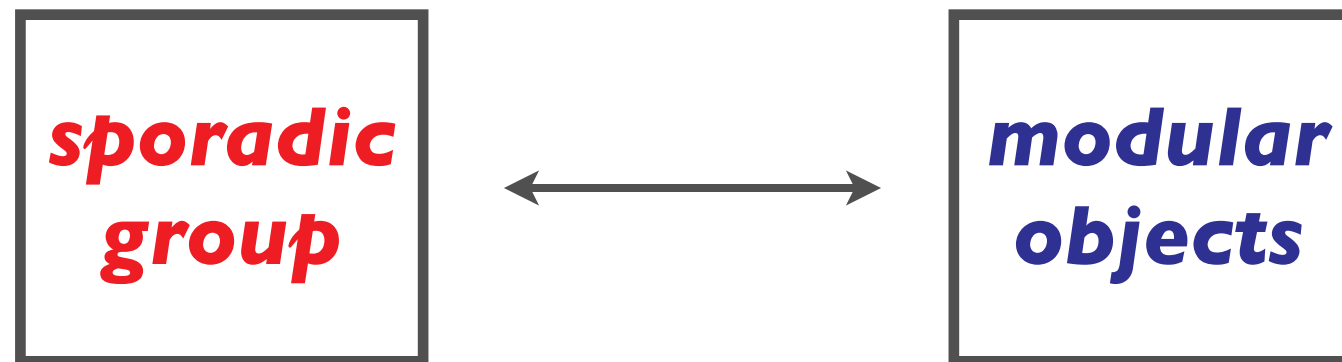


***K3, M₂₄, and
Holographic Moonshines***

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Introduction

The Moonshine Phenomenon



$$J(\tau) = J(\tau + 1) = J(-\frac{1}{\tau}) \quad (q = e^{2\pi i\tau})$$

$$= q^{-1} + 196884q + 21493760q^2 + \dots$$

||

1 + 196883

dim of irreps of Monster

||

1 + 196883 + 21296876

[McKay late 70's]

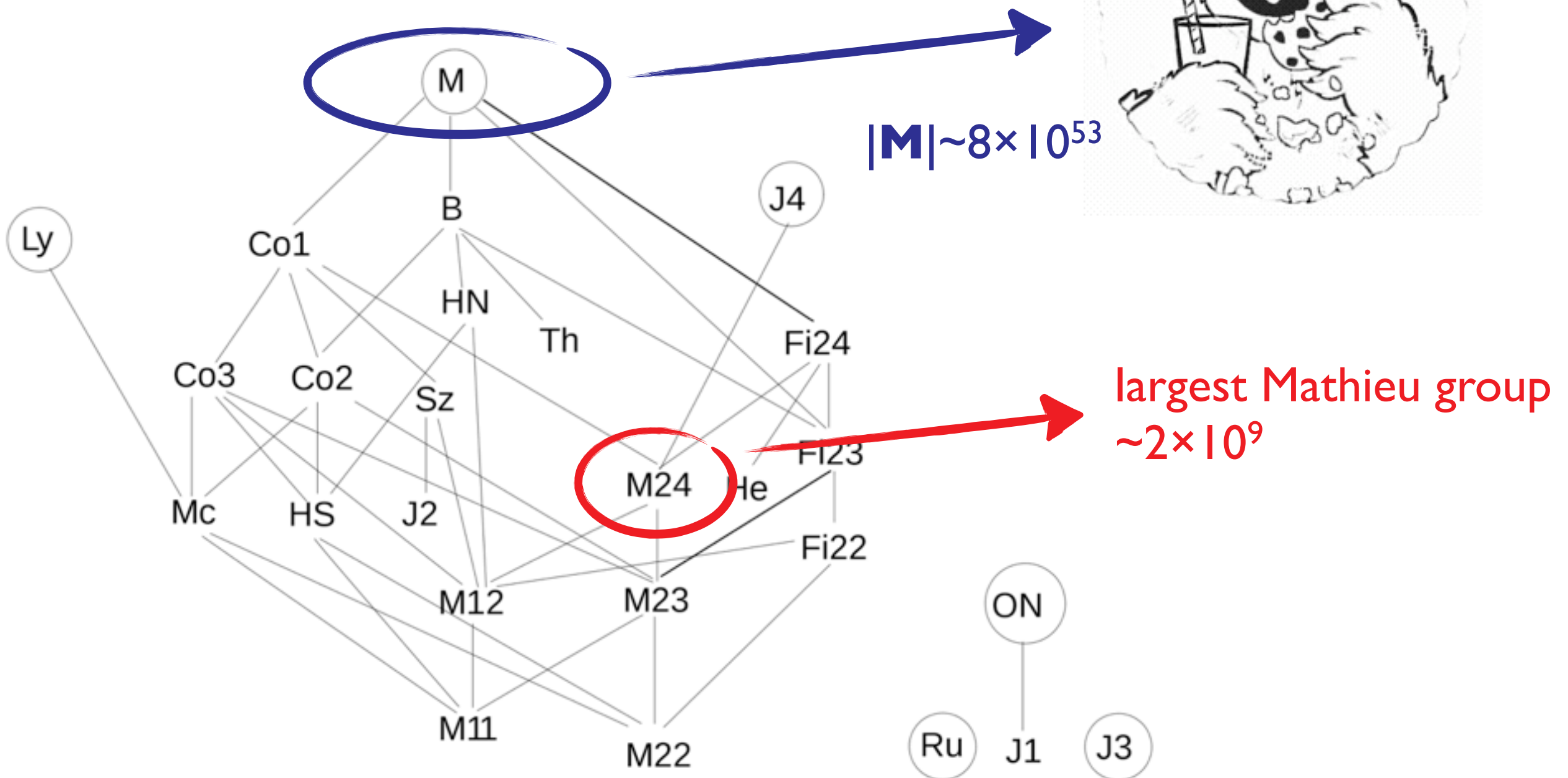
The Largest
Sporadic Group



The Most Natural
Modular Function

Sporadic Groups

The 26 finite simple groups that don't come in ∞ -families.



