

Calibrating Branes, or PhreMology

- Calibrations & the su₃ algebra
- Generalized Calibrations
- Kahler \rightarrow ω -associative - Cayley
- WZ realizations: Domain walls & their intersections
- WZ 'wallpaper'

- J. Gutowski, G. Papadopoulos & P.K.T.
"Supersymmetry & generalized calibrations"
hep-th/9905156

- G. Gibbons and P.K.T.
"A Bogomol'nyi equation for intersecting domain walls"
hep-th/9905196

Also

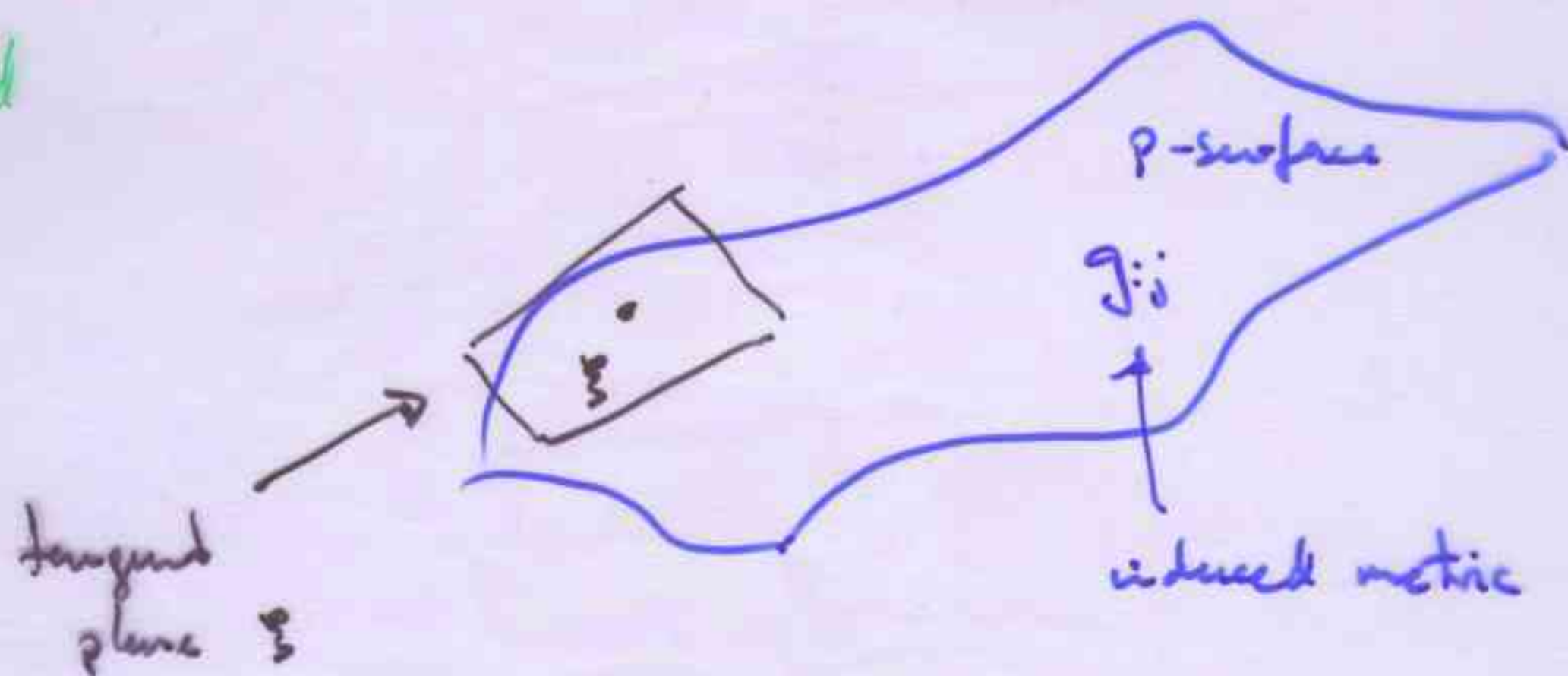
P. Saftin

"Tiling with almost-BPS junctions"
hep-th/9907066

P. Townsend
Strings 99

Calibrations & Minimal Surfaces (Harvey & Lawson)

n -dim Riemannian
manifold



Tangent plane characterized by p coforms
where $\left\{ \frac{\partial}{\partial \sigma^i} ; i=1, \dots, p \right\}$ span plane

$$\xi = \frac{\partial}{\partial \sigma^1} \wedge \dots \wedge \frac{\partial}{\partial \sigma^p}$$

A calibration is a p -form ϕ such that

(i) $\xi \cdot \phi \leq \text{vol}_\xi$ \leftarrow vol. element on ξ

(ii) $d\phi = 0$

Defn: 'Contact set' $G = \{ \xi ; \xi \cdot \phi = \text{vol}_\xi \}$

Theorem: (Harvey & Lawson) $G =$ set of tangent planes
of minimal p -surface

Derivation from super p-brane superalgebra

(Lytowski,
Papadopoulos
& PKT)

