Taming Defects in N=4 Super-Yang-Mills

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Based on 2003.11016, 2004.09514 [with Shota Komatsu], 2005.07197 and works in progress

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Motivations

 $\mathcal{N} = 4$ SYM is perhaps the most well-studied interacting CFT

- Share many features of strongly coupled gauge theories in four dimensions
- A variety of methods: e.g. supersymmetric localization, integrability, bootstrap
- Close connection to quantum gravity via AdS/CFT
- In recent years, lots of progress in determining correlation functions of local operators

Motivations

Rich family of defects of different co-dimensions

- Sensitive to global structures
- New structure constants for the full operator algebra
- Refined consistency conditions on a full-fledged CFT
- Most works focus on (order-type) Wilson loops. More general disorder-type defects. See also Charlotte's talk

Love defects!

A Program for Defects in $\mathcal{N} = 4$ SYM

Classification of (conformal) defects

Many SUSY examples

(e.g. Wilson-'t Hooft loops, Gukov-Witten surface operators, Gaiotto-Witten walls)

Holographic descriptions

[Karch-Randall, Aharony-Karch, Dewolfe-Freedman-Ooguri, Drukker-Gomis-Matsuura, D'Hoker-Estes-Gutperle,...]

Tools to extract defect observables

Analytic methods + Bootstrap philosophy: formulate consistency conditions, constraints from unitarity and crossing

Implications for the fully extended CFT

See Ibou's talk

Higher form symmetries, anomalies through world-volume couplings, complete data to specify extended CFT?

[Gaiotto-Witten, Gukov-Witten Kapustin, Kapustin-Witten,...]

[Liendo-Rastelli-van Rees, Liendo-Meneghelli, Billo-Goncalves-Lauria-Meineri,...]

Bootstrap

Definition

See also Clay's talk

[Aharony-Seiberg-Tachikawa, Gaiotto-Kapustin-Seiberg-Willet...]

A Program for Defects in $\mathcal{N} = 4$ SYM

Classification of (conformal) defect

Many SUSY examples[Gaiotto-Witten, Gukov-Witten Kapustin,
Kapustin-Witten,...](e.g. Wilson-'t Hooft loops, Gukov-Witten surface operators,
Gaiotto-Witten walls)Holographic descriptions[Karch-Randall, Aharony-Karch, Dewolfe-Freedman-Ooguri,
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Analytic • Tools to extract defect observables

Bootstrap: Formulate consistency condit

Implications for the fully extended CFT

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Bootstrap Definition

[Liendo-Rastelli-van Rees, Liendo-Meneghelli, Goo-Goncalves-Lauria-Meineri,...]

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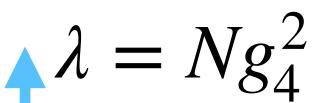
[Aharony-Seiberg-Tachikawa, Gaiotto-Kapustin-Seiberg-Willet...]

Toolkit for SYM

Results (e.g. interface defect) match at common regimes

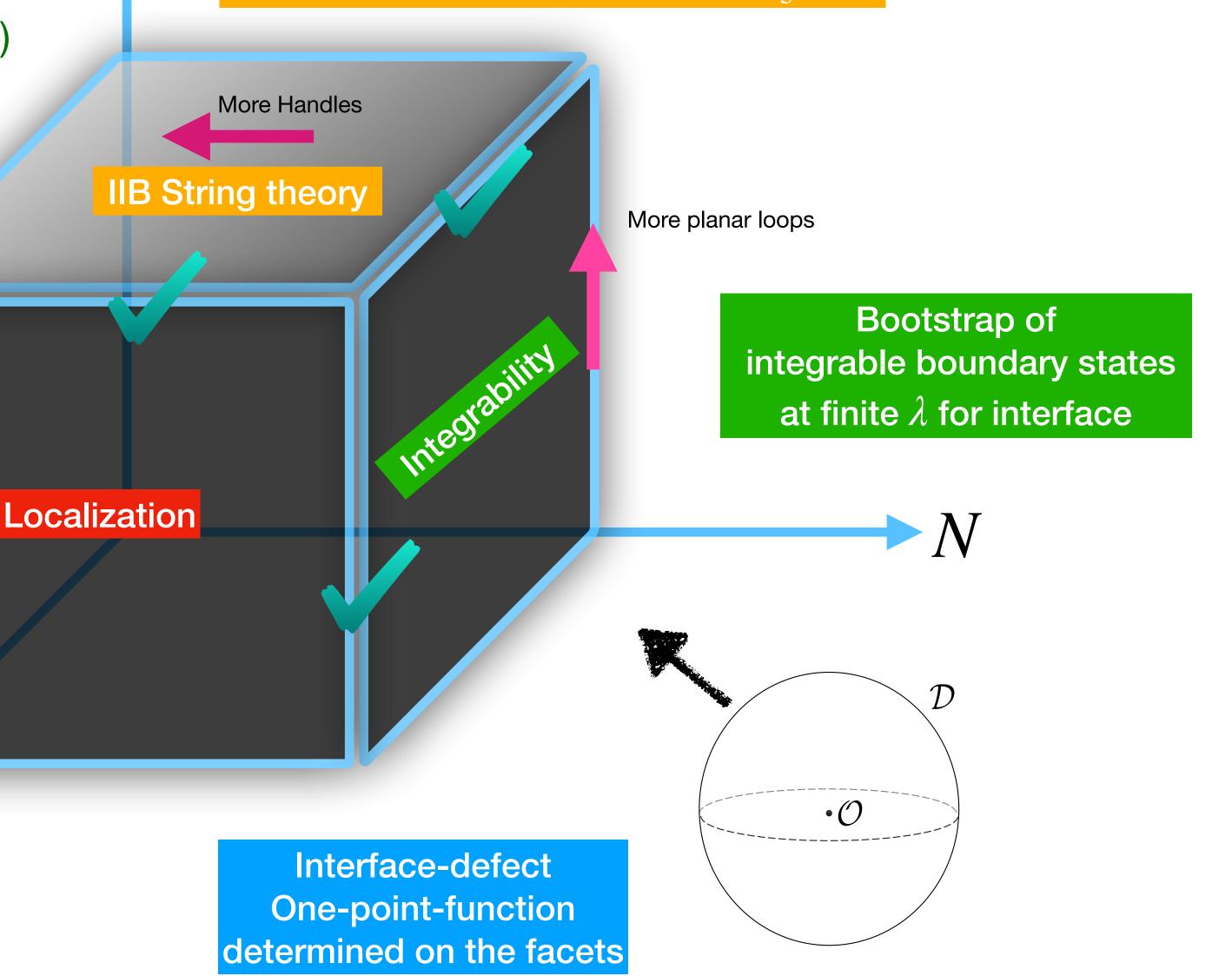
Localization of general 1/16 BPS defect networks Emergent 2d/1d description

