

GEOMETRY

QUANTIZATION

from

SUPERGRAVITY

VYACHESLAV RYCHKOV - U. AMSTERDAM

STRINGS 05 - TORONTO -

- GRANT, MARSANO, MAOZ, PAPADOPOULOS, V.R.
- MAOZ, V.R. to appear hep-th/2505079

THANKS TO: SHIRAZ MINWALLA
AVINASH DHAR

GOAL

MODULI SPACE QUANTIZATION OF
SUGRA SOLUTIONS

↳ "BUBBLING ADS"

PLAN

1. MINISUPERSPACE QUANTIZATION.
CRNKOVIĆ - WITTEN - ZUCKERMAN METHOD

↳ SYMPLECTIC CURRENT

2. "BUBBLING ADS" EXAMPLE

- WAVY LINE APPROXIMATION
(AROUND PLANE WAVE)

- GENERAL DROPLET CASE

QUESTION

• HOW TO QUANTIZE FAMILIES OF GRAVITY SOLUTIONS?

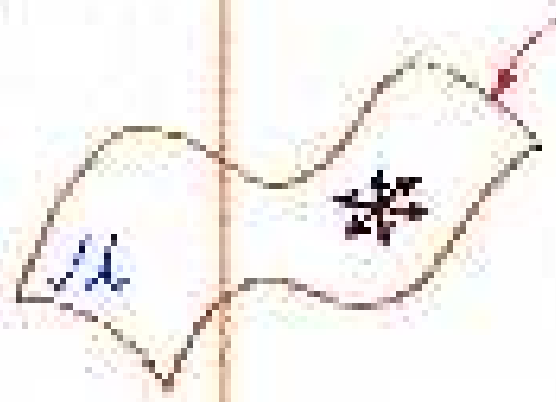


FLUXES ARE QUANTIZED

$$F = F_0 n, \quad n \in \mathbb{Z}$$

(DIRAC QUANTIZATION)

2.