In simple cases, e.g. p-brane = $\mathbb{E}^{(n,1)}$

$$\{Q, Q\} = \int \left[\text{vol}(g) \pm \underbrace{\rho_0 \rho_{i_1 \dots i_p}}_{\substack{\text{closed p-form} \\ \text{'central' charge}}} dX^{i_1} \dots dX^{i_p} \right]$$

\uparrow
p-brane
vol form
= induced
worldspace metric

Let e be constant spinor with $e^\dagger e = 1$

$$(e^\dagger Q)^2 = \int \left[\text{vol}(g) \pm \phi \right] \quad \phi = \bar{e} \rho_{i_1 \dots i_p} e dX^{i_1} \dots dX^{i_p}$$

p-brane

Valid for any open set on worldspace so positivity of

$(e^\dagger Q)^2$ implies

(i) $\int \phi \leq \text{vol}_g$

Topological nature of 'central' charge implies

(ii) $d\phi = 0$

$\Rightarrow \phi$ is a
calibration

Generalized calibrations

(Gutowski & Papadopoulos)
G, P & PKT

The 'simple' cases are those for which

$$\mathcal{L} = \text{vol}(g)$$

↗
p-brane energy density

More generally one has

$$\mathcal{L} = \nu \text{vol}(g) + i_k A$$

↗
'redshift'
factor

↖ (p+1)-form potential
↖ contraction with timelike Killing
vector field k.

In this case one finds, as before, that

$$(i) \quad \mathcal{L} \leq \text{vol}_g$$

but now

$$(ii) \quad d\mathcal{L} = d i_k A \neq 0$$

— defines a 'generalized'
calibration

Theorem: The contact set of a generalized calibration is the set of p-planes tangent to an \mathcal{L} -minimizing p-surface

Supersymmetry & calibrations

Becker², Morrison, Oguchi, Oz, Yin
Ciblaris & Papadoulas
Cantlett, Lambert & West
Acharya, Figueroa-O'Farrill, Spence

Calibrations deduced from susy algebra are

$$\phi = \text{vol}(g) \epsilon^\dagger \Gamma \epsilon$$

for some constant spinor ϵ , where

$$\Gamma = \frac{1}{p! |\text{vol}(g)|} \epsilon^{i_1 \dots i_p} \partial_{i_1} X^{I_1} \dots \partial_{i_p} X^{I_p} \Gamma_{I_1 \dots I_p}$$

$$\left. \begin{array}{l} \Gamma^2 = 1 \\ \epsilon^\dagger \epsilon = 1 \end{array} \right\} \Rightarrow \int \phi \leq \text{vol}_g$$

saturated iff

$$\Gamma_g \epsilon = \epsilon$$

evaluation of Γ on tangent
 p -plane ξ

$$\left. \begin{array}{l} \Gamma_g^2 = 1 \\ \text{tr} \Gamma_g = 0 \end{array} \right\} \Rightarrow +1 \text{ eigenspace of } \Gamma_g \text{ is } \underline{\text{half-maximal}}$$