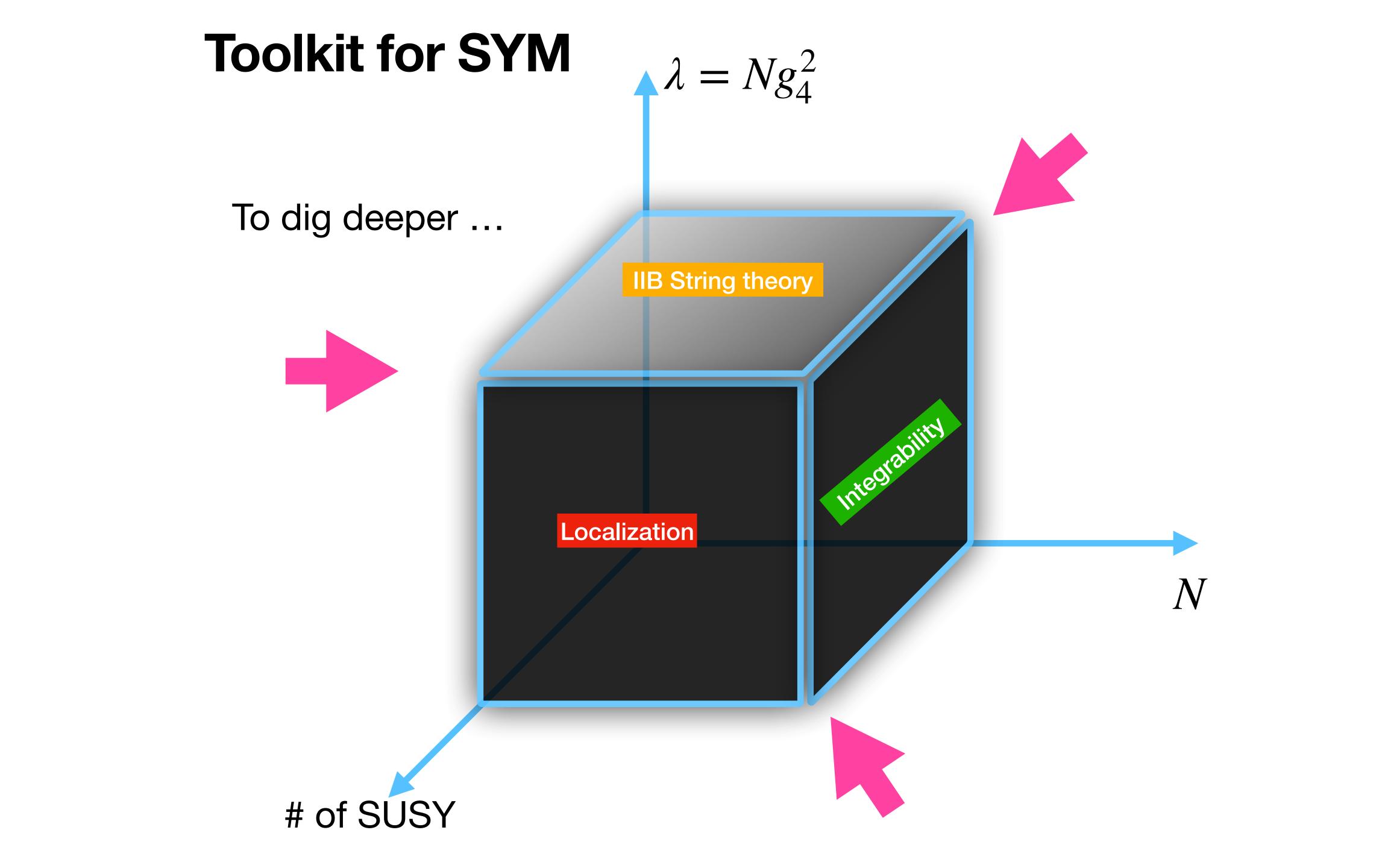


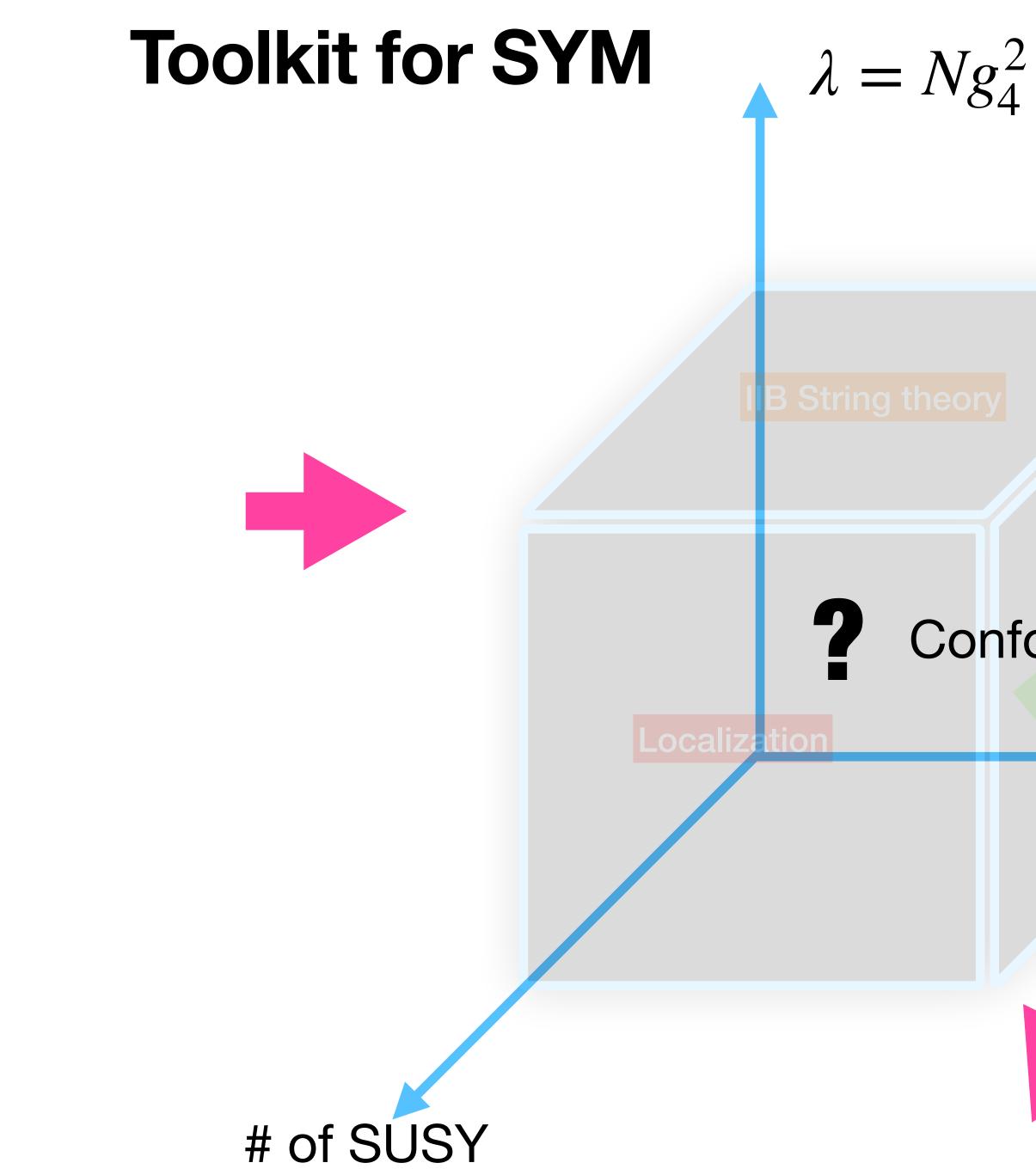


[YW '20, Komatsu-YW '20]

Probe brane analysis in IIB on $AdS_5 \times S^5$









Conformal bootstrap

[Liendo-Rastelli-van Rees, Liendo-Meneghelli, Billo-Goncalves-Lauria-Meineri,...]

