

# AdS branes and Holography

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Strings 2002 - Cambridge

based mainly on:

CB, J. de Boer, R. Dijkgraaf, H. Ooguri  
hep-th/011210

CB, hep-th/0205115

\* Rich extension of holography:

a  $AdS_m$  brane in  $AdS_m$  bulk  
reaches the boundary ( $\infty$  blueshift)  
where it appears as

↳ external source ( $m < m-1$ )

↳ domain wall, or  
defect line ( $m = m-1$ )

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\* Relevant for warped Brane World  
Scenarios without perfect fine  
tuning of brane tension

Karch, Randall  
Porrati

⋮

Can these be realized in string  
theory?

see talk by Giddings  
H. Verlinde, ...

Interesting because product-space  
compactifications so far failed to  
address stability + cosmo. const. problems

\* Can take pp-wave limits

- Skenderis, Taylor
- Bergman, Gaberdiel, Green
- Bain, Meessen, Zamaklar
- Mateos, Ng
- Seki
- ⋮

Here I will concentrate on first point, and comment briefly on second one.

## General set up

AdS<sub>m</sub> bulk supported by flux:

$$ds^2 = L^2 \left\{ \frac{du^2}{u^2} + u^2 (-dt^2 + dx^2 + dy^2 + dz^2) \right\}$$

$$G = L^2 u^2 dt_{\wedge} dx_{\wedge} u^2 dy_{\wedge} dz$$

$$(m=3,5)$$

Consider static p-brane, with embedding  $u(x)$ ;  $y$  &  $z$  are spectator coordinates (either longitudinal or transverse). The brane energy is:

$$\mathcal{E} = T_{\text{eff}} \int \sqrt{u^4 + u'^2} u^{p-1}$$

depends on  
compact part

$$- \rho \int u^{m-1} + \dots$$

↑ only present for  
codimension 1

$$(p=m-2)$$

Extrema given by first-order eqns:

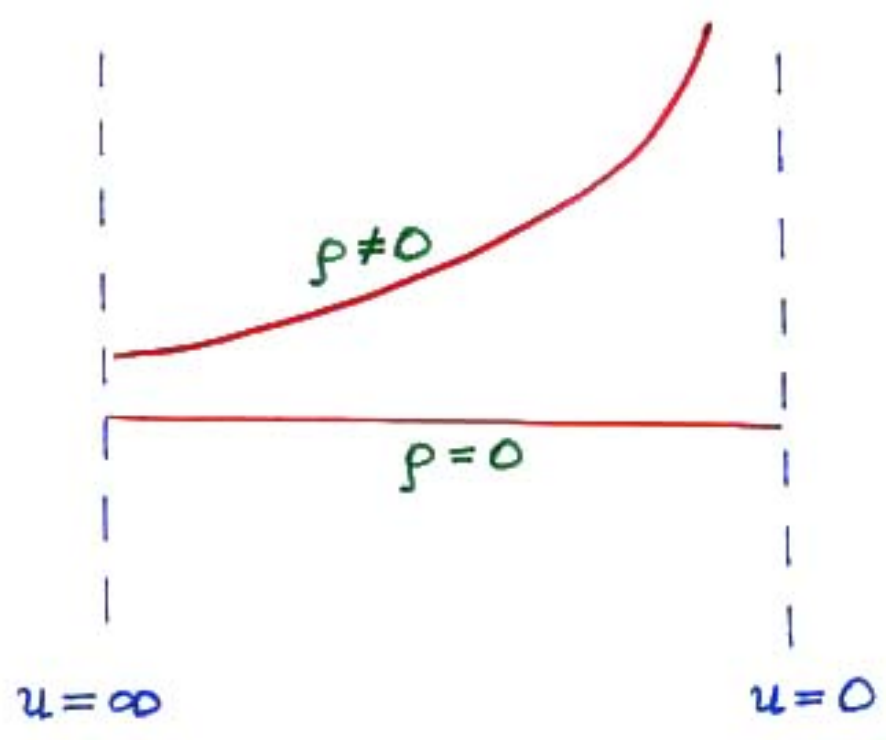
$$\frac{T_{\text{eff}} u^{p+3}}{\sqrt{u^4 + u'^2}} - \rho u^{m-1} = -\Theta_{xx} \text{ constant}$$

if  $p = m - 2$

worldvolume energy-moment. tensor

Free endpoint b.c.s (no applied force)

$$\Rightarrow \boxed{x = \frac{C}{u}} \text{ with } \boxed{C = \frac{\rho}{\sqrt{T^2 - \rho^2}}}$$



$$L_{\text{brane}} = \sqrt{1 + C^2} L \rightarrow \infty \text{ for } T \rightarrow \rho$$