i.e. locally, calibrated surface preserves $\frac{1}{2}$ susy

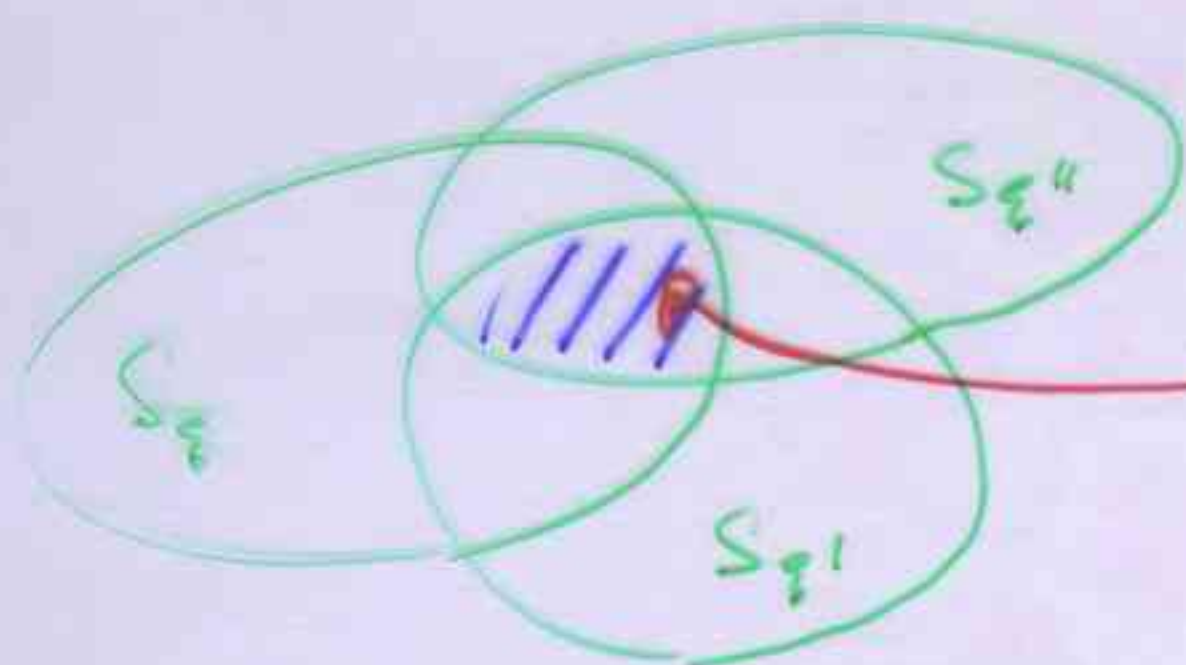
Classification

As ξ varies in contact set G

$$\hat{\Pi}_\xi \rightarrow \hat{\Pi}_{\xi'} = R^{-1} \hat{\Pi}_\xi R$$

rotation in some subgroup of $SO(n)$

Let S_ξ be +1 eigenspace of $\hat{\Pi}_\xi$



solution space of $\hat{\Pi}_\xi \epsilon = \epsilon$
for all $\xi \in G$

Thus, in general we preserve $\leq \frac{1}{2}$ fraction of sym

$$G = \frac{SU(m)}{S[U(p) \times U(m-p)]}$$

(Kähler) ($n=2m$)

$$v = \frac{1}{2^m}$$

$$G = \frac{SU(m)}{SO(m)}$$

(SLAG)

$$v = \frac{1}{2^m}$$

$$G = \frac{G_2}{SO(4)}$$

(Associative
& Co-associative)

$$v = \frac{1}{16}$$

$$G = \frac{Sp \ddot{7}}{[SO(2) \times SU(2) \times SU(2)] / \mathbb{Z}_2}$$

(Cayley)