N

Emergent 2d/1d effective theory

Defect-Yang-Mills

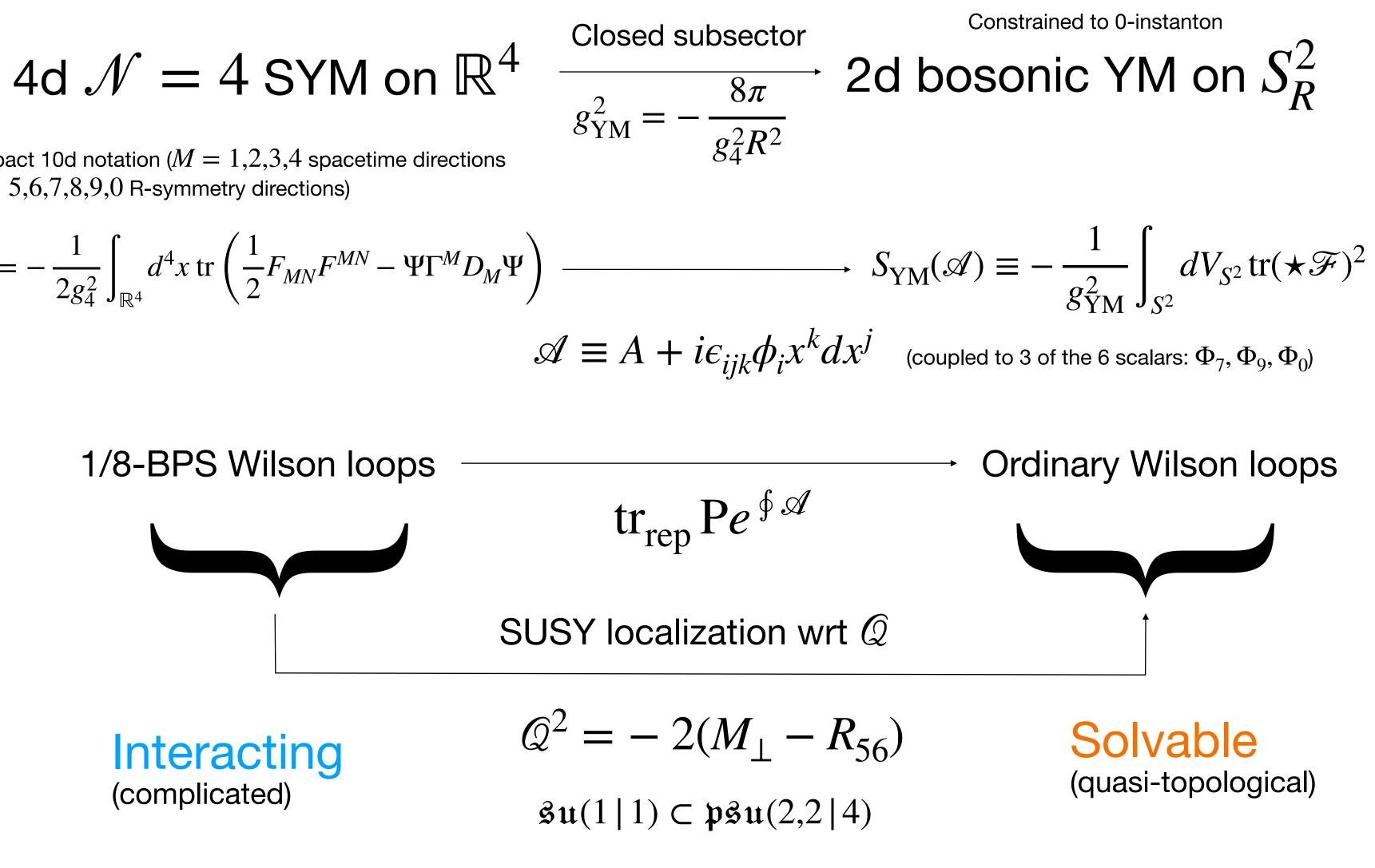
for 1/16 BPS defect observables in $\mathcal{N} = 4$ SYM

Two important ingredients: 2d YM

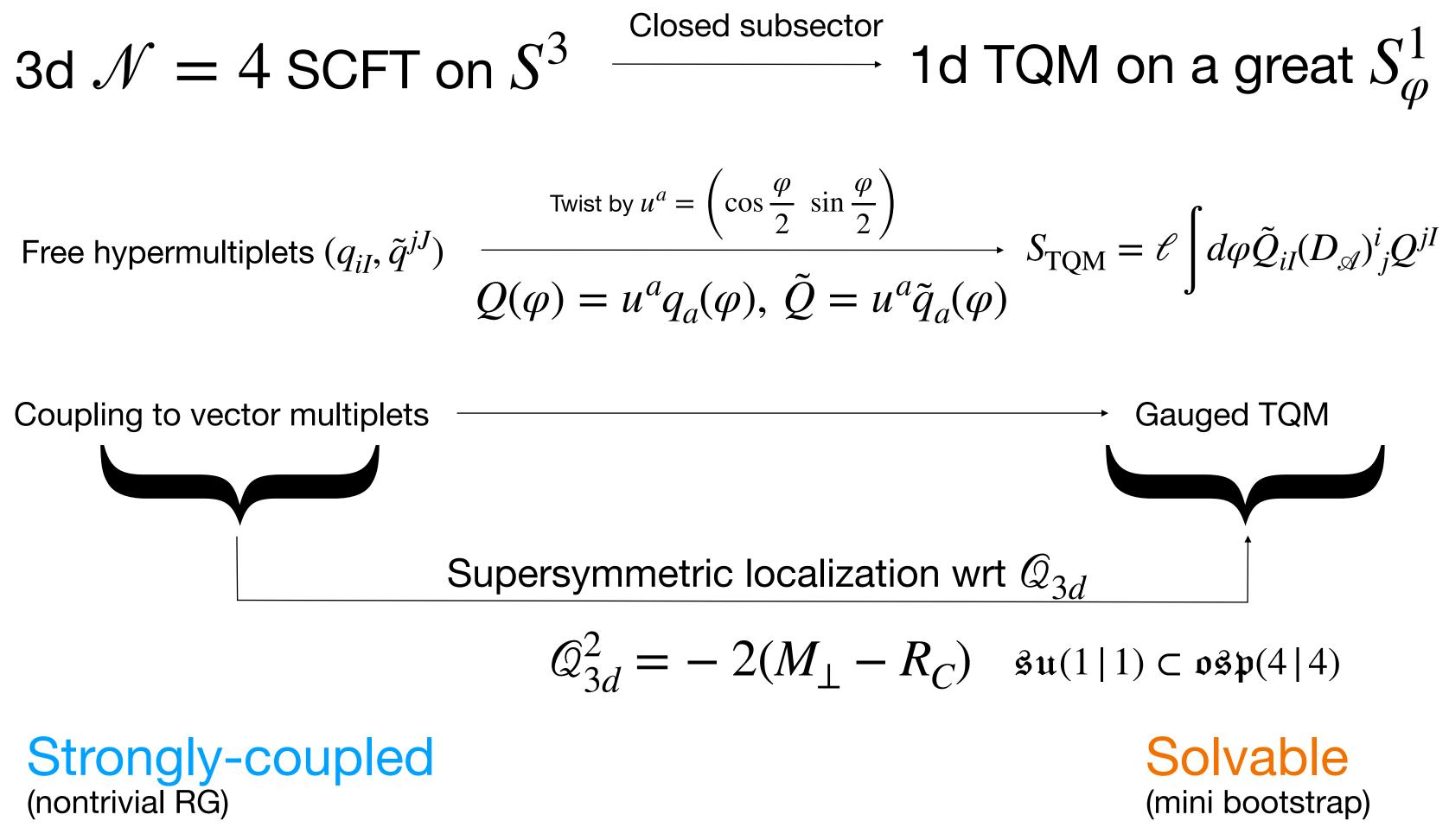
Compact 10d notation (M = 1, 2, 3, 4 spacetime directions M = 5, 6, 7, 8, 9, 0 R-symmetry directions)