Mathieu 24

It has a natural 24-dim representation, on which

$$M_{24} \subset S_{24}.$$

$$\begin{array}{c} N : \text{a 24-dim Niemeier lattice} \\ \curvearrowright \\ M_{24} \end{array}$$

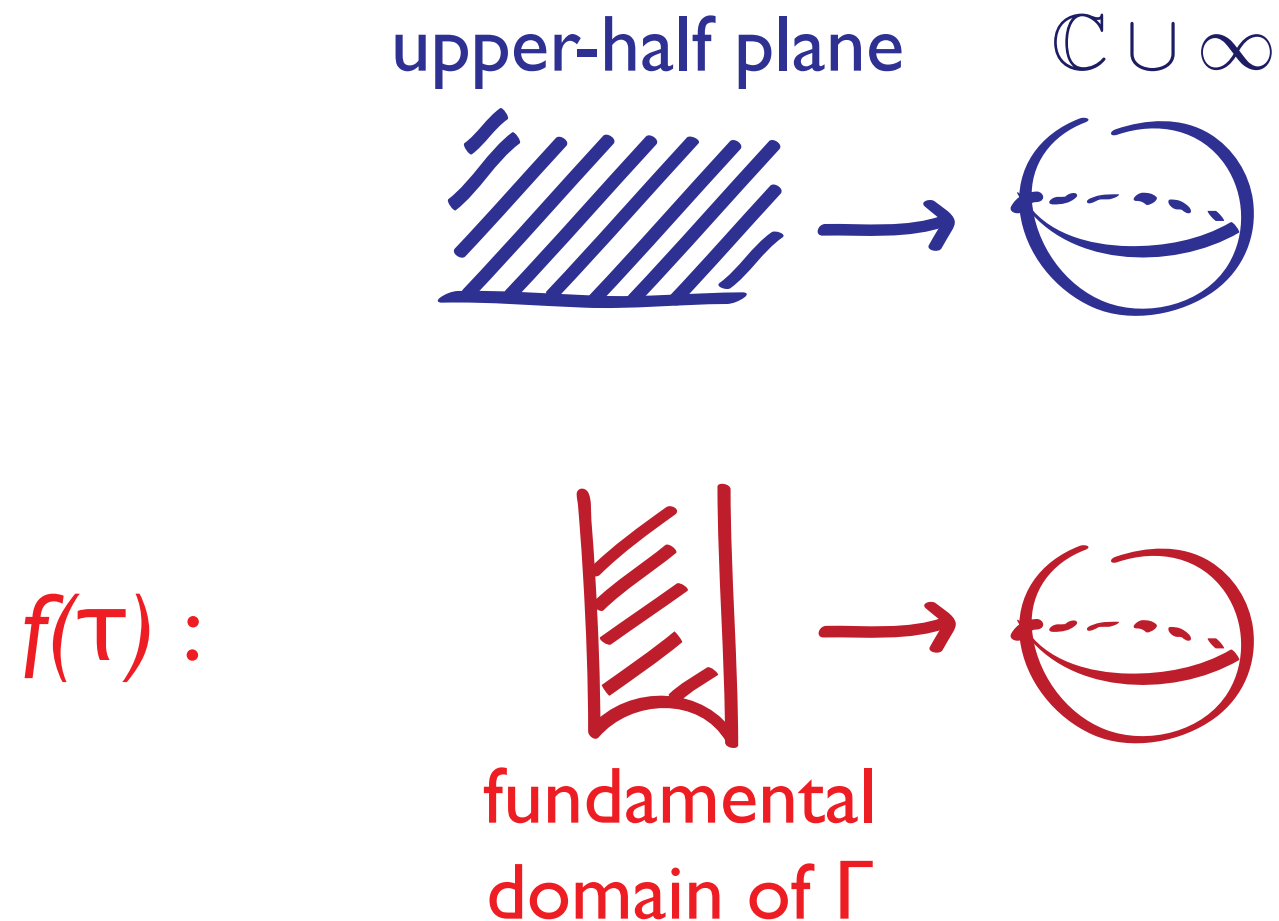
All symmetries G of $K3$ manifold preserving the hyper-Kähler structure satisfy

$$G \subset M_{23} \subset M_{24}.$$

[Mukai '88, Kondo '98]

Modular \subset Automorphic Forms

A modular form $f(\tau)$ transforms “covariantly” under a subgroup Γ of $SL(2, \mathbb{R})$.



An “automorphic form” transforms under more general groups (a higher dimensional version of a modular form).

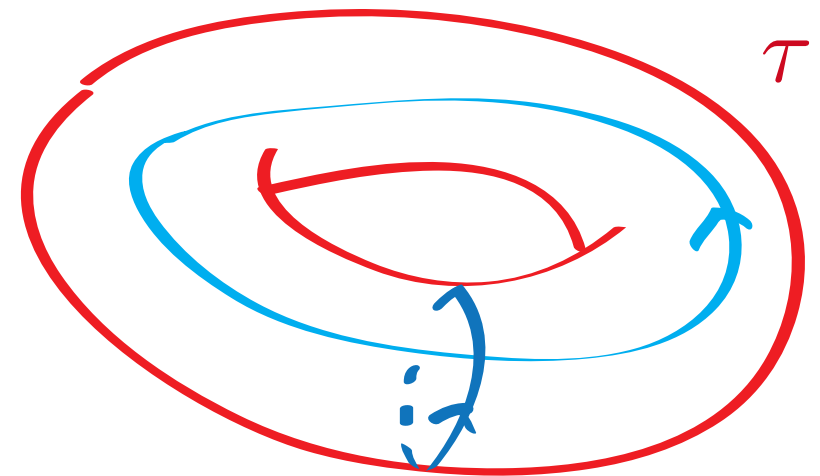
In physics, this “automorphism” reflects an underlying symmetry of the problem.

String theory is good at producing automorphic forms!

World-sheet Symmetries
(mapping class group of Σ)

eg. $Z_{2d} \text{CFT}(\tau)$ transforms under $SL(2, \mathbb{Z})$.