$$v = \frac{1}{32}$$

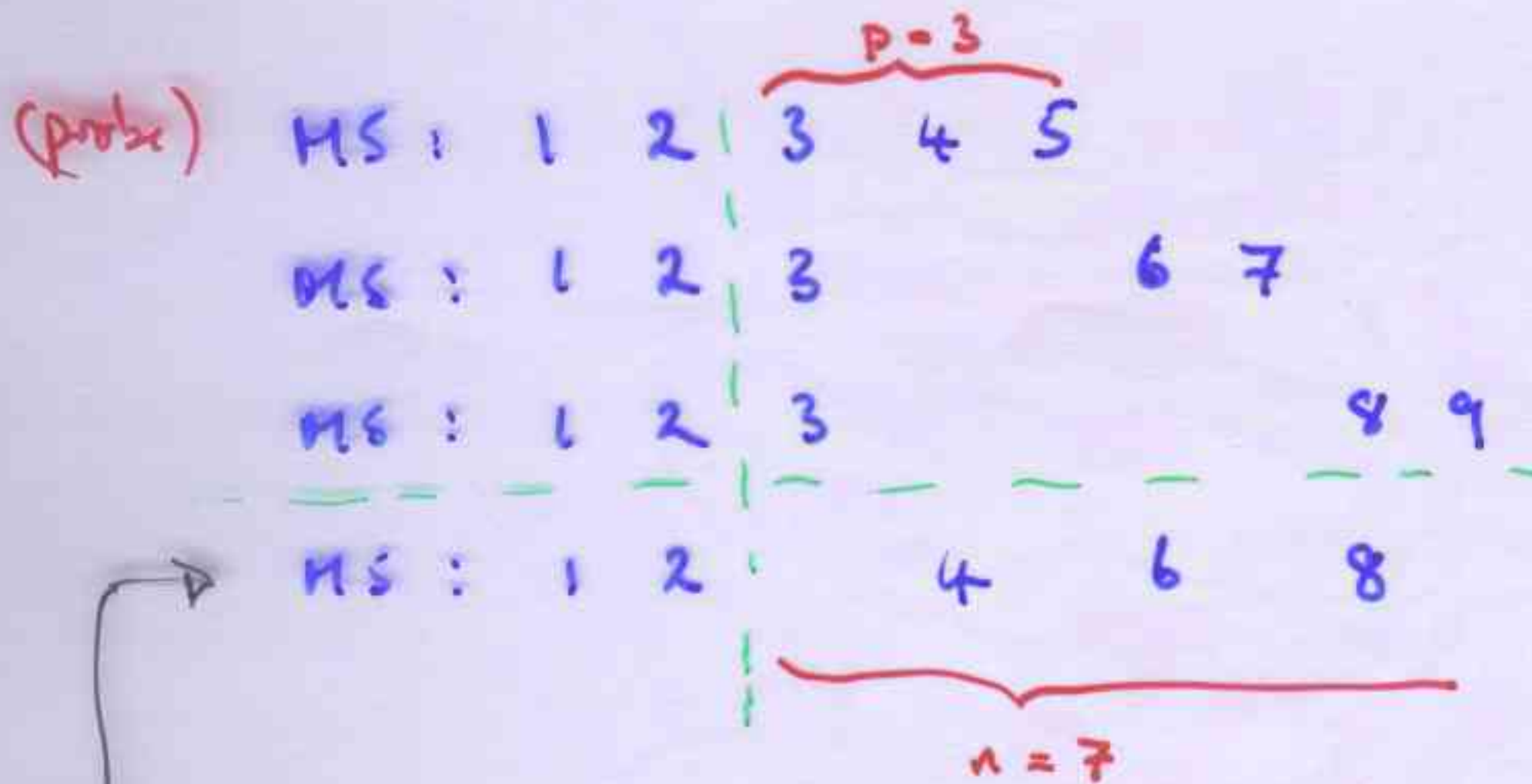
Examples

(taken from Casati, Lambert, West)

(i) Kähler $n=6$ $p=2$ $v=1/8$



(ii) Co-associative ($v=1/16$)



interpret as domain wall in previous $N=1$ $D=4$ theory

eg. in M2CD (Witten, Vasiliev, ...)

N=1 D=4 susy analysis

(Abelian, PGT)
(non-Abelian, PGT)

$$\{\mathcal{Q}, \mathcal{Q}\} = H + \rho^i P_i + \rho^{ij} U_{ij} + \rho^{ij} \gamma_5 V_{ij}$$

↑
4 component spinor
charge

↑ ↑
'electric' x 'magnetic'
2-form charges

Domain walls we associated with constraint

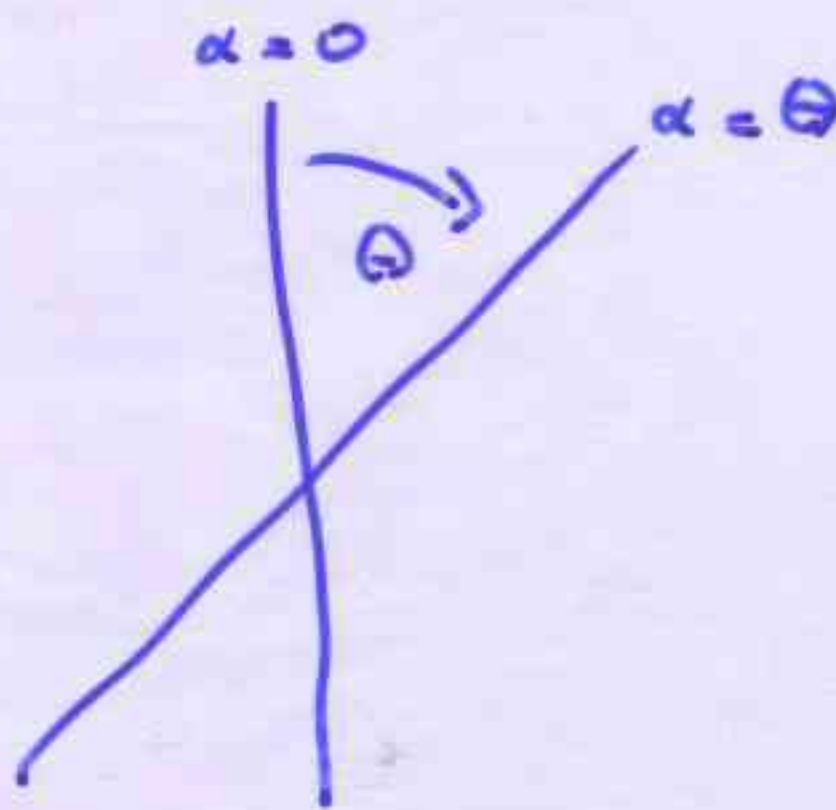
$$\rho^{ij} e^{\alpha \gamma_5} \epsilon = \epsilon \Rightarrow \frac{1}{2} \text{ susy } (\nu = \frac{1}{16})$$

↑
same angle

Take 2 parallel domain walls. Get $\frac{1}{4}$ susy by simultaneous rotation in space and charge space (α)