$$S_{\text{SYM}} = -\frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) - \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x \operatorname{tr} \left(\frac{1}{2} F_{MN} F^{MN} - \Psi \Gamma^M D_M \Psi \right) + \frac{1}{2g_4^2} \int_{\mathbb{R}^4} d^4 x$$









A general picture

BPS interface?

 $\mathfrak{su}(1 \mid 1) \subset \mathfrak{osp}(4 \mid 1)$

- Study \hat{Q} cohomology among general supersymmetric defect observables in the SYM A rich zoo!
- *Q* localization of the defect-enriched SYM gives rise to defect-enriched 2d YM.

• The 3d supercharge Q_{3d} is the restriction of Q to a half-

$$4) \subset \mathfrak{psu}(2,2|4) \qquad \checkmark$$



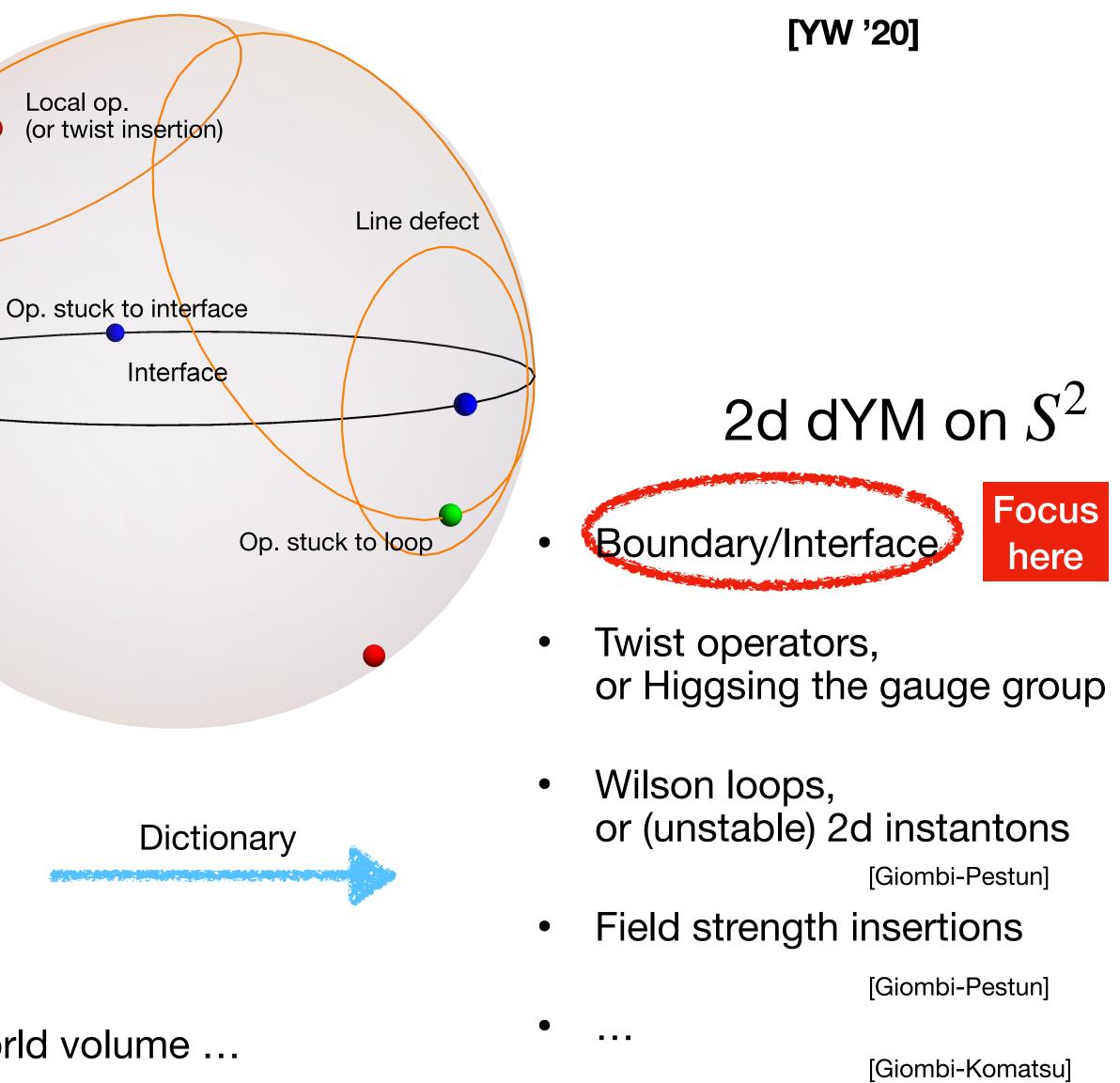
1/16 BPS defect networks

Local op.

In general can have nontrivial topologies (sensitive to global structures): e.g. link of surface operators with line operators.



