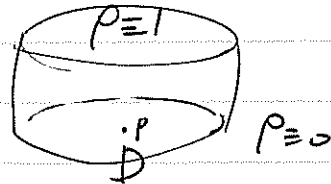


Ward identity

$$\langle \delta \mathcal{O}(p) \mathcal{O}'_1 \dots \mathcal{O}'_s \rangle = \frac{1}{2\pi} \int_{\Sigma} \sqrt{g} d\sigma \epsilon^\nu \nabla^\mu \langle T_{\mu\nu}^c \mathcal{O}(p) \mathcal{O}'_1 \dots \mathcal{O}'_s \rangle$$

if $\epsilon(p'_1) = \dots = \epsilon(p'_s) = 0$.

Take $\epsilon = \rho \epsilon^z(z) \frac{\partial}{\partial z}$
 ↑
 holomorphic



$\rho = \begin{cases} 1 & \text{in } D \\ 0 & \text{outside} \end{cases}$

operator identity

$$\Rightarrow \delta_{\epsilon} \mathcal{O}(p) = \frac{1}{2\pi} \int_D \underbrace{\sqrt{g} d\sigma}_{i g_{z\bar{z}} dz d\bar{z}} \epsilon^z \underbrace{\nabla^\mu T_{\mu z}^c}_{g^{z\bar{z}} \partial_{\bar{z}} T_{z\bar{z}}^{\text{hol}}} \mathcal{O}(p) = T_{z\bar{z}}$$

$$= \frac{i}{2\pi} \int_D dz d\bar{z} \epsilon^z(z) \partial_{\bar{z}} T_{z\bar{z}} \mathcal{O}(p)$$

$$= \frac{-i}{2\pi} \int_D d (dz \epsilon^z(z) T_{z\bar{z}}(z) \mathcal{O}(p))$$

$$= \frac{1}{2\pi i} \oint_{\partial D} dz \epsilon^z(z) T_{z\bar{z}}(z) \mathcal{O}(p)$$

$$= \frac{1}{2\pi i} \oint_p dz \epsilon^z(z) T_{z\bar{z}}(z) \mathcal{O}(p)$$

↑
 small contour that encircles p only.

For $\epsilon = \rho \epsilon^{\bar{z}}(\bar{z}) \frac{\partial}{\partial \bar{z}}$

$$\delta_{\bar{\epsilon}} \mathcal{O}(p) = \frac{i}{2\pi} \int_D d\bar{z} d\bar{z} \underbrace{\epsilon^{\bar{z}}(\bar{z}) \partial_{\bar{z}} T_{\bar{\epsilon}\bar{z}} \mathcal{O}(p)}_{d(d\bar{z} \epsilon^{\bar{z}}(\bar{z}) T_{\bar{\epsilon}\bar{z}} \mathcal{O}(p))}$$

$$= \frac{i}{2\pi} \oint_{\partial D} d\bar{z} \epsilon^{\bar{z}}(\bar{z}) T_{\bar{\epsilon}\bar{z}} \mathcal{O}(p)$$

$$= -\frac{1}{2\pi i} \oint_p d\bar{z} \epsilon^{\bar{z}}(\bar{z}) T_{\bar{\epsilon}\bar{z}} \mathcal{O}(p)$$

Thus Ward ids for holomorphic and antiholomorphic vectors

are

$$\delta_{\epsilon} \mathcal{O}(p) = \frac{1}{2\pi i} \oint_p d\bar{z} \epsilon^{\bar{z}}(\bar{z}) T_{\bar{\epsilon}\bar{z}}(z) \mathcal{O}(p)$$

$$\delta_{\bar{\epsilon}} \mathcal{O}(p) = -\frac{1}{2\pi i} \oint_p d\bar{z} \epsilon^{\bar{z}}(\bar{z}) T_{\bar{\epsilon}\bar{z}}(z) \mathcal{O}(p)$$

This defines an action of hol / antihol vector fields

on the space of local operators.