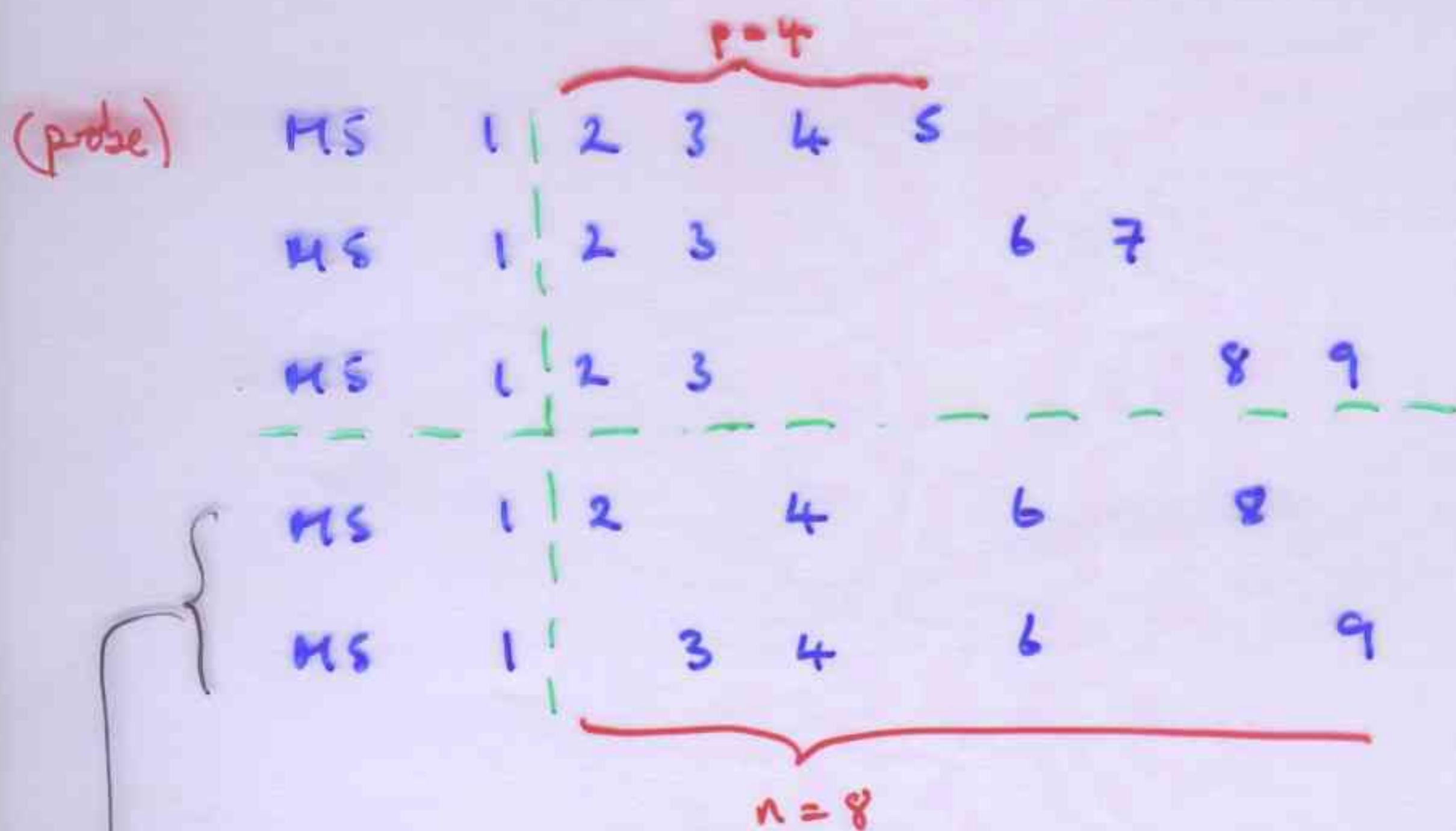
$\frac{1}{2}$ susy
($\nu = \frac{1}{16}$)



$\frac{1}{4}$ susy
($\nu = \frac{1}{32}$)

Relation to calibrations (Giant Hett)

Consider Cayley calibration via intersecting MS-branes



Interpret as domain walls intersecting on 1-axis,

preserving $v = \frac{1}{32}$ susy $\Rightarrow \frac{1}{4}$ of $N=1 D=4$

Domain walls in $N=1 D=4$ can be realized in

WZ model with superpotential $W(\phi)$

(Abraham, PKT)
(Cvetic, Garavito, Rey)
(Dvali & Shifman)

$$\mathcal{H} = \frac{1}{4} \nabla_{\mu} \phi \cdot \nabla^{\mu} \bar{\phi} + |W'(\phi)|^2$$

What about intersecting domain walls?