# Lift to (IIB) string theory

\*  $AdS_3 \times S^3 \times T^4$  (near-horizon of  $n_1$  F1,  $n_5$  NS5 branes)  
 ↳ up to U-duality

Brane is a D3/D1/F1 bound state  
 ↳ couple to background flux

with  $AdS_2 \times S^2$  geometry.

Has been constructed as exact conformal boundary state (in Euclidean theory)

- \* Ponsot, Schomerus, Teschner
- also: Rajaraman, Rozali
- Parnachev, Sahakyan
- Giveon, Kutasov, Schwimmer
- Hikida, Sugawara
- Petroopoulos, Ribault
- Lee, Ooguri, Park,
- Tannenhauser
- ⋮
- Stanciu

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\*  $AdS_5 \times S^5$  (near horizon of D3 branes)

probe: D5/D3 bound state  
with  $AdS_4 \times S^2$  geometry.

Many possible generalizations ...

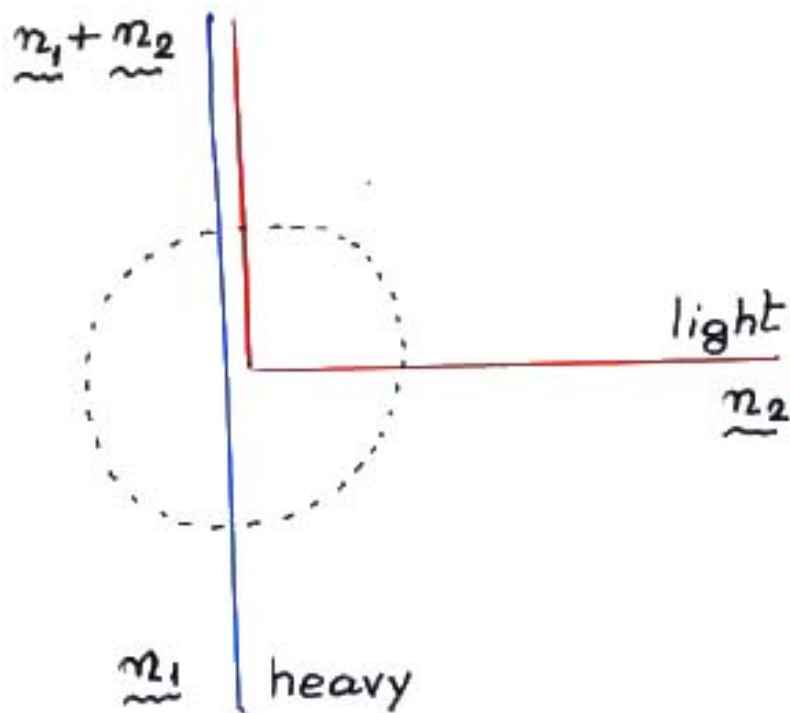
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## Holographic duals

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↳ Consider first  $AdS_3$  case.

- Type IIB theory compactified on  $T^4$  has variety of strings (D5, D3's, D1, NS5, F1) in fundamental rep. of  $U$ -duality group  $SO(5,5|\mathbb{Z})$ .
- These can form supersymmetric string junctions.
- What we describe is the  $nh$  blow-up of a junction between a heavy (background) string and a light (probe) one.



