

# Yang Mills Integrals

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+ ongoing, unpublished work with

Werner Krauth + Jan Plefka (AEI - Potsdam)

What are Yang-Mills integrals?

Answer: D-dimensional Euclidean gauge field theory reduced to zero dimensions.

$$Z_{0,N}^{\mathcal{N}} = \left( \prod_{A=1}^{\mathcal{N}} \int \prod_{\mu=1}^D \frac{dX_{\mu}^A}{\sqrt{2\pi}} \right) \left( \prod_{\alpha=1}^{\mathcal{N}} d\psi_{\alpha}^A \right) e^{-S[X, \Psi]}$$

gauge group:  $SU(N)$

Euclidean action:

$$S[X, \Psi] = \frac{1}{2} \text{Tr} [X_{\mu} X_{\nu}] [X_{\mu} X_{\nu}] + \text{Tr} \Psi_{\alpha} [ [\gamma_{\beta}^{\mu} X_{\mu} ] \gamma_{\beta}^{\nu} \Psi_{\nu} ]$$

Susy:  $\mathcal{N} = 2, 4, 8, 16 \Rightarrow D = 3, 4, 6, 10$

"Bosonic":  $\mathcal{N} = 0 \Rightarrow D = 2, 3, 4, 5, \dots$

• Integrating out fermions gives a

Pfaffian: Homogeneous polynomial in  $X$

Why should we study them?

- "D-branes":  $p = -1$ : D-instanton



- "Pretentious attempts" to define M-theory

"Matrix String Theory":  $10 \rightarrow 2$

"BFSS" or "matrix QM":  $10 \rightarrow 1$

"IKKT":  $10 \rightarrow 0$   $Z_{IKKT} = \sum_{N=2}^{\infty} Z_{10, N}^{16} e^{-\beta N}$

- SUSY index calculations:  $(-1)^F$

- $\mathcal{N} = 4, D = 4$   $SU(\infty)$  multi-instanton calculus

K-instanton sector related to  $Z_{D=10, N=K}^{\mathcal{N}=16}$

(Dovey et al. hep-th/9901128)

- An old idea: "Eguchi-Kawai reduction" (1982)

(see Lovas + Kitazawa 82)



Do YM integrals make sense?

- They look very singular ("valleys").
- back-of-the-envelope calculation:  $D=2$  diverges.

However, for (say)  $SU(2)$  it was ("accidentally") discovered (T'Hooft, Jethava-Stevenson) that convergence is possible:

special gauge:  $X_{\mu}^A = b_{\mu} \delta_{\mu}^A$  (no sum.)  $A=1,2,3$ ,

using  $SU(2) \times SO(D)$  invariance:

$$Z_{0, N=2}^{\text{off}} \sim \int_{0 \leq b_1 \leq b_2 \leq b_3} db_1 db_2 db_3 \Delta(b_i^2) (b_1 b_2 b_3)^{D-3 + \frac{N}{2}} e^{-\frac{1}{2} b_1^2 b_2^2 - \frac{1}{2} b_1^2 b_3^2 - \frac{1}{2} b_2^2 b_3^2}$$

Vandermonde:  $\Delta(b_i^2) = (b_1^2 - b_2^2)(b_1^2 - b_3^2)(b_2^2 - b_3^2)$

(convergence properties) can be seen after a "Nicolai map":

$$\gamma_1 = b_2 b_3 \quad \gamma_2 = b_1 b_3 \quad \gamma_3 = b_1 b_2$$

$$Z_{0, N=2}^{\text{off}} \sim \int_{0 \leq \gamma_1 \leq \gamma_2 \leq \gamma_3} d\gamma_3 d\gamma_2 d\gamma_1 \Delta(\gamma_i^2) (\gamma_1 \gamma_2 \gamma_3)^{\frac{D}{2} - 3 + \frac{N}{4}} e^{-\frac{1}{2} \gamma_1^2 - \frac{1}{2} \gamma_2^2 - \frac{1}{2} \gamma_3^2}$$

$\rightarrow = D-4$  (Susy:  $N > 0$ )

$SU(2)$  convergence:

Bosonic:

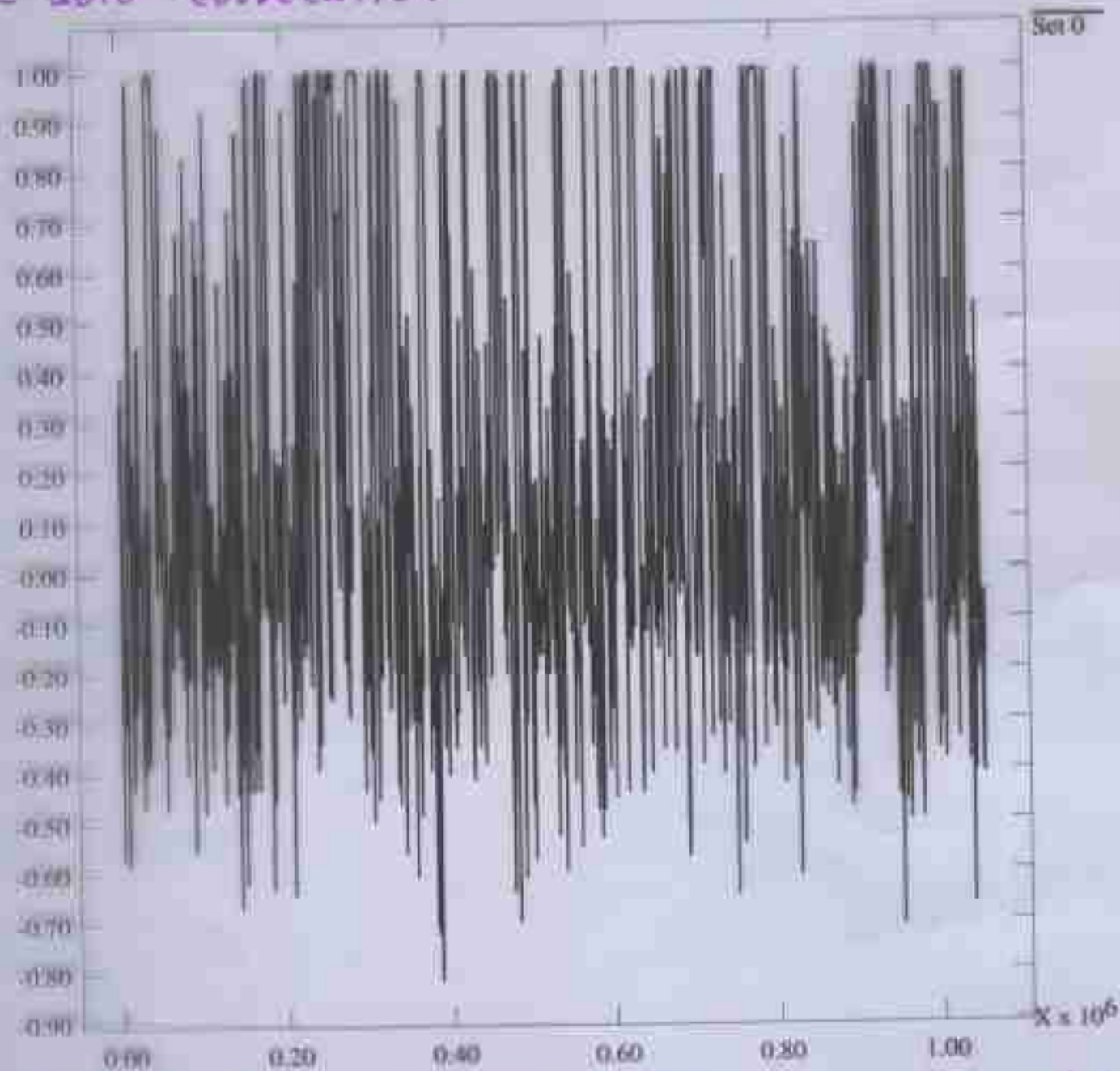
$$D \geq 5$$

Susy:

$$D = 4, 6, 10$$

$N \geq 3$  convergence??

