

"OPEN/CLOSED STRING  
DUALITY"

OR

"HOW GRAVITY COUPLES  
TO MATTER"

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BASED ON

WORK

w/

E. VERLINDE

J. KHORRAM

J. de BOER

C. S. CHAN, P. PANK

+  
WORK IN PROGRESS:

w/ ARKANI-HANED, DIMOPOULOS,  
KALOGER

+ DISCUSSIONS w/

BERKOOB, LARSEN, KACHRU,

E. SILVORSTEIN

# INGREDIENTS

① STRING THEORY:

TYPE I

MATTER



OPEN

GRAVITY



CLOSED

② GAUGE HIERARCHY:

$M_W$

:

$M_{Pl}$

1

$10^{16}$

③ HOLOGRAPHY:

RG - SCALE = EXTRA DIMENSION

① + ②



"BRANE WORLD"

ADD

RS

BACK REACTION →



WARPED COMPACTIFICATION

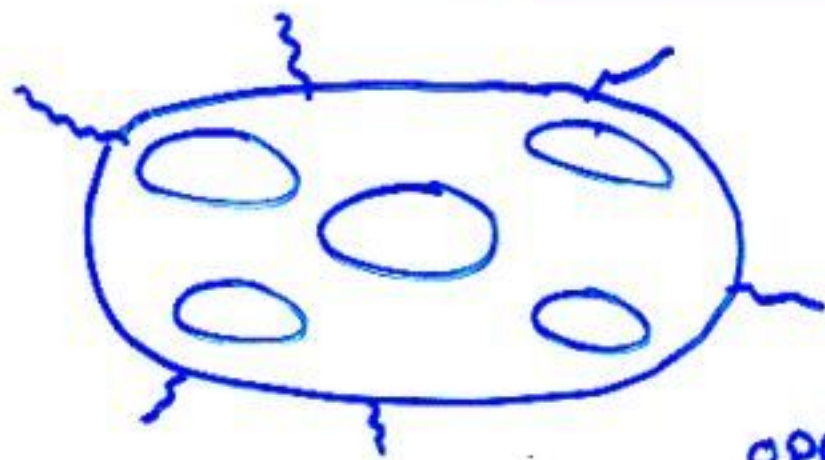
H.I.

AdS / CFT →



③

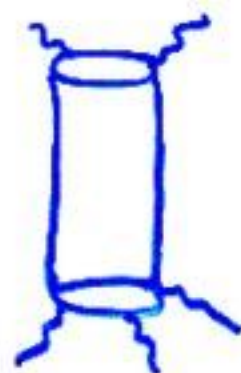
"GEOMETRIZATION OF THE RENORMALIZATION GROUP"



