

# D1/D5 SYSTEMS IN THEORIES WITH 16 SUSY

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hep-th/

(Ads/CFT & D1/D5 systems

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(in Preparation)

32, 16 SUSY

- D1/D5 system in  $\text{IIB} / T^4$  or  $K^3$  is reasonably well understood.
- But for Type I (16 susy) has not been studied yet.
- In general U-duality would map

$\text{IIB} / K_3 \rightarrow$  Type I like models

$\downarrow$   
 $(4,4)$

$\downarrow$   
 $(4,0)$

Effective world volume theory

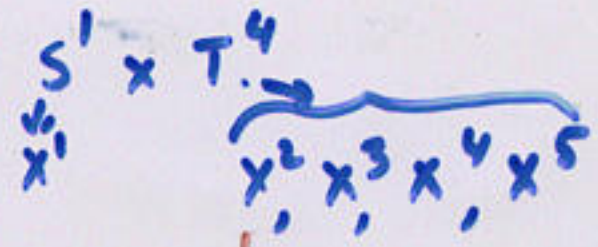
→ As a starting point we study here  $\mathbb{Z}_2$  orbifolds of IIB theory, with  $\mathcal{V}_2$  being  $I_4, \Omega, \Omega I_4$  together with a shift.

→ Applying Adiabatic argument → deduce CFT from Parent CFT.

## Outline

- 1) U-duality chain between  $\Omega$ ,  $I_4$ ,  $\Omega I_4$  models
- 2) Effective field theory of  $D1/D5$  system
- 3) Test : ground state of 2).  
 $\rightarrow$  perturbative spectra in a U-dual theory
- 4) Discussion on  $D1/D5/KK$  3-charge system.
- 5) AdS/CFT correspondence.

IB ON



|              | $S$              | $T_{15}$         | $S$          | $T_{2345}$   | $S$              |
|--------------|------------------|------------------|--------------|--------------|------------------|
| $D_1$        | $F_1$            | $P_1$            | $P_1$        | $P_1$        | $P_1$            |
| $D_{12345}$  | $NS_{12345}$     | $NS_{12345}$     | $D_{12345}$  | $D_1$        | $F_1$            |
| $P_1$        | $P_1$            | $F_1$            | $D_1$        | $D_{12345}$  | $NS_{12345}$     |
| $\Omega$     | $(-1)^{F_1}$     | $(-1)^{F_1}$     | $\Omega$     | $\Omega I_4$ | $(-1)^{F_1} I_4$ |
| $I_4$        | $I_4$            | $(-1)^{F_1} I_4$ | $\Omega I_4$ | $\Omega$     | $(-1)^{F_1}$     |
| $\Omega I_4$ | $(-1)^{F_1} I_4$ | $I_4$            | $I_4$        | $I_4$        | $I_4$            |
| $X$          | $X$              | $C_{15}$         | $B_{15}$     | $G_{15}$     | $G_{15}$         |
| $C_{2345}$   | $C_{12345}$      | $C_{1234}$       | $C_{1234}$   | $C_{15}$     | $B_{15}$         |

