Knot polynomials, homological invariants & topological strings

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Plan:

- (i) Knot polynomials from Chern-Simons gauge theory
- (ii) Our results on mutant knots
- (iii) Homological invariants
- (iv) Large N Chern-Simons & their closed topological string duals
- (v) Conclusions & challenging problems .

Knot polynomials from Chern-Simons theory





Torus Knot :Trefoil 31

Hyperbolic Knot: Figure-Eight knot 41

Chern-Simons theory provides a natural framework for the study of knots

• Chern-Simons action on S^3 based on gauge group G:

$$S_{CS}[A] = \frac{k}{4\pi} \int_{S^3} Tr\left(A \wedge dA + \frac{2}{3}A^3\right)$$

k is the coupling constant, A's are the gauge connections.

• Any knot *K* carrying representation *R* are described by **expectation** value of Wilson loop operators $W_R(K) = Tr[Pexp \oint A]$:

$$V_R^G[K] = \langle W_R(K) \rangle = \frac{\int_{S^3} [\mathcal{D}A] W_R(K) \exp(iS_{CS}[A])}{Z[S^3]}$$

where $Z[S^3] = \int_{S^3} [\mathcal{D}A] \exp(iS_{CS}[A])$ (partition function)

 $V_R^G[K]$ are the **knot invariants**.

Knot invariant computations

•These knot invariants $(V_R^G[K])$ can be directly evaluated using two inputs:(*Kaul*, *Govindarajan*, *PR* (1992))

1) Relation between Chern-Simons theory to G_k Wess-Zumino conformal field theory (Witten 1989).

2) Any knot can be obtained as a closure or platting of braid(Alexander, Birman)

For example, the trefoil can be redrawn as





$$V_R[3_1] = \langle \psi_0 | \psi_3 \rangle = \langle \psi_0 | \mathcal{B}^3 | \psi_0 \rangle$$

where $\ensuremath{\mathcal{B}}$ is the braiding operator.

To write the polynomial form of the knot invariant:

Expand the state $|\psi_0\rangle$ in a suitable basis in which ${\cal B}$ is diagonal.

For the four-punctured S^2 boundary, the conformal block bases are:



where $t \in R_1 \otimes R_2 \cap \overline{R}_3 \otimes \overline{R}_4$ and $s \in R_2 \otimes R_3 \cap \overline{R}_1 \otimes \overline{R}_4$. $a_{st} \begin{bmatrix} R_1 & R_2 \\ R_3 & R_4 \end{bmatrix}$ is the **duality matrix** (fusion matrix) relating these two bases which is proportional to the quantum Wigner 6j symbols:

$$a_{ts} \begin{bmatrix} R_1 & R_2 \\ R_3 & R_4 \end{bmatrix} \propto \begin{cases} R_1 & R_2 & t \\ R_3 & R_4 & s \end{cases}$$

For knots, two of the R_i 's will be R and the other two will be conjugate \overline{R} depending on the <u>orientation</u>.

In the braid diagram for trefoil, middle two strands are parallely oriented and they are braided.

$$|\Psi_0\rangle = \sum_{s \in R \otimes R} \mu_s |\hat{\Phi}_s(\bar{R}, R, R, \bar{R})\rangle$$

where $\mu_s = \sqrt{S_{0s}/S_{00}} \equiv \sqrt{dim_q t}$ (unknot normalisation)

$$V_R[3_1] = \langle \Psi_0 | \mathcal{B}^3 | \Psi_0 \rangle = \sum_s dim_q s (\lambda_s(R, R))^3$$

where braiding eigenvalue for parallelly oriented right-handed half-twists is

$$\lambda_t^{(+)}(R,R) = (-1)^{\epsilon_t} q^{2C_R - C_t/2}, \ q = e^{\frac{2\pi i}{k + C_v}} \text{ where } \epsilon_t = \pm 1$$

Antiparallel braiding eigenvalue will be $\lambda_s^{(-)}(R, \bar{R}) = (-1)^{\epsilon_s} q^{C_s/2}$. We require them for figure-eight drawn as quasi-plat.



$$V_R[4_1] = \sum_{t,s \in R \otimes \bar{R}} \sqrt{\dim_q t \, \dim_q s} \, a_{ts} \begin{bmatrix} \bar{R} & R \\ \bar{R} & R \end{bmatrix} (\lambda_t^{(-)})^2 (\lambda_s^{(-)})^{-2}$$