OPEN STRING  
AMPLITUDES



OPEN LOOP  
STRING



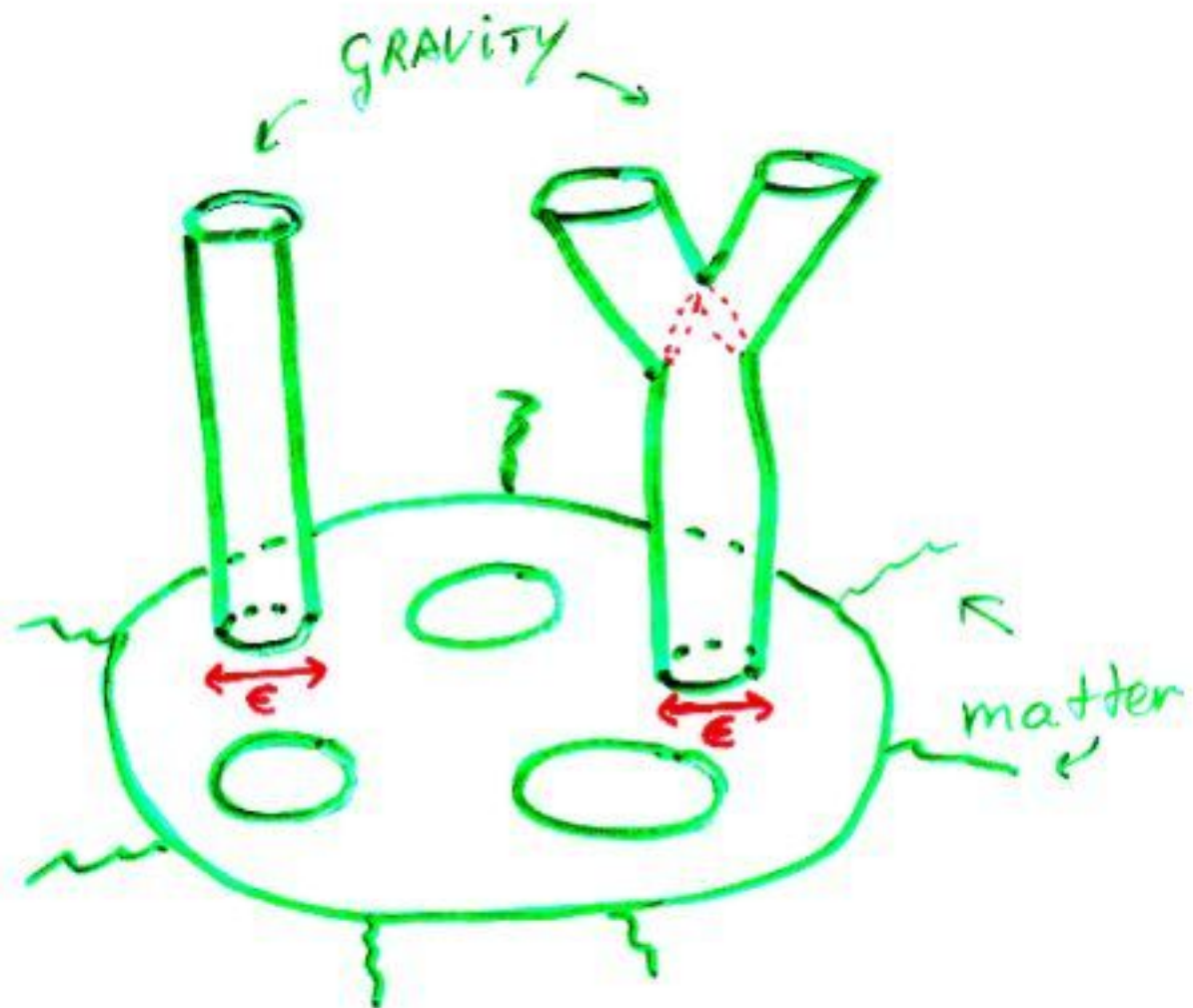
CLOSED STRING  
PROPAGATOR

UV  $\leftrightarrow$  IR

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IN TYPE I, OPEN + CLOSED  
STRINGS ARE IN PEACEFUL ...