A convergent Yang-Mills integral  
 $\gamma = \text{auto-correlation}$

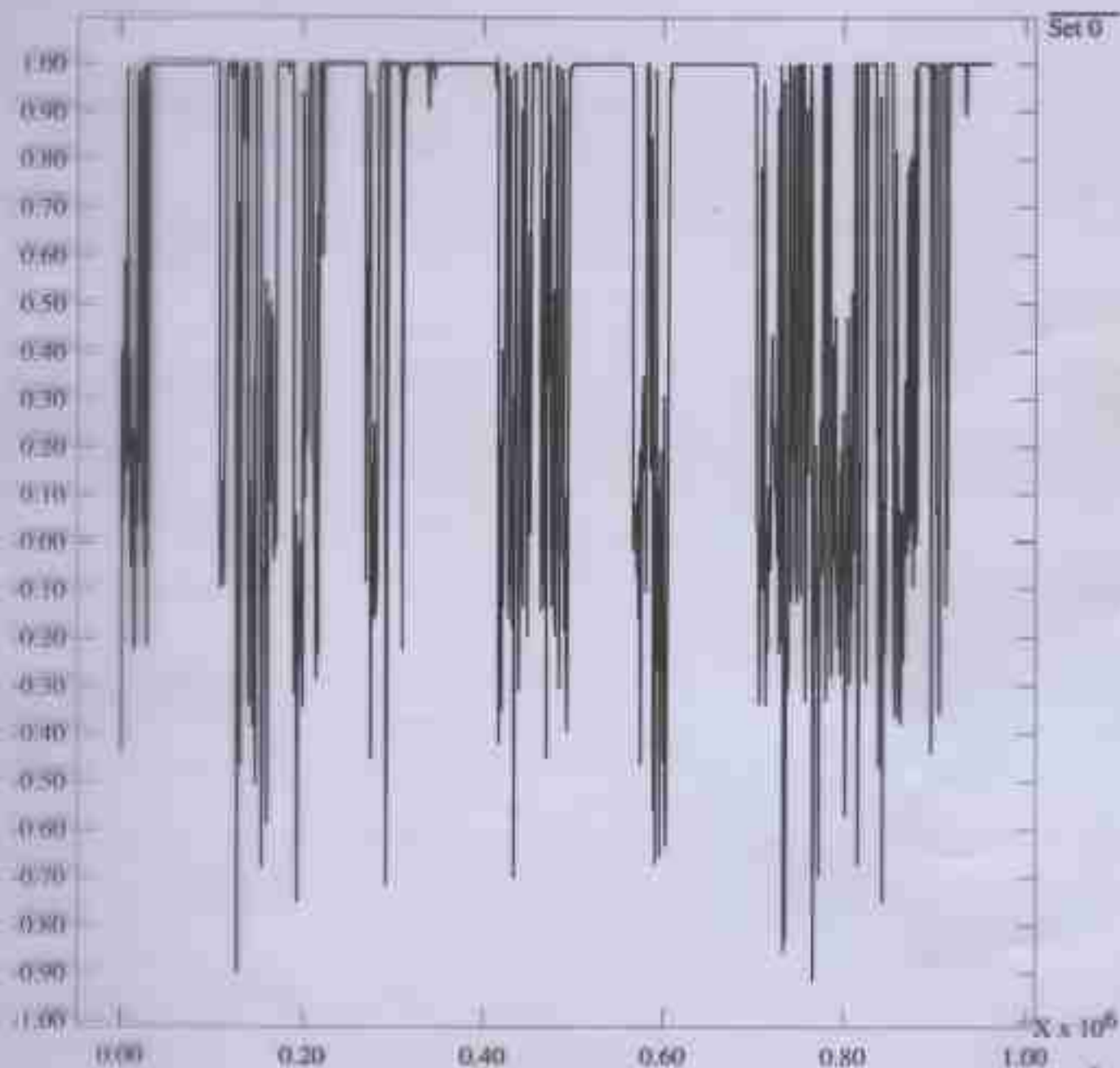


Example: Bosonic  $SU(2)$   $D=5$

# updates

# A (marginally) divergent Yang-Mills integral

$\gamma$  = auto-correlation

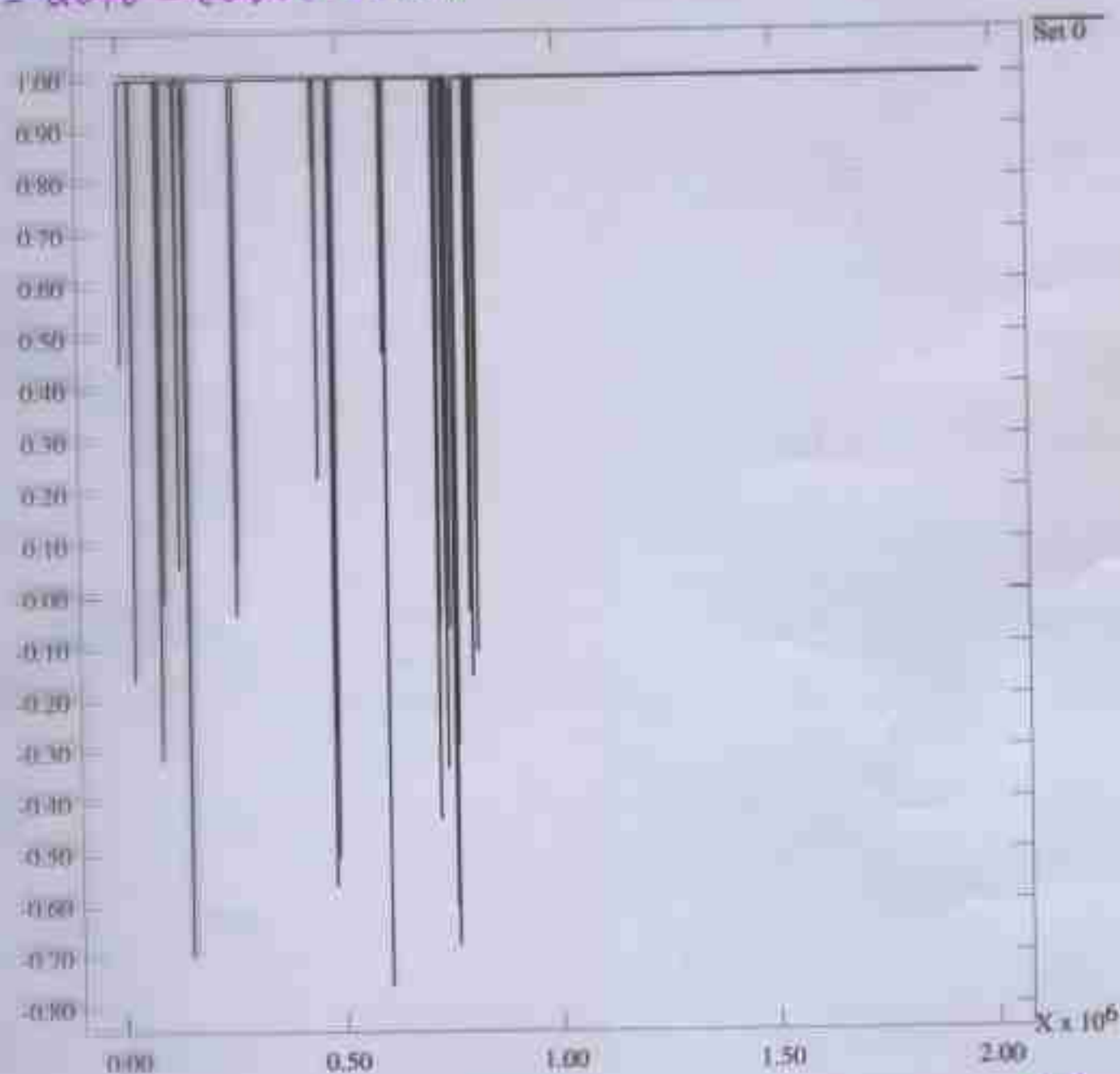


Example: Bosonic  $SU(2)$   $D=4$

#updates

A (strongly) divergent Yang-Mills integral

$\gamma$  = auto-correlation



Example: Bosonic  $SU(2)$   $D=3$  #updates

(convergence of  $SU(N)$  Yang-Mills integrals)

Bosonic

$$D=3 \quad N \geq 4$$

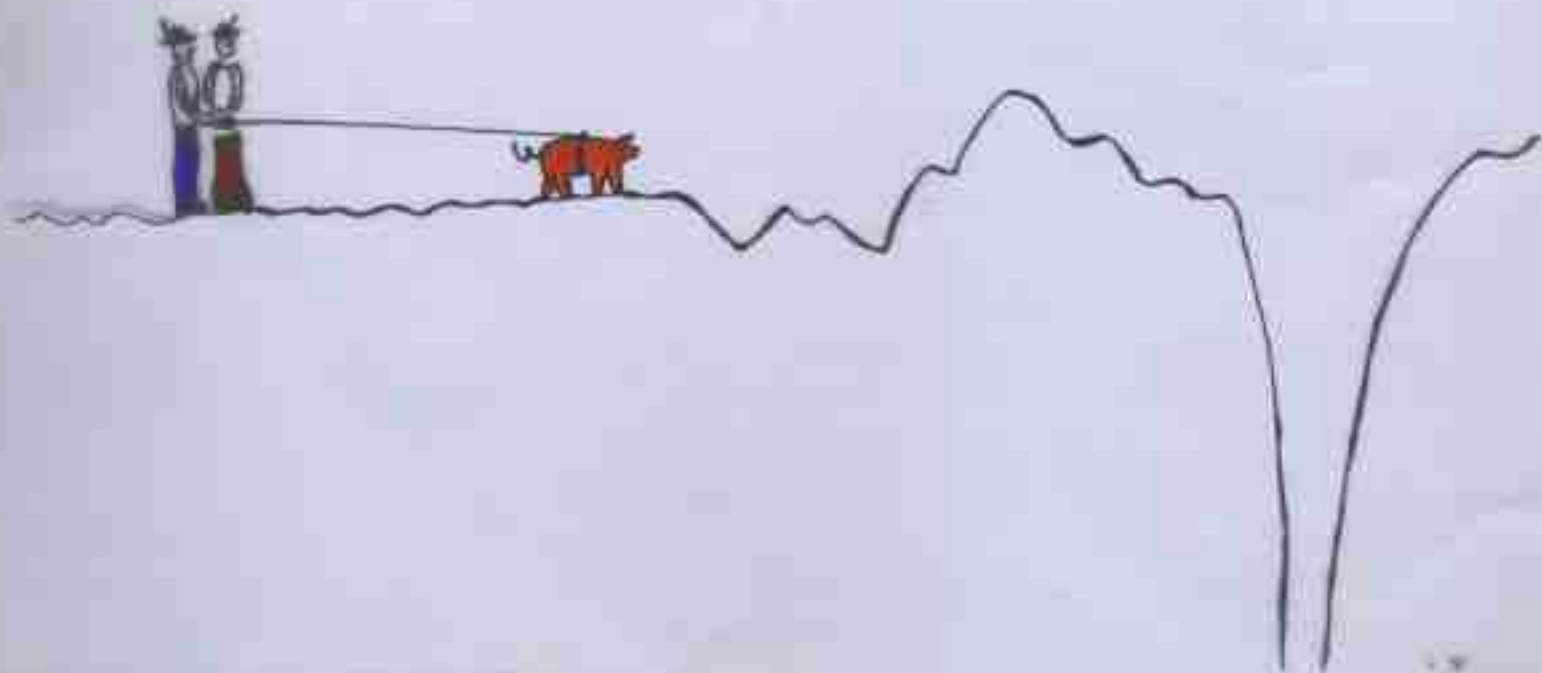
$$D=4 \quad N \geq 3$$

$$D \geq 5 \quad N \geq 2$$

Supersymmetric

$$D=4, 6, 10 \quad N \geq 2$$

All other cases are divergent. (in part: Susy  $D=3$ )





# One loop perturbative estimates

split  $X_\mu = \begin{pmatrix} x_\mu^1 & & 0 \\ & \ddots & \\ 0 & & x_\mu^N \end{pmatrix} + \frac{1}{N}$

$\mathcal{P}(\gamma^2)$  effective ensembles:

SUSY: (Aoki et al. hep-th/9802085)

$$\mathbb{Z}_{D,N} \sim \int \prod_{i,\mu}^{N,D} dx_\mu^i \left[ \prod_{\mu} \delta \left( \sum_i x_\mu^i \right) \right] \times$$
$$\times \sum_{G: \text{maximal tree}} \prod_{(i,j): \text{links of } G} \frac{1}{|x^i - x^j|^{2(D-2)}} + \dots$$



Bosonic: (Hotta et al. hep-th/9811220)

$$\mathbb{Z}_{D,N}^{W=0} \sim \int \prod_{i,\mu}^{N,D} dx_\mu^i \left[ \prod_{\mu} \delta \left( \sum_i x_\mu^i \right) \right] \cdot \prod_{i < j} \frac{1}{|x^i - x^j|^{2(D-2)}}$$

- Power counting reproduces convergence condition
- Far from a proof, breaks down if any two  $x^i, x^j$  get close.