•Knot invariants involves braiding eigenvalues & fusion matrices

•Fusion matrices proportional to quantum Wigner 6j (completely known for SU(2) (Kirillov, Reshetikhin) but not for other groups)

• For few R's, we determined using knot equivalence and properties of Wigner 6j-Kaul, Govindarajan, PR (1992), Zodinmawia, PR(2012)

• Hence the knot invariants $V_R^G[K]$ can be written in variables dependent on k and rank of the group

Well-known knot polynomials match with knot invariants (normalized by unknot) when R is fundamental representation



Jones' polynomial for trefoil:

$$J[T;q] = V_{\Box}^{SU(2)}[T]/V_{\Box}^{SU(2)}[U] = q + q^3 - q^4.$$

We can place any representation R (higher spins) of SU(2) on the knot and obtain colored Jones' polynomials $J_n[K;q]$ where subscript n means spin (n-1)/2 or Young diagram single row with n-1 boxes. The colored Jones for $3_1\&4_1$ are

$$J_n[3_1;q] = \sum_{k=0}^{n-1} (-1)^k q^{k(k+3)/2+nk} (q^{-n-1},q)_k (q^{-n+1},q)_k$$

where $(z;q)_k = \prod_{j=0}^{k-1} (1 - zq^j)$ is called q-Pochhammer symbol.

$$J_n[4_1;q] = \sum_{k=0}^{n-1} (-1)^k q^{nk} (q^{-n-1}, q^{-1})_k (q^{-n+1}, q)_k$$

These two colored Jones polynomial is sufficient to determine $J_n[K_p; q]$ for twist knots K_p



p full twists

p	_4	-3	-2	-1	0	1	2	3	4
knots	10_1	8 ₁	61	4_1		3_1	5_2	7_2	9_2

• [n]-colored HOMFLY $P_n[K_p; a, q]$ (Nawata,Zodin, PR),2012 gave more data leading to conjecture a closed form expression for $U_q(sl_N)$ quantum Wigner 6j symbols for a class of R_i 's (Nawata,Zodin,PR),2013 There are two types of Wigner 6j for SU(N): Type I:



where $n_2 \le n_1 \le n_3$, $k_1 \le n_2$ and $k_2 \le n_1$.

Type II



where $n_1 \leq n_2$, $k_2 \leq \min(n_1, n_3)$ and $k_1 \leq \min(n_1, n_3, n_4)$. The fusion rule requires $n_1 + n_2 = n_3 + n_4$.

$$\begin{cases} \lambda_{1} \quad \lambda_{2} \quad \lambda_{12} \\ \lambda_{3} \quad \lambda_{4} \quad \lambda_{23} \end{cases} = \Delta_{(1,2,12)} \Delta_{(3,4,12)} \Delta_{(1,4,23)} \Delta_{(2,3,23)} [N-1]! \sum_{z \ge 0} (-)^{z} \\ [z+N-1]! C_{z} \{ [z - \frac{1}{2} \langle \lambda_{1} + \lambda_{2} + \lambda_{12}, \alpha_{1}^{\vee} + \alpha_{N-1}^{\vee} \rangle]! \\ \left[z - \frac{1}{2} \langle \lambda_{3} + \lambda_{4} + \lambda_{12}, \alpha_{1}^{\vee} + \alpha_{N-1}^{\vee} \rangle \right]! \\ \left[z - \frac{1}{2} \langle \lambda_{1} + \lambda_{4} + \lambda_{23}, \alpha_{1}^{\vee} + \alpha_{N-1}^{\vee} \rangle \right]! \\ \left[z - \frac{1}{2} \langle \lambda_{2} + \lambda_{3} + \lambda_{23}, \alpha_{1}^{\vee} + \alpha_{N-1}^{\vee} \rangle \right]! \\ \left[\frac{1}{2} \langle \lambda_{1} + \lambda_{2} + \lambda_{3} + \lambda_{4}, \alpha_{1}^{\vee} + \alpha_{N-1}^{\vee} \rangle - z \right]! \\ \left[\frac{1}{2} \langle \lambda_{1} + \lambda_{3} + \lambda_{12} + \lambda_{23}, \alpha_{1}^{\vee} + \alpha_{N-1}^{\vee} \rangle - z \right]! \\ \left[\frac{1}{2} \langle \lambda_{2} + \lambda_{4} + \lambda_{12} + \lambda_{23}, \alpha_{1}^{\vee} + \alpha_{N-1}^{\vee} \rangle - z \right]! \right]^{-1}, \end{cases}$$