IB on  $S^1 \times T^4 \times S^1$   
 $x^1 \quad x^2 \dots x^5 \quad x^6$

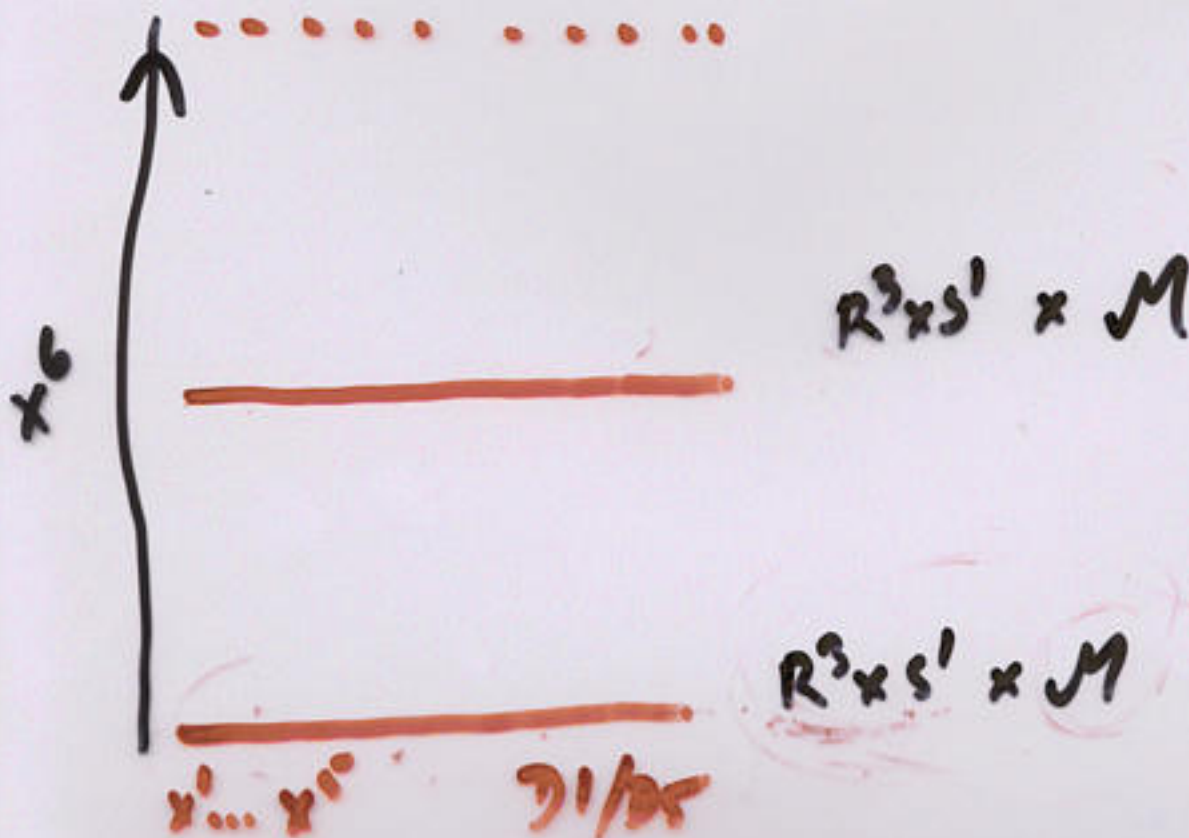
$Z_2$  ORBIFOLD

$g, \sigma_6$

$g = \Omega, \Omega I_{2345}, I_{2345}$

$\sigma_6 = \text{HALF SHIFT ON } x^6$

EFFECTIVE THEORY



$\mathfrak{g}$  will induce an action  
 $\hat{\mathfrak{g}}$  on  $M$

$\hat{\mathfrak{g}}$  identification

$\Rightarrow$  ONLY ONE SET

$$\mathbb{R}^3 \times S^1 \times M$$

DESCRIBES THE SYSTEM

$\Rightarrow$  IF ONE SHIFTS THE ENTIRE  
 SYSTEM  $x^i \rightarrow x^i + \pi R^i$

$$M \rightarrow \hat{\mathfrak{g}} M$$

THEN IT DESCRIBES THE  
 ORIGINAL SYSTEM.

# EFFECTIVE FIELD THEORY

$$(R^3 \times S^1 \times M) / \mathbb{Z}_2$$

$$\downarrow$$
$$\hat{g}, \sigma$$

↳ SHIFT ALONG  
 $x^6$

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STARTING POINT  $\mathbb{R}^3 / T^4$

$$M = T^4 \times (T^4)^N / S_N$$

(Maldacena, Moore  
strominger)  
TATA GROUP

$$Q_5 = 1 \quad Q_1 = N$$

$$\chi = \frac{1}{2} \quad \& \quad C^{(4)} = \frac{2}{N}$$

(Witten  
Seiberg, Witten  
Larsen, Martinec)  
TATA GROUP

(4,4) Supermultiplets

|                  | Boson       | Left fermion   | Right Fermion          |
|------------------|-------------|----------------|------------------------|
| $R^3 \times S^1$ | $X_{AY}$    | $\psi_A$       | $\tilde{\psi}_Y$       |
| $T^4$            | $A_i$       | $\psi_Y$       | $\tilde{\psi}_A$       |
| $(T^4)^N$        | $X_i^{(l)}$ | $\psi_Y^{(l)}$ | $\tilde{\psi}_A^{(l)}$ |