Space-Time Symmetries
(Such as T -, S -dualities)



All symmetries have to be reflected in suitable partition functions.

String Theory and Moonshine

*e.g. Monstrous moonshine (partially) explained by CFT
and proven using bosonic string theory.*

$$J(\tau) = q^{-1} + 196884 q + 21493760 q^2 + \dots$$
$$= \text{Tr}_{\mathcal{H}} q^{L_0 - \frac{c}{24}} \left(\begin{array}{l} \text{partition function of a } c=24 \text{ 2d} \\ \text{chiral CFT with a Monster symmetry} \end{array} \right)$$

[’88 Frenkel-Lepowsky-Meurmann
/Dixon-Ginsparg-Harvey]

The “moonshine conjecture” was proven by introducing
generalised Kac-Moody algebras and considering the lift of $J(\tau)$
into an automorphic form.

[R. Borcherds ’92]

Physically, this “lift” corresponds to considering the
full 26-dim bosonic string.

spectrum generating
algebra

Borcherds-
Kac-Moody
Algebra

symmetry

denominator

Sporadic
Groups

Automorphic
Forms

symmetry

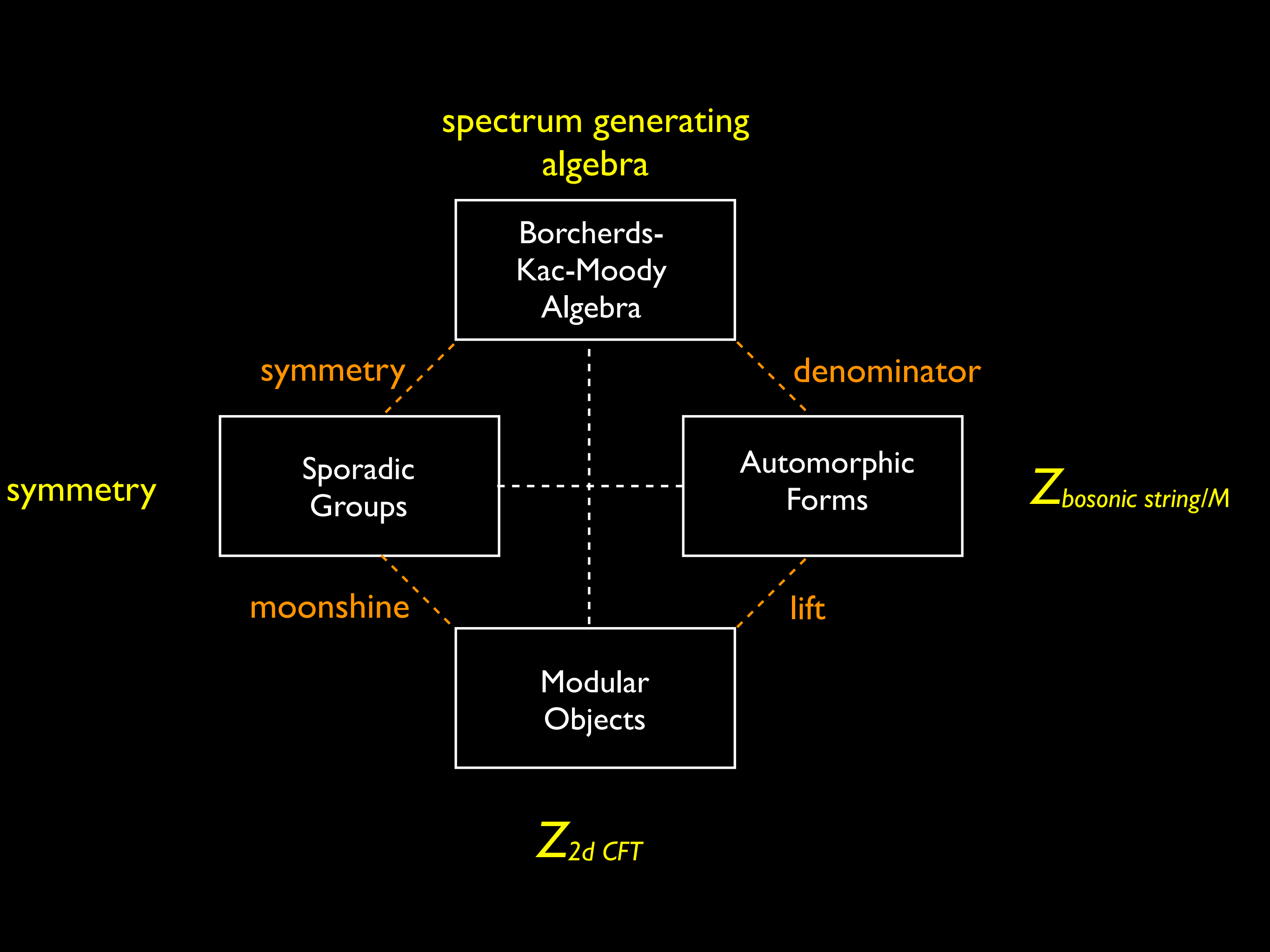
$Z_{\text{bosonic string}/M}$

moonshine

lift

Modular
Objects

$Z_{2d \text{ CFT}}$



K3 and M₂₄

Last year, Eguchi-Ooguri-Tachikawa made a striking observation on the BPS-sector of the 2d SCFT with *K3* target:

$\mathcal{Z}(\tau, z; K3) = K3$ elliptic genus

[Eguchi-Ooguri-Taormina-Yang '89]

$$= \frac{\theta_1^2(\tau, z)}{\eta^3(\tau)} \left(24 \mu(\tau, z) + 2q^{-1/8} \left(-1 + \textcircled{45}q + \textcircled{231}q^2 + \textcircled{770}q^3 + \dots \right) \right)$$

number of massive $N=4$ multiplets

also dimensions of irreps of M_{24} !!!

A New M_{24} Moonshine?

WHY?

WHAT?



- What precisely should we conjecture? In particular, what should be the moonshine diagram in this case?
- What are the consequences? Are they consistent with what we know?
- What are the evidences?