Bogomolnyi eq for WZ domain wall junctions

(Cahlon, P, KT)

(Carroll, Hellerman)
& Trnka

Reduce to 2D $\nabla = (\partial_x, \partial_y)$

Let $\underline{z} = x + iy$. Then

$$\mathcal{H} = \left| \frac{\partial \phi}{\partial \underline{z}} \mp \bar{w}' \right|^2 \pm 2 \operatorname{Re} \left(\frac{\partial W}{\partial \underline{z}} \right) + \frac{1}{2} \mathcal{J}(\underline{z}, \bar{\underline{z}})$$

$$\mathcal{J} = \frac{\partial \phi}{\partial \underline{z}} \frac{\partial \bar{\phi}}{\partial \bar{\underline{z}}} - \frac{\partial \bar{\phi}}{\partial \bar{\underline{z}}} \frac{\partial \phi}{\partial \underline{z}}$$

Define

$$\mathcal{T} = i \int d\underline{z} d\bar{\underline{z}} \frac{\partial W}{\partial \underline{z}}$$

↑
complex domain wall charge

$$\mathcal{Q} = \frac{i}{4} \int d\bar{\phi} \wedge d\phi$$

↑
real domain wall junction charge

Then

$$\mathcal{H} \geq |\mathcal{T}| + |\mathcal{Q}|$$

with equality when

$$\boxed{\frac{\partial \phi}{\partial \underline{z}} = \bar{w}'}$$

→ have to solve this to find intersecting WZ domain walls preserving $\frac{1}{4}$ susy

WZ junctions

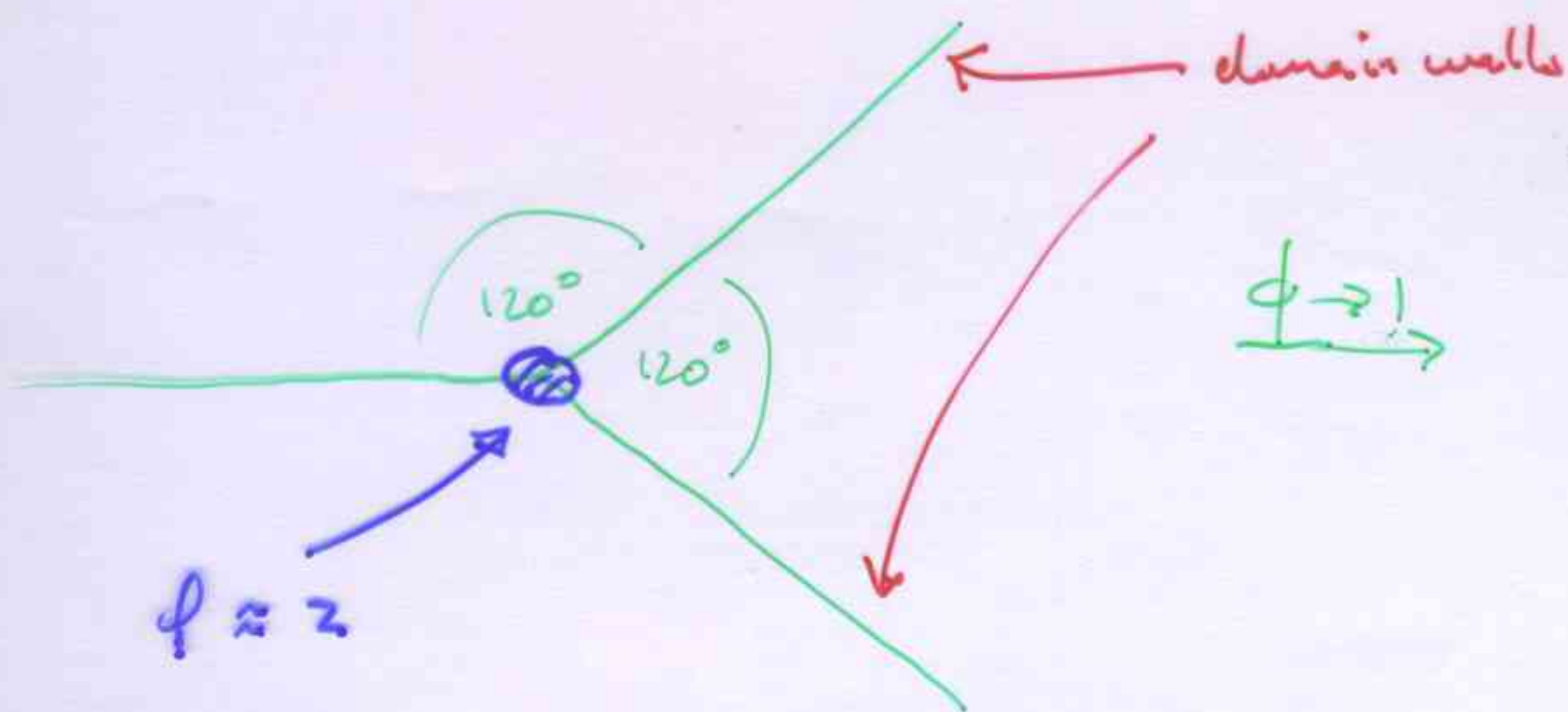
Take $W = \phi - \frac{1}{\bar{\phi}} \phi^4 \Rightarrow W' = 1 - \phi^3$