An exact result for all  $N$  (for  $SU(2)$ )

$$Z_{0,N}^{\mathcal{N}} = \frac{2^{\frac{N(N+1)}{2}} \pi^{\frac{N-1}{2}}}{2\sqrt{N} \prod_{i=1}^{N-1} i!} \times \left\{ \begin{array}{lll} - & D=3 & \mathcal{N}=2 \\ \frac{1}{N^2} & D=4 & \mathcal{N}=4 \\ \frac{1}{N^2} & D=6 & \mathcal{N}=8 \\ \sum_{n|N} \frac{1}{n^2} & D=10 & \mathcal{N}=16 \end{array} \right.$$

$D=10$  conjectured by Green + Gutperle

$D=13, 4, 6, 10$  calculated by Moore et al. [hep-th/9803265](https://arxiv.org/abs/hep-th/9803265)

$D=4, 6, 10$  checked to  $\sim N=5$  by Monte Carlo

Bosonic integrals?

Only the  $SU(2)$  result is known:

$$D=5: \quad Z_{0,N=2}^{\mathcal{N}=0} = 2^{-\frac{3}{4}D-1} \frac{\Gamma(\frac{D}{4}) \Gamma(\frac{D-2}{4}) \Gamma(\frac{D-4}{4})}{\Gamma(\frac{D}{2}) \Gamma(\frac{D-1}{2}) \Gamma(\frac{D-2}{2})}$$

$$Z_{0,N>2}^{\mathcal{N}=0} = ?$$

# Coh FT Approach to susy integrals

(Moave, Nekrasov + Shatashvili)

- Beautiful, but incomplete:

Euclidean light cone:  $\phi = X_1 + iX_2$   $\bar{\phi} = X_1 - iX_2$

$\phi$  is a complex matrix, and  $\bar{\phi} = \phi^\dagger$ .

In the Coh FT method,  $\phi, \bar{\phi}$  are independent

hermitean; corresponding to a "Wick rotation":

$$\phi = X_1 - X_0 \quad \bar{\phi} = X_1 + X_0$$

After elimination of all fields except  $\phi$ :

e.g. susy  $D=4$ :  $Z_{D=4, N}^{\omega=4} \sim \frac{1}{(\omega-1)!} \frac{1}{E^{\omega-1}} \oint \prod_{i=1}^{\omega-1} \frac{d\phi_i}{2\pi i} \prod_{\substack{i,j \\ i \neq j}} \frac{\phi_i - \phi_j}{\phi_i - \phi_j + E \tau i \epsilon}$

- It would be nice to derive contour prescriptions from Wick rotation.
- $D=3$  ??? Euclidean Yang-Mills integral is ill-defined!
- Correlation functions:

$\langle T_\nu \phi^K \rangle = 0$  in original undeformed integral

YM integrals  
(correlators etc.)

$\cong$

Coh FT

"deformed" correlators etc.

?



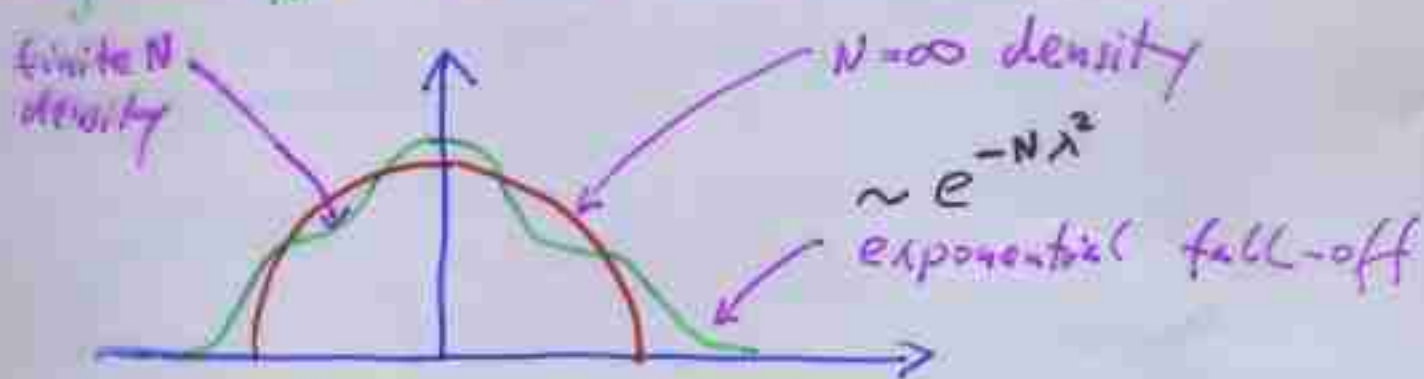
# One-matrix correlators and the density of eigenvalues

- simplest correlators:  $\langle \frac{1}{N} \text{Tr} X_D^{2K} \rangle = \text{"moments"}$
- direct relation to eigenvalue density:

$$X_D = \begin{pmatrix} \lambda_1 & & \\ & \ddots & \\ & & \lambda_N \end{pmatrix} \rightarrow \rho(\lambda) \equiv \langle \frac{1}{N} \sum_{i=1}^N \delta(\lambda - \lambda_i) \rangle$$

$$\langle \frac{1}{N} \text{Tr} X_D^{2K} \rangle = \int_{-\infty}^{\infty} d\lambda \rho(\lambda) \lambda^{2K}$$

- "Wigner-type" matrix models:



- compute moments of YM integrals for  $SU(2)$ :  
surprise! E.g.  $D=4$  susy: all moments are  $\infty$ !
- Does density exist? Yes, but

$$\rho_{D=4, N=2}(\lambda) \sim \frac{1}{\lambda^3} \text{ as } \lambda \rightarrow \infty$$

# Exact results for $SU(2)$

Explicit  $SU(2)$  eigenvalue densities and expectation values (Plefka, unpublished)

$$\rho_{D=4}^{SUSY}(\lambda) = \frac{3 \cdot 2^{5/4}}{\sqrt{\pi}} \lambda^2 U\left(\frac{5}{4}, \frac{1}{2}, 8\lambda^4\right)$$

$$\rho_{D=6}^{SUSY}(\lambda) = \frac{105}{2^{3/4} \sqrt{\pi}} \lambda^2 \left[ U\left(\frac{9}{4}, \frac{1}{2}, 8\lambda^4\right) - \frac{33}{16} U\left(\frac{13}{4}, \frac{1}{2}, 8\lambda^4\right) \right]$$

$$\rho_{D=10}^{SUSY}(\lambda) = \frac{1287}{64 \cdot 2^{3/4} \sqrt{\pi}} \lambda^2 \left[ 546 U\left(\frac{17}{4}, \frac{1}{2}, 8\lambda^4\right) - 147 \frac{17 \cdot 19}{8} U\left(\frac{21}{4}, \frac{1}{2}, 8\lambda^4\right) \right. \\ \left. + 45 \frac{17 \cdot 19 \cdot 21 \cdot 23}{256} U\left(\frac{25}{4}, \frac{1}{2}, 8\lambda^4\right) - \frac{17 \cdot 19 \cdot 21 \cdot 23 \cdot 25 \cdot 27}{2048} U\left(\frac{29}{4}, \frac{1}{2}, 8\lambda^4\right) \right]$$

$U$  is the Kummer- $U$  function:

$$U\left(\nu + \frac{1}{2}, \frac{1}{2}, pz\right) = \frac{\sqrt{p}}{\Gamma(\nu + 1)} \int_0^\infty dx \frac{x^\nu}{(x+z)^{\nu+1/2}} e^{-px}$$

• One has  $\rho_{D=4} \sim \lambda^{-3}$ ,  $\rho_{D=6} \sim \lambda^{-7}$ ,  $\rho_{D=10} \sim \lambda^{-15}$