# .... COEXISTENCE !



OPEN + CLOSED STRING DIAGRAM

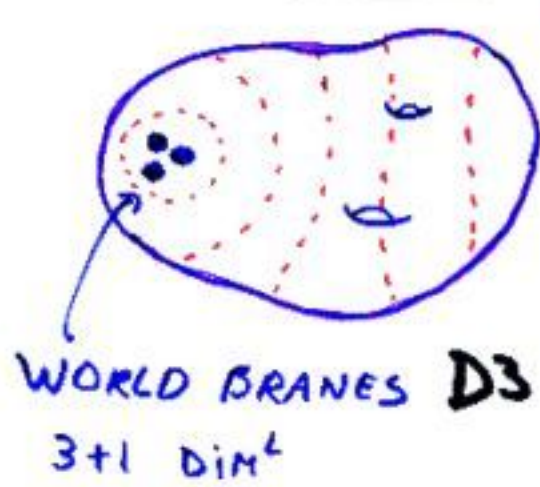


RG RELATION  
 $\epsilon \rightarrow \epsilon + \delta\epsilon$

OPEN STRING  
DIVERGENCES



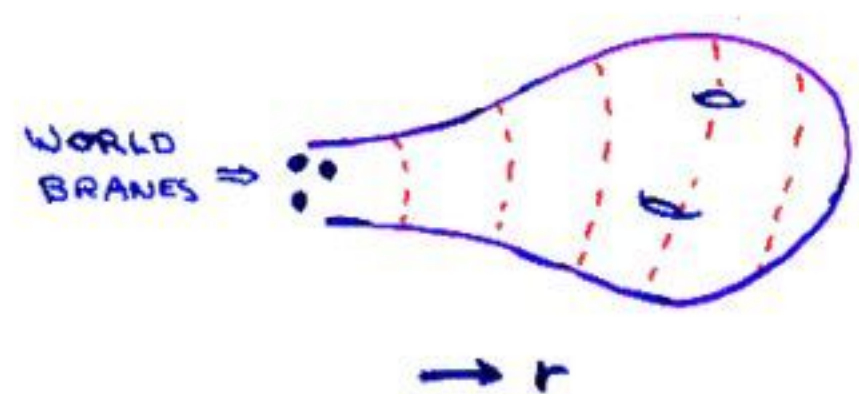
SOURCE TERMS  
FOR CLOSED  
STRING BACKGD



$$K_6 \times \mathbb{R}^4$$

# D3-BRANES  
 $\parallel$   
 $\chi(K_6)$   
 SOME TOPOLOGICAL  
 EULER NUMBER  
F-THEORY

$\Downarrow$   
 BACK REACTION  
 $\Downarrow$



$$K_6 \times \mathbb{R}^4$$

↑  
 WARPED  
 PRODUCT

$$ds^2 = a^2(r) dx_\mu^2 + dr^2 + h_{mn}(r, y) dy^m dy^n$$

DEFORMATION OF COMPACT SLICE  
 OF  $AdS_5$   $\times$   $K_5$

$$\Delta_{\perp} a^2(x_{\perp}) = \sum_i^N \delta(x_{\perp} - y_i) - R^3 \text{-term}$$

||  
EULER  
DENSITY

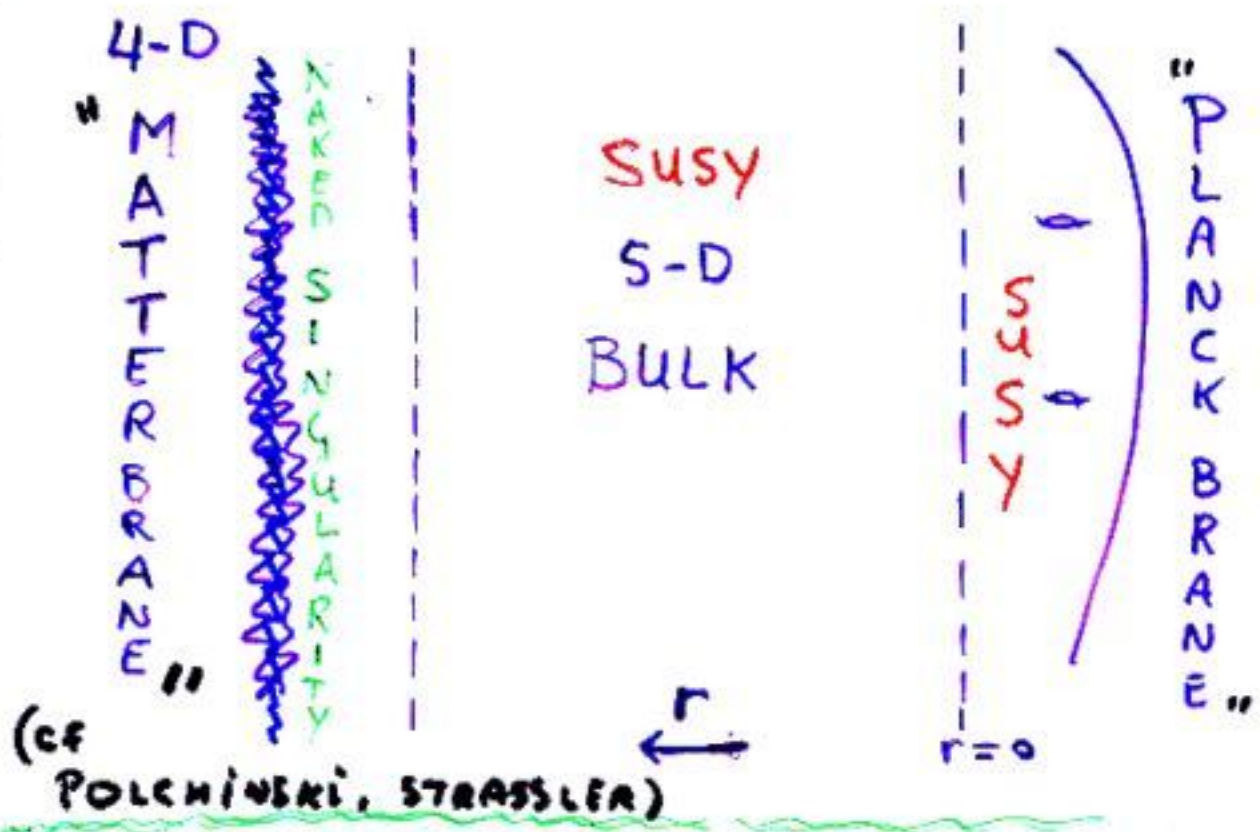
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"PLANCK BRANE"