↑
infinitesimal conformal transf.

eg.

$$\delta \epsilon T_{ww}(w) = \frac{1}{2\pi i} \oint_w dz \epsilon^2(z) \underbrace{T_{zz}(z) T_{ww}(w)}_{\substack{\frac{c}{2} \\ (z-w)^4 + \frac{2}{(z-w)^2} T_{ww} + \frac{1}{z-w} \partial_w T_{ww} \\ + \text{reg}}}}$$

$$= \frac{1}{2\pi i} \oint_w dz \left\{ \frac{\epsilon(z) \frac{c}{2}}{(z-w)^4} + \frac{2\epsilon(z)}{(z-w)^2} T_{ww} + \frac{\epsilon(z)}{z-w} \partial_w T_{ww} \right\}$$

$$\epsilon(w) + (z-w)\epsilon'(w) + \frac{1}{2}(z-w)^2 \epsilon''(w) + \frac{1}{6}(z-w)^3 \epsilon'''(w) + O((z-w)^4)$$

$$\epsilon(w) + (z-w)\epsilon'(w) + O((z-w)^2)$$

$$= \frac{c}{12} \epsilon'''(w) + 2\epsilon'(w) T_{ww}(w) + \epsilon(w) \partial_w T_{ww}$$

↑ Infinitesimal form of $T_{ww}^{(w)} = (\bar{z}'(w))^2 T_{zz}(z) + \frac{c}{12} \{z, w\}$

Similarity

$$\delta \bar{\epsilon} T_{\bar{w}\bar{w}}(\bar{w}) = \frac{c}{12} \bar{\epsilon}'''(\bar{w}) + 2\bar{\epsilon}'(\bar{w}) T_{\bar{w}\bar{w}} + \bar{\epsilon}(\bar{w}) \partial_{\bar{w}} T_{\bar{w}\bar{w}}$$

We denote the action of $z^{n+1} \frac{\partial}{\partial z}$, $\bar{z}^{n+1} \frac{\partial}{\partial \bar{z}}$ on operators at $z=0$ as L_n, \tilde{L}_n :

$$(L_n \mathcal{O})(0) = \frac{1}{2\pi i} \oint_0 dz z^{n+1} T_{zz}(z) \mathcal{O}(0)$$

$$(\tilde{L}_n \mathcal{O})(0) = -\frac{1}{2\pi i} \oint_0 d\bar{z} \bar{z}^{n+1} T_{\bar{z}\bar{z}}(\bar{z}) \mathcal{O}(0)$$

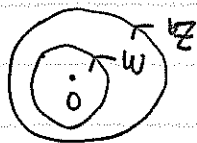
L_n, \tilde{L}_n can also be defined on operators at any point p :

$$(L_n \mathcal{O})(p) = \frac{1}{2\pi i} \oint_p dz (z - z(p))^{n+1} T_{zz}(z) \mathcal{O}(p)$$

$$(\tilde{L}_n \mathcal{O})(p) = -\frac{1}{2\pi i} \oint_p d\bar{z} (\bar{z} - \bar{z}(p))^{n+1} T_{\bar{z}\bar{z}}(\bar{z}) \mathcal{O}(p)$$

The algebra of L_n 's & \tilde{L}_n 's

$$(L_n(L_m \mathcal{O}))(0) = \oint \frac{dz}{2\pi i} z^{n+1} T_{zz}(z) \oint \frac{dw}{2\pi i} w^{m+1} T_{ww}(w) \mathcal{O}(0)$$



$$([L_n, L_m] \mathcal{O})(0) = (L_n(L_m \mathcal{O}))(0) - (L_m(L_n \mathcal{O}))(0)$$

$$= \int \frac{dz}{2\pi i} \frac{dw}{2\pi i} z^{n+1} w^{m+1} T_{zz}(z) T_{ww}(w) \mathcal{O}(0)$$

