he Vernacular of the S-Matrix

Jacob Bourjaily Niels Bohr International Academy and Discovery Center





Thursday, 25th June 2015

Strings 2015, ICTS-TIFR, Bengaluru

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Organization and Outline

- Spiritus Movens: the Discovery of On-Shell Physics
 Using Generalized Unitarity to Compute One-Loop Amplitudes
- Upgrading Unitarity at One-Loop: the *Chiral* Box Expansion
 Chiral Boxes Expansion for One-Loop *Integrands* Making Manifest the Einiteness of All Einite Observables
 - Making Manifest the Finiteness of All Finite Observables
- 3 Generalizing Unitarity for Two-Loop Amplitudes & Integrands
 - The Two-Loop Chiral Integrand Expansion
 - Novel Contributions at Two-Loops and Transcendentality
 - Local, Integrand-Level Representations of All Two-Loop Amplitudes

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition $\int_{\mathbb{R}^{3,1}} d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \ f_{a,b,c,d}$

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition $\int_{\mathbb{R}^{3,1}} d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \ f_{a,b,c,d}$

 $I_{a,b,c,d}$





Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:





the Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



通とくほとくほど

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



通とくほとくほど

Spiritus Movens: One-Loop Generalized Unitarity



Spiritus Movens: One-Loop Generalized Unitarity



Spiritus Movens: One-Loop Generalized Unitarity



Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



A 10

(4) (2) (4) (3) (4)

Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



A 10

(4) (2) (4) (3) (4)

Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



A 10

Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



A 10

Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



A 10

Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



A 10

Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



A 10

Spiritus Movens: One-Loop Generalized Unitarity



Spiritus Movens: One-Loop Generalized Unitarity



Spiritus Movens: One-Loop Generalized Unitarity



Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



(A) (E) (A) (E) (A)

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



直 とう きょう く いう

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



直 とう きょう く いう

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:



直 とう きょう く いう
































































Spiritus Movens: One-Loop Generalized Unitarity

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition $\int_{\mathbb{R}^{3,1}} d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$

Advantages:

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition

$$\int d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

$$\mathbb{R}^{3,1}$$

Advantages:

- each standardized, scalar integral need only be computed once
- all coefficients are easy to compute as on-shell diagrams

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition

$$\int_{\mathbb{R}^{3,1}} d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Advantages:

- each standardized, scalar integral need only be computed once
- all coefficients are easy to compute as on-shell diagrams

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition

$$\int d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

$$\mathbb{R}^{3,1}$$

Advantages:

- each standardized, scalar integral need only be computed once
- all coefficients are easy to compute as on-shell diagrams

Disadvantages:

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition

$$\int d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Advantages:

- each standardized, scalar integral need only be computed once
- all coefficients are easy to compute as on-shell diagrams

Disadvantages:

- all widely-used methods of regularization *severely* obscure the finiteness and dual-conformal invariance of finite observables
- breaks the symmetries of the actual, field-theory loop integrand

Thursday, 25th June 2015

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition

$$\int d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Advantages:

- each standardized, scalar integral need only be computed once
- all coefficients are easy to compute as on-shell diagrams

Disadvantages:

- all widely-used methods of regularization *severely* obscure the finiteness and dual-conformal invariance of finite observables
- breaks the symmetries of the actual, field-theory loop *integrand*

Historically, on-shell functions were first studied in the context of using **generalized unitarity** to determine (*integrated*) one-loop amplitudes:

The Scalar Box Decomposition

$$\int d^4 \ell \ \mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Advantages:

- each standardized, scalar integral need only be computed once
- all coefficients are easy to compute as on-shell diagrams

Disadvantages:

- all widely-used methods of regularization *severely* obscure the finiteness and dual-conformal invariance of finite observables
- breaks the symmetries of the actual, field-theory loop *integrand*

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \,\mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \,\mathcal{A}_n^{(k),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Consider for example the 'MHV' amplitude

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \,\mathcal{A}_n^{(2),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Consider for example the 'MHV' amplitude (k=2)