Then Bog. eq. is

$$\frac{\partial \phi}{\partial z} = 1 - \bar{\phi}^3$$

This has \mathbb{Z}_3 -symmetry: $(z, \phi) \rightarrow (\omega z, \omega \phi)$ $\omega^3 = 1$

so seeks solution of form



- Solution exists in thin-wall limit

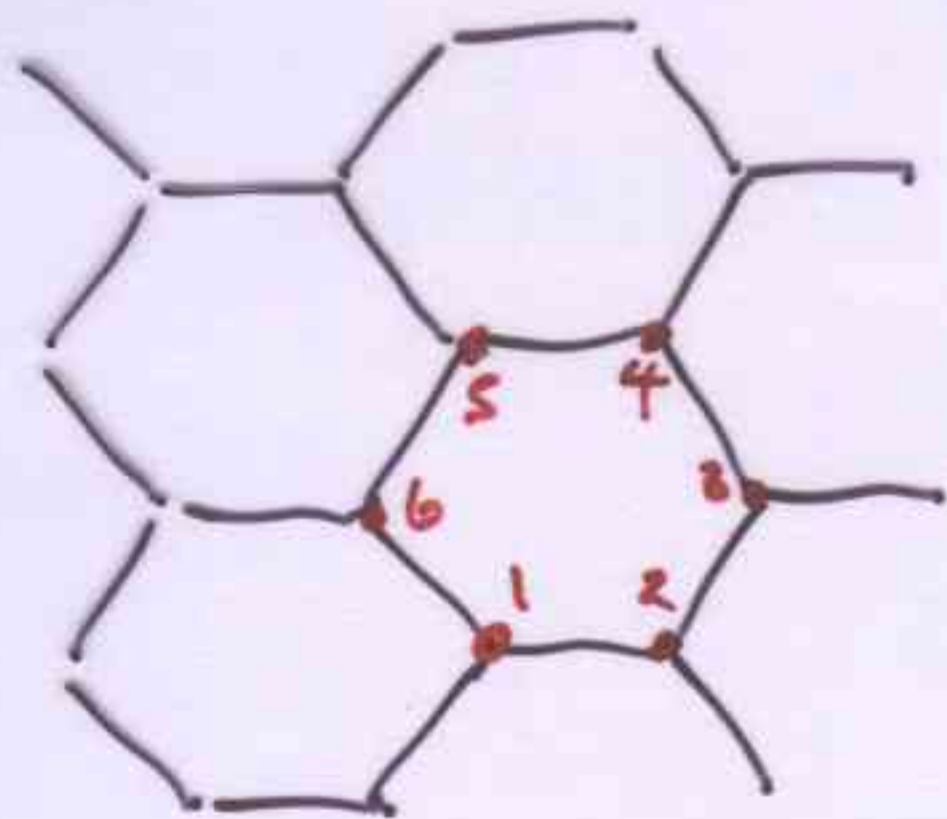
(Stromberg & Zeimer)
(Brouard & Reitich)

- Numerical results confirm existence in general (Saffar)

WZ wallpaper

(Saffin)

WZ junctions can be combined to get 'almost-BPS' networks.