(First part)

A Long-Standing Puzzle

Along the way we were confronted with an old puzzle of a number theoretic nature (the genus zero property).

Surprisingly, considerations from AdS/CFT give a way out by providing natural explanations for some crucial properties of the modular groups appearing in all moonshines known so far.



Holographic Modularity of Moonshines

(Second part)

K3 and M₂₄

K3 Sigma-model

In a 2d $N=(4,4)$ SCFT, BPS states are counted by elliptic genus.

They are modular and are called Jacobi forms.

$$\begin{aligned}\mathcal{Z}(\tau, z; K3) &= \text{Tr}_{\mathcal{H}_{\text{RR}}} \left((-1)^{J_0 + \bar{J}_0} y^{J_0} q^{L_0 - c/24} \bar{q}^{\bar{L}_0 - c/24} \right) = \sum_{\substack{n \geq 0, \\ \ell \in \mathbb{Z}}} c(4n - \ell^2) q^n y^\ell \\ &= \frac{\theta_1^2(\tau, z)}{\eta^3(\tau)} \left(24 \mu(\tau, z) + 2q^{-1/8} (-1 + 45q + 231q^2 + 770q^3 + \dots) \right)\end{aligned}$$

Counting the number of massive $N=4$ multiplets:

$$H(\tau) = 2q^{-1/8} (-1 + 45q + 231q^2 + 770q^3 + \dots)$$

Mock Modular Form

[Zwegers '02/
Eguchi-Hikami'10]

1/4-BPS States in Type II/K3×T²

Further compactifying down to 4d is achieved by “Borcherds-lifting” the elliptic genus.

The resulting automorphic form Φ counts the 1/4-BPS states of the $N=4, d=4$ theory: [Dijkgraaf-Verlinde² '97]

$$\Phi(\tau, z, \sigma) = \frac{1}{pqy} \prod_{n,m,\ell} \left(\frac{1}{1 - p^n q^m y^\ell} \right)^{c(4nm - \ell^2)}$$

Fourier coeff. of $\mathcal{Z}(\tau, z; K3)$

$$= \sum_{P,Q} D(P, Q) p^{Q \cdot Q / 2} q^{P \cdot P / 2} y^{P \cdot Q}$$

This automorphic form also defines an algebraic structure underlying the 1/4-BPS spectrum.

*Dyonic 4d BH
charge (Q,P)
 $S_{BH} \sim \log D(P,Q)$*

See A. Sen, Strings '07, '08, '10

1/2-BPS States in Type II/K3×T²

$$\frac{1}{\eta^{24}(\tau)} = \sum d(Q) q^{\frac{Q \cdot Q}{2}} \Leftarrow \text{heterotic}/T^6$$

$$= \frac{1}{q \prod_{n \geq 1} (1 - q^n)^{24}}$$

[Dabholkar-Harvey '89]

= partition function of 24 chiral bosons

1/2-BPS Hilbert space of Type II/K3×T² forms a M_{24} -representation!

M_{24} maps 1/2-BPS states to each other.

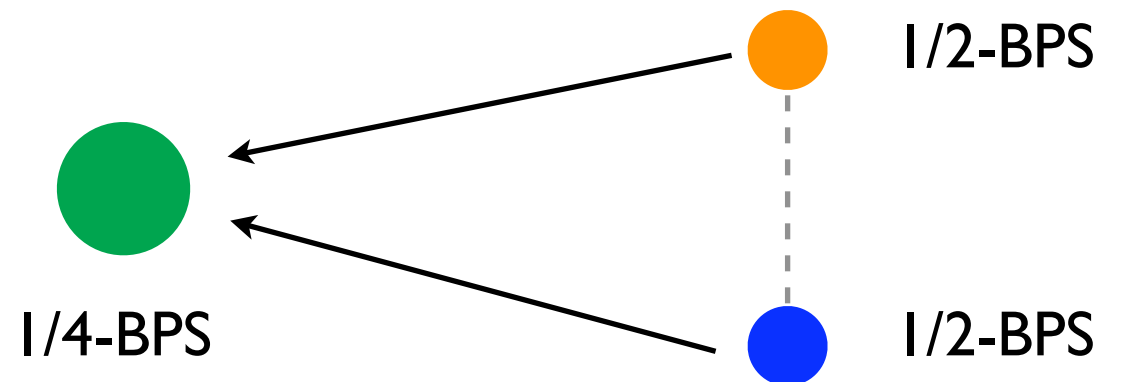
Dyonic Bound States in Type II/K3×T²

I/4-BPS spectrum has to know about the I/2-BPS spectrum too!

