Revisiting soliton contributions to perturbative processes

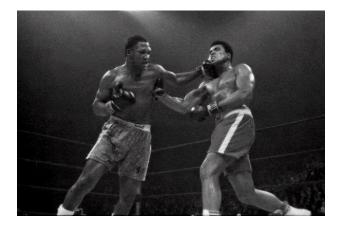
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Strings, Princeton, June 26, 2014

Based on 1403.5017 and 1404.0016 with C. Papageorgakis

Gong Show, Round 16



Even the "Fight of the Century" went only 15 rounds

Q & A

- Q: Do solitons run in loops?
- A: Yes, in the following sense:

$$2 \operatorname{Im} \left\{ \underbrace{\xrightarrow{k}}_{k} \xrightarrow{k} \right\} = \sum_{f} \int d\Pi_{f} \left| \underbrace{\xrightarrow{k}}_{k} \xrightarrow{f} \right|^{2}$$

- Q:
 - Can we compute their contribution to perturbative processes?
 - Is this contribution "exponentially suppressed"? (in what quantity?)¹

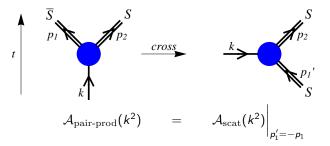
A:

- Yes, under certain assumptions.
- Usually, but not necessarily. (in the ratio R_c/R_s .)

¹Drukier and Nussinov (1982)..., Demidov and Levkov (2011), Banks (2012)

How do we compute?

Use analyticity and crossing symmetry:



compute \mathcal{A}_{scat} perturbatively in soliton sector:

$$\delta(k + p_1' - p_2)\mathcal{A}_{\text{scat}}(k^2) = \langle S(\mathbf{p}_2) | T\{e^{-i\int dt H_l}\} | \mathbf{k}, S(\mathbf{p}_1') \rangle$$

• consider class of scalar models $\phi(x) = \phi_S\left(\frac{\mathbf{x}-\mathbf{x}_0}{R_S};\ldots\right)$

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• find $A_{pair-prod}(k^2)$ vanishes faster than any power in the ratio R_c/R_S for k^2 above[†] threshold $(2M_S)^2$

[†]need to understand threshold effects better

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$$\mathcal{A}_{ ext{pair-prod}}(k^2) \lesssim e^{-2R_{\mathcal{S}}/R_c}$$
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• *provided* R_S bounded away from zero (as a function on the moduli space of soliton solutions)

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Possible lessons for 5D MSYM

- first take off the EFT glasses and suppose that 5D MSYM is a microscopic theory
- \Rightarrow integrate over all loop momenta (*i.e.* $\Lambda_{UV} \rightarrow \infty$)
- but then, for $k^2 > (2M_S)^2$, will produce $S \overline{S}$ pairs
- $R_S = \rho \rightarrow 0$ so argument for exponential suppression breaks down, suggesting soliton contributions may compete
- unfortunately approx. scheme also breaks down...need new methods

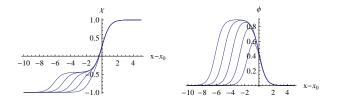
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Thank you!

example with internal modulus³

$$\begin{split} \mathcal{L} &= \frac{1}{2} (\partial_{\mu} \phi \partial^{\mu} \phi + \partial_{\mu} \chi \partial^{\mu} \chi) - \frac{1}{2} (W_{\phi}^{2} + W_{\chi}^{2}) \\ W &= \chi - \frac{1}{3} \chi^{3} - \chi \phi^{2} - \frac{\beta}{3} \phi^{3} \\ \Rightarrow \phi(\chi) &= \frac{\beta \chi + b \sqrt{\beta^{2} + 4} - \sqrt{(\chi \sqrt{\beta^{2} + 4} + \beta b)^{2} + 4(b^{2} - 1)}}{2} , \\ \chi(\chi) &= \sqrt{\frac{b^{2} - 1}{\beta^{2} + 4}} \left[\frac{2 \tanh(\chi - \chi_{0}) + 2b}{\sqrt{b^{2} - 1}(\sqrt{\beta^{2} + 4} - \beta)} - \frac{\sqrt{b^{2} - 1}(\sqrt{\beta^{2} + 4} - \beta)}{2 \tanh(\chi - \chi_{0}) + 2b} - \frac{\beta b}{\sqrt{b^{2} - 1}} \right] \end{split}$$



³BNRT (1997); Brito and de Souza Dutra (2014)