

# String field theory vertex from integrability

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and work in progress...

# Outline

**Motivation for the physical problem**

**Motivation for our approach**

Lessons from the spectral problem

Lessons from form factors

**Light-cone String Field Theory Vertex – the pp-wave**

**Functional equations for the string vertex**

**What happens in  $AdS_5 \times S^5$ ?**

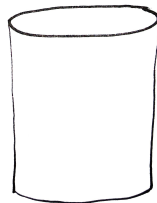
**Conclusions**

# Motivation

- ▶ We have a very good understanding of **the spectrum** of a string on  $AdS_5 \times S^5$ 
  - $\equiv$  quantized energy levels of a string for any value of  $\lambda \equiv g_{YM}^2 N_c$
  - $\equiv$  equivalently (anomalous) dimensions of operators in  $\mathcal{N} = 4$  SYM
- ▶ This is due to **the integrability** of the worldsheet theory

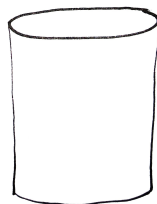
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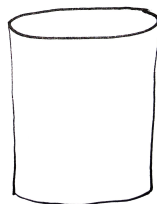
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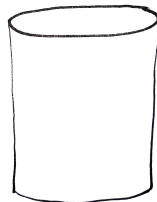
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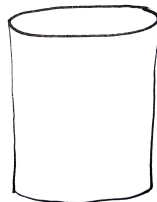
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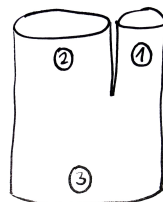
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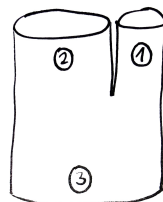
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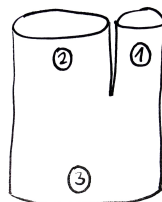
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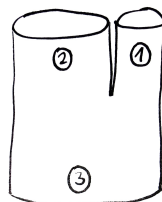
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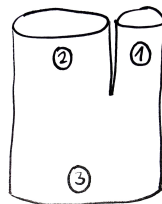
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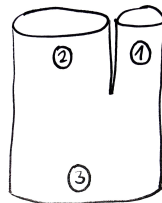
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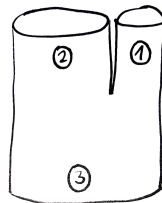
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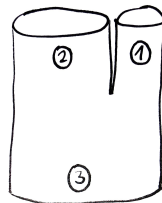




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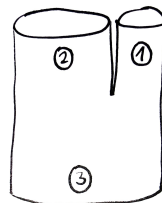
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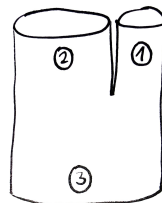
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Of course all this has direct AdS/CFT motivation...

Energy levels of a single string  $\equiv$  Anomalous dimensions

String interactions  $\longrightarrow$  OPE coefficients  
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1) solve the worldsheet theory on an infinite plane

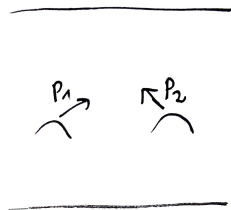
1. Particle momenta/rapidities can be analytically continued into the complex plane...
2. We get **crossing** equation + unitarity
3. ... and symmetry + Yang-Baxter equation
4. determines analytically the **S-matrix**

Works equally for relativistic and  $AdS_5 \times S^5$  case...

**Key role** of the infinite plane  $\rightarrow$  only there do we have crossing+analyticity which allows for solving for the S-matrix (functional equations for the S-matrix)

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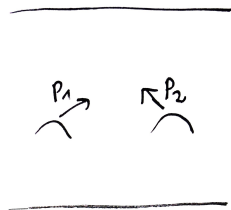
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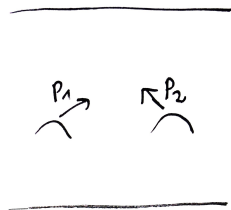
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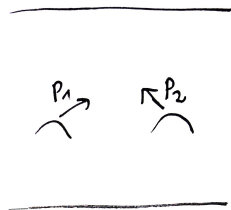
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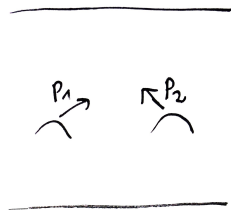
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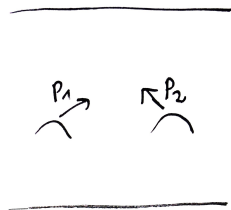
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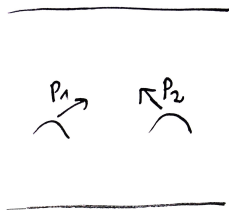
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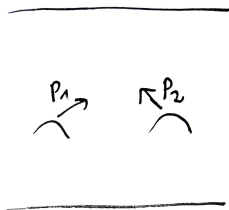
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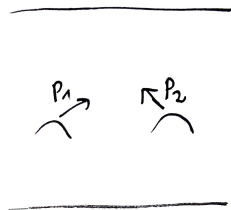
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II) solve the theory on a (large!) cylinder