Little group  $SO(4) \cong SU(2)_A \times SU(2)_Y$

Tangent space  $\mathfrak{g}$

|             | $R^3 \times S^1$ | $T^4$ | $\Omega$ | $\Omega I_4$ |
|-------------|------------------|-------|----------|--------------|
| $X_{AY}$    | +                |       | +        | +            |
| $A_i$       |                  | -     | -        | +            |
| $X_i^{(l)}$ |                  | -     | +        | -            |

U(1) Wilson line

Position of D1 inside D5

Action on the fermions is obtained by

(4,4) SUSY and the fact that

$\hat{g}$  commutes with (4,4) SUSY for  $I_4$

and " " (4,0) SUSY for  $\Omega, \Omega I_4$



|             |                                  |                        |                              |
|-------------|----------------------------------|------------------------|------------------------------|
|             | $I_4$                            | $\Omega$               | $\Omega I_4$                 |
| $\hat{g} =$ | $I_4^{cm} \cdot I_4^{sym Prod.}$ | $(-1)^{F_L} I_4^{cm.}$ | $(-1)^{F_L} I_4^{Sym Prod.}$ |
|             | $(4,4)$                          | $(4,0)$                | $(4,0)$                      |

Longitudinal Shift case

$\sigma$  along  $x'$  (common world volume)

→ Periodic or anti-periodic B.C. along  $x'$  depending on the eigenvalues of  $\hat{g}$ .

# GENERAL PROBLEM

$$\left( (\text{CFT})^N / S_N \right) / \mathbb{Z}_2$$

$$Z = \sum_{\mathcal{H}_{\text{CFT}}} q^{L_0} \bar{q}^{\bar{L}_0} y^{J_3} \bar{y}^{\bar{J}_3}$$

$J_3, \bar{J}_3$   $SU(2)_A \times SU(2)_Y$  Quantum #'s

$$Z_{ab} = \sum_{ab} C(\Delta, \bar{\Delta}, \theta, \tilde{\theta}) q^\Delta \bar{q}^{\bar{\Delta}} y^\theta \bar{y}^{\tilde{\theta}}$$

$a, b = \pm$  twists along  $\sigma, t$  directions

Put  $\bar{\Delta} = 0$  (To sector in SuperSymm)  
in  $(4,4)$  or  $(4,0)$   
case)

Ellipticize The Four Sectors of the  
Symmetric Product CFT give

$$E_{ab} = \sum_N p^N E_{N; ab} (q, y, \tilde{y}) ; \quad \bar{\Delta} = 0$$

$$E_{++} = (1 - p^n q^m y^l \tilde{y}^{\tilde{l}})^{-c_{++}(nm, l, \tilde{l})}$$

$$E_{+-} = (1 - p^n q^m y^l \tilde{y}^{\tilde{l}})^{-\frac{1}{2}(c_{++} + c_{+-})(nm, l, \tilde{l})} \cdot (1 + p^n q^m y^l \tilde{y}^{\tilde{l}})^{-\frac{1}{2}(c_{+-} - c_{++})(nm, l, \tilde{l})}$$

$$E_{-+} = (1 - p^{2n} q^m y^l \tilde{y}^{\tilde{l}})^{-\frac{1}{2}(c_{++} + c_{+-})(2nm, l, \tilde{l})} \cdot (1 \mp p^{2n} q^{m-\frac{1}{2}} y^l \tilde{y}^{\tilde{l}})^{-\frac{1}{2}(c_{+-} - c_{++})(2n(m-\frac{1}{2}), l, \tilde{l})} \cdot (1 - p^{2n+1} q^m y^l \tilde{y}^{\tilde{l}})^{-\frac{1}{2}(c_{--} + c_{-+})(2n+1)m, l, \tilde{l}} \cdot (1 \mp p^{2n+1} q^{m-\frac{1}{2}} y^l \tilde{y}^{\tilde{l}})^{-\frac{1}{2}(c_{-+} - c_{--})(2n+1)(m-\frac{1}{2}), l, \tilde{l}}$$