where

$$\begin{aligned} \Delta_{(1,2,3)} &\equiv \Delta(\lambda_1, \lambda_2, \lambda_3) = \left\{ \begin{bmatrix} \frac{1}{2} \langle -\lambda_1 + \lambda_2 + \lambda_3, \alpha_1^{\vee} + \alpha_{N-1}^{\vee} \rangle \end{bmatrix} \\ &\times \begin{bmatrix} \frac{1}{2} \langle \lambda_1 - \lambda_2 + \lambda_3, \alpha_1^{\vee} + \alpha_{N-1}^{\vee} \rangle \end{bmatrix} \\ &\times \begin{bmatrix} \frac{1}{2} \langle \lambda_1 + \lambda_2 - \lambda_3, \alpha_1^{\vee} + \alpha_{N-1}^{\vee} \rangle \end{bmatrix} \\ &\times \left\{ \begin{bmatrix} \frac{1}{2} \langle \lambda_1 + \lambda_2 + \lambda_3, \alpha_1^{\vee} + \alpha_{N-1}^{\vee} \rangle + N - 1 \end{bmatrix} \right\}^{1/2} \end{aligned}$$

$$C_{z}^{(I)} = \begin{cases} \delta_{z,z_{\min}+i} \begin{bmatrix} N-2+k_{2}-i \\ k_{2}-i \end{bmatrix}^{-1} & \text{for } k_{1} > k_{2} ,\\ \delta_{z,z_{\min}+i} \begin{bmatrix} N-2+k_{1}-i \\ k_{1}-i \end{bmatrix}^{-1} & \text{for } k_{1} \le k_{2} , \end{cases}$$

$$C_{z}^{(II)} = \begin{cases} \delta_{z,z_{\max}-i} \begin{bmatrix} N-2+k_{2}-i \\ k_{2}-i \end{bmatrix}^{-1} & \text{for } k_{1} > k_{2} ,\\ \delta_{z,z_{\max}-i} \begin{bmatrix} N-2+k_{1}-i \\ k_{1}-i \end{bmatrix}^{-1} & \text{for } k_{1} \le k_{2} , \end{cases}$$

OBSERVATIONS

• [n]-colored HOMFLY for knots drawn as quasi-plat of 4-strand braids can be obtained using this data.

 Our data is not sufficient to write the polynomial for knots obtained from quasi-plat of braids with more than 4-strands.
 (dual basis of 6-point or higher point conformal blocks requires Wigner 6-j beyond our conjectured class) How do we obtain [n]-colored HOMFLY for 9_{42} knot involving six braids?





Knots 10_{71} , 10_{152} can be drawn gluing 3-manifolds involving three S^2 boundaries each with four-punctures:



Requires the following building blocks to compute knot polynomials



 $u_r = \sum_t (dim_q t)^{(1-r/2)} |\phi_t^{(1)}\rangle \dots |\phi_t^{(r)}\rangle$



 v_8



 v_7

2

• We have redrawn many knots using these building blocks enabling evaluation of [n]-colored HOMFLY polynomials. These polynomials do not distinguish mutants.



Kinoshita-Terasaka & Conway mutants

- Need to work out polynomials for mixed representations.
- The two types of $U_q(sl_N)$ Wigner 6j has been recently determined for [2, 1] (Gu,Jockers),2014-first mixed representation
- Wanted to check the power [2,1] colored HOMFLY for the mutant pair
- The two knots (mutant pairs) are indeed distinguished (Nawata,Singh,PR,2015)

•Enumerated class of mutants which can be distinguished but some pretzel mutants with antiparallel odd-braids cannot be distinguished (Mironov, Morozov, Morozov, Singh, PR),2015

• Crucial input in the context of mixed representation: multiplicity

$$(21;0) \otimes (21;0) = (42;0)_0 \oplus (2^3;0)_0 \oplus (31^3;0)_0 \oplus (321;0)_0 \\ \oplus (321;0)_1 \oplus (41^2;0)_0 \oplus (3^2;0)_0 \oplus (2^21^2;0)_0$$