MODULI SPACE,
ALL FLUXES
CONSTANT



RECENT EXAMPLES

③

BUBBLING ADS GEOMETRIES

LIN, LUNIN, MALDACENA



MICROSCOPIC

FREE FERMION DROPLETS



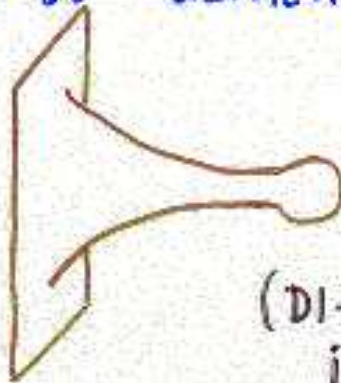
CORLEY, JEVICKI, RAMGOOLAN
BERENSTEIN

D1-D5 GEOMETRIES WITH ANGULAR MOMENTUM

LUNIN, MALDACENA, MAOZ



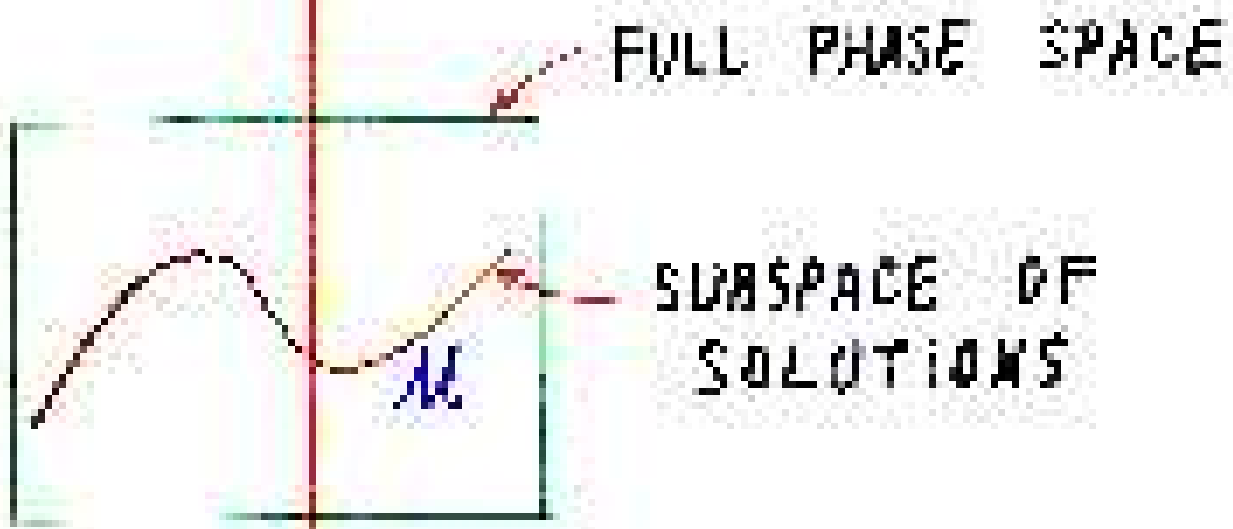
CHIRAL FUND. STRING



$F_i(\nu)$
 $i=1, \dots, 4$

(D1-D5 BH MICROSTATES
IN MATHUR'S PROGRAM)

GENERAL IDEA



RESTRICT :

$$\omega = \Omega / M$$

Ω - SYMPLECTIC FORM OF GIVEN THEORY

EXAMPLE

CHIRAL SECTOR OF A BOSON

$$S = \frac{1}{2} \int dt dx (\partial\varphi)^2$$

• SYMPLECTIC FORM:

$$\Omega = \int_{t=\text{const}} dx \delta\dot{\varphi} \wedge \delta\varphi$$

δ - EXTERIOR DIFFERENTIAL

• MODULI SPACE OF SOLUTIONS:

$$\mathcal{M} = \{ \varphi = f(t+x), f \text{ arbitrary} \}$$

• RESTRICTION:

$$\omega = \Omega|_{\mathcal{M}} = \int du \delta f'(u) \wedge \delta f(u)$$

⇒ POISSON BRACKETS

$$\{ f(u_1), f'(u_2) \} = \delta(u_1 - u_2)$$

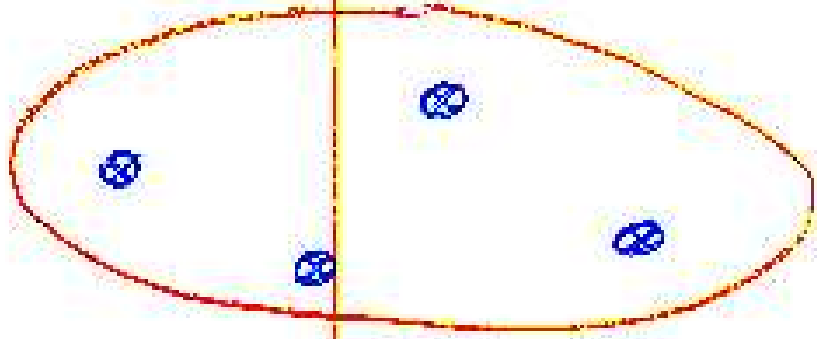
⇒ QUANTIZATION

THIS IS NOT MODULI SPACE

QUANTIZATION à la MANTON

MANTON STORY:

MANTON
ATYAH, HITCHIN
GIBBONS, RUBACK
FERRELL, EARDLEY
...



3n-DIM
MODULI SPACE

• MULTI-CENTERED STATIC EXTREMAL
BH SOLUTIONS

• SLOW SCATTERING DESCRIBED
BY GEODESICS IN (CALCULABLE)
METRIC $ds^2 = g_{ij}(q) dq^i dq^j$

• PHASE SPACE IS 6n-DIM

• SYMPLECTIC FORM DEGENERATE
ON THE ORIGINAL 3n-DIM MODULI
SPACE

→ NO BH CENTER QUANTIZATION

SYMPLECTIC FORM OF GRAVITY ⑦

- REGULAR SOLUTIONS (NO HORIZON)
- CAUCHY SURFACE Σ



CANONICAL FORMALISM

- 3+1 SPLITTING

$$ds^2 = -N^2 dt^2 + h_{ij} (dx^i + N^i dt)(dx^j + N^j dt)$$

ARNOWITT, DESER, MISNER

- PHASE SPACE VARIABLES (h_{ij}, Π^{ij})

$$\Pi^{ij} = \sqrt{h} (K \cdot h^{ij} - K^{ij})$$

- CONSTRAINTS
- SYMPLECTIC FORM

$$\Omega = \int_{\Sigma} d^3x \delta \Pi^{ij}(x) \wedge \delta h_{ij}(x)$$

COVARIANT APPROACH

CRNKOVIĆ - WITTEN CURRENT

⑧

$$\Omega = \int d\Sigma_\mu J^\mu$$

SYMPLECTIC CURRENT:

$$J_{CW}^\mu = \delta\Gamma_{\nu\lambda}^\mu \wedge \delta[\sqrt{g} g^{\nu\lambda}] - \delta\Gamma_{\nu\lambda}^\nu \wedge \delta[\sqrt{g} g^{\mu\lambda}]$$