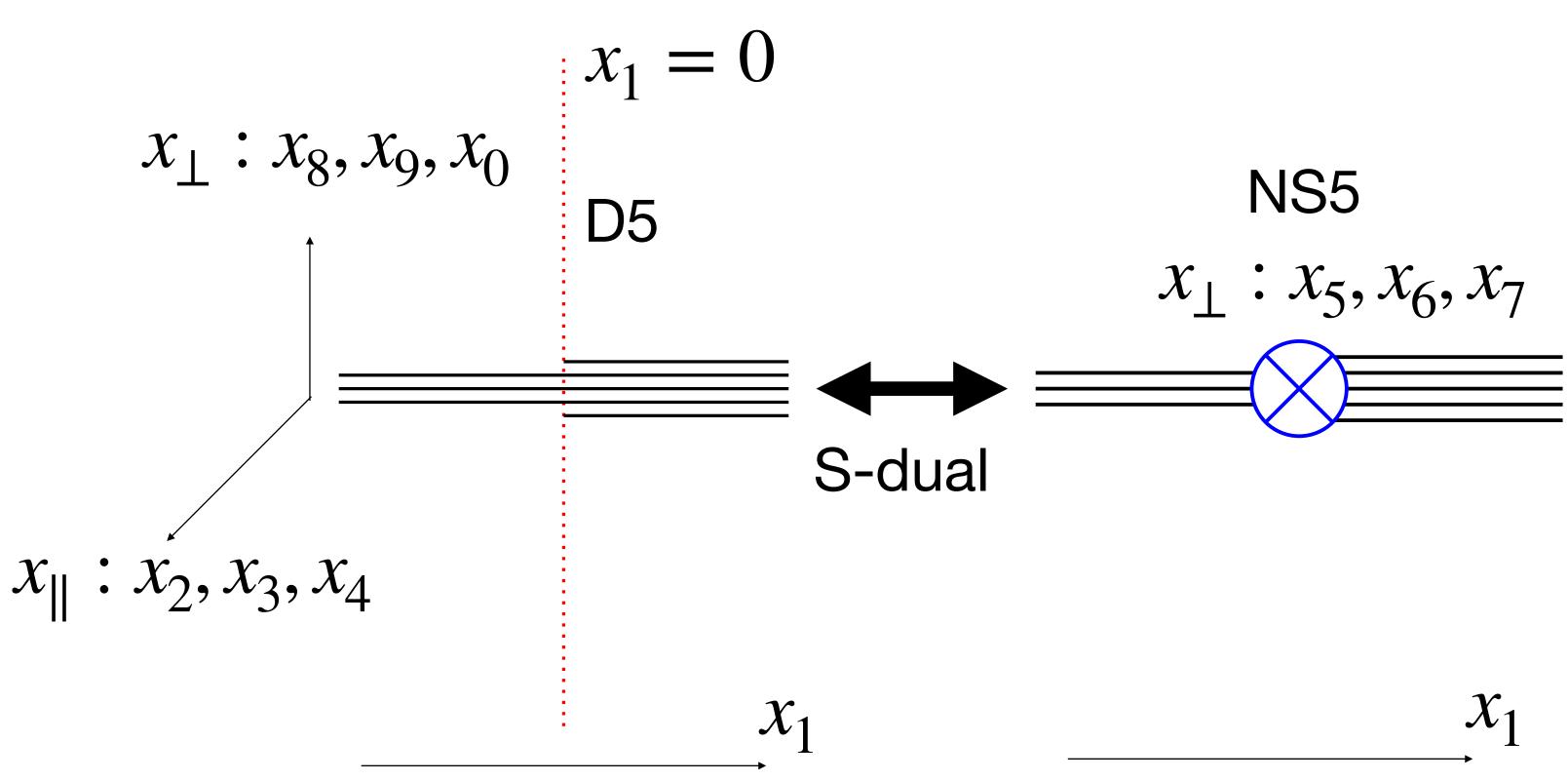
- Boundary/Interface lacksquare
- Surface operators: transverse \bullet or longitudinal to S^2
- Line operators: ulletWilson, 't Hooft lines
- Local operators \bullet
- Local operators on the defect world volume ... \bullet



Application D5-brane interface defects

D5-brane interface for SYM: brane picture

where k D3 branes can end



• Engineered by intersection of one D5 brane with D3 branes

D5-brane interface for SYM: field theory

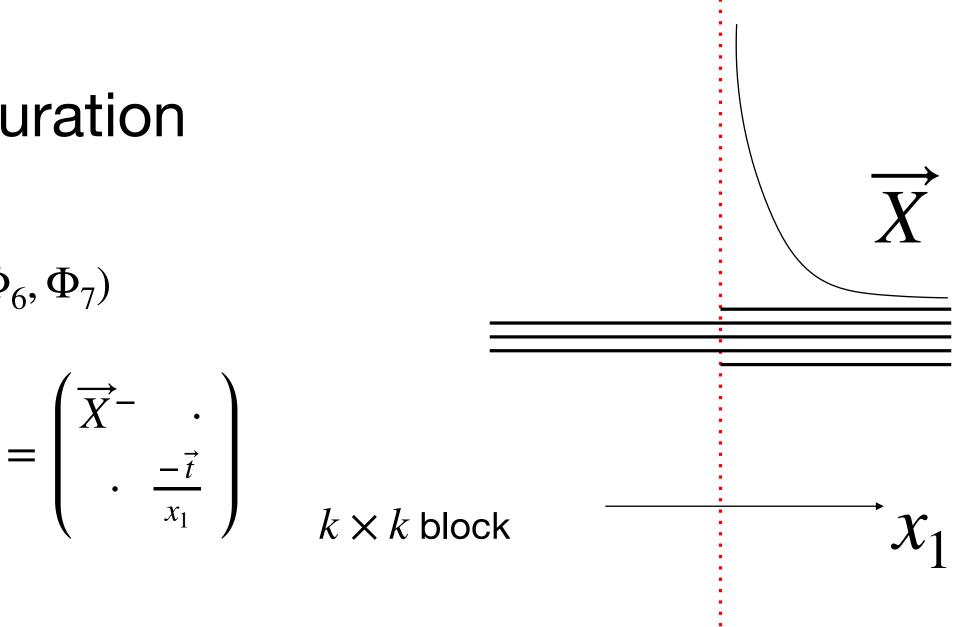
- Singular Nahm pole configuration (related to usual N.P. bc by folding trick)

$$\overrightarrow{X} = (\Phi_8, \Phi_9, \Phi_0), \quad \overrightarrow{Y} = (\Phi_5, \Phi_6,$$

$$A^+_{\mu} = \begin{pmatrix} A^-_{\mu} & \\ & \cdot \\ & \cdot \end{pmatrix}, \quad \overrightarrow{Y^+} = \begin{pmatrix} \overrightarrow{Y^-} & \\ & \cdot \\ & \cdot \end{pmatrix}, \quad \overrightarrow{X^+} =$$

Nahm
$$\begin{cases} [t_i, t_j] = \epsilon_{ijk} t_k \\ \text{Pole} \\ \text{Config} \end{cases} \begin{cases} t_3^{k \times k} = -\frac{i}{2} \text{Diag} \end{cases}$$