- On holographic screen see domain wall between CFTs with different central charge or different moduli.
  - ↳ modified attractor mechanism  
Ferrara, Kallosh, Strominger
- Probe brane is AdS and supersymmetric  
⇒ wall is superconformal
- Similar story in higher dimensions, but more complicated since wall can have interacting CFT on its worldvolume

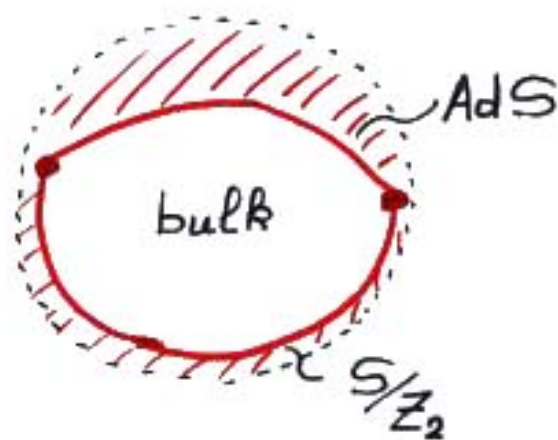
De Wolfe, Freedman, Oguri  
Erdmenger, Guralnik, Kirsch  
Mateos, Ng, Townsend



## Important clarification:

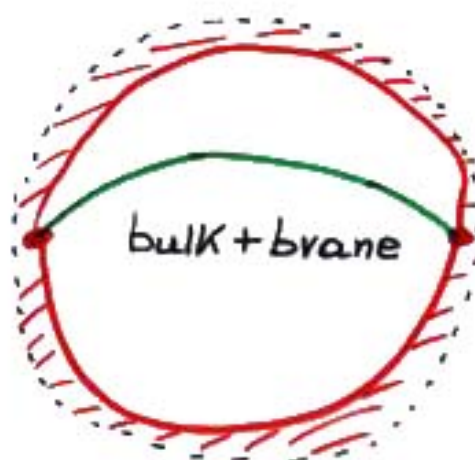
↳ many papers on subject  
take holographic screen  
to be  $AdS \times S/Z_2$

Karch, Randall  
Porrati  
Bousso, Randall  
Duff, Liu, Sati



gravity localized  
on AdS brane further  
mapped on  $\partial(S/Z_2)$

↳ Here holographic screen will be  $S$



# Confirm with analysis à la Brown-Henneaux

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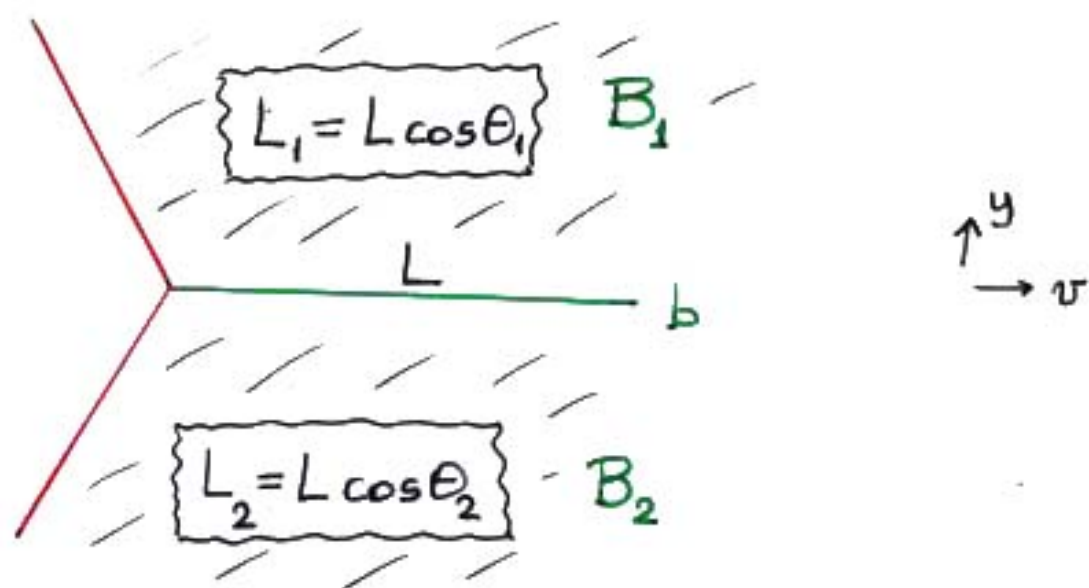
or in 'modern version'  
Henningson + Skenderis  
Balasubramanian + Krauss  
⋮

They showed that asymptotic symmetries of gravity in  $AdS_3$  are two (left + right) Virasoro algebras with  $c = 3L/2G_N$

