Toda Theory From Six Dimensions

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Strings, Princeton June 25th, 2014 Daniel Jafferis & C.C. ArXiv: 1406.XXXX

General Motivation

• Reductions of 6d (2,0) ---> geometric perspective on SQFTs

– s-dualities, mirror symmetries, ...

Supersymmetric localization ---> new partition functions

- instanton sums, sphere partition functions, indices, ...

 Fusion ---> novel interpretations of partition functions for QFTs with 6d parent

Specific Motivation

6d (2,0) on $S^4 \times \Sigma$

6d conformal invariance + twisting on Σ + supersymmetry

Ζ

---> Z independent of $vol(\Sigma)$ and $vol(S^4)$

Specific Motivation

6d (2,0) on S⁴×Σ ~

[Witten, Gaiotto]

[Nekrasov, Pestun]

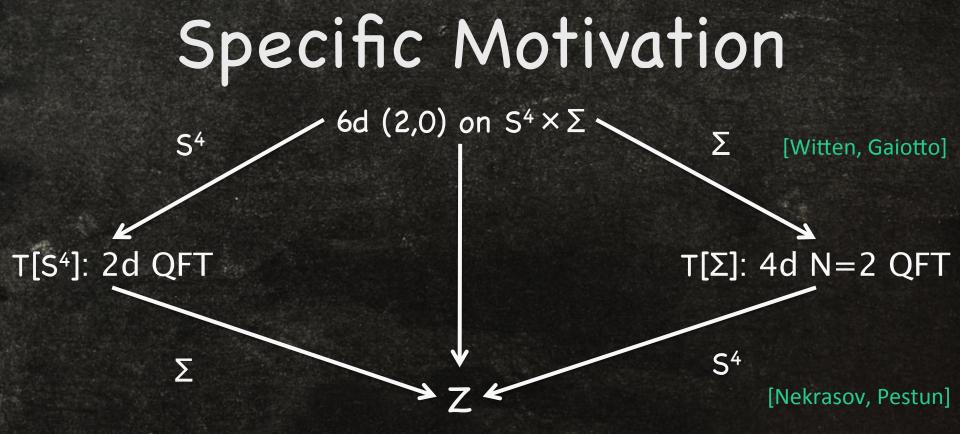
T[Σ]: 4d N=2 QFT

Σ

S⁴

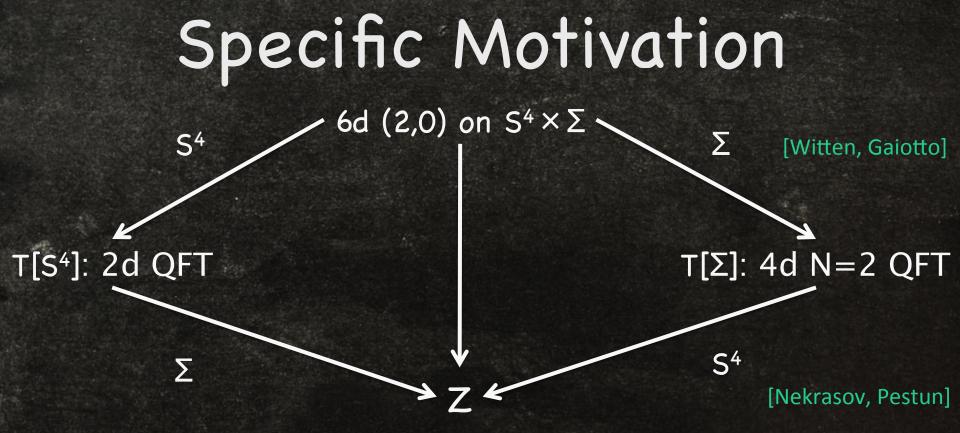
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6d conformal invariance + twisting on Σ + supersymmetry

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AGT conjecture: $T[S^4] = Toda \ CFT \longrightarrow N-1 \ real \ scalars \ \Phi^i$

