Integrability in AdS₃/CFT₂

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based on work in collaboration with R. Borsato, O. Ohlsson Sax, B. Stefański jr. & A. Torrielli

see in particular arXiv:1403.4543 and 1406.2971







Integrability in AdS₃/CFT₂

AdS_3/CFT_2 holography

- Many interesting properties: CFT₂, black-hole physics, higher-spin theories, rich dual gauge theory flowing to SCFT...
- In string theory, we can obtain it from RR and/or NSNS fluxes.
- For pure-NSNS, CFT techniques can be used on the worldsheet.

[Maldacena, Ooguri '00]

• RR fluxes are problematic in this approach.

Integrability and massless modes

• Classical integrability for maximally supersymmetric backgrounds

 $\mathsf{AdS}_3\times\mathsf{S}^3\times\mathsf{T}^4\qquad\text{and}\qquad\mathsf{AdS}_3\times\mathsf{S}^3\times\mathsf{S}^3\times\mathsf{S}^1$

with pure-RR and mixed flux.

[Babichenko, Stefański, Zarembo '10] [Cagnazzo, Zarembo '13]

- \bullet We would like quantum integrability, like for $AdS_5 \times S^5.$
- However, **massless modes** seemingly spoil usual approach. Integrable massless scattering can be subtle.

[Zamolodchikov, Zamolodchikov '92] [Fendley, Saleur '93]

 \longrightarrow major obstacle for integrability.

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Pure-RR AdS $_3 \times S^3 \times T^4$ light-cone symmetries

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\mathfrak{psu}(1,1|2)_{\mathsf{L}} \oplus \mathfrak{psu}(1,1|2)_{\mathsf{R}}
\downarrow (\mathsf{l.c. gauge})
\mathfrak{psu}(1|1)^4_{\mathsf{centr.ext.}}
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$$\downarrow (\mathsf{I.c. gauge})$$

$$\mathfrak{psu}(1|1)_{\mathsf{centr.ext.}}^{4}$$

Central charges:

H Hamiltonian, M Mass,

$$\mathbf{C} = +i \frac{h}{2} (e^{+i\mathbf{P}} - 1), \qquad \mathbf{\bar{C}} = -i \frac{h}{2} (e^{-i\mathbf{P}} - 1)$$

Pure-RR $AdS_3 \times S^3 \times T^4$ light-cone symmetries

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Exact statements about massless modes

• All one-particle representations, including massless are short

$$\mathbf{H}^2 = \mathbf{M}^2 + 4\mathbf{C}\mathbf{ar{C}}$$
 .

Masslessness is protected at all-loops

 $\mathbf{M} | \mathsf{massless} \rangle = 0$ protected by $\mathfrak{su}(2) \subset \mathfrak{so}(4)$.

• All-loop massless dispersion relation and group velocity

$$E(p) = \pm 2h \sin \frac{p}{2}, \qquad v(p) = \frac{\partial E}{\partial p} = \pm h \cos \frac{p}{2}.$$

The complete all-loop S matrix

• Write down irreducible representations of symmetries,

• = massive,
$$\circ$$
 = massless.

Impose invariance

$$\begin{bmatrix} \mathbf{S}, \mathbf{Q} \end{bmatrix} = \mathbf{0},$$

and find

$$\mathbf{S} = \left(\begin{array}{cc} \mathbf{S}^{\bullet\bullet} & \mathbf{S}^{\bullet\circ} \\ \mathbf{S}^{\circ\bullet} & \mathbf{S}^{\circ\circ} \end{array} \right)$$

• Yang-Baxter equation holds.

Results and outlook

• For pure-RR AdS $_3 \times S^3 \times T^4$, complete exact S matrix was found. [Borsato, Ohlsson Sax, Stefański, AS '14]

- Validation: world-sheet perturbative calculations up to two loops.
 [Sundin, Wulff '12] [Beccaria, Levkovich-Maslyuk, Macorini, Tseytlin '12] [Abbott '13]
 [Engelund, McKeown, Roiban '13] [Babichenko, Dekel, Ohlsson Sax '14] [...]
- Mixed fluxes and $AdS_3 \times S^3 \times S^3 \times S^1$: massive sector already known, full S matrix should follow similarly.

[Borsato, Ohlsson Sax, AS '12] [Hoare, Stepanchuk, Tseytlin '13]