Repeat analysis for gravity in two patches of  $AdS_3$  spacetime glued along common  $AdS_2$  boundary:

$$ds^2 = \frac{1}{f^2} (du^2 + dy^2 - dt^2)$$

with  $f(u,y) = \begin{cases} \frac{u + y \tan \theta_1}{L} & y > 0 \\ \frac{u + y \tan \theta_2}{L} & y < 0 \end{cases}$



Solves eqns. deriving from

$$S_{\text{bulk}} + S_{\text{boundary}} + S_{\text{brane}}$$

$\downarrow$  discontinuous cosmo. constant       $\downarrow$  Gibbons-Hawking & counterterm

$$S_{\text{brane}} = T \int \sqrt{-\hat{g}} - \frac{1}{8\pi G_N} \int \sqrt{\hat{g}} [K]$$

due to thin brane approx.

iff  $\tan \theta_1 - \tan \theta_2 = 8\pi G_N L \cdot T$

Change coordinates to

$$ds^2 = L^2 \left\{ \frac{du^2}{u^2} + u^2(-dt^2 + dx^2) + 2 \sin \theta_r dx du \right\}$$

in  $B_r$

↳ One Virasoro algebra of asymptotic symmetries at  $u \rightarrow \infty$ :

$$\begin{cases} \delta x^\pm = -\xi^\pm - \frac{1}{2u^2} \partial_\mp^2 \xi^\mp + \frac{\sin \theta_r}{2u} (\partial_+ \xi^\pm - \partial_- \xi^\mp) \\ \delta u = \frac{u}{2} (\partial_+ \xi^\pm + \partial_- \xi^\mp) - \frac{\sin \theta_r}{2} (\partial_+^2 \xi^\pm - \partial_-^2 \xi^\mp) \end{cases}$$

with  $\xi^\pm = g(x^\pm)$

↳ same function for continuity at brane

↳ Brown-York tensor

$$T_{ab} = \frac{2}{\sqrt{-g}} \frac{\delta S}{\delta g^{ab}} \quad \rightarrow \text{metric at boundary}$$

$$\delta T_{\pm\pm} = \frac{L \cos \theta_r}{16\pi G_N} \partial_\pm^3 \xi^\pm \quad \text{in } \partial B_r$$

$$\delta T_{+-} = 0$$

i.e. consistent with expected jump of central charge.

# Permeable CFT walls

Can two CFTs interact non-trivially along conformal wall?

Is this generic RG fixed point?

Folding trick:

Affleck, Oshikawa  
⋮



$$T_{xt}^1 = T_{xt}^2 \quad (\text{no net flow of energy to boundary})$$

fold ↘



$$T_{xt}^{\text{tot}} = T_{++}^{\text{tot}} - T_{--}^{\text{tot}} = 0$$

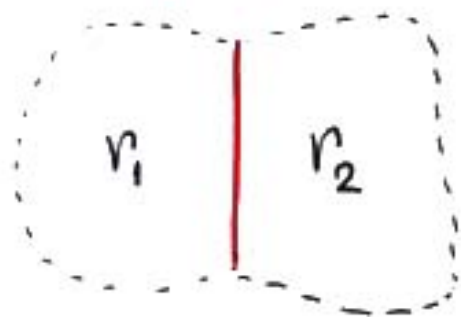
CFT1  $\otimes$  CFT2

∴ Conformal boundary state in tensor product theory.

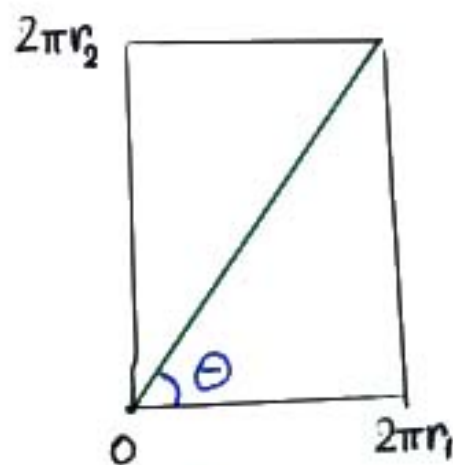
Permeable  $\Leftrightarrow$  not  $\sum \otimes$  (separate Ishibashi states)