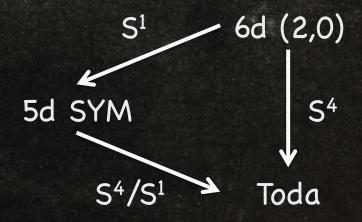
 $L \sim \Sigma_{ij} C_{ij} d_{\mu} \Phi^{i} d_{\mu} \Phi^{j} - \Sigma_{i} \exp(\frac{1}{2} \Sigma_{j} C_{ij} \Phi^{j})$ ($C_{ij} = SU(N)$ Cartan matrix)

Result & Method

Result: derivation of Toda theory via reduction from 6d

Method: factorize reduction:

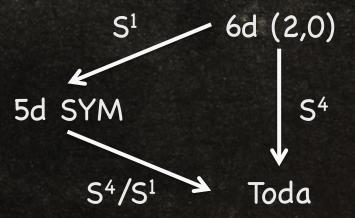
[Kim-Kim-Kim, Fukada-Kawana-Matsumiya, Lee-Yamazaki, Jafferis-C.C.]



Result & Method

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 Method: factorize reduction:
 [Kim-Kim, Fukada-Kawana-Matsumiya, Lee-Yamazaki, Jafferis-C.C.]
 why it works:



- Higher derivative corrections to 5d SYM unimportant

- suppressed by small r(S¹)
- Q exact

- S^4/S^1 not smooth but at singularity, $g^2_{ym} \sim r(S^1) \rightarrow 0$

---> understand with weakly coupled 5d physics

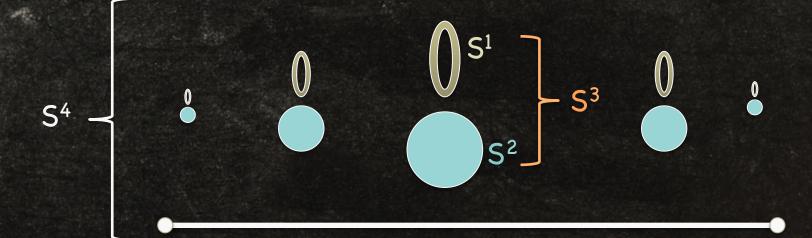


- Compactification on $S^4 \times R^{1,1}$ has OSP(2|4) symmetry
- 5d SUSY ---> reduce on Hopf circle of equatorial S³ in S⁴

S⁴ Geometry

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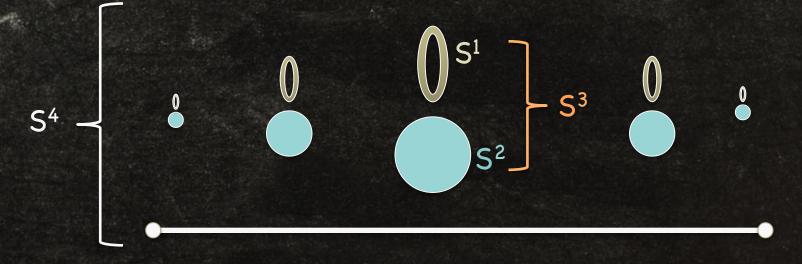
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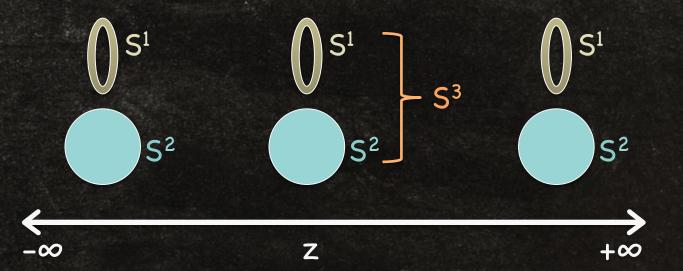


Use 6d Weyl invariance to stretch interval to ∞ length

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Use 6d Weyl invariance to stretch interval to ∞ length
 ---> S³ now constant radius
 ---> ds² = (dΩ₃)² + dz² + cosh²(z/r)(-dt²+dx²)

Plan of Attack

- Reduce from 6d (2,0) to 5d SYM on Hopf circle of S^3
- Reduce 5d SYM on S^2 with one unit of RR-flux
- Place resulting 3d theory on manifold R^{1,2} with non-trivial metric:

$$ds^2 = dz^2 + \cosh^2(z/r)(-dt^2 + dx^2)$$

- Add suitable boundary conditions at $|z| = \infty$
- Determine effective boundary theory

Relation To Chern-Simons

 5d SYM on S² with 1 unit of RR-flux ---> complex CS in 3d [Lee-Yamazaki, Jafferis-C.C.]