The susy one-matrix correlators for  $D = 6$

$$\langle \text{Tr} X_D^2 \rangle_{D=6} = \frac{1}{2} \sqrt{\frac{2}{\pi}} \quad \langle \text{Tr} X_D^4 \rangle_{D=6} = \frac{25}{64}$$

as well as the susy  $D = 10$  ones

$$\begin{aligned} \langle \text{Tr} X_D^2 \rangle_{D=10} &= \frac{8}{25} \sqrt{\frac{2}{\pi}} & \langle \text{Tr} X_D^4 \rangle_{D=10} &= \frac{9}{80} \\ \langle \text{Tr} X_D^6 \rangle_{D=10} &= \frac{3}{32} \sqrt{\frac{2}{\pi}} & \langle \text{Tr} X_D^8 \rangle_{D=10} &= \frac{297}{4096} \\ \langle \text{Tr} X_D^{10} \rangle_{D=10} &= \frac{1089}{8192} \sqrt{\frac{2}{\pi}} & \langle \text{Tr} X_D^{12} \rangle_{D=10} &= \frac{184041}{655360} \end{aligned}$$

all other  $SU(2)$  susy one-matrix correlators do not exist.

• Bosonic  $SU(2)$  case:

$$\rho_{D, N=2}^{W=0}(\lambda) \sim \frac{1}{\lambda^{D-3}}$$



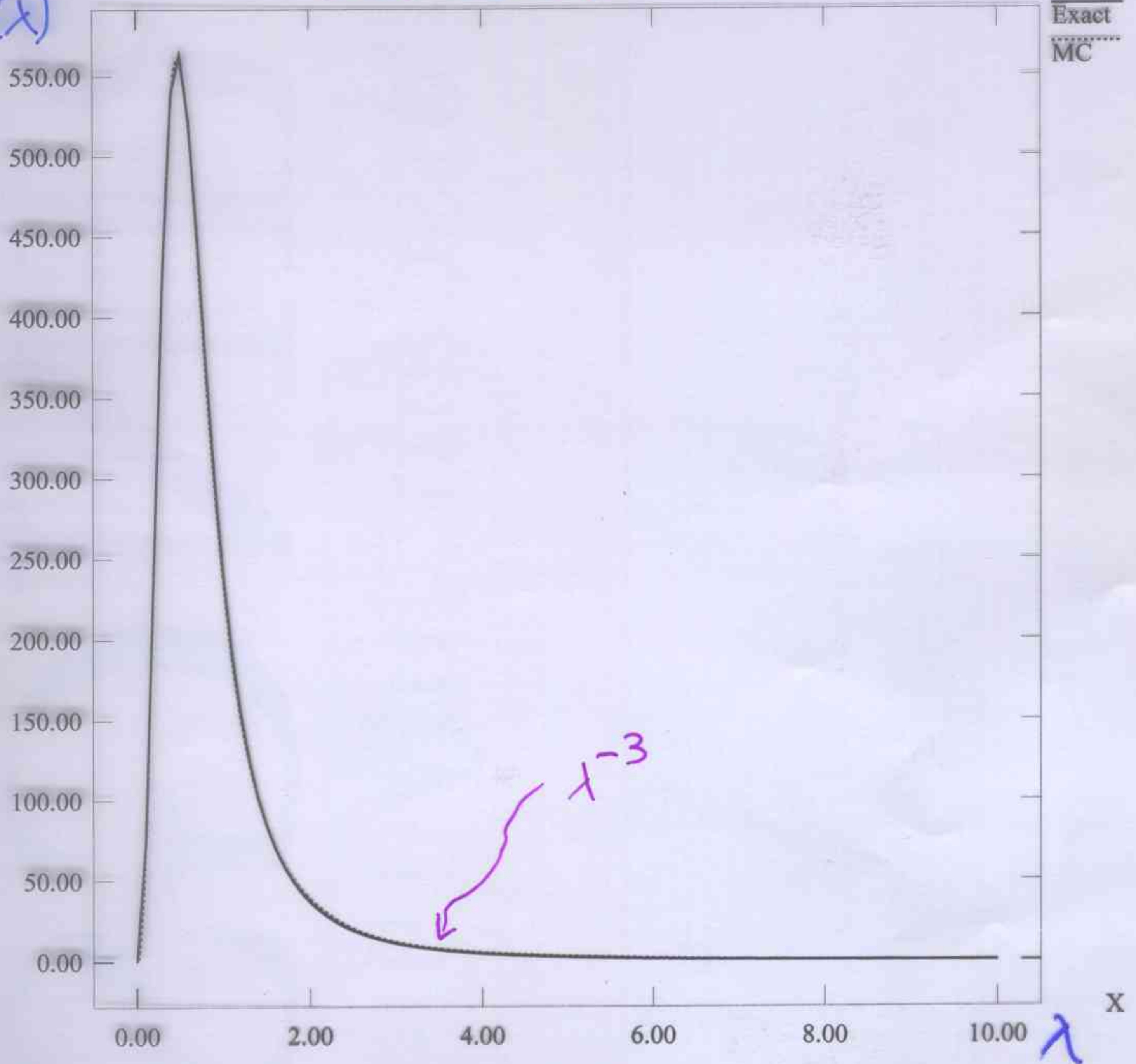
# Accuracy of Monte Carlo computation

SU(2) D=4 Susy

Susy Distribution n\_2d\_4

(half-density)