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \,\mathcal{A}_n^{(2),1} = \sum_{a,b,c,d} I_{a,b,c,d} \left(f_{a,b,c,d}^1 + f_{a,b,c,d}^2 \right)$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,b,c,d} \, I_{a,b,c,d} \, f_{a,b,c,d}^1$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,b,c,d} \, I_{a,b,c,d} \, f_{a,b,c,d}^1$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,b,c,d} \, I_{a,b,c,d} \, f_{a,b,c,d}^1$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$, and the only non-vanishing $f_{a,b,c,d}^1$ are:

$$f_{a,a+1,c,c+1}^1 = \begin{array}{c} c+1 & Q_1 \\ & Q_1 \\ & \ddots \\ & & a+1 \end{array}$$

 < □ > < □ > < □ > < ≡ > < ≡ > ≡ < ○ Q (?)</td>

 The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$, and the only non-vanishing $f_{a,b,c,d}^1$ are:



the Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$, and the only non-vanishing $f_{a,b,c,d}^1$ are:



the Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$, and the only non-vanishing $f_{a,b,c,d}^1$ are: d^a

$$f_{a,a+1,c,c+1}^{1} = \underbrace{c+1}_{c} \underbrace{\mathcal{Q}_{1}}_{c} a+1$$

$$c+1 \underbrace{\mathcal{Q}_{1}}_{c} a+1 \Leftrightarrow \int d^{4}\ell \underbrace{\frac{(a,c)(a,a+1)-(a,c+1)(c,a+1)}{(\ell,a)(\ell,a+1)(\ell,c)(\ell,c+1)}}_{\mathcal{I}_{a,a+1,c,c+1}^{1}} \underbrace{\ell,\mathcal{Q}_{2}(X,\mathcal{Q}_{1})}_{\mathcal{I}_{a,a+1,c,c+1}^{1}}$$

 < □ > < □ > < □ > < ≡ > < ≡ > ≡ < ○ Q (?)</td>

 The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

The Scalar Box Expansion for the One-Loop Amplitude

$$\int d^4 \ell \, \mathcal{A}_n^{(2),1} = \sum_{a,c} \, I_{a,a+1,c,c+1} \, f^1_{a,a+1,c,c+1}$$

Consider for example the 'MHV' amplitude (k=2), for which $f_{a,b,c,d}^2 = 0$, and the only non-vanishing $f_{a,b,c,d}^1$ are: d^a

$$f_{a,a+1,c,c+1}^{1} = \underbrace{c+1}_{c} \underbrace{\mathcal{Q}_{1}}_{c} a+1$$

$$c+1 \underbrace{\mathcal{Q}_{1}}_{c} a+1 \Leftrightarrow \int d^{4}\ell \underbrace{\frac{(a,c)(a,a+1)-(a,c+1)(c,a+1)}{(\ell,a)(\ell,a+1)(\ell,c)(\ell,c+1)}}_{\mathcal{I}_{a,a+1,c,c+1}^{1}} \underbrace{\ell,\mathcal{Q}_{2}(X,\mathcal{Q}_{1})}_{\mathcal{I}_{a,a+1,c,c+1}^{1}}$$

 < □ > < □ > < □ > < ≡ > < ≡ > ≡ < ○ Q (?)</td>

 The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(2),1} \stackrel{\underline{i}^2}{=} \sum_{a,c} \mathcal{I}_{a,a+1,c,c+1}^1 f_{a,a+1,c,c+1}^1$$

$$f_{a,a+1,c,c+1}^{1} = \underbrace{\begin{array}{c} c+1 \\ c \end{array}}_{c} \underbrace{\begin{array}{c} Q_{1} \\ Q_{1} \\ c \end{array}}_{a+1} \Leftrightarrow \int d^{4}\ell \underbrace{\frac{(a,c)(a,a+1)-(a,c+1)(c,a+1)}{(\ell,a)(\ell,a+1)(\ell,c)(\ell,c+1)}}_{\mathcal{I}_{a,a+1,c,c+1}^{1}} \underbrace{\begin{array}{c} (\ell, Q_{2})(X, Q_{1}) \\ (\ell, X)(Q_{2}, Q_{1}) \\ \mathcal{I}_{a,a+1,c,c+1}^{1} \end{array}}_{\mathcal{I}_{a,a+1,c,c+1}^{1}}$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{\flat}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

$$f_{a,a+1,c,c+1}^{1} = \overset{c+1}{\underbrace{\begin{array}{c} & & \\ & &$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}^?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}^?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on all co-dimension four residues

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*.