1. Bethe Ansatz Quantization

$$e^{ip_k L} \prod_{l \neq k} S(p_k, p_l) = 1$$

2. Get the energies from

$$E = \sum_k E(p_k) = \sum_k \sqrt{1 + \frac{\lambda}{\pi^2} \sin^2 \frac{p_k}{2}}$$

This gives the spectrum up to wrapping corrections...

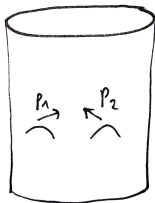
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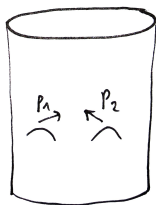
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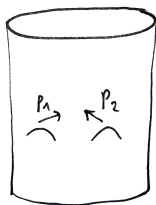
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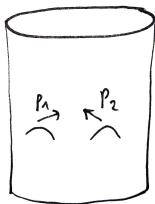
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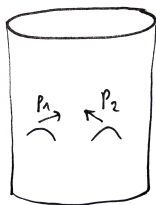
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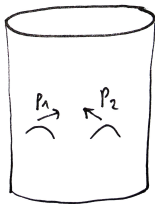
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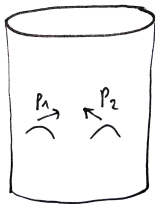
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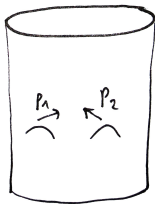
This gives the spectrum up to wrapping corrections...

relativistic  $\sim e^{-mL}$     weak coupling  $\sim \lambda^L$     strong coupling  $\sim e^{-\frac{2\pi L}{\sqrt{\lambda}}}$

**Main message:** Simple passage to finite volume (up to wrapping)

## How to solve the spectral problem?

### II) solve the theory on a (large!) cylinder



#### 1. Bethe Ansatz Quantization

$$e^{ip_k L} \prod_{l \neq k} S(p_k, p_l) = 1$$

#### 2. Get the energies from

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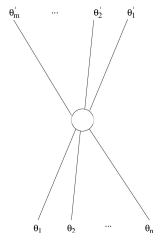


**Proceed to form factors...**

# Form factors

- Form factors are expectation values of a local operator sandwiched between specific multiparticle *in* and *out* states

$$p_k = m \sinh \theta$$



- Form factors in infinite volume (on an infinite plane) satisfy a definite set of **functional equations**

$$f(\theta_1, \theta_2) = S(\theta_1, \theta_2) f(\theta_2, \theta_1)$$

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Pozsgay, Takacs

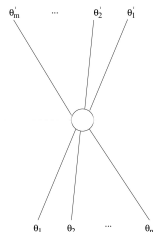
$$\langle \emptyset | \mathcal{O}(0) | \theta_1, \theta_2 \rangle_L = \frac{1}{\sqrt{\rho_2 \cdot S(\theta_1, \theta_2)}} \cdot f(\theta_1, \theta_2) + O(e^{-mL})$$

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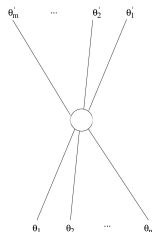
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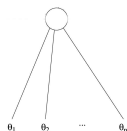
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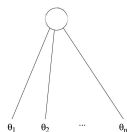
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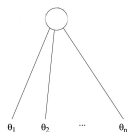
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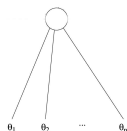
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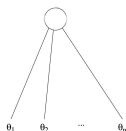
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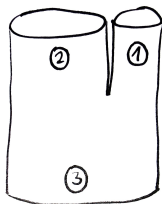
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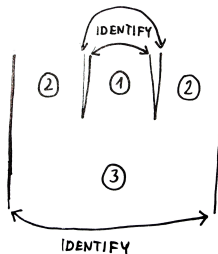
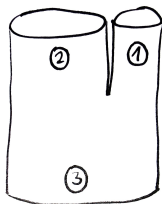


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- ▶ impose continuity conditions for  $\phi$  and  $\Pi \equiv \partial_t \phi$
- ▶  $\phi$  expressed in terms of  $\cos \frac{2\pi n}{L_r}$  and  $\sin \frac{2\pi n}{L_r}$  modes...