$Y \times 10^{-3}$   
 $P(x)$



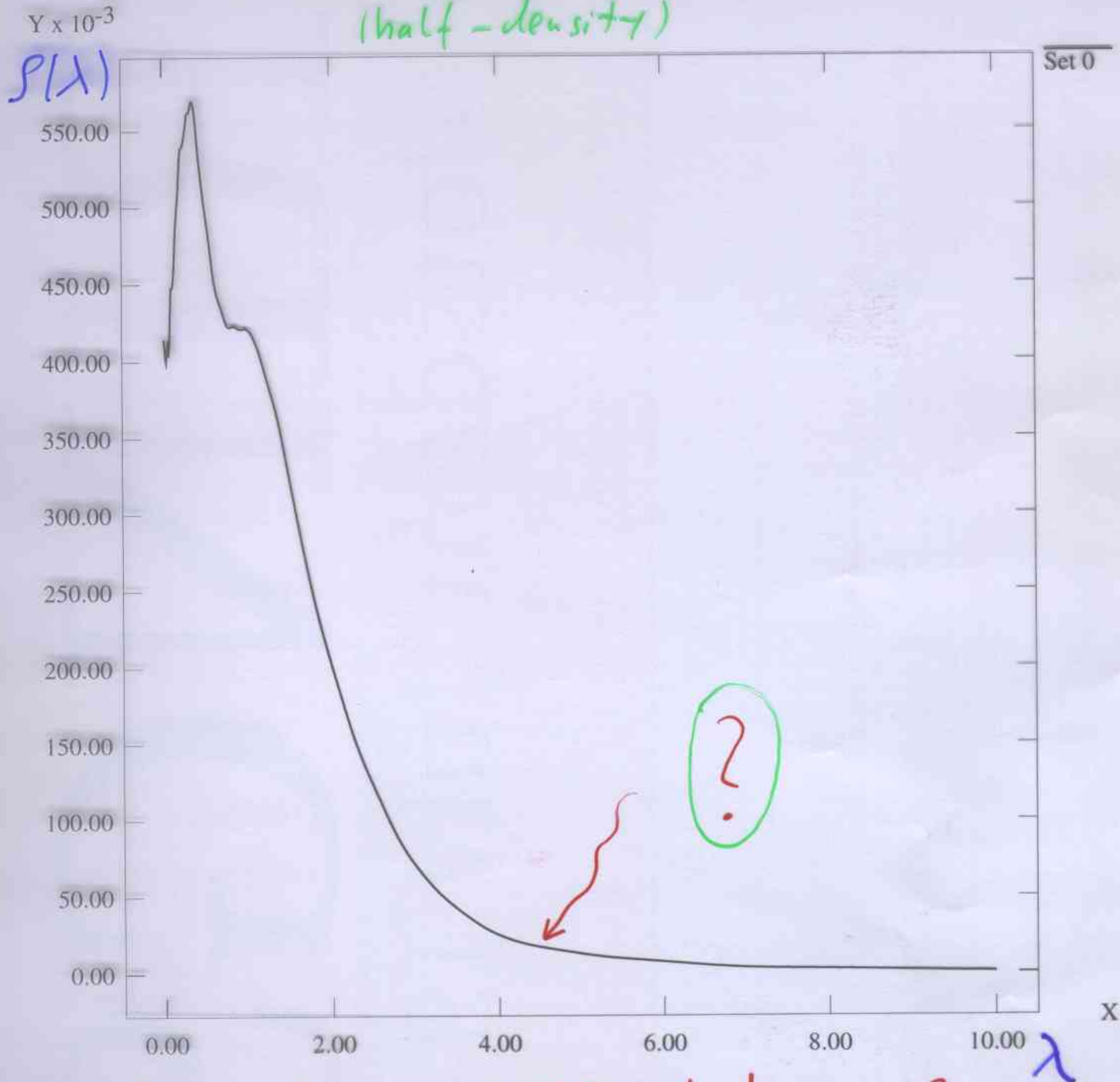
"Kummer - U a' la Monte Carlo"

$N > 2$  computations

$SU(4)$   $D=4$  susy

Susy Distribution n\_4d\_4

(half-density)



What is the asymptotic behavior?



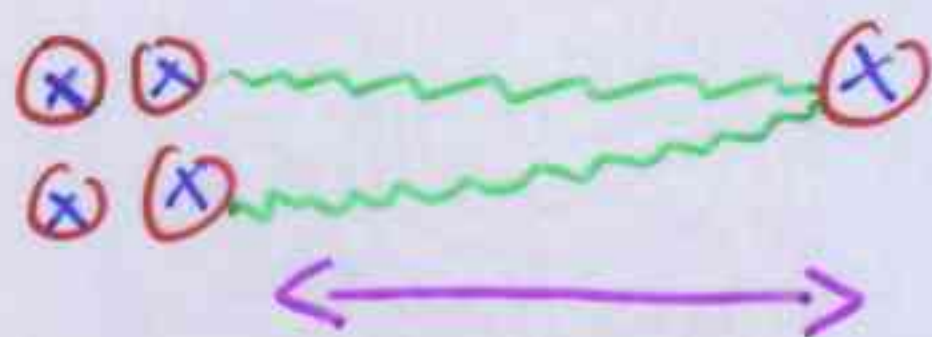
Intuition: One-loop estimates

We had:

$$Z_{D,N}^{\text{susy}} \sim \int \prod_i^{N-1} \prod_\mu^D dx_\mu^i \cdot \sum_{\text{trees}} \sum_{\substack{\text{links} \\ (i,j)}} \frac{1}{(x^i - x^j)^{3(D-2)}}$$

$$Z_{D,N}^{\text{bosonic}} \sim \int \prod_i^{N-1} \prod_\mu^D dx_\mu^i \cdot \prod_{i < j} \frac{1}{(x^i - x^j)^{2(D-2)}}$$

Most dangerous "infrared" configuration:



This suggests:

$$\left\langle \frac{1}{N} \text{Tr} X_D^{2k} \right\rangle < \infty$$

if and only if

$$k < D - 3 \quad \text{susy}$$

$$k < N(D-2) - \frac{3}{2}D + 2 \quad \text{bosonic}$$

In bosonic case, in marked difference to susy, all moments exist as  $N \rightarrow \infty$ !

Is this true non-perturbatively?