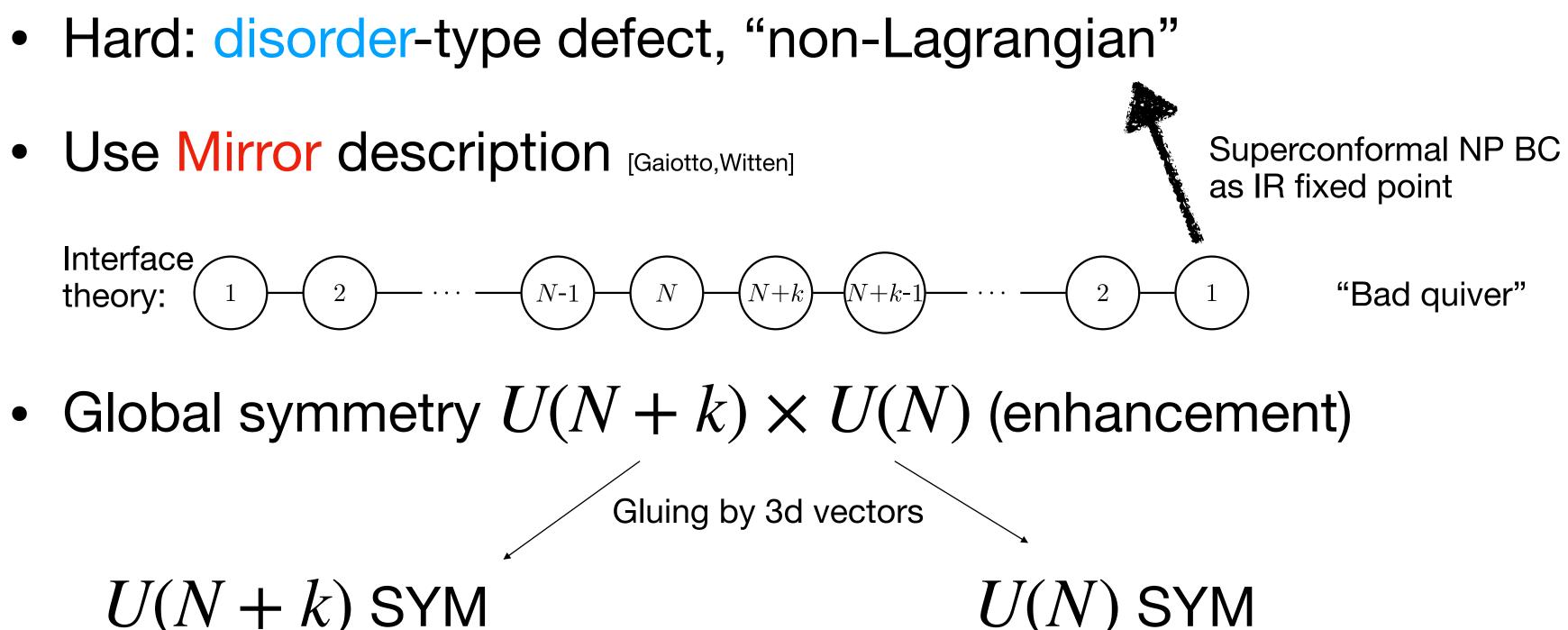
• Interpolate between U(N) and U(N + k) SYM at the same coupling



given by SU(2) reps

g[k-1, k-3, ..., 1-k]

D5-brane interface for SYM: field theory



U(N+k) SYM with Dirichlet BC

For SUSY observables this suffices by non-renormalization

with Dirichlet BC

• A UV Lagrangian for SYM on S^4 coupled to the D5 interface at equator that preserves $\mathcal{Q} \in \mathfrak{ogp}(4|4,\mathbb{R}) \subset \mathfrak{pgu}(2,2|4)$

[YW '20, Komatsu-YW '20]

Defect-Yang-Mills for D5-brane interface

• Run localization by \hat{Q} to find 2d/1d description

$$Z_{\rm dYM} = \int_{S_{\rm YM}^2} D\mathscr{A} |_{HS_{-}^2} D\mathscr{A}' |_{HS_{+}^2} L$$

$$S_{\rm YM}(\mathscr{A}) \equiv -\frac{1}{g_{\rm YM}^2} \int_{HS_{-}^2} dV_{S^2} \operatorname{tr}(\star \mathscr{F})^2$$

in the presence of the D5-brane interface

 $DQD\tilde{Q}|_{S_{\text{TOM}}^1} e^{-S_{\text{YM}}(\mathscr{A}) - S_{\text{YM}}(\mathscr{A}') - S_{\text{TQM}}(\mathscr{A}, \mathscr{A}', Q, \tilde{Q})}$

Ready to use to compute general 1/16 BPS SYM observables

[YW '20, Komatsu-YW '20]

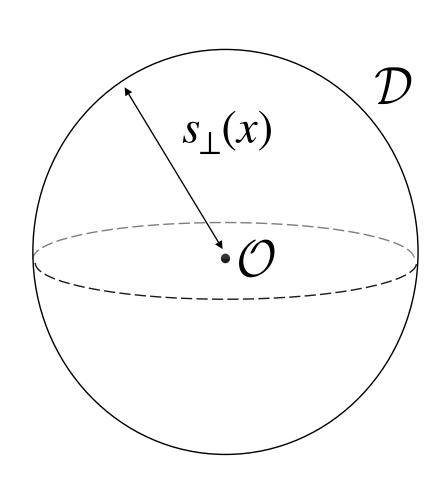
Defect One-point function

Basic defect observables: higher point correlators determined by bulk OPE

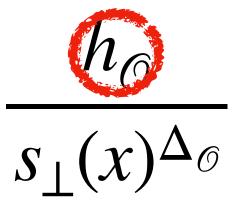
$$\langle \mathcal{O}(x) \rangle_{HS^4} =$$

$$\langle \mathcal{O}_J \rangle_{\mathrm{dSYM}} = \frac{(-i)^J}{Z_{\mathrm{dYM}}} \int D \mathscr{A} D$$

- Compute by first integrate out 1d fields (Q, Q)
- techniques



[YW '20, Komatsu-YW '20]



Fixed by conformal Ward identity

Take $\mathcal{O}_I \equiv \operatorname{tr}(\Phi_7 + i\Phi_8)^J$

half-BPS

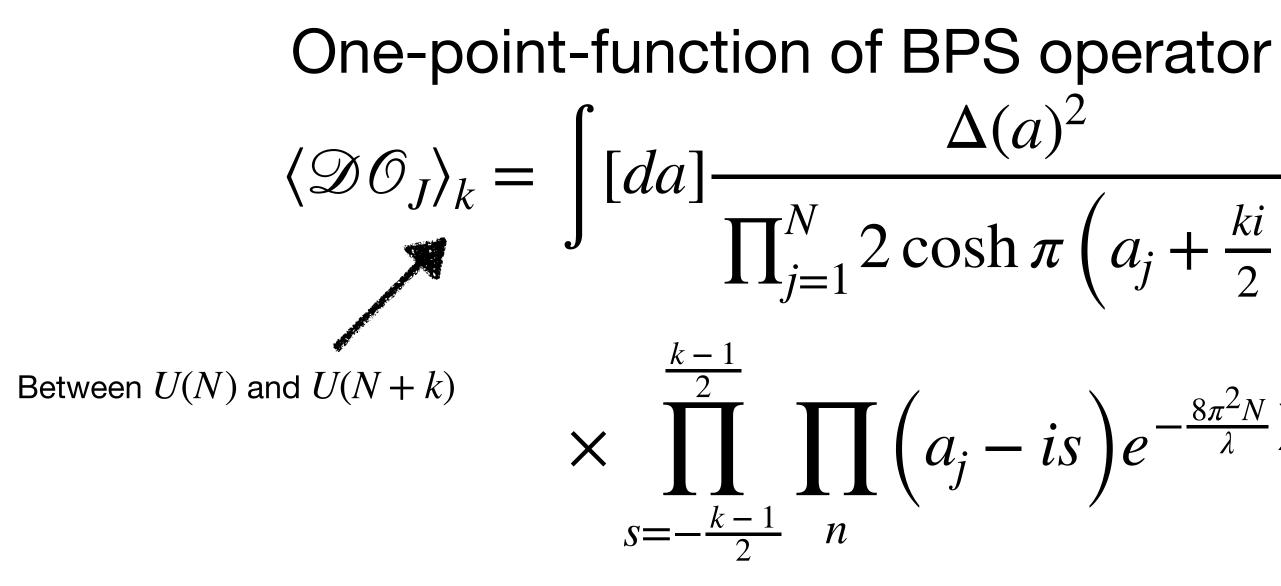
 $OQD\tilde{Q} \operatorname{tr}(\star \mathscr{F})^{J} e^{-S_{\mathrm{YM}}(\mathscr{A}) - S_{\mathrm{TQM}}(Q,\tilde{Q},\mathscr{A})}$