- $\partial_\mu J^\mu = 0 \Rightarrow \Omega$ INDEPENDENT OF Σ VARIATIONS

- Ω INV. UNDER

$$\delta g_{\mu\nu} \rightarrow \delta g_{\mu\nu} + \nabla_{(\mu} \xi_{\nu)}$$

ON SOLUTIONS

CRNKOVIĆ, WITTEN 86

GENERALIZATION TO ANY LAGRANGIAN THEORY

$$\mathcal{L} = \mathcal{L}(\varphi_A, \partial_\mu \varphi_A)$$

$$\Omega = \int d\Sigma_\mu \mathcal{J}^\mu$$

$$\mathcal{J}^\mu = \delta \left[\frac{\partial \mathcal{L}}{\partial \partial_\mu \varphi^A} \right] \wedge \delta \varphi^A$$

- $\partial_\mu \mathcal{J}^\mu = 0$
 - Ω GAUGE INV.
- } ON SOLUTIONS OF E.O.M.

ZUCKERMAN
LEE, WALD

FIXED FLUXES RESTRICTION

10

GAUGE FIELD QUANTIZATION
IN NONTRIVIAL TOPOLOGY



E. g. ELECTROMAGNETISM

$$S = \int \sqrt{g} F_{\mu\nu}^2 d^4x$$

SYMPLECTIC CURRENT:

$$J^M = \delta[\sqrt{g} F^{M\nu}] \wedge \delta A_\nu$$

USES δA_ν !

→ MUST CHOOSE $\delta A_\nu(x)$ REGULAR
EVERYWHERE ON Σ

→ IMPOSSIBLE UNLESS
 $\delta F = d(\delta A)$ HAS
ZERO FLUX ON ALL 2-CYCLES

WORKED EXAMPLE:

(11)

"BUBBLING ADS"

MODULI SPACE:



PROBLEM:

FIND SYMPL. FORM RESTRICTED ON MODULI SPACE

• WHICH VARIATIONS TO CONSIDER?

ALL FLUXES FIXED \Leftrightarrow

ALL BLACK & WHITE AREAS FIXED

RELEVANT ACTION

$(g_{\mu\nu}, F_5)$ SECTOR

$$S = \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{g} (R - 4|F_5|^2)$$

SYMPL. CURRENT:

$$J^M = - \delta \Gamma_{\nu\lambda}^M \wedge \delta [\sqrt{g} g^{\nu\lambda}] + \delta \Gamma_{\lambda\nu}^{\nu} \wedge \delta [\sqrt{g} g^{M\lambda}]$$

(GRAVITY PART)

$$- 8 \delta [\sqrt{g} F^M \dots] \wedge \delta A \dots$$

(GAUGE PART)

$$(F = dA)$$

$$\begin{aligned}
 ds^2 = & -h^{-2} (dt + V_i dx^i)^2 + h^2 (dy^2 + dx_i^2) \\
 & + y e^G d\Omega_3^2 + y e^{-G} d\tilde{\Omega}_3^2 \\
 & i=1,2
 \end{aligned}$$

$$h^{-2} = \frac{y}{\sqrt{1/4 - z^2}} ; \quad e^{2G} = \frac{1/2 + z}{1/2 - z}$$

$$z = \frac{1}{\pi} \frac{y^2}{(x^2 + y^2)^2} * Z$$

$$V_i = \frac{\epsilon_{ij}}{\pi} \frac{x_j}{(x^2 + y^2)^2} * Z$$

$$Z = \begin{cases} -1/2, & x \in \bullet \\ 1/2, & x \notin \bullet \end{cases}$$

$$\begin{aligned}
 F_5 &= F \wedge d\Omega + \tilde{F} \wedge d\tilde{\Omega} \\
 F &= dB \quad \tilde{F} = d\tilde{B}
 \end{aligned}$$

$$B_i = -\frac{y^2 V_i}{4(1/2 - z)} - \frac{\epsilon_{ij}}{4\pi} \frac{x_j}{x^2 + y^2} * Z + \frac{x_1}{4} \delta_{i,2}$$

