Anomalies and Entanglement Entropy

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based on a work with A. Yarom (Technion) (to appear)



Conformal anomalies in CFT_d:

$$S_A = rac{c_{d-2}}{\epsilon^{d-2}} \cdots + c_0 \log \epsilon + \cdots, \qquad c_0 \sim ext{central charges}$$

Gravitational anomaly in CFT₂ with $c_L \neq c_R$ [Wall 11, Castro-Detournay-Iqbal-Perlmutter 14]



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How about chiral and (mixed-)gravitational anomalies in other dimensions?



- Use a thermodynamic partition function W reproducing flavor and (mixed-)gravitational anomalies
- Evaluate W on the n-fold cover M_n that is an S¹ fibration over an entangling region A



Calculate the variation of the entanglement entropy with W



- Use a thermodynamic partition function W reproducing flavor and (mixed-)gravitational anomalies
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1 Entanglement entropy in QFT

2 Anomalies in CFT₂

3 Thermodynamic partition function method





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Definition of entanglement entropy

Divide a system to A and $B = \overline{A}$: $\mathcal{H}_{tot} = \mathcal{H}_A \otimes \mathcal{H}_B$



Definition

 $S_A = -\mathrm{tr}_A \rho_A \log \rho_A$

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Replica trick

Entanglement entropy

$$S_A = \lim_{n \to 1} \left(\partial_n - 1 \right) W_n$$

- $W_n = -\log Z[\mathcal{M}_n]$: the Euclidean partition function
- \mathcal{M}_n : the *n*-fold cover with a deficit angle $2\pi(1-n)$ around the entangling surface $\Sigma \equiv \partial A$

The *n*-fold cover \mathcal{M}_n

Suppose A is a semi-infinite line in two dimensions

$$A = \{0 \le r < \infty, \ \phi = 0\}$$





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CFT_2 with gravitational anomaly

• CFT₂ with central charges $c_L \neq c_R$ has a gravitational anomaly

$$\nabla_{\mu}T^{\mu\nu} \sim (c_L - c_R)X^{\nu} \neq 0$$

The boost (rotation) of the line A is not a symmetry



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Conformal transformation



• \mathcal{M}_n conformally mapped to a cylinder of circumference $2\pi n$

Cylinder partition function

$$Z[\mathcal{M}_n] = \langle 0 | \exp\left[-\ell H + i\tilde{\theta} P\right] | 0 \rangle$$

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Entanglement entropy

$$S_A = \frac{c_L + c_R}{12} \log(\Lambda/\epsilon) - i \frac{c_L - c_R}{24} \theta$$



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The imaginary part depends on the rotation angle!

(Similar calculation for an interval in CFT_2 [Wall 11,

Castro-Detournay-Iqbal-Perlmutter 14])



The chiral anomaly in 2d takes the form

$$abla_{\mu}J^{\mu} \sim c_s \,\epsilon_{\mu
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Can it be seen through EE in not only CFT but also QFT?



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Thermodynamic partition function

Consider a manifold $\mathcal M$ with U(1) isometry

$$ds_{\mathcal{M}}^2 = e^{2\sigma(x)} (d\tau + a_i(x)dx^i)^2 + g_{ij}(x)dx^i dx^j$$

When au has a period eta, the local temperature is

$$T = e^{-\sigma(x)}/\beta \ ,$$

$$T^{-1}$$

 The thermodynamic partition function
 [Banerjee-Bhattacharya-Bhattacharyya-Jain-Minwalla 12, Jensen-Kaminski-Kovtun-Meyer-Ritz-Yarom 12, Jain-Sharma 12]

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Anomalous partition function

Decompose W into gauge-invariant and anomalous parts

$$W = W_{\text{gauge-inv}} + W_{\text{anom}}$$

■ W_{anom} can be fixed by the anomaly inflow mechanism [Jensen-Loganayagam-Yarom 12]

$$\delta W_{\text{anom}} = \delta \int_{\hat{\mathcal{M}}} i I_{CS} , \qquad \partial \hat{\mathcal{M}} = \mathcal{M}$$

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In two dimensions

• The *n*-fold cover \mathcal{M}_n with U(1) symmetry

$$ds_{\mathcal{M}_n}^2 = dr^2 + r^2 d\tau^2$$

The local temperature

$$T = 1/(2\pi nr) \; ,$$



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$$S_A = \lim_{n \to 1} \left(\partial_n - 1 \right) W[\mathcal{M}_n]$$

Anomalous partition function in 2d

Anomaly polynomial in 2d

$$\mathcal{P}_{2d} = c_s F \wedge F + c_g \operatorname{tr}(\mathcal{R} \wedge \mathcal{R})$$

Chern-Simons form, $\mathcal{P}_{2d} = dI_{CS}$:

$$I_{CS} = c_s A \wedge F + c_g \operatorname{tr} \left[\Gamma \wedge d\Gamma + \frac{2}{3} \Gamma \wedge \Gamma \wedge \Gamma \right]$$

Anomalous partition function:

$$\delta W_{\text{anom}} = i \int_{\partial \mathcal{M}} [c_s \chi F + c_g (\partial_\mu \xi^\nu) d\Gamma^\mu_{\ \nu}] (\delta g_{\mu\nu} = (\mathcal{L}_{\xi} g)_{\mu\nu} , \quad \delta A_\mu = (\mathcal{L}_{\xi} A)_\mu + \partial_\mu \chi)$$

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Anomalies in EE

Rotation
$$\tau \to \tau + \theta$$
, $\xi^{\tau} = \theta$

Gravitational anomaly for a rotation by angle θ

 $\delta_{\theta} S_A \sim i \, c_g \, \theta$

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$$U(1)$$
 gauge variation $A_{\mu} \rightarrow A_{\mu} + \partial_{\mu} \chi$

Chiral anomaly for a gauge variation

$$\delta_{\chi}S_A \sim i \, c_s \, \chi \, \int I$$

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Anomalies in 4d

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 $I_{CS} = A \wedge [c_A F \wedge F + (1 - \alpha)c_m \operatorname{tr}(\mathcal{R} \wedge \mathcal{R})] + \alpha c_m F \wedge j_{GCS}(\Gamma)$

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Mixed anomaly

Take a plane entangling surface



• Magnetic field $B \equiv F_{xy}$ through the entangling surface

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Mixed anomaly

Take a plane entangling surface

$$\mathcal{M}_n = \boxed{\begin{array}{c} \tau \sim \tau + 2\pi n \\ & &$$

• Magnetic field $B \equiv F_{xy}$ through the entangling surface

Mixed anomaly for a rotation by angle θ $\delta_{\theta}S_A \sim i\,\theta\,\alpha\,c_m\,B\,\mathrm{vol}(\mathbb{R}^2)$



Similarly calculation to the mixed anomaly yields

$U(1)^3$ anomaly for a gauge variation

$$\delta_{\chi} S_A \sim i \, \chi \, \left[c_A \, N_{\text{inst}} + 3 c_m (1 - \alpha) \tau[\mathcal{M}] \right]$$

 $N_{\rm inst}:$ Instanton number , $-\tau[{\mathcal M}]:$ Hirtzbruch signature

It is fixed by topological numbers



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- Systematic treatment of anomalies in entanglement entropy using the thermodynamic partition function
- The partition function method does not need conformal symmetry
- Can be applied as well in higher even dimensions (including mixed anomalies), but how about other anomalies such as parity anomaly in odd dimensions?