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**!
Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\Leftrightarrow \int d^4\ell \, \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

 Image: Constraint of the S-Matrix: Revisiting Generalized Unitarity

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\mathcal{I}_{\rm div}^a \equiv : \qquad \Longleftrightarrow \int d^4\ell \, \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

 Image: Constraint of the S-Matrix: Revisiting Generalized Unitarity

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\mathcal{I}_{\rm div}^a \equiv : a = d^4\ell \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

and the residue about the point $\ell \rightarrow x_a$ must be the tree amplitude: $\mathcal{A}_n^{(k),0}$

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_n^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^1 f_{a,b,c,d}^1 + \mathcal{I}_{a,b,c,d}^2 f_{a,b,c,d}^2 \right)$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\mathcal{I}_{\rm div}^a \equiv : a = d^4\ell \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

and the residue about the point $\ell \rightarrow x_a$ must be the tree amplitude: $\mathcal{A}_n^{(k),0}$

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_{n}^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^{1} f_{a,b,c,d}^{1} + \mathcal{I}_{a,b,c,d}^{2} f_{a,b,c,d}^{2} \right) + \mathcal{A}_{n}^{(k),0} \sum_{a} \mathcal{I}_{\text{div}}^{a}$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\mathcal{I}_{\rm div}^a \equiv \frac{1}{2} \left(\frac{d^2 \ell}{d^2 \ell} \frac{(a-1,a+1)(a,X)}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,X)} \right)$$

and the residue about the point $\ell \rightarrow x_a$ must be the tree amplitude: $\mathcal{A}_n^{(k),0}$

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_{n}^{(k),1} \stackrel{\underline{i}?}{=} \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^{1} f_{a,b,c,d}^{1} + \mathcal{I}_{a,b,c,d}^{2} f_{a,b,c,d}^{2} \right) + \mathcal{A}_{n}^{(k),0} \sum_{a} \mathcal{I}_{\text{div}}^{a}$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\mathcal{I}_{\rm div}^a \equiv \frac{1}{2} \left(\frac{d^2 \ell}{d^2 \ell} \frac{(a-1,a+1)(a,X)}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,X)} \right)$$

and the residue about the point $\ell \rightarrow x_a$ must be the tree amplitude: $\mathcal{A}_n^{(k),0}$

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_{n}^{(k),1} = \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^{1} f_{a,b,c,d}^{1} + \mathcal{I}_{a,b,c,d}^{2} f_{a,b,c,d}^{2} \right) + \mathcal{A}_{n}^{(k),0} \sum_{a} \mathcal{I}_{\text{div}}^{a}$$

This ansatz matches the correct integrand on **all** co-dimension four residues *involving four distinct propagators*. **However**, each chiral box is **IR-finite**! There are **also** co-dimension four residues involving only three propagators:

$$\mathcal{I}_{\rm div}^a \equiv : \qquad \Leftrightarrow \int d^4\ell \, \frac{(a-1,a+1)(a,\mathbf{X})}{(\ell,a-1)(\ell,a)(\ell,a+1)(\ell,\mathbf{X})}$$

and the residue about the point $\ell \rightarrow x_a$ must be the tree amplitude: $\mathcal{A}_n^{(k),0}$

the Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_{n}^{(k),1} = \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^{1} f_{a,b,c,d}^{1} + \mathcal{I}_{a,b,c,d}^{2} f_{a,b,c,d}^{2} \right) + \mathcal{A}_{n}^{(k),0} \sum_{a} \mathcal{I}_{\text{div}}^{a}$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand

$$\mathcal{A}_{n}^{(k),1} = \sum_{a,b,c,d} \left(\mathcal{I}_{a,b,c,d}^{1} f_{a,b,c,d}^{1} + \mathcal{I}_{a,b,c,d}^{2} f_{a,b,c,d}^{2} \right) + \mathcal{A}_{n}^{(k),0} \sum_{a} \mathcal{I}_{\text{div}}^{a}$$

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand



Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands



Because the divergences are universal, the ratio function is manifestly finite!