$$\tilde{B}_i = -\frac{y^2 V_i}{4(1/2 + z)} - \frac{\epsilon_{ij}}{4\pi} \frac{x_j}{x^2 + y^2} * Z - \frac{x_1}{4} \delta_{i,2}$$

$$B_y = \tilde{B}_y = 0$$

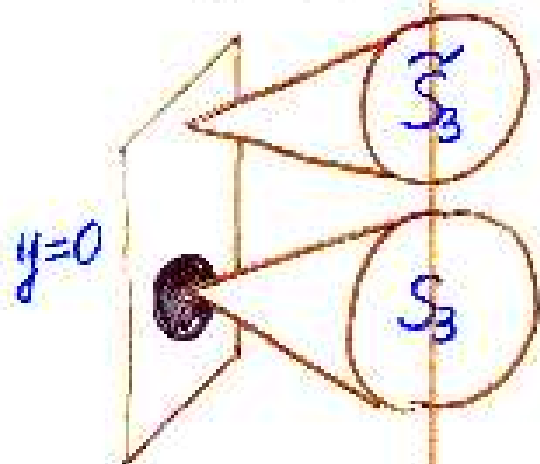
SUBTLETY

(13)

- $A_4 = B \wedge d\Omega + \tilde{B} \wedge d\tilde{\Omega}$

IS SINGULAR ON $y=0$ PLANE:

$$B, \tilde{B} \rightarrow \text{CONST} (y \rightarrow 0)$$



S_3 (\tilde{S}_3) SHRINKS

- δA_4 CAN BE MADE REGULAR BY GAUGE TRANSF. (DROPLET DEP.)

$$\delta B_i = \partial_i \lambda \quad (y=0, Z = -1/2)$$

$$\delta \tilde{B}_i = \partial_i \tilde{\lambda} \quad (y=0, Z = 1/2)$$

FULL SYMPLECTIC FORM

$$\Omega = \int_{t=\text{const}} J^t \quad (\text{BULK})$$

$$+ 8 \int_{y=0, Z = -1/2} \lambda \wedge \delta \tilde{F}_{12} \quad (\text{BOUNDARY COMPENSAT.})$$

$$- 8 \int_{y=0, Z = 1/2} \tilde{\lambda} \wedge \delta F_{12}$$

SIMPLEST COMPUTATION :

PERTURBATION AROUND PLANE WAVE

WAVY LINE APPROXIMATION

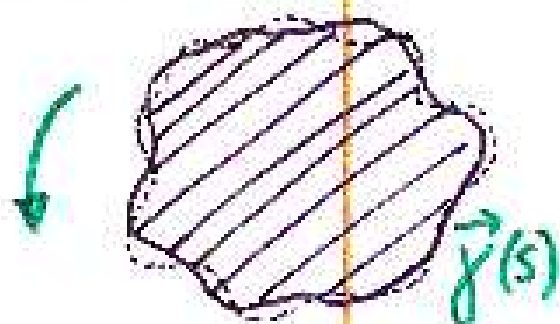


1. EXPRESS VARIATIONS OF ALL FIELDS VIA $E(x_1)$
2. COMPUTE SYMPL. CURRENT
3. INTEGRATE TO GET SYMPL. FORM

ANSWER (IN MOM. REP.)

$$\omega = \frac{(2\pi^2)^2}{K_{10}^2} \int \frac{dp}{2\pi} \frac{i}{p} E(p) \wedge E(-p)$$

DUAL PICTURE - FREE FERMIONS (15)



SEMICLASSICAL
QUANTIZATION
(LARGE N LIMIT)
HYDRODYNAMIC APPROACH

$$H_{\text{TOT}} = \int_{\text{DROPLET}} \frac{dp dq}{2\pi \hbar} \frac{p^2 + q^2}{2} = H_{\text{TOT}}[\vec{\gamma}(s)]$$

1D PLANCK CONSTANT

$$\frac{d\vec{\gamma}}{dt} = \{ \vec{\gamma}, H_{\text{TOT}} \}$$

P.B. GENERATED BY SYMPLECTIC FORM:

$$\omega = \frac{1}{8\pi \hbar} \oint \oint ds ds' \text{Sign}(s-s') \delta\chi_{\perp}(s) \wedge \delta\chi_{\perp}(s')$$

$$\left(\oint ds \delta\chi_{\perp}(s) = 0 \right)$$

POLYCHRONAKOS
DHAR
G.M.M.P.R.