Denote gauge fields on the two hemispheres collectively by \mathscr{A}

• Obtain matrix model for $\langle O_J \rangle_{dSYM}$ by 2d gauge theory



Full D5 brane matrix model at any (N, λ)



One-point-function of the interface

$$\langle \mathcal{D} \rangle_{k} = \int [da] \frac{\Delta(a)^{2}}{\prod_{j=1}^{N} 2\cosh \pi \left(a_{j} + \frac{ki}{2}\right)} \prod_{s=-\frac{k-1}{2}}^{\frac{k-1}{2}} \prod_{j=1}^{N} \left(a_{j} - is\right) e^{-\frac{8\pi^{2}N}{\lambda} \sum_{i=1}^{N} a_{i}^{2}}$$

deg-J polynomial from G.S. procedure $\langle \mathcal{D}\mathcal{O}_J \rangle_k = \int [da] \frac{\Delta(a)^2}{\prod_{j=1}^N 2\cosh \pi \left(a_j + \frac{ki}{2}\right)} \left| \sum_{j=1}^N f_J(a_j) + \sum_{s=-\frac{k-1}{2}}^{\frac{k-1}{2}} f_J(is) \right|$

$$is)e^{-\frac{8\pi^2N}{\lambda}\sum_{i=1}^Na_i^2}$$

Normalize by taking ratio

[Komatsu-YW '20]

Large N One-point functions and checks

$$\begin{split} \langle \mathcal{O}_J \rangle_k &= \frac{i^J}{2^{\frac{J}{2}}\sqrt{J}} \left(\text{integ} \right. \\ & \left\{ \oint \frac{dx}{2\pi i} g\left(1 - \frac{1}{x^2} \right) \right. \\ & \left\{ \oint \frac{dx}{2\pi i} g\left(1 - \frac{1}{x^2} \right) \right. \\ & \left\{ \oint \frac{dx}{2\pi i} g\left(1 - \frac{1}{x^2} \right) \right. \\ & \left. \sup = \sum_{s = -\frac{k-1}{2}}^{\frac{k-1}{2}} \left(x_s^J + \delta_{J,2} \right) \right. \end{split}$$

Large λ : check with IIB string theory answers [Nagasaki-Yamaguchi]

- Probe D5 brane on $AdS_4 \times S^2$ with k units of w.v. flux

Small λ : check perturbation theory [Buhl-Mortensen-de Leeuw-Ipsen-Kristjansen-Wilhelm,Kristjansen-Müller-Zarembo]

[YW '20, Komatsu-YW '20] gral + sum) $g = \frac{\sqrt{\lambda}}{4\pi}$ $x^{J}\pi \operatorname{coth}\left[\pi g\left(x+\frac{1}{x}\right)\right]$ $k\in 2\mathbb{Z}+1$ $x^J \pi \tanh \left| \pi g \left(x + \frac{1}{x} \right) \right| \quad k \in 2\mathbb{Z}$

$$x_u = \frac{u + \sqrt{u^2 - 4g^2}}{2g}$$

Zhukovsky Hint for integrability!

One-point function from w.v. couplings to closed string KK modes

Bootstrap approach to

Integrable boundary states

for interface defects in $\mathcal{N} = 4$ SYM

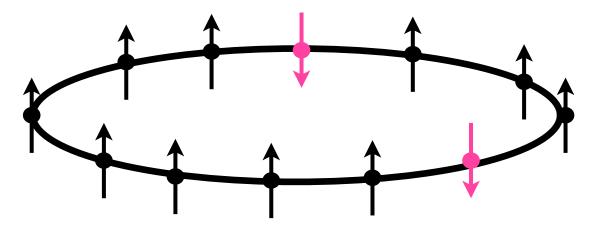
D5 brane and Integrable spin chain

Traditional Spin chain approach [Minahan-Zarembo]: $\mathcal{O} = tr(\dots ZYZ\dots ZYZ\dots)$ "Y magnons"

- $|\mathcal{D}\rangle$ represented by a matrix-product state
- $|O\rangle$ represented by Bethe states, labelled by rapidities (momentum) u_m of the magnons
- One-point function from overlap
- Limitation to weak coupling (nearest-neighbor) interactions, fixed length)

[Buhl-Mortensen, Gimenez Grau, Ipsen, de Leeuw, Linardopoulos, Kristjansen, Pozsgay, Volk, Widen, Wilhelm, Zarembo ...]





D5 brane and Integrable field theory

Weak coupling evidences for defect observable:

2) Overlaps involve ratio of Gaudin determinants



Boundary state \mathscr{D} D5-brane interface defect \leftarrow on IIB worldsheet in $AdS_5 \times S^5$ Large N (lightcone G.S.)

> • Single trace operators (non-BPS) - closed string state $|\Psi\rangle$ • One-point function from overlap $\langle \mathcal{O} \rangle_{\Im} = \langle \mathcal{D} | \Psi \rangle$

1) Rapidities pair up $(u_1, -u_1, u_2, -u_2, \dots, u_{M/2}, -u_{M/2})$, i.e. no particle production

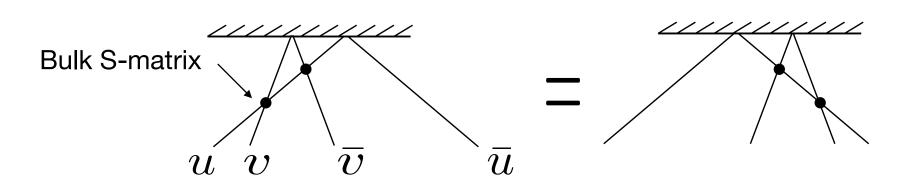
[Jiang-Komatsu-Vescovi]



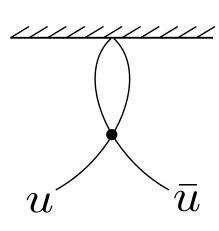
[Komatsu-YW '20]