Two-Centered Bound States



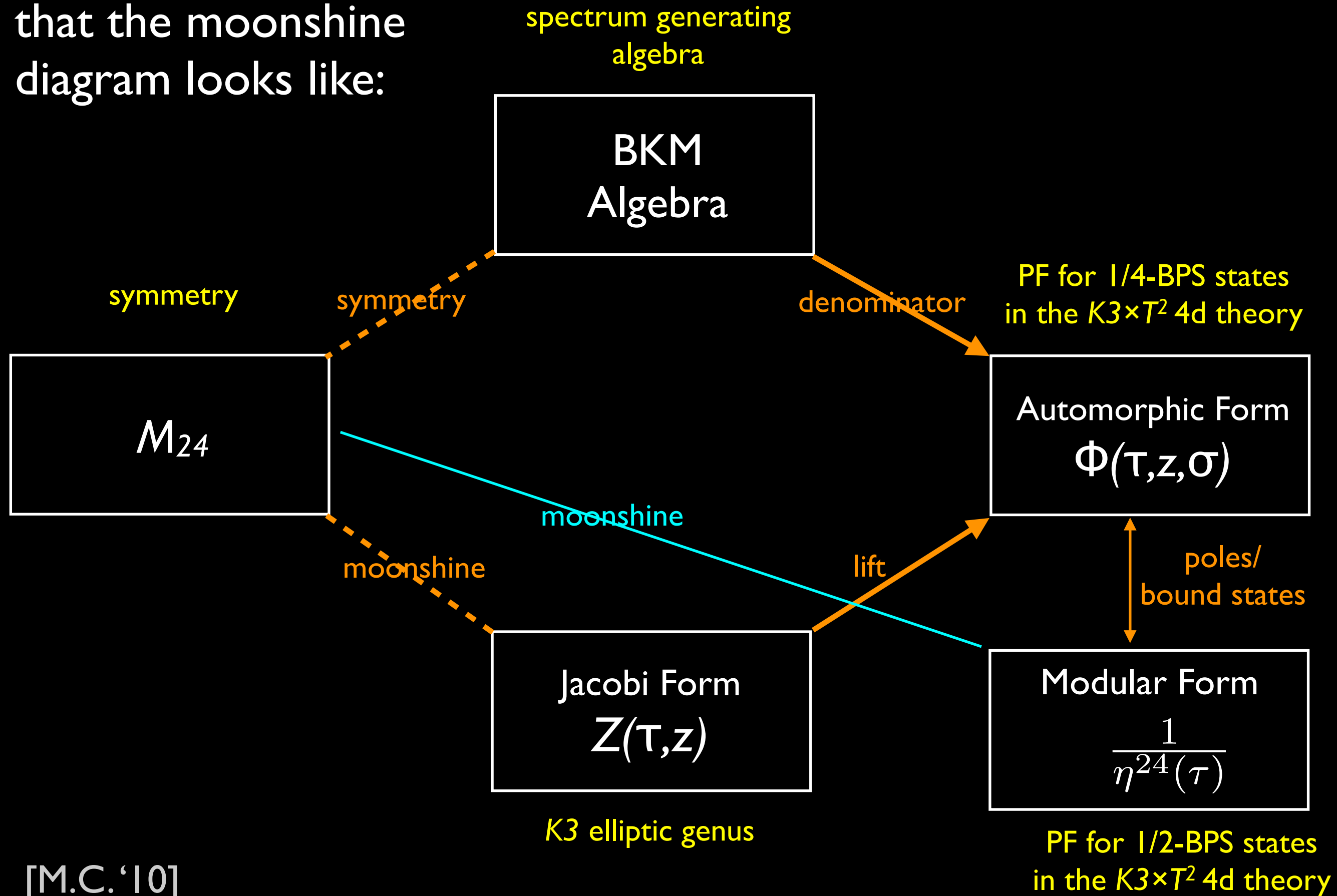
Poles in I/4-BPS P.F.



$$\Phi(\tau, z, \sigma) \xrightarrow{z \rightarrow 0} \frac{1}{(2\pi i z)^2} \frac{1}{\eta^{24}(\tau)} \frac{1}{\eta^{24}(\sigma)}$$

[A. Sen/M.C.-Verlinde '07, '08]

Physics suggests
that the moonshine
diagram looks like:



What exactly are we conjecturing?

The mock modular form containing the M_{24} information in $\mathcal{Z}(\tau, z; K3)$

$$H(\tau) = q^{-1/8} \left(-2 + 90q + 462q^2 + 1540q^3 + \dots \right)$$

$$\stackrel{?}{=} q^{-1/8} \left(-2 + \sum_{n=1}^{\infty} q^n \dim K_n \right) \quad K_n = M_{24}\text{-representation}$$

True. But meaningless.

Instead we need information about the vector (Hilbert) space!

Twisting the Symmetries

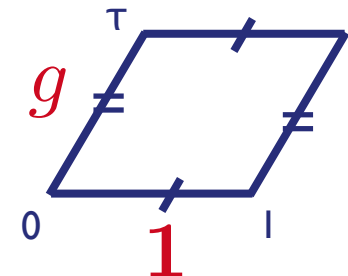
Hilbert space \mathcal{H}



The **twisted** partition function $Z_g = \text{Tr}_{\mathcal{H}}(g \dots)$ gives finer information about \mathcal{H} than just the dimension.

e.g. 2d CFT

$$Z_g(\tau) = \text{Tr}_{\mathcal{H}}(g q^{\hat{H}} \dots)$$



boundary condition: ~~$SL(2, \mathbb{Z})$~~ $\rightarrow \Gamma_g$

e.g. When there is supersymmetry, we can compute

$$\text{Tr}_{\mathcal{H}}(q^{L_0}) \quad \text{or} \quad \text{Tr}_{\mathcal{H}}((-1)^F q^{L_0})$$

PF

index (twisted PF)

The invariance group has to preserve the spin structure.

A Concrete Conjecture

In particular, we want to compute the **twisted** K3 elliptic genus $\mathcal{Z}_g(\tau, z; K3)$ and extract the interesting part $H_g(\tau)$.

Conjecture:

For all $g \in M_{24}$ It should be a (mock) modular form of a discrete subgroup $\Gamma_g \subset SL(2, \mathbb{R})$ and satisfy

$$H_g(\tau) = q^{-1/8} \left(-2 + \sum_{n=1}^{\infty} q^n \text{Tr}_{K_n} g \right)$$

Status: The candidate (mock) modular forms $H_g(\tau)$ have been proposed for all $[g] \in M_{24}$.

[M.C. / Gaberdiel-Hohenegger-Volpato/Eguchi-Hikami '10]

Checks:

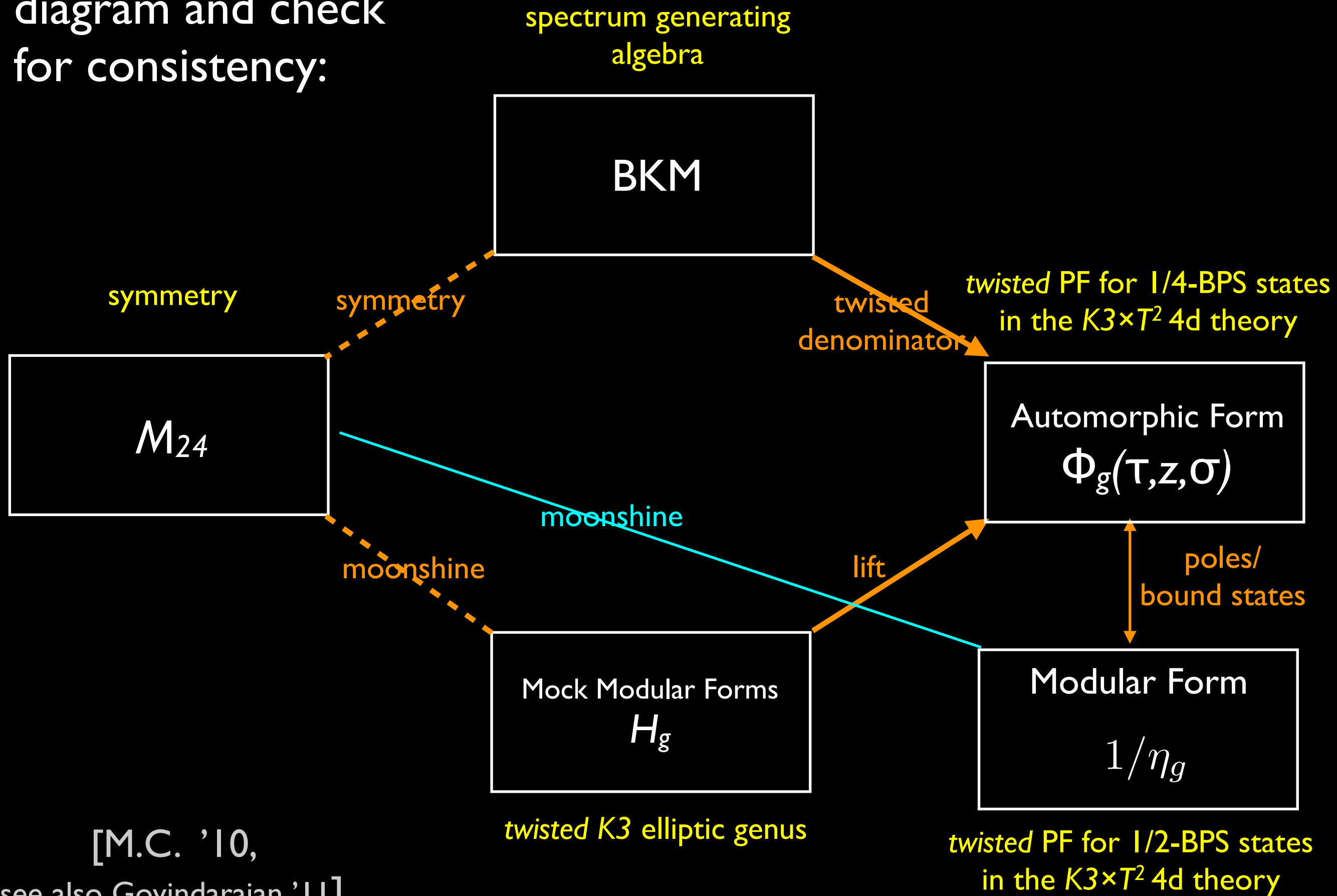
1) $K_{n < 1000}$ constructed!

[Y. Tachikawa]

2) Coincides with geometric calculation when the latter is available.

[David-Jatkar-Sen '06]

Twist the moonshine diagram and check for consistency:



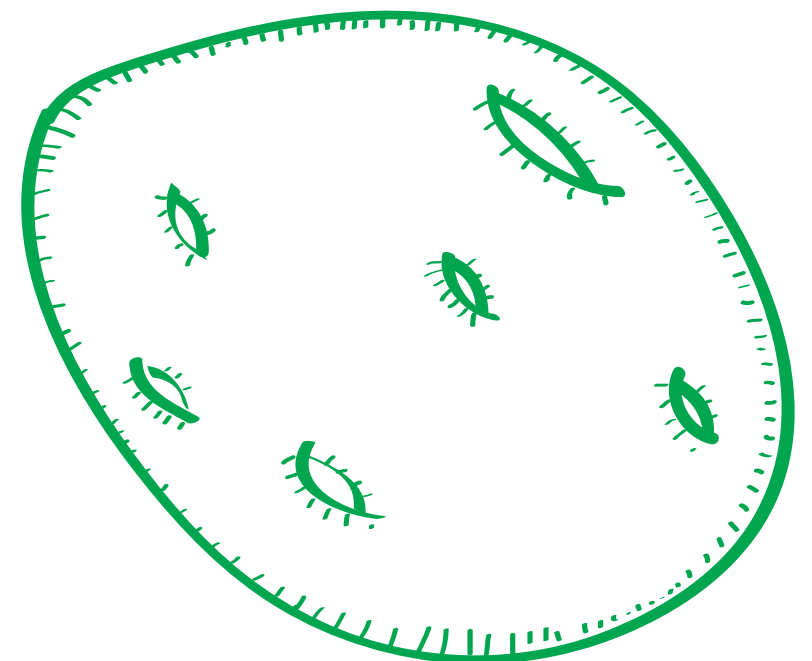
[M.C. '10,
see also Govindarajan '11]

Summary: M_{24} in K3 String Theories

We don't understand, but we are pretty sure!

Challenge:
Stringy K3 Symmetries

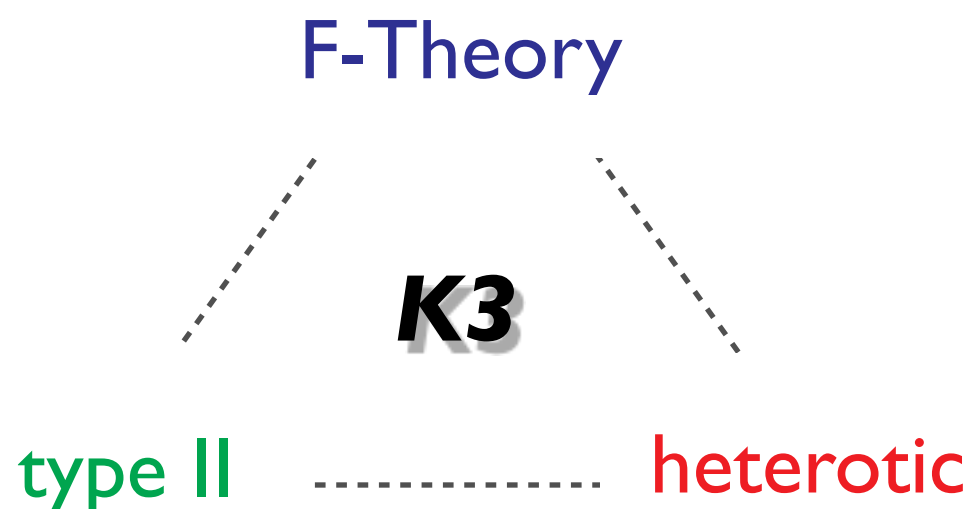
These results are clearly evidence for uncharted symmetries of the stringy geometries of K3. **Note: Classical symmetries are not enough! Stringy symmetries are needed.**



Summary: M_{24} in K3 String Theories

This moonshine is physically interesting.