We see appears twice. Multiplicity incorporated four-point conformal blocks:

$$\begin{array}{c} R_{2} \\ R_{1} \\ R_{4} \end{array}^{R_{4}} = |\phi_{t,r_{3}r_{4}}^{(1)}(R_{1},\ldots,R_{4})\rangle, \quad R_{2} \\ R_{2} \\ R_{3} \\ R_{2} \\ R_{3} \\ R_{4} \end{array} = |\phi_{s,r_{1}r_{2}}^{(2)}(R_{1},\ldots,R_{4})\rangle$$

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Table gives the three-manifolds obtained from surgery of framed knots

Frame link knot	3-Manifold		
\bigcirc \bigcirc	S ² X S ¹		
8 070	S³		
	RP ³		
ODO P	L(P,1)		
$a_{1}=2$ $a_{2}=3$ $p = \frac{a_{1}}{a_{2}-1}$ a_{3}	L(5,3)		
	P ³		

Framed links related by Kirby moves gives the same three-manifold



Three-manifold invariant proportional to Chern-Simons partition function Z[M] respecting Kirby moves is(Kaul,PR),2000

$$Z[M] \propto \sum_{R} \dim_{q} R \, V_{R}^{SU(N)}[K]$$

Questions

- large N expansion $\log Z[M]$?
- For any knot, we observe Laurent series:

$$P[K; a, q] = \sum_{i,j} c_{i,j} a^{i} q^{j}$$

where $c_{i,j}$ are integers. Topological meaning or reason?

Why integers?- two parallel developments

 From physics (topological strings, BPS states counting) Ooguri, Vafa (1999)

•From mathematics(homological chain complex) Khovanov(1999), Khovanov-Rosansky(2004)

$$c_{i,j} = \sum_{k} (-1)^k dim H_{i,j,k}$$

Homological Invariants

Introduce A and B slicing as shown in the diagram. Define $n(s) = n_B$, $j(s) = n_B + n_+ - n_-$



 C_{nj} is the vector space with basis as states with n(s) = n and j(s) = j. $J[K;q] = \sum_{n,j} (-1)^n q^j \dim(C_{nj})$

where the homology chain

$$\partial: C_{n,j} \longrightarrow C_{n+1,j}, \quad \partial^2 = 0$$



$$C_{*j}: C_{0j} \longrightarrow C_{1j} \longrightarrow C_{2j} \longrightarrow \dots$$

The vector space

$$H_n(C_{*j}) = \frac{\ker(\partial : C_{n,j} \longrightarrow C_{n+1}, j)}{\operatorname{Image}(\partial : C_{n-1,j} \longrightarrow C_{n,j})}$$

$$Kh(K;q,t) = \sum_{n,j} t^n q^j \dim(H_{nj}).$$

Taking t = -1 gives the Jones polynomial J[K; q]



Large N Chern-Simons & topological strings

- Gopakumar-Vafa duality gives A-model closed topological string on a resolved conifold from large N expansion of $\log Z[S^3]$.
- Ooguri-Vafa conjecture a form for reformulated knot invariants $f_R(q, \lambda)$.
- We verified Ooguri-Vafa conjecture for non-torus knots (Sarkar, PR 2000) and higher crossing knots(Nawata,Zodin,PR 2013) using our [n]-colored HOMFLY polynomials.
- Making use of these two duality conjectures, we attempted the *N* expansion of logZ[M] for some three-manifolds *M* giving(P. Borhade, T. Sarkar, PR(2003)

$$\log Z_0[M] = \sum_{n=1}^{\infty} \sum_g \frac{1}{n} \left(\sinh \frac{dg_s}{2} \right)^{2g-2} \times \left\{ \sum_Q \sum_{\{\ell_\alpha\}} \sum_{\{s_\alpha\}} \hat{N}_{(R_{\ell_1,s_1},\dots,R_{\ell_r,s_r}),g,Q} (-1)^{\sum_\alpha s_\alpha} \right. \\ \left. (-1)^n \sum_\alpha \ell_\alpha p_\alpha \lambda^{\frac{1}{2}n \sum_\alpha \ell_\alpha p_\alpha} \left(\lambda^n \{Q + \sum_\alpha (\frac{-\ell_\alpha}{2} + s_\alpha)\} \right) \right. \\ \left. - \lambda^n \{Q + 1 + \sum_\alpha (\frac{-\ell_\alpha}{2} + s_\alpha)\} \right) \} \\ = \left. \sum_{g,n,m} \frac{1}{n} (2 \sinh \frac{ng_s}{2})^{2g-2} n_{g,m} e^{-dmt} \right\}$$

