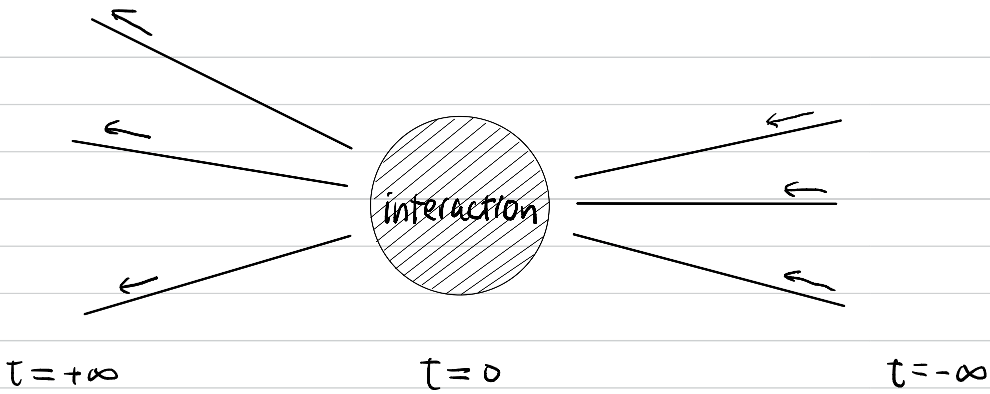
Recap



$$|\psi\rangle_{in} \longleftarrow |\psi, \text{free}\rangle$$

$$|\chi, \text{free}\rangle \longleftarrow |\chi\rangle_{out}$$

$$\langle \chi, \text{free} | S | \psi, \text{free}\rangle := {}_{out}\langle \chi | \psi \rangle_{in}$$

$\mathcal{O} = \mathcal{O}^\dagger$ a scalar opr s.t. $\langle 0 | \mathcal{O}(x) | 0 \rangle = 0$

$$\langle p | \mathcal{O}(x) | 0 \rangle = \sqrt{Z}\, e^{ip\cdot x}$$

- wave packet $f(x)$, $\mathcal{O}_f(t) = \frac{-i}{\sqrt{Z}} \int d^{d-1}\mathbf{x}\, f(t,\mathbf{x}) \overleftrightarrow{\partial_t} \mathcal{O}(t,\mathbf{x})$

$$\mathcal{O}_f(t)|0\rangle \xrightarrow{t\to\pm\infty} |f\rangle = |f\rangle_{in} = |f\rangle_{out}$$

$$\mathcal{O}_f(t)^\dagger |0\rangle \xrightarrow{t\to\pm\infty} 0$$

- $f_1, \dots, f_n$ widely separated at $t \to \pm\infty$:

$$|f_1, \dots, f_n\rangle_{in/out} = \lim_{T\to\mp\infty} \mathcal{O}_{f_1}(T)\cdots\mathcal{O}_{f_n}(T)|0\rangle$$

$$\overset{!}{=} \lim_{T_1\to\mp\infty} \cdots \lim_{T_n\to\mp\infty} \mathcal{O}_{f_1}(T_1)\cdots\mathcal{O}_{f_n}(T_n)|0\rangle$$

LSZ reduction formula

$$\langle g_1, \cdots, g_n \text{ free} \,|\, S \,|\, f_1, f_2 \text{ free} \rangle = {}_{\text{out}}\langle g_1, \cdots, g_n \,|\, f_1, f_2 \rangle_{\text{in}} = ?$$

$$\mathcal{O}_f(-T) - \mathcal{O}_f(T) = -\int_{-T}^{T} dt \, \frac{d}{dt} \mathcal{O}_f(t)$$

$$= \int_{-T}^{T} dt \, \frac{i}{\sqrt{Z}} \int_{\mathbb{R}^{d-1}} d^{d-1}\mathbf{x} \, \underbrace{\partial_t \left(f \partial_t \mathcal{O} - \partial_t f \, \mathcal{O} \right)}_{f \partial_t^2 \mathcal{O} - \underbrace{\partial_t^2 f}_{= (\nabla^2 - m^2) f} \mathcal{O}} (t, \mathbf{x})$$

[as $f(t, \mathbf{x}) \to 0$ as $|\mathbf{x}| \to \infty$,

spatial partial integration is allowed.

$$= \int_{-T}^{T} dt \int_{\mathbb{R}^{d-1}} d^{d-1}\mathbf{x} \, f(t, \mathbf{x}) \frac{i}{\sqrt{Z}} \underbrace{\left(\partial_t^2 - \nabla^2 + m^2 \right)}_{\partial_x^2 + m^2} \mathcal{O}(t, \mathbf{x})$$

$$=: \int_{-T}^{T} d^d x \, f(x) \frac{i}{\sqrt{Z}} (\partial^2 + m^2) \mathcal{O}(x)$$

Taking its adjoint

$$\mathcal{O}_f(T)^{\dagger} - \mathcal{O}_f(-T)^{\dagger} = \int_{-T}^{T} d^d x \, f(x)^* \frac{i}{\sqrt{Z}} (\partial^2 + m^2) \mathcal{O}(x)$$