$$= \oint_0 \frac{dw}{2\pi i} w^{m+1} \int \frac{dz}{2\pi i} z^{n+1} T_{zz}(z) T_{ww}(w) \mathcal{O}(0)$$

$$= \oint_0 \frac{dw}{2\pi i} w^{m+1} \left\{ \frac{c}{12} (w^{n+1})''' + 2(w^{n+1})' T_{ww}(w) + w^{n+1} \partial_w T_{ww}(w) \right\} \mathcal{O}(0)$$

$$= \oint_0 \frac{dw}{2\pi i} \left\{ \frac{c}{12} w^{m+1} (w^{n+1})''' + 2 w^{m+1} (w^{n+1})' T_{ww} - (w^{m+n+2})' T_{ww} \right\} \mathcal{O}(0)$$

$$\frac{c}{12} w^{m+n-1} (n+1)n(n-1) \quad (2(n+1) - (m+n+2)) w^{m+n+1} T_{ww}$$

$$= \frac{c}{12} \delta_{m+n,0} (n^3 - n) \mathcal{O}(0) + (n-m) \underbrace{\oint_0 \frac{dw}{2\pi i} w^{m+n+1} T_{ww} \mathcal{O}(0)}_{(L_{n+m} \mathcal{O})(0)}$$

$$(L_{n+m} \mathcal{O})(0)$$

$$\boxed{[L_n, L_m] = (n-m) L_{n+m} + \frac{c}{12} \delta_{n+m,0} (n^3 - n)}$$

This is called the Virasoro algebra of central charge C .

"Vir_C".

Similarly we find

$$[\tilde{L}_n, \tilde{L}_m] = (n-m) \tilde{L}_{n+m} + \frac{C}{12} \delta_{n+m,0} (n^3 - n)$$

$$[L_n, \tilde{L}_m] = 0$$

The infinitesimal conformal transformations of operators realized a representation of the Lie algebra

$$\begin{array}{cc} \text{Vir}_C \oplus \text{Vir}_C & \\ \uparrow & \uparrow \\ \{L_n\} & \{\tilde{L}_n\} \end{array}$$

$$\underline{L_{-1}, \tilde{L}_{-1}}$$

$T_{\mu\nu}$ can also be regarded as Noether current for translations. Thus,

$$L_{-1} \mathcal{O}(p) = \frac{1}{2\pi i} \oint_p dz T_{zz}(z) \mathcal{O}(p) = \frac{\partial}{\partial z} \mathcal{O}(p)$$

$$\tilde{L}_{-1} \mathcal{O}(p) = -\frac{1}{2\pi i} \oint_p d\bar{z} T_{\bar{z}\bar{z}}(\bar{z}) \mathcal{O}(p) = \frac{\partial}{\partial \bar{z}} \mathcal{O}(p).$$

$$\underline{L_0, \tilde{L}_0}$$

$$\text{Suppose } L_0 \mathcal{O} = \Delta \mathcal{O}, \quad \tilde{L}_0 \mathcal{O} = \tilde{\Delta} \mathcal{O}$$

(i.e. \mathcal{O} is an eigenoperator for L_0, \tilde{L}_0 .)

$$\text{recall } z \frac{\partial}{\partial z} + \bar{z} \frac{\partial}{\partial \bar{z}} \quad \dots \text{ dilatation}$$

$$z \frac{\partial}{\partial z} - \bar{z} \frac{\partial}{\partial \bar{z}} \quad \dots \text{ rotation.}$$

$\Delta + \tilde{\Delta}$ is the dimension of the operator \mathcal{O}

$\Delta - \tilde{\Delta}$ is the spin of the \mathcal{O} .

or \uparrow angular momentum

$\Delta, \tilde{\Delta}$ play the analogous role as energy in ordinary QFT. (called conformal weights)

The algebra \supset

$$[L_0, L_{-n}] = nL_{-n}, \quad [\tilde{L}_0, \tilde{L}_{-n}] = n\tilde{L}_{-n}$$

L_{-n} raises the (right) conformal weight by n

\tilde{L}_{-n} raises the (left) " " by n .