*looks like an inherently finite-volume computation...*

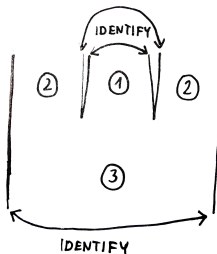
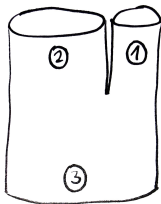
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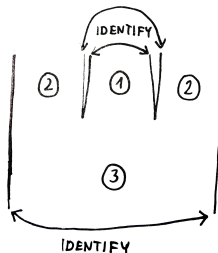
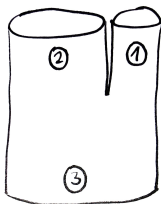
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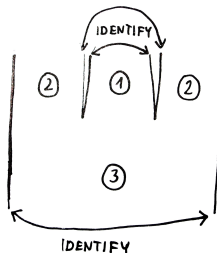
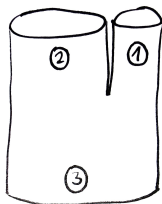


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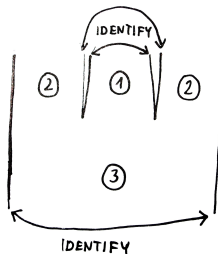
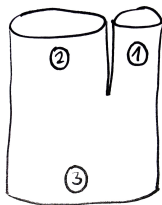


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- ▶ In the pp-wave times, people used simplified expressions for  $N_{nm}^{rs}$  neglecting exponential  $e^{-\mu\alpha_r}$  terms  $\alpha_r = L_r/L_3$  (these are exactly wrapping terms  $e^{-ML_r}$ !!)

- ▶ Going to an exponential basis (BMN basis) one got in this limit

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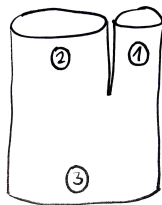
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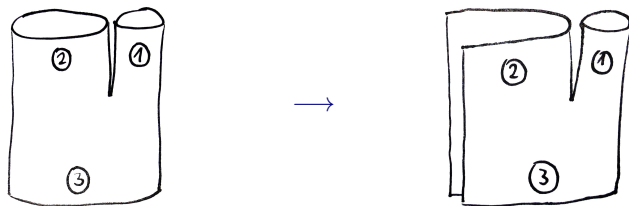
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Formulate functional equations...

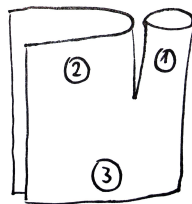
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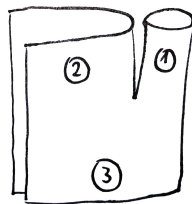
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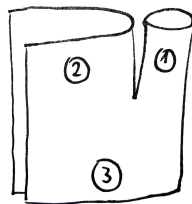
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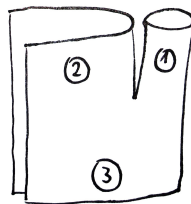
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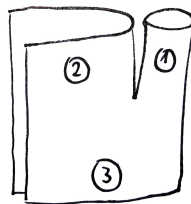
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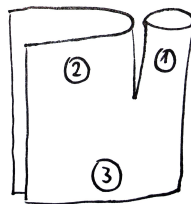
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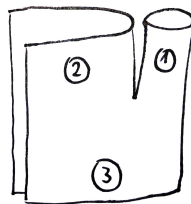