For  $\Delta = 0$  (i.e. pure D1/D5 system)

$E_{++}$  &  $E_{+-}$  agree with perturbative results in the  $U$ -dual theory

# 3-CHARGE SYSTEM & a "PUZZLE"

Suppose  $\rightarrow$  u-duality

$D_1$   
 $D_{12345}$   
 $P_1$

u-duality  $\rightarrow$

$P_1$   
 $D_{12345}$   
 $D_1$

$E(p, q, r)$

$\rightarrow$

$E'(r, p, q)$

(Recall  $p^{Q_1} q^{K_1}$ )

( $\gamma = 1$ )

(A)

(B)

$I_4$

$\rightarrow$

$\Omega I_4$

$D_1/D_5$  system

$D_1/D_5$  system

is  $(4, 4)$

is  $(4, 0)$

$\downarrow$

Elliptic genus

$E_{CM}(q, r) E_{Int}(p, r)$

$$E_{Int.}(p, z, \gamma) = E_F(p, \gamma) \cdot \hat{E}_{Int.}(p, z, \gamma)$$

(where  $E_F(p, \gamma) = E_{Int.}(p, 0, \gamma)$ )

⇒ U-duality implies:

$$E_{ch}^A = E_F^B$$
$$E_F^A = E_{ch}^B$$

} TRUE FOR CFT'S PROPOSED HERE

$$\hat{E}_{Int.}^A = \hat{E}_{Int.}^B$$

→ Not true

WHAT IS GOING WRONG !

THIS IS A GENERIC PHENOMENON  
WHEN MODEL (A) IS (4,4) & sym. Prod CFT  
and MODEL (B) IS (4,0)

(ALSO FOR IB / K3)

Indeed take model (A) which is (4,4)

Sym. Prod. CFT.

$$E^A = E_{cm}^A(q, \gamma) E_F^A(p, \gamma) \underbrace{\hat{E}_{\text{Sym. Prod.}}^A(p, q, \gamma)}_{\text{Sym in } p \leftrightarrow q}$$

1) for  $q^0$ ;  $E^A$  has second order zero in  $(\gamma-1)^2$  ( $\frac{1}{2}$  BPS)

2) for  $q^m$ ;  $m > 0$   $E^A$  has fourth order zero  $(\gamma-1)^4$  ( $\frac{1}{4}$  BPS)

↓ U-duality

$$E_{cm}^A(p, \gamma) E_F^A(q, \gamma) \hat{E}_{\text{Sym. Prod.}}^A(p, q, \gamma)$$

expand  $p^1$  from the U-dual  
Internal contribution

$$\rightarrow \boxed{E_{cm}^A(p, \gamma) \hat{E}_{\text{Sym. Prod.}}^A(p, q, \gamma) = E_{cm}^A(q, \gamma) - E_{cm}^A(0, \gamma) + f_1^A(\gamma)}$$

where

→  $E_1(2, \gamma)$  is the elliptic genus for single copy of  $(CFT)^A$  (4, 4)

→  $f_1(\gamma)$  is the coeff of  $p'$  in  $E_{CM}^A(p, \gamma)$

This expression has no well defined modular property under some subgroup of  $SL(2, \mathbb{Z})$

$q \quad \tau = \frac{\ln q}{2\pi i} \rightarrow -\frac{1}{c}$

$\tau = \frac{\ln \gamma}{2\pi i} \rightarrow -\tau/c$

UNLESS  $E_1^A(0, \gamma) = f_1^A(\gamma)$

BUT THIS IS IN GENERAL NOT TRUE

$E_1^A(0, \gamma)$  counts the # of D1/D5 BOUND STATES IN MODEL (A) ((4, 4) Theory)

In GENERAL  $E_1^A(0,1) \neq 0$

But  $f_1^A(\gamma) \rightarrow 0$  at  $\gamma \rightarrow 1$

since it counts the  $1/4$  BPS states  $D1/D5$  in MODEL (B)

(4,0) Theory

What is the possible resolution?