Subtle Issues

In the large k limit

$$Z[M] = \sum_{c} Z_{c}[M]$$

t' Hooft proposal requires

 $\ln Z_c[M] = \text{Closed String expansion}$

whereas we find

$$\ln\left(\sum_{c} Z_{c}[M]\right) = \text{Closed String partition function}$$

Hence we cannot predict duality between Chern-Simons gauge theory on M with the A-model string theory with the $n_{g,m}$'s we have determined

Generalisation of the duality to *SO* **gauge groups**

A-model closed strings on an orientifold of the resolved conifold is dual to SO/Sp Chern-Simons theory (Sinha and Vafa)

$$lnZ_{(CS)}^{(SO)}[S^3] = \frac{1}{2}\mathcal{Z}^{(or)} + \mathcal{Z}^{(unor)}$$

Incorporating Wilson loop observables

$$\ln\langle Z(\{U_{\alpha}\},\{V_{\alpha}\})\rangle = \mathcal{F}_{\mathcal{G}}(V) = \frac{1}{2}\mathcal{F}_{\mathcal{R}}^{(or)}(V) + \mathcal{F}^{(unor)}(V)$$

Not clear how to seperate, we showed LHS (Pravina Borhade and PR)

$$\langle Z(\{U_{\alpha}\},\{V_{\alpha}\})\rangle = \exp\left[\sum_{n=1}^{\infty}\sum_{\{R_{\alpha}\}}g_{R_{1},R_{2},\ldots,R_{r}}(q^{n},\lambda^{n})\frac{1}{n}\prod_{\alpha=1}^{r}Tr_{R_{\alpha}}V_{\alpha}^{n}\right]$$

where $g_{R_1,...R_r}(q,\lambda) = \sum_{Q,s} \frac{1}{(q^{1/2}-q^{-1/2})} N_{(R_1,R_2...R_r),Q,s} q^s \lambda^Q$ $N_{(R_1,...R_r),Q,s}$ are integers-how to find oriented contribution?

•Composite representation invariants to extract oriented contribution (Marino)

$$\mathcal{F}_{\mathcal{R}}^{(or)}(V) = \sum_{R,S} \langle W_{(R,S)}[C] \rangle (Tr_R V) (Tr_S V) = \sum_t \mathcal{R}_t[C] Tr_t V$$
$$= \sum_{d=1}^{\infty} \sum_R h_R(q^d, \lambda^d) Tr_R V^d$$

where R and S are representation on the knot C in two Lagrangian submanifolds \mathcal{N}_{ϵ} and $\mathcal{N}_{-\epsilon}$ related by orientifolding action. The composite representation (R, S) has highest weight $\Lambda_R + \Lambda_{\overline{S}}$.

• Obtained invariants of torus knots & links for some composite representations and verified the form of $h_R(q, \lambda)$ (2010, C. Paul, P. Borhade, PR) • Using [r]-colored Kauffman polynomials for figure-eight knot from the study of structural properties of colored Kauffman homologies of knots (Nawata,Zodin,PR,2013) and Morton's result for (\Box, \Box) composite invariant for figure-eight, we verified the integrality structure for h and



Summary

• Elegant computation of [n]-colored HOMFLY for many knots from Chern-Simons approach

• [2,1]-colored HOMFLY distinguishes many mutants including Kinoshita-Terasaka & Conway mutant pair.

• These invariants are definitely useful to verify integrality structures predicted by U(N) and SO topological string duality conjectures Challenging problems

(i) Obtaining closed form expression for $SU(N)_q$ fusion matrices for general *R*'s which could be rectangular or mixed representation.

(ii) Can we get a better understanding of polynomial variable t in homological invariants?

(iii) The results on refined CS for [2, 2p + 1]-torus knots suggests that there could be a combinatorial definition of colored Khovanov invariants for non-torus knots. So far no success!

(iv) Any idea on getting # of solutions to elliptic Nahm equations accounting for integer coefficients in Khovanov invariant for trefoil? Gaitto,Witten (2011)

(v) Appears Kevin Costello work may be applicable to the determine spectral parameter dependent R-matrix for N-vertex models using 4-d twisted supersymmetric field theory.

Due to time constraint, I have not discussed volume conjecture, A-polynomials and their relations to (i) SL(2, C) 3-d gravity(Gukov,2003) and (ii) their appearance in mirror Calabi-Yau geometry (Aganagic,Vafa(2012)-interesting developments.

Thank You