- complex SL(N,C) gauge field $\mathcal{B} = A + i X$

 $- L = \frac{1}{8\pi} \left[\operatorname{Tr} \left(\mathcal{B} \, \mathrm{d}\mathcal{B} + \frac{2}{3} \, \mathcal{B}^3 \right) + \operatorname{Tr} \left(\overline{\mathcal{B}} \, \mathrm{d}\overline{\mathcal{B}} + \frac{2}{3} \, \overline{\mathcal{B}}^3 \right) \right]$

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 Puzzle: How does SUSY reduction of SU(N) gauge theory result in bosonic SL(N,C) gauge theory?

• Answer: the two concepts are equivalent!

SU(N) covariant Lorenz gauge condition D^aX_a = 0
 breaks SL(N,C) to SU(N)

- fermions reinterpreted as Faddeev-Popov ghosts

Boundary Data – One Side

D6

D4

From IIA perspective --> D4 ending on D6

- scalars have a Nahm pole X_a ~ T_a /w
 [Diaconescu]
- T_a valued in SU(2), $[T_a, T_b] = \varepsilon_{abc} T_c$
- A chosen so that SL(N,C) field, \mathcal{B} , is flat
- ---> $\mathcal{B} = (iT_3) dw/w + (T_+) dx_+/w$
- Fermions lifted by Dirichlet condition

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Fermions lifted by Dirichlet condition

The terms ... are less singular in w, and are fluctuating fields.
 They give rise to a chiral Toda theory

Executive Summary

Question: What are the Toda fields?

• Answer:

The Toda fields are modes of the 5d scalars X_a localized at the poles of the S⁴

[Nekrasov-Witten]

Map to Toda

CS Theory --> boundary theory of currents (WZW-model) [Witten]

 $\mathcal{B} = F^{-1} d F + F^{-1} (H^{-1} dH) F$ (pure gauge HF)

F is background giving Nahm pole, H is dynamical

$\begin{array}{l} & \textbf{Map to Toda} \\ & \textbf{CS Theory --> boundary theory of currents (WZW-model)} \\ & \textbf{B} = F^{-1} d F + F^{-1} (H^{-1} dH) F & (pure gauge HF) \end{array}$

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<u>Properties of H</u>:

- Gauge $D^aX_a = 0 \rightarrow H$ is SU(N) valued not SL(N,C) valued

- $H = H(x_{+}, x_{-})$ depends only on boundary coordinates

- Flatness of $H^{-1} dH \rightarrow H = H(x_{+})$ is chiral

- Regularity of ... --> H fixed by N-1 real scalars = Toda fields

Map to Toda

 More Briefly: Nahm boundary conditions provide constraints on WZW currents which reduce it to Toda [Balog-Fehér-Forgács-O'Raifeartaigh-Wipf]

 Each boundary (region near a pole in S⁴) gives a chiral half of Toda. Together they form the full non-chiral Toda.

[Elitzer-Moore-Schwimmer-Seiberg]

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 <u>Central Charge</u>: [Elitzer-Moore-Schwimmer-Seiberg]

• Toda central charge, $c = N-1 + N(N^2-1)(b+b^{-1})^2$. S⁴ gives b = 1

 $() S^{1}$

squashed S³

Recover b by squashing geometry

Future Directions

Understand the dictionary between Toda operators, and 6d defect operators

 Use similar techniques to study 6d (2,0) on other geometries. An interesting case is S⁶ which should lead to direct information about 6d correlation functions

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Thanks for Listening!