- We have black holes.
- Since $K3$ is ubiquitous, it is important to understand these new symmetries.
- We should also explore where else M_{24} shows its presence.



Sporadic Quantum
Black Holes

Holographic Modularity of the Moonshines

with **John Duncan** 1107.xxxx

A Long-Standing Mystery

For positive every integer n we can define a certain subgroup $\Gamma_0(n)^+ \subset SL(2, \mathbb{R})$.

Q: What are the prime numbers p such that $\Gamma_0(p)^+$ is **genus zero**?

A: $\Gamma_0(p)^+$ is genus zero iff p divides $|\mathbf{M}|$!



[Ogg '73]

WHY??



WHY??

An “Explanation” from Monstrous Moonshine

$$J(\tau) = q^{-1} + 196884 q + 21493760 q^2 + \dots$$

$$= \text{Tr}_{\mathcal{H}} q^{L_0 - \frac{c}{24}}$$

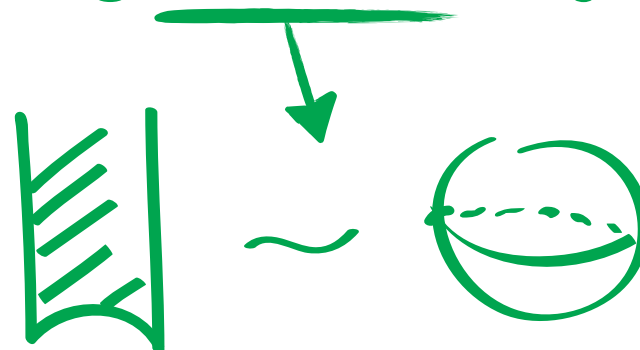


As before, we consider the twisted partition function

$$J_g(\tau) = \text{Tr}_{\mathcal{H}}(g q^{L_0 - \frac{c}{24}}) = \sum_{n \geq -1} c_g(n) q^n \quad \text{for every } g \in \mathbf{M}$$

Moonshine Conjecture (Conway-Norton '79):

$J_g(\tau)$ is invariant under some genus zero $\Gamma_g \subset SL(2, \mathbb{R})$.



Genus Zero Property

Genus zero groups $\Gamma \subset SL(2, R)$ are rare.

But only genus zero groups seem to appear in moonshines.



Not just in Monstrous moonshine, this has also been extended to

- The “generalised moonshine”
- Groups other than the Monster

see for instance
Norton '84/Carnahan '08
Höhn '03/Duncan '05, '06

BUT WHY GENUS ZERO??

NO Genus Zero for the New M_{24} Moonshine

Heresy!

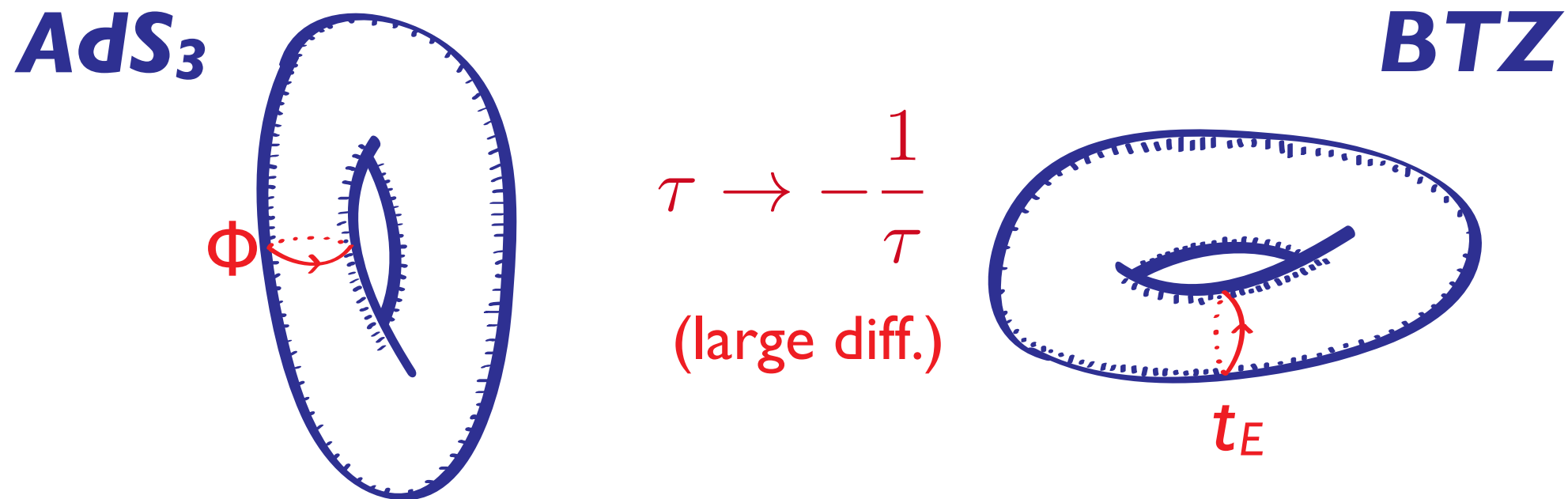
But true, by inspecting

$$H_g(\tau) = q^{-1/8} \left(-2 + \sum_{n=1}^{\infty} \underbrace{c_{g,n}}_{\downarrow} q^n \right)$$