An operator that cannot be lowered by L_n, \tilde{L}_n ($n \geq 1$)

ie.

$$L_n \mathcal{O} = 0 \quad \tilde{L}_n \mathcal{O} = 0 \quad \forall n \geq 1$$

is called a primary operator (or primary field)

The basic idea (of representation theory of $\text{Vir}_c \oplus \text{Vir}_c$):

Classify all operators into primary operators

and their descendants.

[An operator of the form $L_{-n_1} \dots L_{-n_k} \mathcal{O}$ ($n_1, \dots, n_k \geq 1$)
are called a descendant of \mathcal{O} .]

Example

$$\mathcal{O} = \mathbb{1}$$

$$(L_n \mathbb{1})(0) = \frac{1}{2\pi i} \oint_0 d\tilde{z} \underbrace{\tilde{z}^{n+1} T_{z\tilde{z}}(z)}_{\text{regular as } z \rightarrow 0 \text{ if } n \geq -1} \mathbb{1}(0) = 0 \quad n \geq -1$$

$$(\tilde{L}_n \mathbb{1})(0) = 0 \quad n \geq -1.$$

$\therefore \mathcal{O} = \mathbb{1}$ is a primary operator of conformal weight $(\Delta, \tilde{\Delta}) = (0, 0)$.

$$(L_{-2} \mathbb{1})(0) = \frac{1}{2\pi i} \oint_0 d\tilde{z} \tilde{z}^{-1} T_{z\tilde{z}}(z) = T_{z\tilde{z}}(0)$$

$$(\tilde{L}_{-2} \mathbb{1})(0) = -\frac{1}{2\pi i} \oint_0 d\tilde{z} \tilde{z}^{-1} T_{\tilde{z}\tilde{z}}(\tilde{z}) = T_{\tilde{z}\tilde{z}}(0)$$

$\therefore T_{z\tilde{z}}, T_{\tilde{z}\tilde{z}}$ are descendants of $\mathcal{O} = \mathbb{1}$.

• Note: L_n, \tilde{L}_n depends on the coordinate system.

Suppose z and w are holomorphic coordinates

both $z(p) = w(p) = 0$ at a point p .

Then, they are related by $w = cz + O(z^2) = w(z)$

$$(L_n^{\{w\}} \mathcal{O})_{(0)} = \frac{1}{2\pi i} \oint_0 dw \underbrace{w^{n+1}}_{w'(z) dz} \underbrace{T_{ww}(w)}_{(w'(z))^{-2} (T_{zz}(z) - \frac{c}{12} \{w, z\})} \mathcal{O}(0)$$

$$= \frac{1}{2\pi i} \oint_0 dz w(z)^{n+1} (w'(z))^{-1} (T_{zz}(z) - \frac{c}{12} \{w, z\}) \mathcal{O}(0)$$

$w'(z)^{-1}$ is analytic at $z=0$ $\therefore w(z)^{n+1} w'(z)^{-1} \{w, z\}$ is analytic at $z=0$ (\rightarrow no pole)

$$= \frac{1}{2\pi i} \oint_0 dz \underbrace{w(z)^{n+1} (w'(z))^{-1}}_{z^{n+1} \cdot \left(\frac{w(z)}{z}\right)^{n+1} (w'(z))^{-1}} T_{zz}(z) \mathcal{O}(0)$$

$$z^{n+1} \cdot \left(\frac{w(z)}{z}\right)^{n+1} (w'(z))^{-1} = z^{n+1} \left(c^n + \sum_{m=1}^{\infty} c_m^n z^m \right)$$

$$= \left(\left(c^n L_n^{\{z\}} + \sum_{m=1}^{\infty} c_m^n L_{n+m}^{\{z\}} \right) \mathcal{O} \right)_{(0)}$$

• The definition of primary does not depend on coordinate system you use.

• For a primary operator, conformal weights do not depend on coordinate system you use.