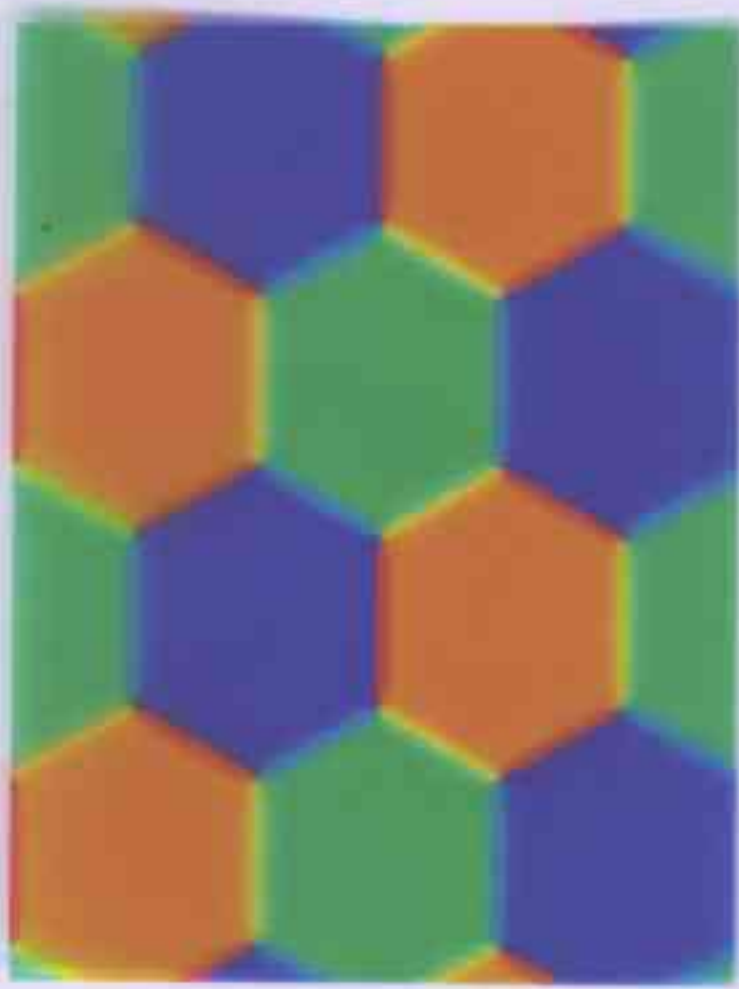


6 vertices are $\left\{ \begin{array}{l} \mathbb{Z}_3 \text{ pairs of junctions} \\ \mathbb{Z}_2 \text{ pairs of } \underline{\text{anti-junctions}} \end{array} \right.$
↑
so not BPS

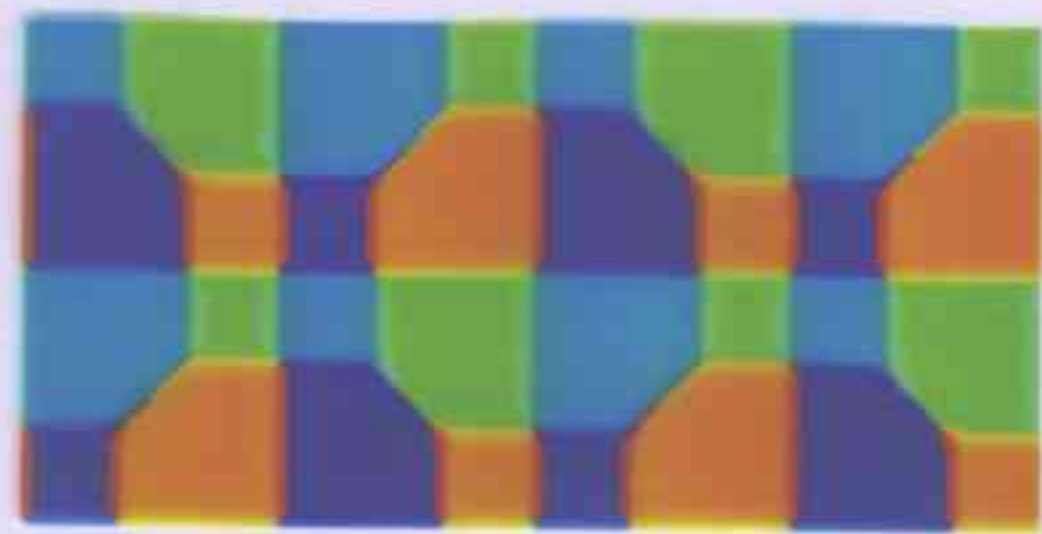
But network is meta-stable - decays by a
cell-nucleation tunnelling (so classically stable)

For different choices of $W(\phi)$ different 'wallpaper'
patterns are possible \rightarrow tilings of plane

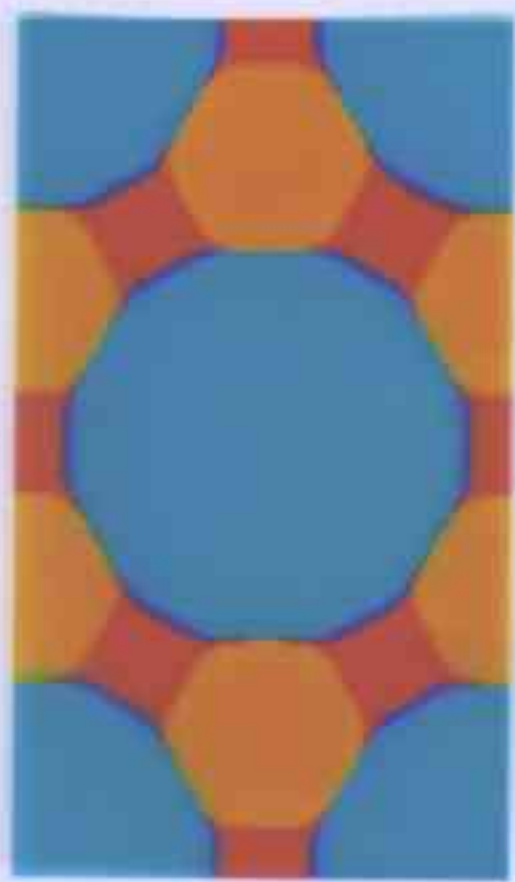
$$W' = 1 - \phi^3$$



$$W' = 1 - \phi^4$$



$$W' = 1 - \phi^6$$



$$W' = 1 - \phi^6$$