$\Rightarrow$  Monte Carlo random walk

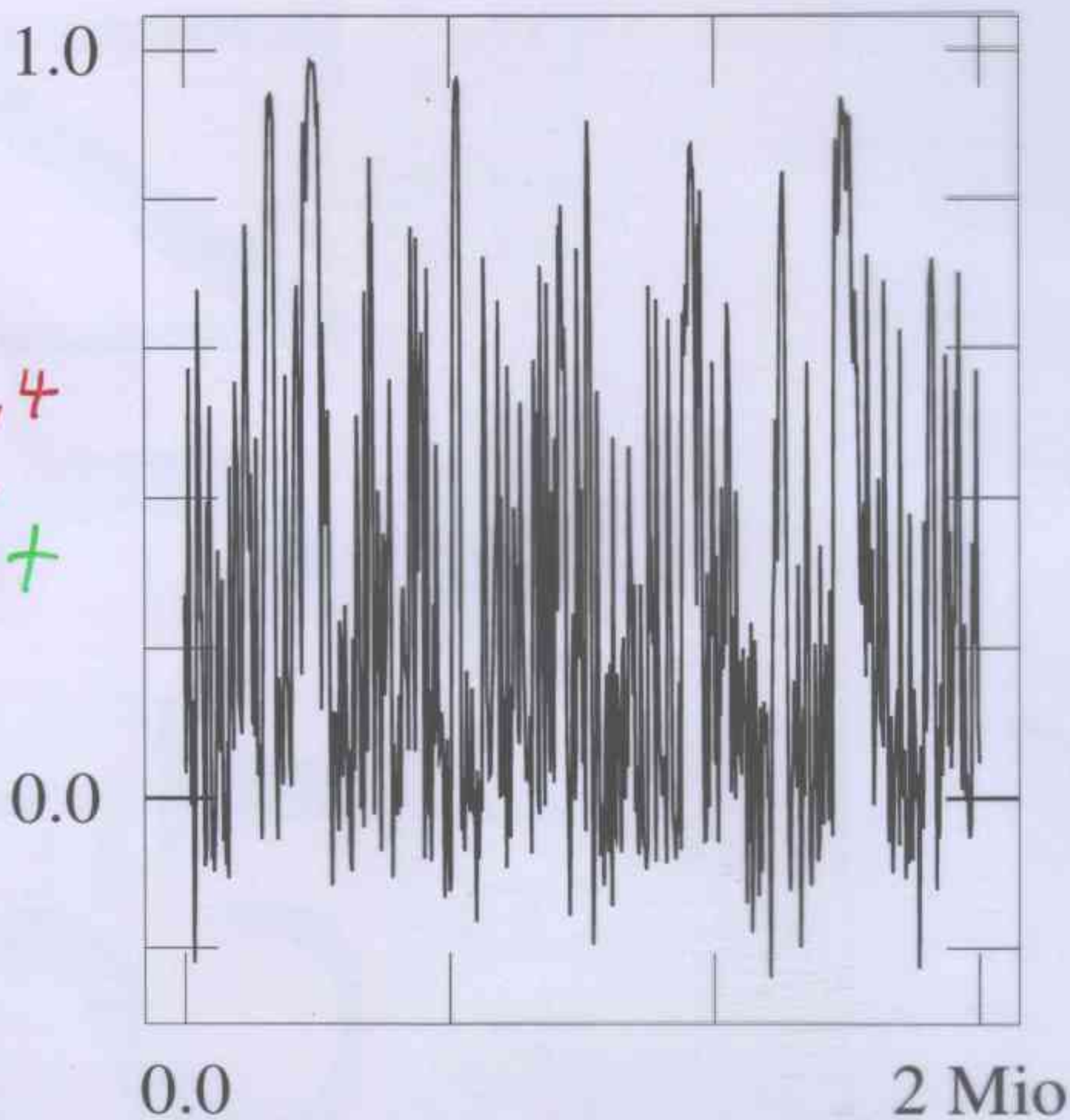


Auto correlation function:

Susy  $D=6$   $SU(4)$

$\text{Tr } X^4$   
convergent

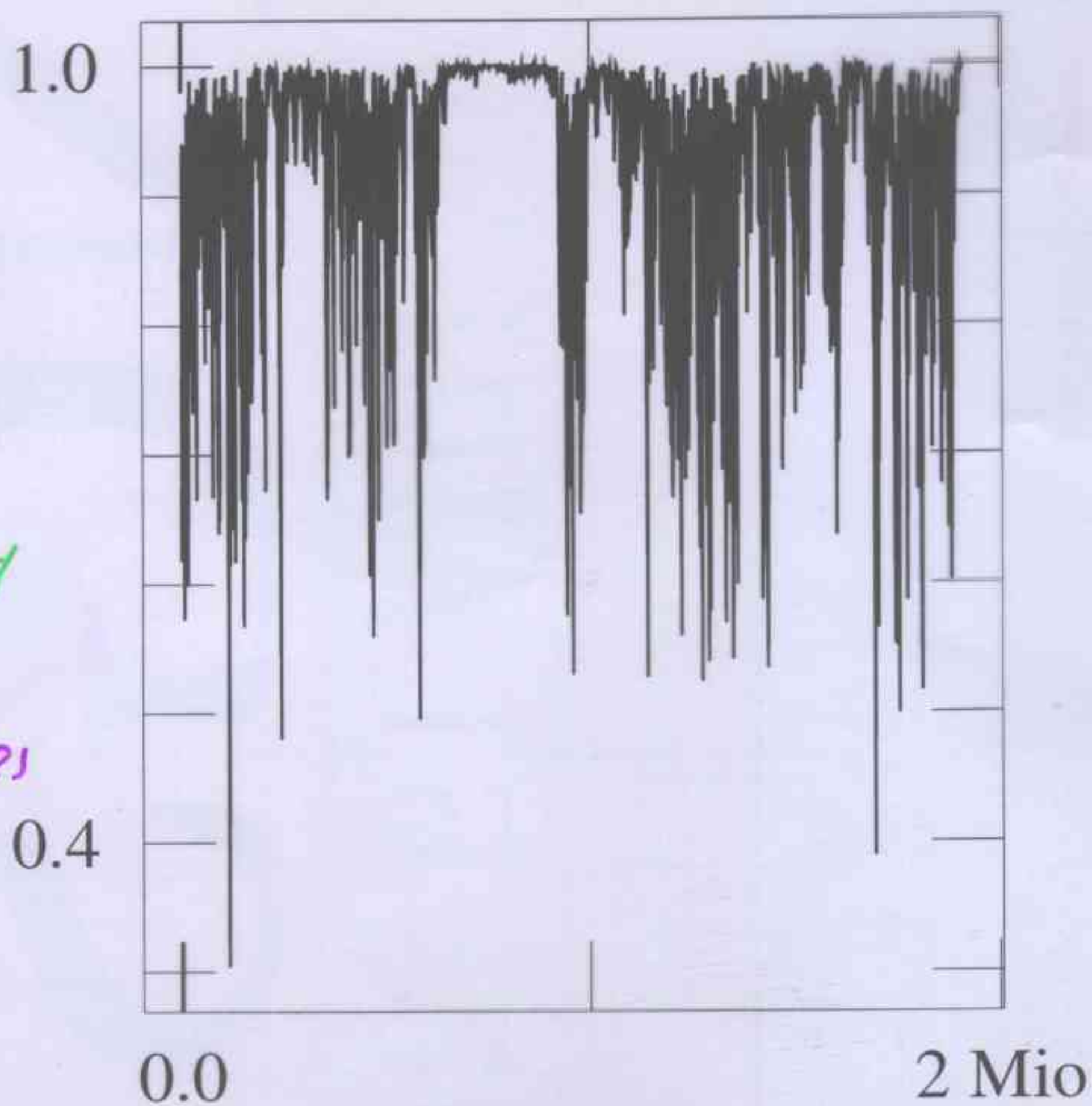
Moment  
exists!



Autocorrelation function:

Susy  $D=6$   $SU(4)$

$\text{Tr } X^6$   
marginally  
divergent  
moment does  
not exist!