Note: $(L_n \mathcal{O})(0) = \oint_0 \frac{dz}{2\pi i} z^{n+1} T_{zz}(z) \mathcal{O}(z)$

$$\Leftrightarrow T_{zz}(z) \mathcal{O}(0) = \sum_{n \in \mathbb{Z}} \bar{z}^{-n-2} (L_n \mathcal{O})(0)$$

If \mathcal{O} is a primary operator, $n \geq 1$ terms vanish:

$$\begin{aligned} T_{zz}(z) \mathcal{O}(0) &= \sum_{n \leq 0} \bar{z}^{-n-2} (L_n \mathcal{O})(0) \\ &= \bar{z}^{-2} \underbrace{(L_0 \mathcal{O})(0)}_{\substack{\text{if conformal} \\ \text{weight } \Delta}} + \bar{z}^{-1} \underbrace{(L_{-1} \mathcal{O})(0)}_{\parallel} + \bar{z}^0 \underbrace{(L_{-2} \mathcal{O})(0)}_{\parallel} + \dots \\ &\quad \parallel \quad \parallel \\ &\quad \Delta \mathcal{O}(0) \quad \partial_{\bar{z}} \mathcal{O}(0) \end{aligned}$$

\therefore For a primary operator \mathcal{O} of conformal weight $\Delta, \tilde{\Delta}$:

$$T_{zz}(z) \mathcal{O}(w) = \frac{\Delta}{(z-w)^2} \mathcal{O}(w) + \frac{1}{z-w} \partial_w \mathcal{O}(w) + \text{regular as } z \rightarrow w$$

$$T_{\bar{z}\bar{z}}(z) \mathcal{O}(w) = \frac{\tilde{\Delta}}{(\bar{z}-\bar{w})^2} \mathcal{O}(w) + \frac{1}{\bar{z}-\bar{w}} \partial_{\bar{w}} \mathcal{O}(w) + \text{regular}$$

Note: $\mathcal{O}(w) = T_{ww}(w)$ is indeed NOT a primary operator.

Consequence :

$$\begin{aligned} (\delta_\epsilon \mathcal{O})(z) &= \oint_z \frac{dw}{2\pi i} \epsilon(w) \underbrace{T_{ww}(w)}_{\left(\leftarrow \text{primary of weight } \Delta \right)} \mathcal{O}(z) \\ &= \frac{\Delta}{(w-z)^2} \mathcal{O}(z) + \frac{1}{2w-z} \partial_z \mathcal{O}(z) + \text{reg} \\ &= \Delta \epsilon'(z) \mathcal{O}(z) + \epsilon(z) \partial_z \mathcal{O}(z) \end{aligned}$$

$$(\delta_{\bar{\epsilon}} \mathcal{O})(z) = \tilde{\Delta} \bar{\epsilon}'(\bar{z}) \mathcal{O}(z) + \bar{\epsilon}(\bar{z}) \partial_{\bar{z}} \mathcal{O}(z)$$

→ Finite version : under $z \rightarrow w(z)$

$$\mathcal{O}(z) \rightarrow (w'(z))^\Delta (\bar{w}'(\bar{z}))^{\tilde{\Delta}} \mathcal{O}(w(z))$$

$$\mathcal{O} = \mathcal{O}(z) (dz)^\Delta (d\bar{z})^{\tilde{\Delta}} \rightarrow \mathcal{O}(w(z)) (dw)^\Delta (d\bar{w})^{\tilde{\Delta}} = w^* \mathcal{O}.$$

∴ primary operator of conformal weight $(\Delta, \tilde{\Delta})$ defines a " $(\Delta, \tilde{\Delta})$ -differential".