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands



Because the divergences are universal, the ratio function is manifestly finite!

$$\mathcal{R}_n^{(k),1} \equiv \mathcal{A}_n^{(k),1} - \mathcal{A}_n^{(k),0} \times \mathcal{A}_n^{(2),1}$$

(k), 1

Chiral Boxes Expansion for One-Loop Integrands Making Manifest the Finiteness of All Finite Observables

A 'Box'-Expansion for One-Loop Integrands

A Chiral 'Box'-Expansion for the One-Loop Amplitude Integrand



Because the divergences are universal, the ratio function is manifestly finite!

$$\begin{aligned} \mathcal{R}_n^{(k),1} &\equiv \mathcal{A}_n^{(k),1} - \mathcal{A}_n^{(k),0} \times \mathcal{A}_n^{(2),1} \\ &= \mathcal{A}_{n,\text{fin}}^{(k),1} - \mathcal{A}_n^{(k),0} \times \mathcal{A}_{n,\text{fin}}^{(2),1} \end{aligned}$$

Thursday, 25th June 2015

Strings 2015, ICTS-TIFR, Bengaluru

the Vernacular of the S-Matrix: Revisiting Generalized Unitarity

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders.

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders.

$$\mathcal{A}_{n}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell}$$

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell}$$

Manifesting the Exponentiation of Divergences to All Orders

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell}$$

with

$$\mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

Manifesting the Exponentiation of Divergences to All Orders

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell}$$

with

$$\mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

And this separation makes *manifest* the finiteness of **all** finite observables

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

$$\mathcal{R}_n^{(k)} \equiv rac{\mathcal{A}_n^{(k)}}{\mathcal{A}_n^{(2)}} \equiv \sum_{\ell=0}^\infty g^\ell \mathcal{R}_n^{(k),\ell}$$

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

$$\mathcal{R}_{n}^{(k)} \equiv \frac{\mathcal{A}_{n}^{(k)}}{\mathcal{A}_{n}^{(2)}} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{R}_{n}^{(k),\ell} \quad \text{where} \quad \mathcal{A}_{n}^{(k)} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{A}_{n}^{(k),\ell}$$

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

$$\mathcal{R}_{n}^{(k)} \equiv \frac{\mathcal{A}_{n}^{(k)}}{\mathcal{A}_{n}^{(2)}} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{R}_{n}^{(k),\ell} \quad \text{where} \quad \mathcal{A}_{n}^{(k)} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{A}_{n}^{(k),\ell}$$

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

And this separation makes *manifest* the finiteness of **all** finite observables *e.g.* the ℓ -loop ratio function:

 $\mathcal{R}_{n}^{(k)} \equiv \frac{\mathcal{A}_{n}^{(k)}}{\mathcal{A}_{n}^{(2)}} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{R}_{n}^{(k),\ell} \quad \text{where} \quad \mathcal{A}_{n}^{(k)} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{A}_{n}^{(k),\ell}$ Using the separation of $\mathcal{A}_{n}^{(k),\ell}$ together with the form of $\mathcal{A}_{n,\text{div}}^{(k),\ell}$ given above, it can be shown that:

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

And this separation makes *manifest* the finiteness of **all** finite observables *e.g.* the ℓ -loop ratio function:

$$\mathcal{R}_{n}^{(k)} \equiv \frac{\mathcal{A}_{n}^{(k)}}{\mathcal{A}_{n}^{(2)}} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{R}_{n}^{(k),\ell} \quad \text{where} \quad \mathcal{A}_{n}^{(k)} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{A}_{n}^{(k),\ell}$$

the separation of $\mathcal{A}_{n}^{(k),\ell}$ together with the form of $\mathcal{A}_{n,\text{div}}^{(k),\ell}$ given above,