$\ell \backslash [g]$	1A	2A	2B	3A	3B	4A	4B	4C	5A	6A	6B	7A	$\overline{7A}$	8A	10A	11A
1	90	-6	10	0	6	-6	2	2	0	0	-2	-1	-1	-2	0	2
2	462	14	-18	-6	0	-2	-2	6	2	2	0	0	0	-2	2	0
3	1540	-28	20	10	-14	4	-4	-4	0	2	2	0	0	0	0	0
4	4554	42	-38	0	12	-6	2	-6	-6	0	4	4	4	-2	2	0
5	11592	-56	72	-18	0	-8	8	0	2	-2	0	0	0	0	2	-2
6	27830	86	-90	20	-16	6	-2	6	0	-4	0	-2	-2	2	0	0
7	61686	-138	118	0	30	6	-10	-2	6	0	-2	2	2	-2	-2	-2
8	131100	188	-180	-30	0	-4	4	-12	0	2	0	-3	-3	0	0	2
9	265650	-238	258	42	-42	-14	10	10	-10	2	6	0	0	-2	-2	0

AdS_3/CFT_2

Recall: In a Euclidean 3d gravity, the only smooth solution with aAdS boundary conditions and torus boundary is the **solid torus**:



⇒ Saddle points of the gravity path integral are labeled by the **contractible cycle**.

$$\gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_\infty \setminus SL(2, \mathbb{Z}) \quad (\tau \sim \tau + 1)$$

[Maldacena-Strominger '98]

Twisting the boundary condition: $SL(2, \mathbb{Z}) \rightarrow \Gamma_g$.

AdS_3/CFT_2

Assuming a CFT has a dual description given by **semi-classical-like AdS gravity**

⇒ The (twisted) partition function $Z_g(\tau)$ can also be computed from the gravity side by summing over saddle point contributions

$$\sum_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_\infty \setminus \Gamma_g} f\left(\frac{a\tau + b}{c\tau + d}\right)$$

⇒ $Z_g(\tau)$ has to be **Rademacher-summable!**

See also the “Farey Tail” papers:

Dijkgraaf-Maldacena-Moore-Verlinde '00,

Kraus-Larsen/Dijkgraaf-de Boer-M.C.-Manschot-Verlinde/Denef-Moore/Manschot-Moore '06

Rademacher-Summability

$$Z_g(\tau) = \text{Reg} \left[\sum_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_\infty \setminus \Gamma_g} f\left(\frac{a\tau + b}{c\tau + d}\right) \right] \quad \boxed{\text{convergent, anomaly-free}}$$

$$e.g. \quad J(\tau) = e(-\tau) + \lim_{K \rightarrow \infty} \sum_{\substack{0 < c < K \\ -K^2 < d < K^2}} e\left(-\frac{a\tau + b}{c\tau + d}\right) - e\left(-\frac{a}{c}\right) - 12 \quad , \quad e(x) = e^{2\pi i x}$$

$$= \lim_{K \rightarrow \infty} \sum_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in (\Gamma_\infty \setminus SL(2, \mathbb{Z}))_{<K}} \left(q^{-1} \Big|_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \cdot \tau} - q^{-1} \Big|_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \cdot i\infty} \right) - 12$$

$$= \text{Reg} \left[\sum_{\gamma \in (\Gamma_\infty \setminus SL(2, \mathbb{Z}))} q^{-1} \Big|_{\gamma \cdot \tau} \right]$$

Rademacher-Summability

puts stringent constraints on Γ_g .

e.g. If we want to have a **modular function** (weight 0, weakly holomorphic),

Rademacher-summability

\Leftrightarrow

Γ_g is genus zero

[Duncan-Frenkel '09]

In particular, $J_g(\tau) = \text{Reg} \left[\sum_{\gamma \in (\Gamma_\infty \setminus \Gamma_g)} q^{-1} \Big|_{\gamma \cdot \tau} \right]$ for all $g \in \mathbf{M}$!

Note: this holds for higher central charge cases (q^n) as well.

Rademacher-Summability of the M_{24} Moonshine

Here we have $H_g(\tau)$ = weight 1/2 Mock modular form.

Rademacher-summability ~~\iff~~ Γ_g is genus zero

Instead, with the help of previous results of Bringmann-Ono ('06) and Eguchi-Hikami ('09), we show

$$H_g(\tau) = -2 \operatorname{Reg} \left[\sum_{\gamma \in (\Gamma_\infty \setminus \Gamma_g)} q^{-1/8} \Big|_{\gamma \cdot \tau} \right] \quad \text{for all } [g] \in M_{24}.$$

To Summarise

We propose a new pattern
Rademacher Summability

$Z_g(\tau)$ = modular functions
e.g. Monster moonshine

$g=0$

$Z_g(\tau)$ = *mock* modular forms
for the new M_{24} moonshine

verified

Motivation for the new pattern:
All known theories of moonshine have a CFT interpretation.

Assuming the existence of a good dual description
(for a higher c version of the sporadic CFT)
 \Rightarrow All twisted PF have to be Rademacher-summable.



A glass of whiskey for physicists?

AdS/CFT $\rightarrow \rightarrow$ AdS/CMT, AdS/QCD,
AdS/Hydrodynamics..... +AdS/NT??

Thank You!

To Summarise

Our result raises many physical questions:

- For M_{24} moonshine the dual situation is (more or less) clear: $AdS_3 \times S^3 \times K3$. How about other sporadic CFT's? Our results can be interpreted as an evidence that there exists 2d chiral CFT's with sporadic symmetries with semi-classical-like (saddle-dominated) AdS_3 duals.
- Relatedly, this brings back the question: What are the criteria for a CFT to have a (weakly coupled) AdS_3 duals?

[see for instance, Heemskerk-Penedones-Polchinski-Sully '09, El-Showk-Papadodimas '11]

- How to understand the twisting from a bulk viewpoint in general?
[some progress reported in A. Sen Strings 2011]