constraints preliminary
  - Current uncertainties reported, not central values

## SDSSIII BOSS "CMASS" (z<sub>eff</sub> = 0.57) galaxy sample in perspective



Eisenstein et al. arXiv:1101.1529

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**SDSSIII** 

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## BAO fits in $P(k)/\xi(r)$ consistent

X. Xu et al. (in prep; DR7) BOSS Galaxy Clustering (in prep.)

#### BAO fit plot was here

2-3% uncertainty on BAO position in angle-averaged  $P(k)/\xi(r)$ 

Constrains  $D_V \propto (D_A^2/H)^{1/3}$ 



#### The CMASS measurements

#### • 26 log bins in s for $\xi_0$ and $\xi_2 = 52$ DOF

Measurement of  $\xi_0/\xi_2$  was here

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#### Model Fits

• We test the LCDM hypothesis in 4 models, always marginalizing over P(k) shape and  $\sigma^{2}_{FOG}$ :

- LCDM (bσ<sub>8</sub>)
- LCDM + fσ<sub>8</sub>: (bσ<sub>8</sub>, fσ<sub>8</sub>)
  - LCDM + geometry:  $(b\sigma_8, D_V, D_A*H)$
- LCDM++:  $(b\sigma_8, f\sigma_8, D_V, D_A*H)$

### SDSS

Current status

- $D_V/D_{V,fid} = x \pm 0.019$  (i.e., minimal information gain on  $D_V$  compared to BAO only!)
- Geometry LCDM:  $f\sigma_8 = xx \pm 0.03$  (7%) [WMAP7 LCDM: 0.45 ± 0.025]
- $f\sigma_8 LCDM: \eta = xx \pm 0.04$  (4%) [WMAP7 LCDM: 1.00 ± 0.012]
  - Fit both:  $f\sigma_8 = xx \pm 0.07$ ,  $\eta = xx \pm 0.07$

#### SDSSIII Testing alternative models with amplitude of peculiar velocities



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#### Expansion rate at z=0.57



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#### Error Budget/Future Prospects



s<sub>min</sub> (Mpc/h)

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#### Error Budget/Future Prospects



s<sub>min</sub> (Mpc/h)

## SDSSII

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#### Summary/Conclusions

#### • DR9 CMASS results:

- high significance detection of BAO in  $\xi_0(r)$ ,  $P_0(k)$  (~2% constraint on  $Dv \propto D_A^2/H$ )
- 7% (4%) measurement of  $f\sigma_8$  (D<sub>A</sub> \* H) at z=0.57
- Two "easy" ways to improve our precision:
  - use information on small scales to constrain  $\sigma^2_{FOG}$ .
  - Push modeling of halo clustering to smaller scales