= REGION OF  $K_6$  OUTSIDE  
OF ADS - THROAT

AFTER DIMENSIONAL REDUCTION OVER  $K_5$

10-D  $\rightarrow$  5-D



$$S_{\text{TOTAL}} = S_{\text{MATTER}}^{4-D} + S_{\text{BULK}}^{5-D} + S_{\text{PLANCK}}^{4-D}$$

$$S_{\text{BULK}} = \int \sqrt{-g} [R_5 + (\nabla\phi)^2 + V(\phi)]$$

+ FERMIONS

$\phi =$  MODULI OF  $K_5$ , ETC.

# MATTER BRANE & HOLOGRAPHY

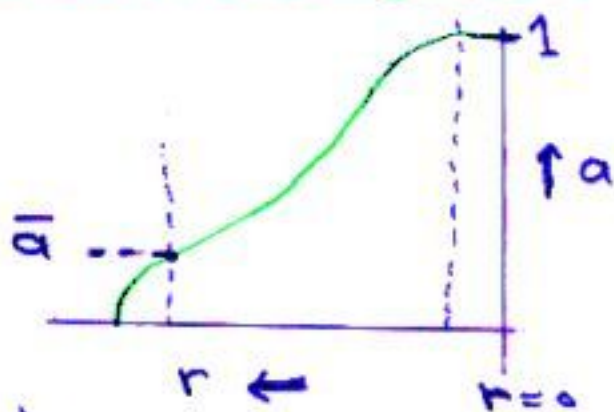
$$e^{\frac{i}{\hbar} \int_{\text{MATTER}} (\bar{a}, \phi)} = \int_{\text{Energy} \leq \bar{a}} d[\text{MATTER}] e^{\frac{i}{\hbar} S(\text{MATTER}, \phi)}$$

$\phi$  = COUPLING CONSTANTS

$\bar{a}$  = UV CUT-OFF ENERGY SCALE

$$ds_s^2 = a^2(r) ds_v^2 + dr^2$$

$$\phi = \phi(r)$$



$$a \frac{\partial}{\partial a} \phi(a) = \beta_\phi(\phi, a)$$

$\Rightarrow$  RADIAL EVOLUTION OF  $a$  &  $\phi$

REPRESENTS RG-FLOW OF THE

QFT COUPLINGS. ■



## CLAIM:

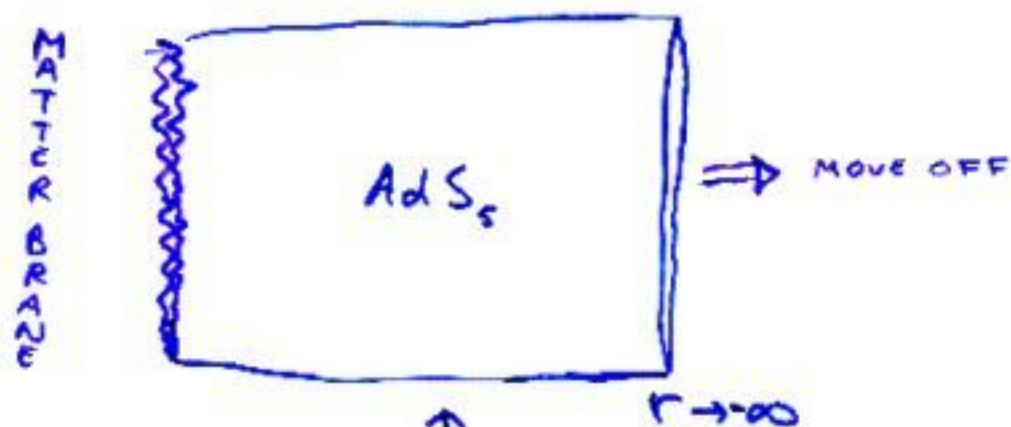
4-D COSMOLOGICAL CONSTANT  $\Lambda_4$  IS NOT  
FIXED BY THE MATTER BRANE EQN. OF MOTION.  
INSTEAD, IT'S AN UNDETERMINED INTEGRATION CNST.