Interesting classification problem  
for bCFT. Can construct explicit  
examples:

\* Radius jump for free field



described after folding  
by diagonally-embedded  
D1-brane



$$\tan \theta = r_2 / r_1$$

\* Jump in Kähler moduli  
of CY CFTs

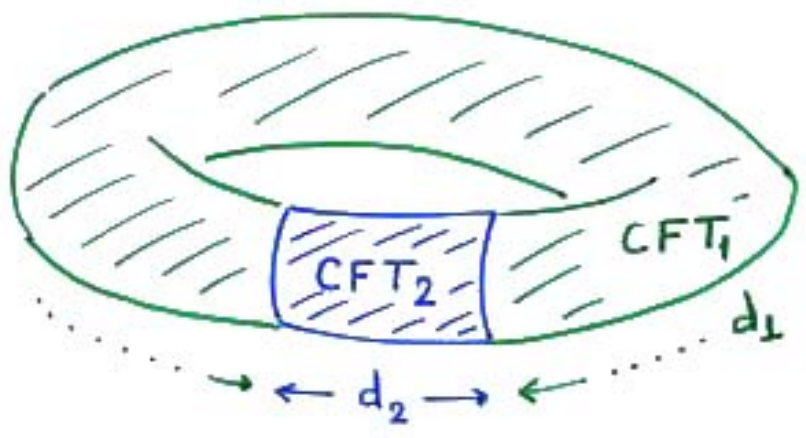
CY, CY': same complex structure  
 $\leadsto$  diagonal brane  $z_i = z_i'$

\* Some algebraic constructions

Quella, Schomerus  
Recknagel

# Casimir energy of conformal bubble

↳ CFT calculation:



$$\bar{Z} = \langle\langle B | q_1 L_0^{(1)} - \frac{c^{(1)}}{24} q_2 L_0^{(2)} - \frac{c^{(2)}}{24} | B \rangle\rangle$$

with  $q_i = \exp(-2\pi d_i/T)$

$$E_{\text{casimir}} = \lim_{T \rightarrow \infty} \lim_{d_i \rightarrow \infty} -\frac{1}{T} \log \langle\langle \quad \rangle\rangle$$

\* For free scalar with radius jump

$$\lim_{q_i \rightarrow 0} \langle\langle \quad \rangle\rangle \propto \prod_{n=1}^{\infty} (1 - \cos^2 2\theta q_2^{2n})^{-1}$$

quantum dilogarithm

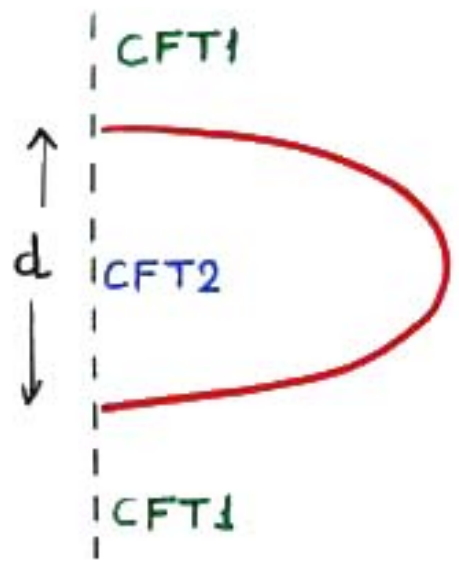
$$\Rightarrow E_{\text{Casimir}} = -\frac{1}{8\pi d} \text{Li}_2(\cos^2 2\theta)$$

↳ classical dilogarithm

$-\frac{\pi}{48d}$  when  $\frac{n_1}{n_2} \rightarrow 0, \infty$   
(perfect reflection)

$-\frac{1}{8\pi d} \left(\frac{\delta n}{n}\right)^2$   
for  $n_1 \approx n_2$   
(near perfect transmission)