→ AGREES WITH WAVY LINE
GRAVITY COMPUTATION

$$\left(\hbar = \frac{\kappa_{10}}{4\pi^{5/2}} \right)$$

GENERAL DROPLET CASE

(16)

SYMPLECTIC CURRENT VIA FIELD VARIATIONS:

$$J^t = y^3 \frac{3/4 + z^2}{(1/4 - z^2)^2} \delta(V_i \partial_i z) \wedge \delta z$$

$$+ y^3 \delta(h^{-4} V_j \partial_{[i} V_{j]}) \wedge \delta V_i$$

$$+ 8 \epsilon^{ij} (\delta \tilde{B}_i \wedge \partial_y \delta B_j - \delta B_i \wedge \partial_y \delta \tilde{B}_j)$$

NEED TO INTEGRATE:

$$\omega \sim \int dy dx_1 dx_2 J^t + \int_{y=0} (\dots)$$

CAN SHOW THAT

$$J^t = \partial_y \left[\epsilon_{ij} \delta \left(\frac{y^4 (1/4 + z^2) z V_i}{(1/4 - z^2)^2} \right) \wedge \delta V_j \right]$$

$$+ \partial_i \left[\frac{2 y^3 z}{1/4 - z^2} \delta z \wedge \delta V_i \right]$$

BOUNDARY CONTRIBUTION

DOES NOT CONTRIBUTE

(17)

ADDING $y=0$ BOUNDARY CONTRIBUTIONS,
ONE GETS

$$\omega \sim \int_{y=0, z=-1/2} \epsilon^{ij} \delta C_i \wedge \delta C_j$$

$$\delta C_i = \epsilon_{ij} \frac{x_j}{x^2} * \delta Z = \partial_i \lambda$$

$$\omega \sim \int_{\partial \mathcal{D}} \lambda \wedge \delta C_{||}$$

\leftarrow DROPLET BDRY

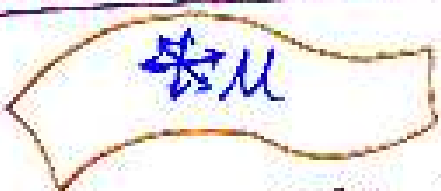
$$\begin{array}{l} \delta C_{||} \rightarrow \delta \chi_{\perp}(s) \\ \lambda \rightarrow \int^s \delta \chi_{\perp}(s') ds' \end{array} \quad \left. \begin{array}{l} \text{on } \partial \mathcal{D} \\ \text{sp. TOP} \end{array} \right\}$$

$$\Rightarrow \omega \sim \oint ds \delta \chi_{\perp}(s) \wedge \int^s \delta \chi_{\perp}(s') ds'$$

AGREEMENT
(ALSO COEFFICIENT)

CWZ VS. EFFECTIVE ACTION

Itzhak Kleban
Walter Kerner



QUANTIZING QUADRATIC FLUCTUATIONS

$S_I (\delta g, \delta A, \dots)$ - quadratic action

- COULD FIRST QUANTIZE THEN SPECIALIZE.

S_{II} ;

-TRACTABLE IN CERTAIN CASES (AdS, PLANE WAVE)

-HARD (IMPOSSIBLE?) IN GENERAL

- TRUNCATING ACTION TO MODES FROM M IS IMPOSSIBLE IN GENERAL.

E.g. FOR LLM AROUND $AdS_5 \times S^5$

$\Delta = J$ MODES COUPLE TO $\Delta = -J$

IN S_{II} (ANGULAR MOMENTUM CONSERVATION)

CWZ BYPASSES THESE PROBLEMS

CONCLUSIONS

- MINISUPERSPACE QUANTIZATION
→ CHECK/FIND MICROSCOPIC DUALS OF SUGRA SOLUTION FAMILIES
- COVARIANT CNZ METHOD
- QUANTUM EQUIVALENCE
"BUBBLING ADS" \Leftrightarrow FREE FERMIONS

FUTURE

- D1-D5 SOLUTIONS OF LUNIN-MALDACENA-NAOZ
→ EXPECT $\omega \sim \int dt \delta F_i'(v) \wedge \delta F_i(v)$
- FURTHER APPLICATIONS TO BH MICROSTATE GEOMETRY COUNTING
→ SEE BENA'S TALK
- 1/4 BPS SOLS OF LIU-VAMAN
→ FRACTIONAL STATISTICS?
- BETTER UNDERSTAND CORRECTIONS/
ROLE OF SUSY