## REASON / MECHANISM

- ① AdS / CFT
- ② ADJUSTMENT OF WARP FACTOR
- ③ OPEN STRING PERTURBATION THEORY
- ④ RG - SCALE DEPENDENCE OF  $S_{\text{MATTER}}(\phi, a)$

ALL ① - ④ ARE IN FACT THE SAME.

① SINCE 5-D EQNS ARE LOCAL, WE CAN SEND THE "PLANCK BRANE" OFF TO  $\infty$ , AND JUST LOOK AT THE MATTER BRANE SEPARATELY, EMBEDDED INSIDE A NON-COMPACT AdS-LIKE 5-D SPACE.



AdS/CFT:

EQUIVALENCE

4-D QFT WITHOUT 4-D GRAVITY!

$$M_{PL} \rightarrow \infty$$

CAN LIVE ON ANY 4-D GEOMETRY

$\Rightarrow R^{(4)}$  IS DETERMINED BY BOUNDARY COND.  
 $r \rightarrow \infty$

## ② 5-D EINSTEIN EQUATION:

$$(\partial_r g_{\mu\nu})^2 - (\partial_r \phi)^2 + V(\phi) + R_4 = 0$$

$$\downarrow ds^2 = a^2(r) ds_4^2 + dr^2$$

$$\left(\frac{\partial_r a}{a}\right)^2 - (\partial_r \phi)^2 + V(\phi) + \frac{k}{a^2} = 0$$

cf. FRW

$k = \Lambda_4$ , THE 4-D COSMOLOGICAL CONSTANT

$\Rightarrow \Lambda_4 =$  INTEGRATION CONSTANT OF THE  $r$  EVOLUTION.

$\Rightarrow \Lambda_4$  IS DETERMINED BY INITIAL CONDITIONS IN THE PLANCK REGION



SCHRÖDINGER = RG FLOW EQN:

$$H e^{\frac{i}{\hbar} S_{\text{MATTER}}[\phi, \bar{a}]} = 0 \quad (*)$$

WITH

$$H = \left( \bar{a} \frac{\partial}{\partial \bar{a}} \right)^2 - \left( \frac{\partial}{\partial \phi} \right)^2 + a^4 V(\phi) + a^2 R^{(4)}$$

THIS EQN FOLLOWS FROM THE INVARIANCE UNDER SHIFTING THE DIVIDING LINE BETWEEN THE BULK & THE MATTER BRANE.

(\*)  $\Rightarrow$  SUPPOSE THAT:

$$S_{\text{MATTER}} = \int \sqrt{-g} [U(\phi, \bar{a}) + \Phi(\phi, \bar{a}) R^{(4)} + \dots]$$

THEN:

$$(**) \left( \frac{\partial}{\partial \phi} U \right)^2 - \left( U + \frac{\bar{a}}{4} \frac{\partial}{\partial \bar{a}} U \right)^2 = V(\phi)$$

$$\overset{\hbar \rightarrow 0}{\left[ \bar{a} \frac{\partial}{\partial \bar{a}} - \rho_{\phi} \frac{\partial}{\partial \phi} \right] \Phi} = \frac{V_{\phi}}{U}$$