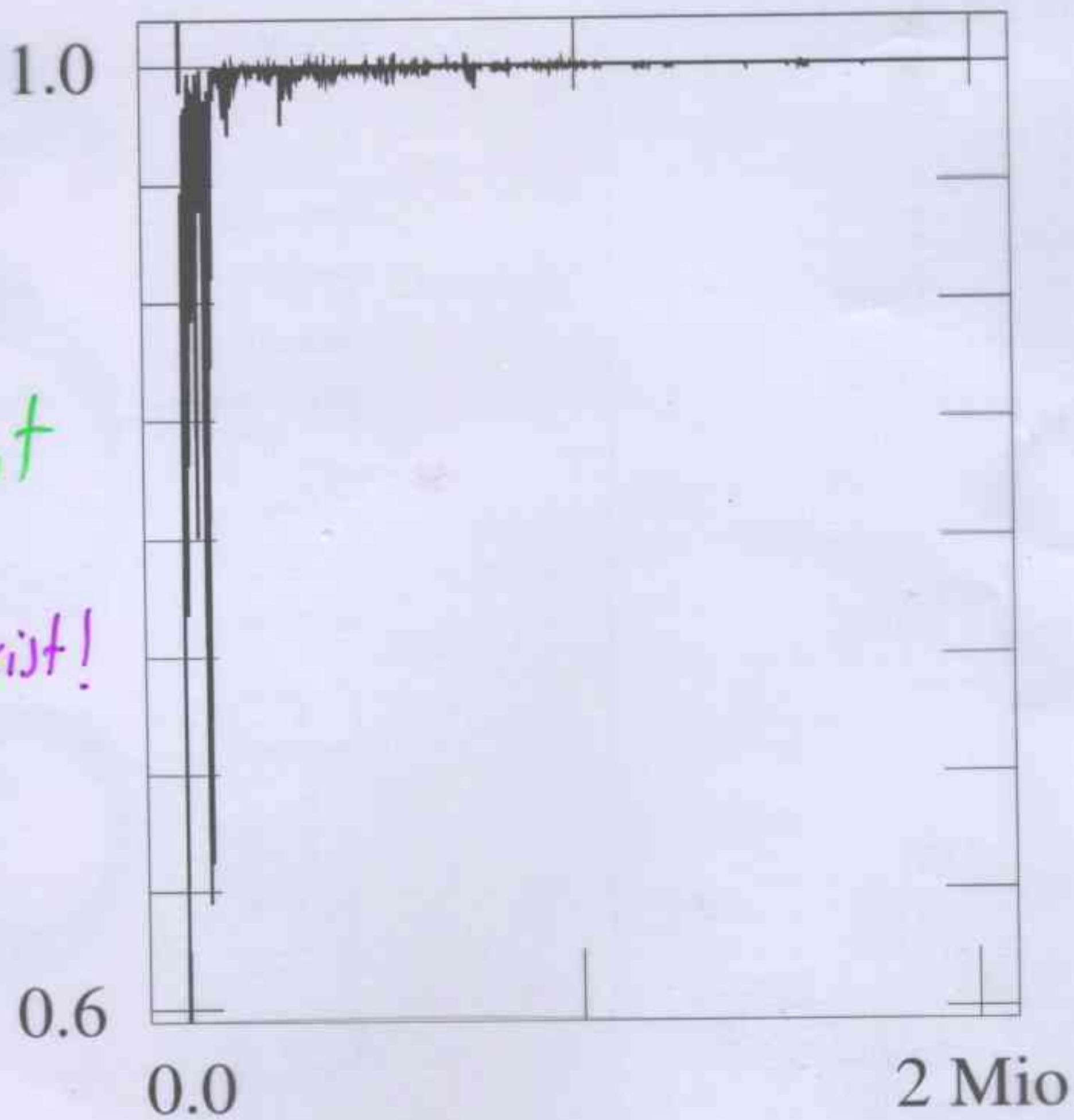




Autocorrelation function:

Susy  $D=6$   $SU(4)$

$\overline{T_r X^8}$   
divergent  
Does not exist!



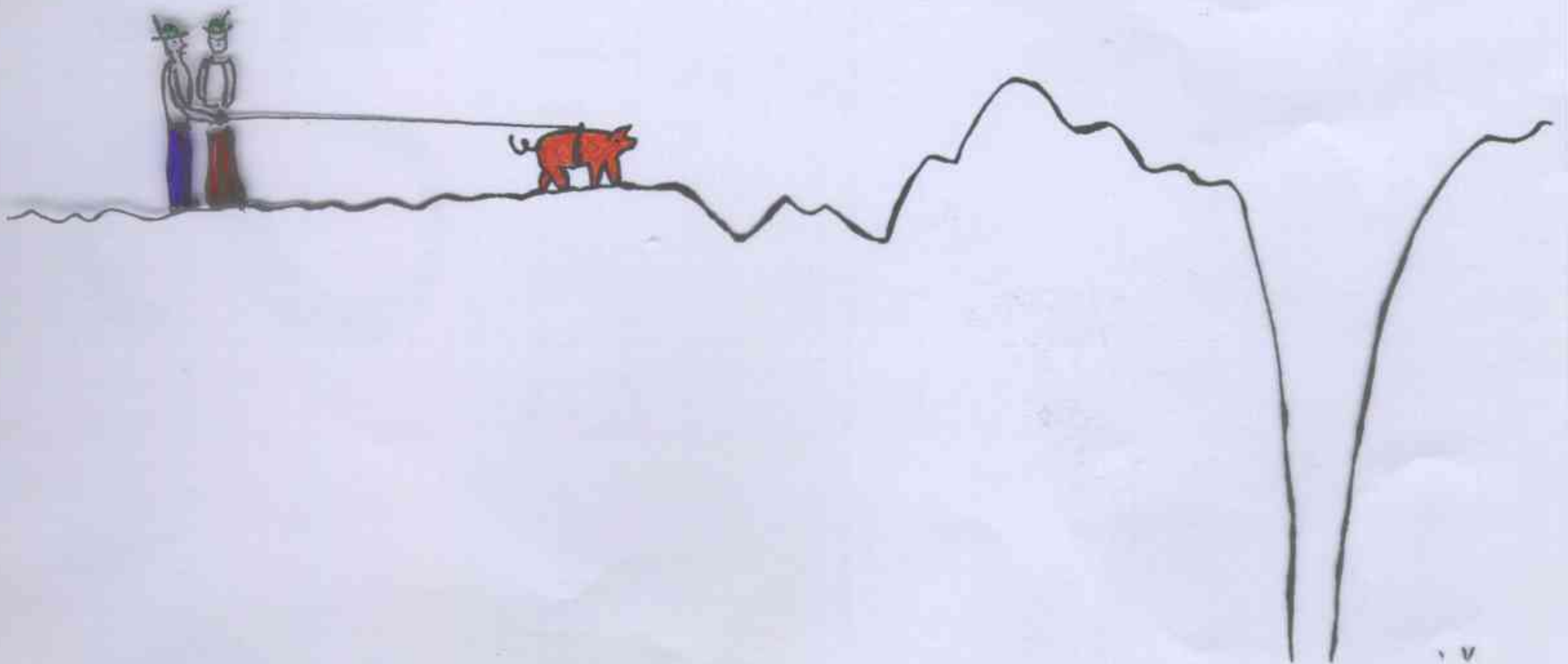
# Asymptotic eigenvalue density

Susy:

$D=4$	$\rho(\lambda) \sim \lambda^{-3}$	all $N$	No clear defined edge!
$D=6$	$\rho(\lambda) \sim \lambda^{-7}$	all $N$	
$D=10$	$\rho(\lambda) \sim \lambda^{-15}$	all $N$	

Bosonic:

$$D \geq 3 \quad \rho(\lambda) \sim \lambda^{-2N(D-2) + 3D - 5}$$





# conclusions

- Established **non-perturbative** convergence conditions for Yang-Mills integrals (w/without susy).
- Ditto for simplest correlators  
⇒ **Non-perturbative eigenvalue densities.**
- Demonstrated, that numerically **precise** results may be computed for "not too large"  $N$ .

**Examples:** partition functions,  
correlation functions,  
Wilson loops,  
spectral distributions ...



## Open Problems I: analytic

Yang-Mills integrals are

- a challenge to our "analytic toolbox",

and

- an ideal "laboratory".

- Rigorously prove convergence properties

- calculate analytically partition functions and "observables":

$$\langle T_\nu A_\mu^k \rangle, \langle T_\nu F^4 \rangle, \langle T_\nu P \exp i \oint A \rangle, \dots$$

## Open Problems II: numerical

- Learn to go to higher  $N$ , especially in the  $SUSY$  case, especially in the  $D=10$  case.
- Can we say more about the full effective action of diagonal elements?
  - "Eguchi-Kawai" mechanism
  - "self-compactification"



## Exciting Prospect:

- Can we make further **precise** "predictions" for **Yang-Mills integrals** from **"dual"** formulations (string theory, **Sugra**)?  
(especially at finite  $N$ !)

$\Rightarrow$

As opposed to, say, Matrix QM,  
here we have a system that  
is amenable to  
**non-perturbative analysis!**