Using the separation of $\mathcal{A}_n^{(\kappa),c}$ together with the form of $\mathcal{A}_{n,\text{div}}^{(\kappa),c}$ given abovit can be shown that:

$$\mathcal{R}_{n}^{(k),\ell} = \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} - \sum_{q=1}^{\ell} \mathcal{R}_{n}^{(k),\ell-q} \mathcal{A}_{n,\mathrm{fin}}^{(2),q}$$

Thursday, 25th June 2015

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

The separation of amplitudes into *manifestly* finite and *manifestly* divergent parts can be done at all loop orders. Moreover, **all** divergences exponentiate:

$$\mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} \equiv \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} + \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \qquad \text{with} \qquad \mathcal{A}_{n,\mathrm{div}}^{(k),\ell} \equiv \sum_{q=1}^{\ell} \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell-q} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)^{q}$$

And this separation makes *manifest* the finiteness of **all** finite observables *e.g.* the ℓ -loop ratio function:

$$\mathcal{R}_{n}^{(k)} \equiv \frac{\mathcal{A}_{n}^{(k)}}{\mathcal{A}_{n}^{(2)}} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{R}_{n}^{(k),\ell} \quad \text{where} \quad \mathcal{A}_{n}^{(k)} \equiv \sum_{\ell=0}^{\infty} g^{\ell} \mathcal{A}_{n}^{(k),\ell}$$

the separation of $\mathcal{A}_{n}^{(k),\ell}$ together with the form of $\mathcal{A}_{n,\text{div}}^{(k),\ell}$ given above,

Using the separation of $\mathcal{A}_n^{(\kappa),c}$ together with the form of $\mathcal{A}_{n,\text{div}}^{(\kappa),c}$ given abovit can be shown that:

$$\mathcal{R}_{n}^{(k),\ell} = \mathcal{A}_{n,\mathrm{fin}}^{(k),\ell} - \sum_{q=1}^{\ell} \mathcal{R}_{n}^{(k),\ell-q} \mathcal{A}_{n,\mathrm{fin}}^{(2),q}$$

Thursday, 25th June 2015

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

$$\mathcal{A}_{n,\operatorname{div}}^{(k),2}$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

$$\mathcal{A}_{n,\operatorname{div}}^{(k),2} =$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\mathrm{fin}}^{(k),2} + \mathcal{A}_{n,\mathrm{div}}^{(k),2}$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \bigotimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) \quad + \quad \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$
The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

All the divergent contributions, $\mathcal{A}_{n,\text{div}}^{(k),2}$, are easy to identify:

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) \quad + \quad \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

All the divergent contributions, $\mathcal{A}_{n,\text{div}}^{(k),2}$, are easy to identify:

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) \quad + \quad \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

All the divergent contributions, $\mathcal{A}_{n,\text{div}}^{(k),2}$, are easy to identify:

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) \quad + \quad \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop *Chiral* Expansion
$$\mathcal{A}_{n}^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

All the divergent contributions, $\mathcal{A}_{n,\text{div}}^{(k),2}$, are easy to identify:

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \bigotimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) \quad + \quad \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \bigotimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right) \quad + \quad \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

 $\mathcal{I}_L(X) \otimes \mathcal{I}_R(X)$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) \quad + \quad \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

 $\mathcal{I}_L(\mathbf{X}) \bigotimes \mathcal{I}_R(\mathbf{X}) \equiv$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right) + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \Big) + \Big(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \Big)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right) + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

Thursday, 25th June 2015

the Vernacular of the S-Matrix: Revisiting Generalized Unitarity

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{n,\mathrm{div}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{\mathrm{uv}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{(k),0} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