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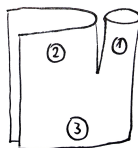
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### Functional equations for the (decompactified) string vertex

written here for two incoming particles and, for the moment, free theory



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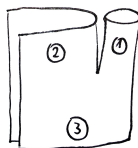
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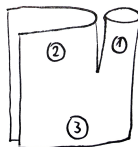
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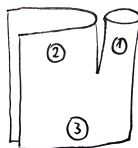
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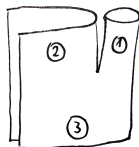
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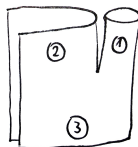
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### Functional equations for the (decompactified) string vertex

written here for two incoming particles and **interacting** worldsheet theory



$$\begin{aligned}N^{33}(\theta_1, \theta_2) &= N^{33}(\theta_2, \theta_1) \cdot \mathbf{S}(\theta_1, \theta_2) \\N^{33}(\theta_1, \theta_2) &= e^{-ip_1 L} N^{33}(\theta_2, \theta_1 - 2\pi i) \\-i \operatorname{res}_{\theta'=\theta} N^{33}(\theta + i\pi, \theta) &= (1 - e^{ipL}) F_0\end{aligned}$$

- ▶ The exact pp-wave solution, involving the  $\Gamma_\mu(\theta)$  special function solves these equations and can be reconstructed from them!

$$n(\theta)n(\theta + i\pi) = -\frac{1}{2\pi^2} ML \sinh \theta \sin \frac{p(\theta)L}{2}$$

- ▶ This includes **all exponential wrapping corrections**  $e^{-\mu\alpha_1} = e^{-ML}$  for the #1 string
- ▶ Straightforward generalization of the axioms to an **interacting** integrable QFT

## Step II) The string vertex — back to finite volume

We considered so far the ‘decompactified string vertex’...

but ultimately we are interested in the finite volume one...

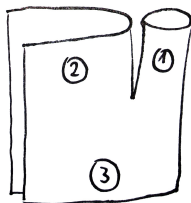
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We neglect exponential corrections for strings #2 and #3 but **keep** all size dependence of string #1...



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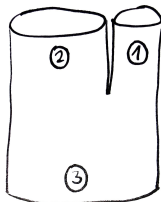
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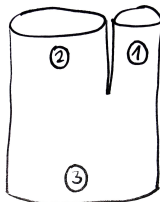
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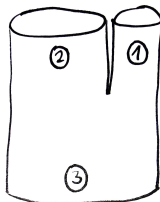
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## What happens in $AdS_5 \times S^5$ ?

### Novel kinematics

- ▶ Complex rapidities are defined on a covering of an elliptic curve
- ▶ Only  $e^{ip}$  (and not the momentum  $p$  itself) is a well defined elliptic function
- ▶ The phase factors  $e^{ipL}$  make sense directly only for integer  $L$  which is nice from the point of view of  $\mathcal{N} = 4$  SYM...

### Complicated dynamics

- ▶ The S-matrix does not depend on the difference of rapidities
- ▶ The S-matrix is nondiagonal which drastically complicates solving form factor axioms (which are a special case of our SFT axioms)

We would like to separate the two problems...

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## The $AdS_5 \times S^5$ Neumann coefficient

Factorize the  $AdS_5 \times S^5$  Neumann coefficient as

$$N_L^{33}(z_1, z_2)_{i_1 i_2} = \underbrace{F(z_1, z_2)_{i_1, i_2}}_{\text{form factor}} \cdot \underbrace{N_L^{33}(z_2, z_1)}_{\text{includes all } L \text{ dependence}}$$

- ▶ The  $N_L^{33}(z_2, z_1)$  includes **all** dependence on the size of the string **L** at any coupling ( $\equiv$  infinite set of wrapping corrections)
- ▶ Currently we have a solution of the functional equations for any even **L** at arbitrary coupling

$$N_L^{33}(z_2, z_1) \sim \text{conventional elliptic functions} \times \text{Elliptic Gamma}$$

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## Conclusions

- ▶ We propose a framework for formulating functional equations for string interactions (light cone string field theory vertex) when the worldsheet theory is **integrable**
  - ▶ This approach should work in particular for strings in the full  $AdS_5 \times S^5$  geometry
  - ▶ A key step is the existence of an infinite volume setup, which allows for formulating functional equations incorporating e.g. crossing
  - ▶ Second step involves reduction to (large) finite size of two of the three strings
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- ▶ We reproduced pp-wave string field theory formulas for the Neumann coefficients
  - ▶ We solved for the 'kinematical' part of the  $AdS_5 \times S^5$  Neumann coefficient describing exact volume dependence (currently for even  $L$ ) at any coupling – may describe all order wrapping w.r.t. one string

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