- 1) Perhaps  $1/4$  BPS states are not stable. (cf. 3-string Junction a la Ashoke Sen in  $N=4$  SYM Theory)

But in model A (4,4) theory the singularity (threshold) appears in a subspace of co-dimension  $> 1$  in the moduli space.

$\Rightarrow$  Elliptic genus does not jump !!



$\Rightarrow$  3-charge (D1, D5, KK)  
must be stable.

2) Sym. Prod. theory is defined at  
 $\chi = 1/2$

Im model (A)  $\chi = 1/2 \rightarrow \chi = 0$

Im model (B)  $\chi = 1/2 \quad \Delta \chi = 0$   
are distinct models.

U-duality on model (A) at  $\chi = 1/2$

$\rightarrow$  model (B) at  $\chi = 0$

This is entirely singular

On the other hand  $\chi = 1/2$  in Model (B)  
at which Sym. Prod. CFT exists  
is mapped to Model (A) with non-zero  
 $G_{15}$  which breaks Lorentz invariance of  
the world sheet.

### 3. AdS/CFT Correspondence

If  $R_6$  is very large the near horizon geometry of D1/D5 system

$$\sim \text{AdS}_3 \times S^2 \times T^4$$

Fluctuation analysis

$h_{m,n}$   $(m+n, m+n)$   $(4,4)$  Super multiplets

$h_{m,n}$  are the hodge #'s of  $T^4$

$$h_{00}^{(1)}$$

$$h_{01}^{(2)} \quad h_{10}^{(2)}$$

$$h_{02}^{(1)} \quad h_{11}^{(4)} \quad h_{20}^{(1)}$$

$$h_{12}^{(2)} \quad h_{21}^{(2)}$$

$$h_{22}^{(1)}$$

$\mathbb{Z}_2$  action on (c.c) primers

$$\hat{g} = (-1)^{k+s} \quad \text{for } I_4 \text{ model}$$

$$= (-1)^s \quad \text{for } \Omega I_4 \text{ model}$$

Furthermore

$$\hat{g} \text{ commutes with } \overset{\text{(Left)}}{a_{-1/2}} \text{ \& } \overset{\text{(Right)}}{\tilde{a}_{-1/2}} \quad \text{for } I_4 \text{ model} \quad (4,4)$$

But  $\hat{g}$  commutes with  $\tilde{a}_{-1/2}$  &  
anti commutes with  $a_{-1/2}$  for  
~~the~~  $\Omega I_4$  model.  $(4,0)$

By Following Jan De Boer  
associate "degree"

Poincare Polynomials give correct  
D1/D5 Bound states

Elliptic genus for multiparticle  
SUGRA states

for model  $I_4 \xrightarrow{\text{agrees}} \text{CFT for } I_4 \text{ model}$

for model  $\Omega I_4 \xrightarrow{\text{agrees}} \text{CFT for } I_4 \text{ model}$   
with  $p \leftrightarrow v$

$$p^N q^0$$

$$\left. \begin{array}{l} \Delta < N/4 \\ N < 0/4 \end{array} \right\}$$

This result seems to be consistent with

$\rightarrow$   $(D1, D5, KK)$  system is stable in  
 $I_4$  model

$\rightarrow$   $\Omega I_4$  model at  $\kappa=0$  is indeed U-dual  
to  $I_4$  model

Question: Can one do SUGRA  
analysis for  $\chi = 1/2$  for  $\mathcal{N}=4$   
model?

AND

Does This agree with the  
(4,0) CFT we have proposed  
here?