General Strategy for non-perturbative boundary state

- two particle overlaps e.g. $\langle \mathcal{D} | Y(u) Y(\bar{u}) \rangle$
- Explore symmetry constraints under
 - $SU(2 | 1) \times SU(2 | 1) \subset SU(2 | 2) \times SU(2 | 2)_{ce}$
- Impose consistency conditions \bullet



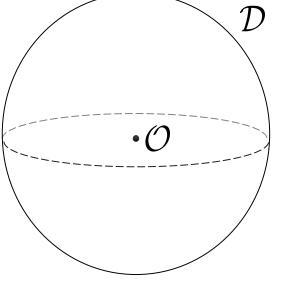
Boundary Young-Baxter

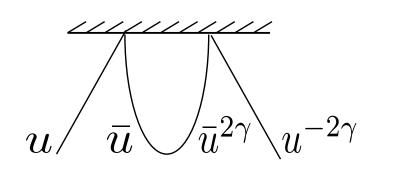


Watson (crossing)

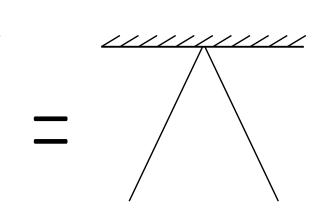
Apply bootstrap philosophy to integrable boundary state [Ghoshal-Zamolodchikov, Jiang-Komatsu-Vescovi]

Integrability implies general overlaps determined by





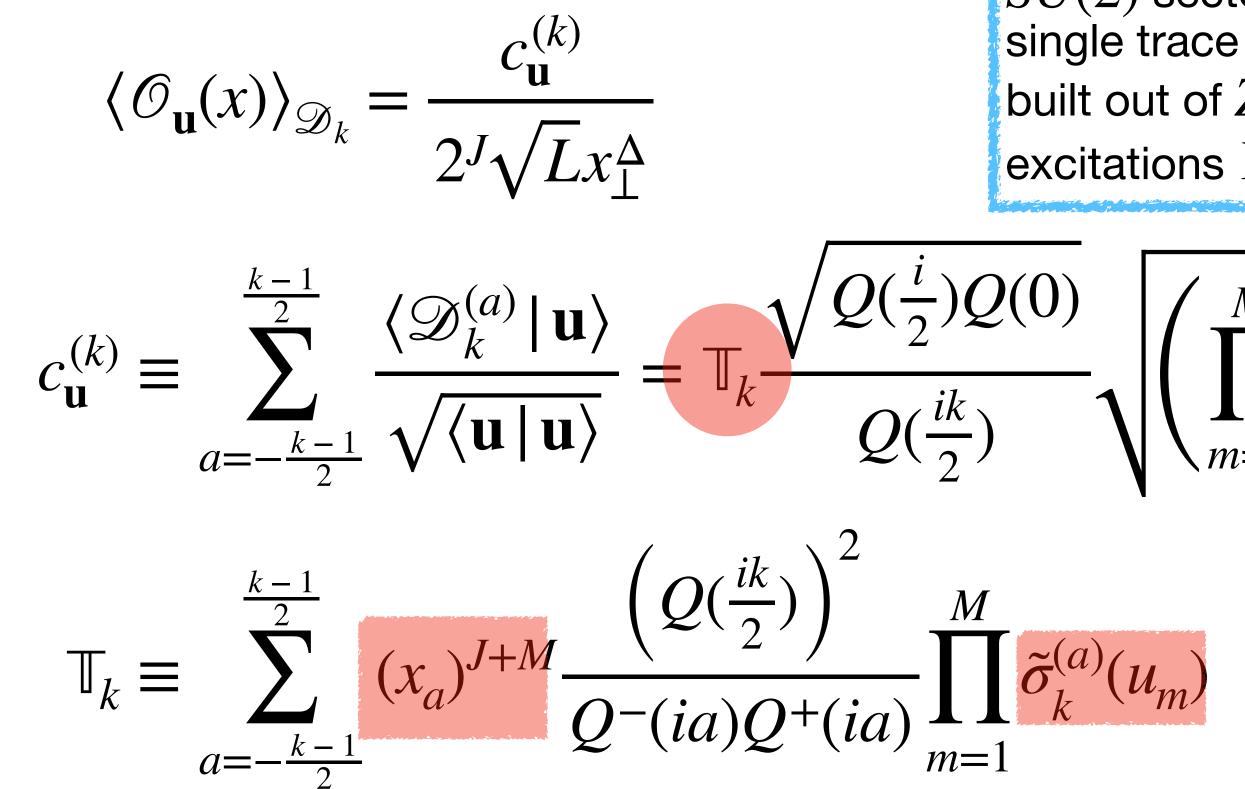
Boundary unitarity



[Komatsu-YW '20]

See also [Gombor-Bajnok]

Large *J* R-charge non-BPS one-point function



- Zero magnon M = 0 case gives large J BPS one point function
- Weak coupling expansion agrees with spin chain method

SU(2) sector: single trace operators of length L built out of $Z = \Phi_7 + i\Phi_8$ and excitations $Y = \Phi_5 + i\Phi_9$

$$\frac{Q(\frac{i}{2})Q(0)}{Q(\frac{ik}{2})} \sqrt{\left(\prod_{m=1}^{M} \sigma_{B}(u_{m})\right)} \frac{\det G_{+}}{\det G_{-}}$$

Extend to SO(6) in [Gombor-Bajnok]

 $Q(u) = \prod^{m} (u - u_m)$ m=

Important: excited boundary stares (bound states with bdy)

More checks in [Kristjansen-Müller-Zarembo]

• Finite J (non-pert.) from wrapping corrections (using TBA and analytic continuation)



Summary

- Localization: 1/16 BPS, arbitrary N, arbitrary λ
- Integrability: non-BPS, large N, arbitrary λ
- Application: Interface-defect one-point-function

A toolkit for defects in $\mathcal{N} = 4$ SYM

More applications and future directions

- Correlation functions of both bulk and defect operators [YW '20]
- Unorientable spacetimes, e.g. SYM \mathbb{RP}^4 governed by YM on \mathbb{RP}^2 [YW '20]
- Surface defects (linked with line defects)
- Discrete theta angles and global structures of gauge theories: from 4d SYM to 2d dYM
- Classification of integrable boundary states in IIB string on $AdS_5 \times S^5$ e.g. D7, D-instanton, more exotic ones?
- Topological defects (bi-branes) on w.s. sigma model are known to generate boundary states [Fuchs-Schweigert-Waldorf, Kojita-Maccaferri-Masuda-Schnabl] Relation to SYM interfaces?
- *TT* deformation: 2d YM has solvable TT [Sentilli-Tierz] learn about $T\bar{T}$ deformations 4d $\mathcal{N} = 4$ SYM [Caetano-Rastelli-Peelaers]
- $SL(2,\mathbb{Z})$ properties:
 - Recent progress in $SL(2,\mathbb{Z})$ invariant four point functions [Chester-Green-Pufu-YW-Wen] including defects which transform nontrivially under $SL(2,\mathbb{Z})$?
- Extend the defect analysis to other AdS/CFT pairs?

e.g. deriving AdS_3/CFT_2 beyond local correlators by incorporating defects [See Rajesh's talk]

Thank you!