 $\frac{(b-1,b+1)(b,X)}{(\ell_1,b-1)(\ell_1,b)(\ell_1,b+1)(\ell_1,X)} \otimes \frac{(X,a)(a-1,a+1)}{(X,\ell_2)(\ell_2,a-1)(\ell_2,a)(\ell_2,a+1)}$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{\mathbf{A}_{n,\mathrm{div}}^{(k),1}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

$$\frac{(b-1,b+1)(b,a)(a-1,a+1)}{(\ell_1,b-1)(\ell_1,b+1)(\ell_1,\ell_2)(\ell_2,a-1)(\ell_2,a)(\ell_2,a+1)}$$

Ш

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{\mathbf{A}_{n,\mathrm{div}}^{(k),1}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

$$\frac{(b-1,b+1)(b,a)(a-1,a+1)}{(\ell_1,b-1)(\ell_1,b+1)(\ell_1,\ell_2)(\ell_2,a-1)(\ell_2,a)(\ell_2,a+1)}$$

Ш

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{\mathbf{A}_{n,\mathrm{div}}^{(2),1}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{\mathbf{A}_{n,\mathrm{div}}^{(2),1}} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{b} + \left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)$$

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

The Two-Loop Chiral Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_n^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1} \right)}_{l} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1} \right)}_{l}$$

7

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1}\right)}_{b} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1}\right)}_{a} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(k),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{($$

Thursday, 25th June 2015

h

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1}\right)}_{b} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1}\right)}_{a} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(k),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{($$

Thursday, 25th June 2015

h

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Constructing Local Integrands for Two-Loop Amplitudes

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\text{fin}}^{(k),2} + \mathcal{A}_{n,\text{div}}^{(k),2}$$

"Merging" One-Loop, Chiral (X-dependent) Integrands

$$\mathcal{I}_{L}(X) \otimes \mathcal{I}_{R}(X) \equiv \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, X)}{(\ell_{1}, X)} \otimes \frac{(X, \mathcal{N}_{R})}{(X, \ell_{2})} \mathcal{I}_{R}' \mapsto \mathcal{I}_{L}' \frac{(\mathcal{N}_{L}, \mathcal{N}_{R})}{(\ell_{1}, \ell_{2})} \mathcal{I}_{R}'$$

$$\mathcal{A}_{n,\mathrm{div}}^{(k),2} = \mathcal{A}_{n}^{(k),0} \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{div}}^{(2),1}\right)}_{b} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(2),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{(k),1}\right)}_{a} + \underbrace{\left(\mathcal{A}_{n,\mathrm{div}}^{(k),1} \otimes \mathcal{A}_{n,\mathrm{fin}}^{($$

Thursday, 25th June 2015

h

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



2. Finite Penta-Boxes:

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:





The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:



2. <u>Finite</u> Penta-Boxes:



 Image: state of the s-Matrix: Revisiting Generalized Unitarity

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

1. "Kissing" Boxes:





The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

3. Finite Double-Boxes:



Thursday, 25th June 2015

Strings 2015, ICTS-TIFR, Bengaluru

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

3. <u>Finite</u> Double-Boxes:





The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

3. <u>Finite</u> Double-Boxes:





The Two-Loop Chiral Integrand Expansion

Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Finite Integrand Contributions to Two-Loop Amplitudes

3. <u>Finite</u> Double-Boxes:





The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:
The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'.

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

 $\mathcal{A}_{10}^{(5)}(\varphi_{12},\varphi_{12},\varphi_{12},\varphi_{23},\varphi_{23},\varphi_{34},\varphi_{34},\varphi_{34},\varphi_{41},\varphi_{41})$

 $\propto (\widetilde{\eta}_1^1 \widetilde{\eta}_1^2) (\widetilde{\eta}_2^1 \widetilde{\eta}_2^2) (\widetilde{\eta}_3^1 \widetilde{\eta}_3^2) (\widetilde{\eta}_4^2 \widetilde{\eta}_4^3) (\widetilde{\eta}_5^2 \widetilde{\eta}_5^3) (\widetilde{\eta}_6^3 \widetilde{\eta}_6^4) (\widetilde{\eta}_7^3 \widetilde{\eta}_7^4) (\widetilde{\eta}_8^3 \widetilde{\eta}_8^4) (\widetilde{\eta}_9^4 \widetilde{\eta}_9^1) (\widetilde{\eta}_{10}^4 \widetilde{\eta}_{10}^1)$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