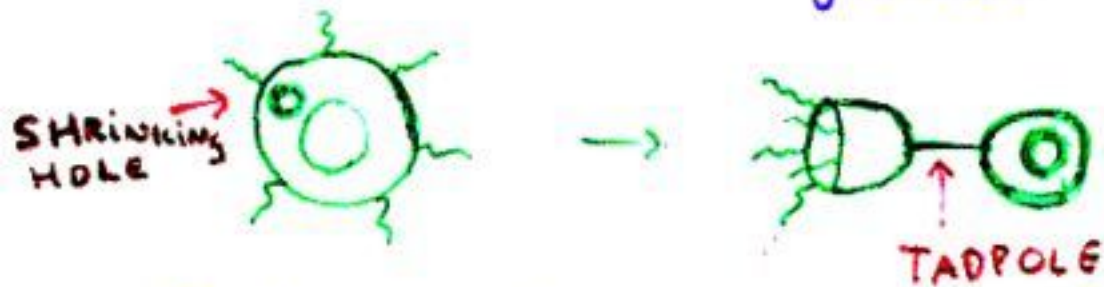
(3) CONSIDER OPEN STRING PERT. SERIES:

$$\begin{aligned}
 \Lambda_4^{\text{OPEN}} &= \text{DISK} + \text{ANNULUS} + \text{ETC} \\
 &= \frac{1}{g_s} \lambda^0 + (g_s)^0 \lambda^1 + \dots
 \end{aligned}$$

↑↑

HOW IS IT CANCELED?

⇒ ~~ONE~~ TWO LOOP DIVERGENCE:



⇒ MUST ADD COUNTERTERM:



RENORMALIZES CLOSED STRING BACKGROUND

SUCH THAT:

$$\Lambda_4^{\text{CLOSED}} = \text{DISK WITH X} = - \Lambda^{\text{OPEN}}$$

"MATTER BRANE"



$$ds_4^2 = \bar{a}^2 dx_n^2$$

"PLANCK SCALE BRANE"



$$a = 1$$



$$\frac{1}{2} \int d^4x [\bar{a}^2 (\partial h)^2 + m_0^2 h^2 \bar{a}^4]$$

$h =$  scalar field on matter brane.

1-loop determinant

$$\frac{\mu^\epsilon}{2} \int \frac{d^{3-\epsilon} k}{(2\pi)^{3-\epsilon}} \sqrt{k^2 + m_0^2 \bar{a}^2}$$



$$\frac{\bar{a}^4 m_0^4}{32 \pi^2} \left[ \frac{1}{\epsilon} - \frac{1}{2} \log\left(\frac{m_0^2 \bar{a}^2}{4\pi \mu^2}\right) - \frac{\gamma}{2} + \frac{3}{2} \right]$$

GOLDBERGER, RUTHSTEIN

AFTER A FEW STANDARD STEPS ONE FINDS THAT BOTH PLANCK + MATTER BRANE ARE STABILIZED, PROVIDED ONE SATISFIES THE MATCHING RELATIONS:

$$\frac{\dot{a}}{a} = \left(1 + \frac{\bar{a}}{4} \frac{\partial}{\partial \bar{a}}\right) U(\bar{\phi}, \bar{a}) = -W(\bar{\phi})$$

$$\dot{\phi} = \partial_{\phi} U(\bar{\phi}, \bar{a}) = -\partial_{\phi} W(\bar{\phi})$$

FOR GENERAL  $U(\bar{\phi}, \bar{a})$  (SOLVING EQNS (\*\*)) THERE EXIST SOLUTIONS TO THESE CONDITIONS. IN FACT, THERE'S A WHOLE "RG-TRAJECTORY" OF SOLUTIONS, GENERATED BY

$$\bar{a} \frac{\partial}{\partial \bar{a}} \bar{\phi}(\bar{a}) = \beta_{\phi}(\bar{\phi})$$

$$\text{WITH } \beta_{\phi} = 6 \frac{\partial_{\phi} W}{W}$$



SUPPOSE WE CHOOSE: (cf. TALK BY  
R. KALLOSH)  
K. STELLER

$$S_{\text{PLANCK}} = \int_{4-D} \sqrt{-g} W(\phi)$$

$\Rightarrow$  BOUNDARY CONDITION AT  $r=0$  BECOMES

$$\partial_r g_{\mu\nu} = W(\phi) g_{\mu\nu}$$

$$\partial_r \phi = \partial_\phi W(\phi)$$

SINCE PLANCK REGION  $r \geq 0$  IS SUPER-SYMMETRIC, IT'S NOT UNREASONABLE TO ASSUME A FINE-TUNED RELATION BETWEEN  $V(\phi)$  AND  $W(\phi)$  SUCH THAT:

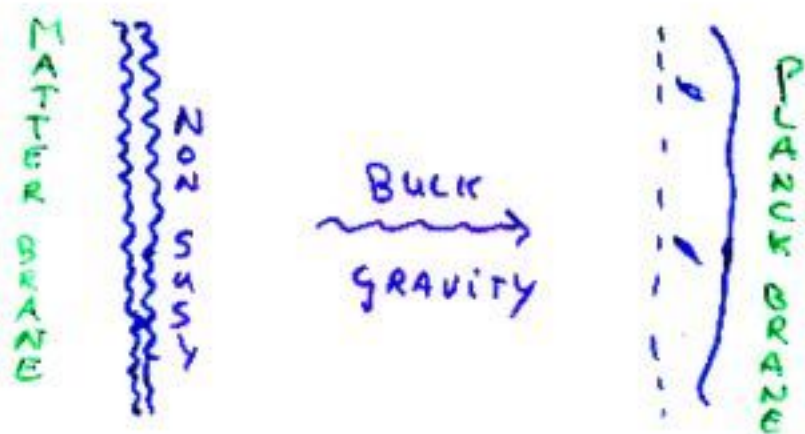
$$\Lambda_4 = V(\phi) + (W(\phi))^2 - (\partial_\phi W)^2 \equiv 0$$

$\Rightarrow$   $W(\phi) =$  SUPERPOTENTIAL  
OF 5-D GAUGED SUGRA

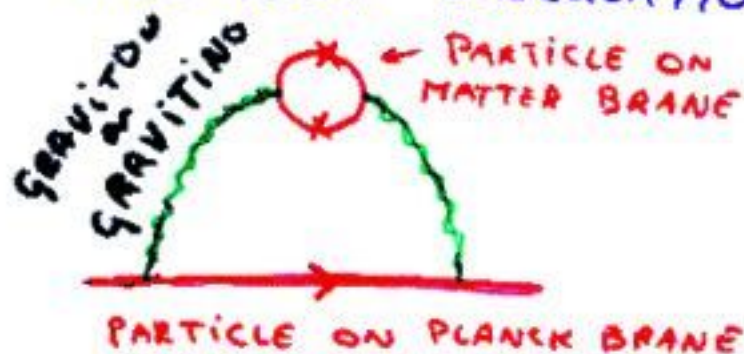
$\Rightarrow$  "PLANCK BRANE" IS FLAT  $\nabla$

FINALLY: How LARGE/SMALL is  $\Lambda_4$ ?

NATURAL SCALE IS SET BY THE AMOUNT OF SUSY VIOLATION IN PLANCK REGION.



ONE LOOP CALCULATION GIVES:



$$M_{3/2} \approx \frac{m_{\text{susy}}^2}{m_{\text{pl}}}$$

$$\Lambda_4 \approx \frac{(M_{3/2})^4}{m_{\text{pl}}^2} \approx \frac{(m_{\text{susy}})^8}{(m_{\text{pl}})^6}$$

(C. T. BANKS, C. SCHMIDHUBER.)

## CONCLUSIONS

- ① WARPED GEOMETRY NATURALLY ARISES IN STRING THEORY. LARGE HIERARCHY STILL REQUIRES FINE-TUNING, HOWEVER.  
⇕  
SUSY
- ② HOLOGRAPHIC INTERPRETATION OF EXTRA DIMENSION AS RG SCALE  $\Rightarrow$  KK PARADIGM MODIFIED.
- ③  $\Lambda_4$  IS DETERMINED LOCALLY IN THE PLANCK REGION. CAN BE PROTECTED BY SUPERSYMMETRY  $\Rightarrow \Lambda_4 \approx (m_{\text{SUSY}})^8?$

## QUESTIONS

- ① EXPLICIT STRING REALIZATIONS. (MAYR) H.V.
- ② STABILITY, FLUCTUATIONS
- ③ GRAVITY LOOPS,  $1/N$ ?
- ④ PHASE TRANSITIONS?
- ⑤ 4-D EFFECTIVE FIELD THY?