↳ Supergravity calculation



solution with  $\Theta_{xx} \neq 0$   
(not free endpoint  $\Rightarrow$  not AdS)

analogous to  $q\bar{q}$  potential  
Maldacena  
Rey



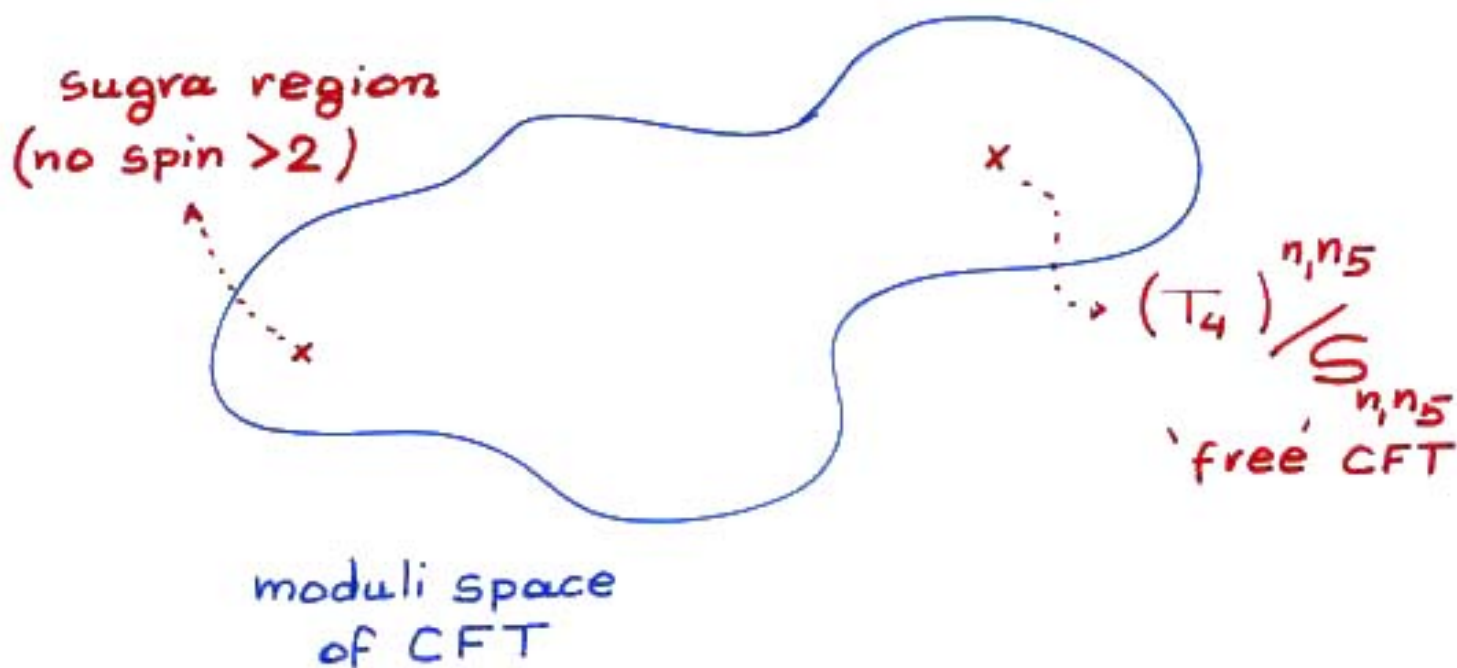
end result :

$$E_{\text{Casimir}} = -\frac{2LT}{d} \left\{ 2E(R) - K(R) \right\}^2$$

complete elliptic integrals

$$R^2 = \frac{T - \rho}{2T}$$

Cannot compare, since different regions of validity



↳ Strong coupling prediction. Need exact Conf. boundary state for asymp. AdS branes to approach symmetric-product orbifold??

## Speculation

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- \* In 'RS-like' scenarios the Planck brane in AdS bulk serves to generate 4d gravity. But where does SM live?
- \* If it were realized holographically then  $N, g_{YM}^2 N$  large.  
Uncomfortable idea  $\Rightarrow$  better to realize SM with probe brane in susy 'RS-like' bulk.
- \* Tantalizing parametric suppression of cosmological constant

$$L_{\text{brane}} \sim \sqrt{\frac{T}{\delta T}} \cdot L_{\text{bulk}}$$

\* Suppose that:

$$L_{\text{bulk}} \sim \text{mm} \sim 10^{+12} / \text{GeV}$$

$$T \sim (10^{18} \text{ GeV})^4$$

$$\delta T \sim (10^3 \text{ GeV})^4$$

↳  $m_{\text{susy}}^4$  from loops on the brane

$$\Rightarrow L_{\text{brane}} \sim 10^{30} \text{ mm}$$

present Hubble radius of our universe

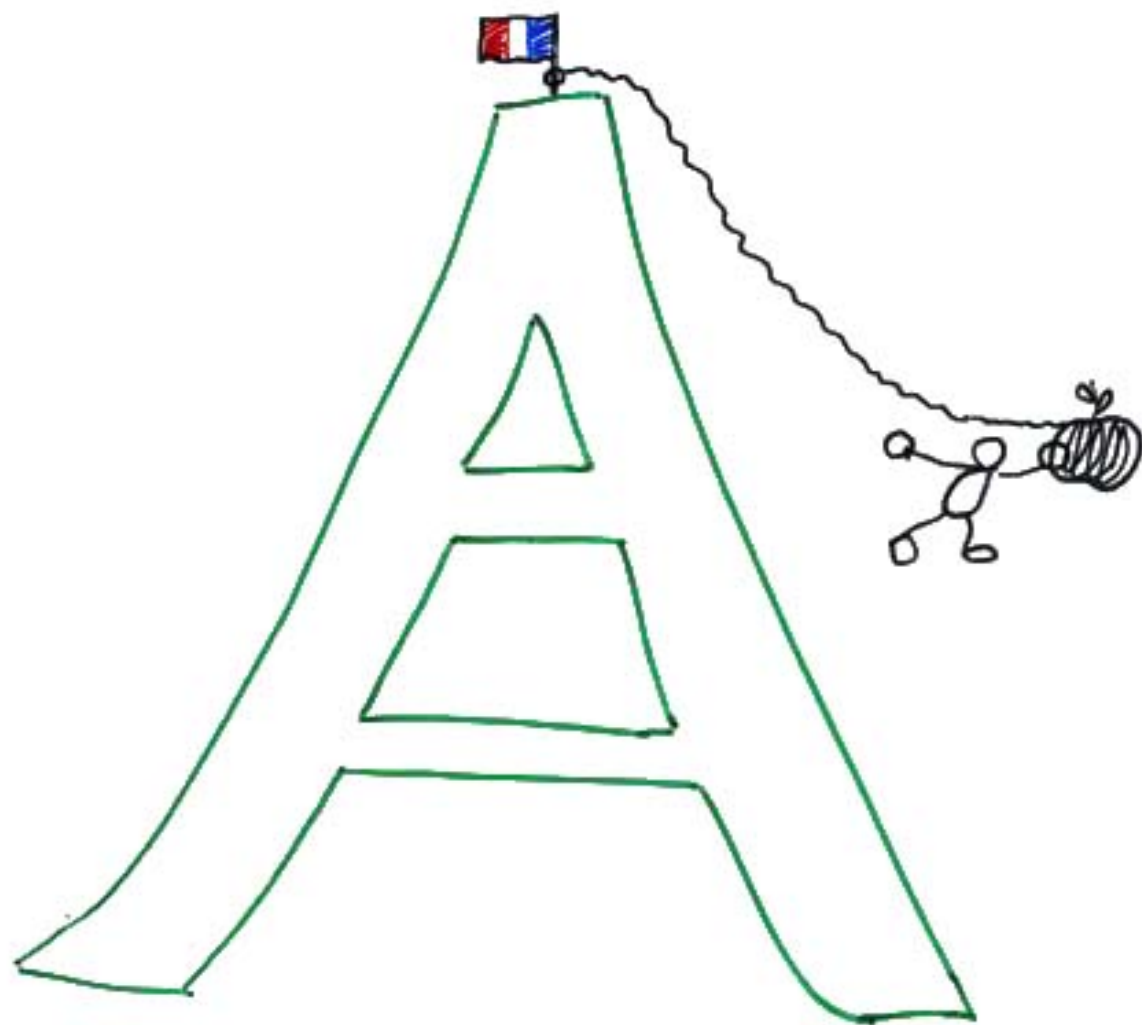
\* Problem: how to generate

$$G_N^{(4)} \sim (10^{19} \text{ GeV})^{-2} \quad ??$$

STRINGS '04

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We hope to welcome  
you all there, in 2 years