 $\mathcal{A}_{10}^{(5)}(\varphi_{12},\varphi_{12},\varphi_{12},\varphi_{23},\varphi_{23},\varphi_{34},\varphi_{34},\varphi_{34},\varphi_{41},\varphi_{41})$

 $\propto (\widetilde{\eta}_1^1 \widetilde{\eta}_1^2) (\widetilde{\eta}_2^1 \widetilde{\eta}_2^2) (\widetilde{\eta}_3^1 \widetilde{\eta}_3^2) (\widetilde{\eta}_4^2 \widetilde{\eta}_4^3) (\widetilde{\eta}_5^2 \widetilde{\eta}_5^3) (\widetilde{\eta}_6^3 \widetilde{\eta}_6^4) (\widetilde{\eta}_7^3 \widetilde{\eta}_7^4) (\widetilde{\eta}_8^3 \widetilde{\eta}_8^4) (\widetilde{\eta}_9^4 \widetilde{\eta}_9^1) (\widetilde{\eta}_{10}^4 \widetilde{\eta}_{10}^1)$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

$$\mathcal{A}_{10}^{(5)}(\varphi_{12},\varphi_{12},\varphi_{12},\varphi_{23},\varphi_{23},\varphi_{34},\varphi_{34},\varphi_{34},\varphi_{41},\varphi_{41})$$

Problem: all (isolated) on-shell functions vanish on this component!



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:

$$\mathcal{A}_{10}^{(5)}ig(arphi_{12},arphi_{12},arphi_{12},arphi_{23},arphi_{23},arphi_{34},arphi_{34},arphi_{34},arphi_{41},arphi_{41}ig)$$

Problem: all (isolated) on-shell functions vanish on this component!



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:


The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:


The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Novel Contributions Required: the Shifted Double-Boxes

4. "Shifted" Double-Boxes:

It turns out that here are contributions to $\mathcal{A}_{n,\text{fin}}^{(k),2}$ which **cannot** be written as 'superfunction'×'integral'. To see this, consider the following 10-particle all-scalar, component amplitude:



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\operatorname{div}}^{(k),2} + \mathcal{A}_{n,\operatorname{fin}}^{(k),2}$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\operatorname{div}}^{(k),2} + \mathcal{A}_{n,\operatorname{fin}}^{(k),2}$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n, ext{div}}^{(k),2} + \mathcal{A}_{n, ext{fin}}^{(k),2}$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\mathrm{div}}^{(k),2} + \mathcal{A}_{n,\mathrm{fin}}^{(k),2}$$



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion (l) = (l

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\mathrm{div}}^{(k),2} + \mathcal{A}_{n,\mathrm{fin}}^{(k),2}$$





The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes





The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n,\mathrm{div}}^{(k),2} + \mathcal{A}_{n,\mathrm{fin}}^{(k),2}$$

The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n, ext{div}}^{(k),2} + \mathcal{A}_{n, ext{fin}}^{(k),2}$$



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes

Local Integrand Expansion

$$\mathcal{A}_n^{(k),2} = \mathcal{A}_{n, ext{div}}^{(k),2} + \mathcal{A}_{n, ext{fin}}^{(k),2}$$





The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes





The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



The Two-Loop Chiral Integrand Expansion Novel Integrand Contributions at Two-Loops and Transcendentality Local, Integrand-Level Representations of All Two-Loop Amplitudes

Local Integrand Expansion for All Two-Loop Amplitudes



GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015
GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO LEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015

GRASSMANNIAN GEOMETRY OF SCATTERING AMPLITUDES

NIMA ARKANI-HAMED JACOB BOURJAILY FREDDY CACHAZO ALEXANDER GONCHAROV ALEXANDER POSTNIKOV JAROSLAV TRNKA



CAMBRIDGE UNIVERSITY PRESS

The Vernacular of the S-Matrix: Revisiting Generalized Unitarity

Thursday, 25th June 2015