$$\mathcal{O} = \mathcal{O}(z) (dz)^\Delta (d\bar{z})^{\tilde{\Delta}}$$

Example massless scalar

$$T_{zz} = - : \partial_z X \partial_z X : , \quad T_{\bar{z}\bar{z}} = - : \partial_{\bar{z}} X \partial_{\bar{z}} X :$$

$$T_{zz}(z) \partial_w X(w) = - : \partial_z X(z) \overbrace{\partial_z X(z) \partial_w X(w)}^{\substack{\text{const } \delta(z-w) \\ -\frac{1}{2(z-w)^2}}}: \times z - : (\partial_z X)^2 \partial_w X : \uparrow_{\text{reg}}$$

$$= \frac{1}{(z-w)^2} \partial_z X(z) + \text{reg}$$

$$= \frac{1}{(z-w)^2} \partial_w X(w) + \frac{1}{z-w} \partial_w(\partial_w X) + \text{reg.}$$

$$T_{\bar{z}\bar{z}}(z) \partial_w X(w) = - : \partial_{\bar{z}} X(z) \overbrace{\partial_{\bar{z}} X(z) \partial_w X(w)}^{\text{const } \delta(z-w)}: \times z - : (\partial_{\bar{z}} X)^2 \partial_w X : \uparrow_{\text{reg}}$$

$$\left(\Rightarrow (\tilde{L}_n \partial_w X) = 0 \quad n \geq 0 \right)$$

$\therefore \partial_z X$ is a primary operator of conformal weight $(\Delta, \tilde{\Delta}) = (1, 0)$.

$$\left(\Rightarrow \text{dimension } 1, \text{ spin } = 1 \right)$$

Similarly

$\partial_{\bar{z}} X$ is a " " " " $(\Delta, \tilde{\Delta}) = (0, 1)$

$$\left(\Rightarrow \text{dimension } 1, \text{ spin } = -1 \right)$$

Example massless Dirac

$$T_{zz} = -\frac{1}{2} : \bar{\Psi}_- \partial_z \Psi_- : + \frac{1}{2} : \partial_z \bar{\Psi}_- \Psi_- :$$

$$T_{\bar{z}\bar{z}} = -\frac{1}{2} : \bar{\Psi}_+ \partial_{\bar{z}} \Psi_+ : + \frac{1}{2} : \partial_{\bar{z}} \bar{\Psi}_+ \Psi_+ :$$

$$T_{zz}(z) \Psi_-(w) = -\frac{1}{2} : \overbrace{\bar{\Psi}_-(z) \partial_z \Psi_-(z)}^{\Psi_-(z)} : + \frac{1}{2} : \partial_z \overbrace{\bar{\Psi}_-(z) \Psi_-(z)}^{\Psi_-(z)} : + \text{reg.}$$

$$= -\frac{1}{2} \frac{1}{z-w} \underbrace{\partial_z \Psi_-(z)}_{(-1) \leftarrow \text{statistics}} + \frac{1}{2} (-1) \frac{-1}{(z-w)^2} \Psi_-(z) + \text{reg.}$$

$$= \frac{\frac{1}{2}}{(z-w)^2} \Psi_-(z) + \frac{1}{z-w} \partial_z \Psi_-(z) + \text{reg.}$$

$$T_{\bar{z}\bar{z}}(z) \Psi_-(w) = \text{reg.}$$

$\therefore \Psi_-$ is a primary operator of weight $(\Delta, \tilde{\Delta}) = (\frac{1}{2}, 0)$

$$\Rightarrow \text{dimension} = \frac{1}{2}, \text{ spin} = \frac{1}{2}$$

$\bar{\Psi}_-$ is also. $(\Delta, \tilde{\Delta}) = (\frac{1}{2}, 0)$, dimension $\frac{1}{2}$, spin $\frac{1}{2}$.

Ψ_+ is a primary op of $(\Delta, \tilde{\Delta}) = (0, \frac{1}{2})$ (dim = $\frac{1}{2}$, spin = $-\frac{1}{2}$)

$\bar{\Psi}_+$ //

exercises (HW)

$$T_{zz}(z) : \bar{\psi}_-(\omega) \psi_-(\omega) : = ?$$

$$T_{zz}(z) : e^{i\hbar X(\omega)} : = ?$$