Consider $X_{T_1 \cdots T_n, T_1', T_2'} :=$

$$\prod_{i=1}^{n} \int_{-T_i}^{T_i} d^4y_i \, g_i(y_i)^* \frac{i}{\sqrt{Z}} (\partial_{y_i}^2 + m^2) \prod_{j=1}^{2} \int_{-T_j'}^{T_j'} d^4x_j \, f_j(x_j) \frac{i}{\sqrt{Z}} (\partial_{x_j}^2 + m^2)$$

$$\langle 0 | T \, \mathcal{O}(y_1) \cdots \mathcal{O}(y_n) \mathcal{O}(x_1) \mathcal{O}(x_2) | 0 \rangle .$$

$$\int_{-T_1}^{T_1} d^4y_1 \, g_1(y_1)^* \frac{i}{\sqrt{Z}} (\partial_{y_1}^2 + m^2) \langle 0 | T \, \mathcal{O}(y_1) \cdots \mathcal{O}(x_2) | 0 \rangle$$

$$\xrightarrow[T_1 \to \infty]{} \langle 0 | \mathcal{O}_{g_1}(\infty)^\dagger \, T(\mathcal{O}(y_2) \cdots \mathcal{O}(x_2)) | 0 \rangle - \langle 0 | T(\mathcal{O}(y_2) \cdots \mathcal{O}(x_2)) \underbrace{\mathcal{O}_{g_1}(-\infty)^\dagger | 0 \rangle}_{= 0}$$

Thus

$$X_{\vec{T}, \vec{T}'} \xrightarrow{T_1 \to \infty} \prod_{i=2}^{n} \int_{-T_i}^{T_i} \cdots \prod_{j=1}^{2} \int_{-T_j'}^{T_j'} \cdots \langle 0 | \mathcal{O}_{g_1}(\infty)^\dagger \, T(\mathcal{O}(y_2) \cdots \mathcal{O}(x_2)) | 0 \rangle$$

Repeating this for $T_2, \cdots, T_n$, we find

$$X_{\vec{T}, \vec{T}'} \xrightarrow{T_1, T_2, \cdots T_n \to \infty} \prod_{j=1}^{2} \int_{-T_j'}^{T_j'} \cdots \underbrace{\langle 0 | \mathcal{O}_{g_1}(\infty)^\dagger \cdots \mathcal{O}_{g_n}(\infty)^\dagger}_{{}_{\text{out}}\langle g_1, \cdots, g_n |} T(\mathcal{O}(x_1) \mathcal{O}(x_2)) | 0 \rangle$$

Further limits :

$$\xrightarrow{T_1' \to \infty} \int_{-T_2'}^{T_2'} d^d x_2\, f_2(x_2) \frac{i}{\sqrt{Z}} \left(\partial_{x_2}^2 + m^2\right)$$

$$\left({}_{out}\langle g_1, \cdots, g_n | \mathcal{O}(x_2)\, \mathcal{O}_{f_1}(-\infty) | 0 \rangle - {}_{out}\langle g_1, \cdots, g_n | \mathcal{O}_{f_1}(\infty)\, \mathcal{O}(x_2) | 0 \rangle \right)$$

$$\xrightarrow{T_2' \to \infty} {}_{out}\langle g_1, \cdots g_n | \mathcal{O}_{f_2}(-\infty)\, \mathcal{O}_{f_1}(-\infty) | 0 \rangle$$

$|f_1, f_2\rangle_{in}$

$|f_1, f_2\rangle_{out}$

$$- {}_{out}\langle g_1, \cdots g_n | \mathcal{O}_{f_2}(\infty)\, \underbrace{\mathcal{O}_{f_1}(-\infty) | 0 \rangle}_{\mathcal{O}_{f_1}(\infty) | 0 \rangle}$$

$|f_1, f_2\rangle_{out}$

$$- {}_{out}\langle g_1, \cdots, g_n | \mathcal{O}_{f_1}(\infty)\, \underbrace{\mathcal{O}_{f_2}(-\infty) | 0 \rangle}_{\mathcal{O}_{f_1}(\infty) | 0 \rangle}$$

$|f_1, f_2\rangle_{out}$

$$+ {}_{out}\langle g_1, \cdots, g_n | \mathcal{O}_{f_1}(\infty)\, \mathcal{O}_{f_2}(\infty) | 0 \rangle$$

$$= {}_{out}\langle g_1, \cdots, g_n | f_1, f_2 \rangle_{in} - {}_{out}\langle g_1, \cdots, g_n | f_1, f_2 \rangle_{out}$$

$$\cancel{- {}_{out}\langle g_1, \cdots, g_n | f_1, f_2 \rangle} \cancel{+ {}_{out}\langle g_1, \cdots, g_n | f_1, f_2 \rangle_{out}}$$

$$= \langle g_1, \cdots, g_n, \text{free} | S | f_1, f_2, \text{free} \rangle - \langle g_1, \cdots, g_n, \text{free} | f_1, f_2, \text{free} \rangle$$

We obtained

$$\langle g_1,\cdots,g_n,\text{free}|S|f_1,f_2,\text{free}\rangle$$

$$= \langle g_1,\cdots,g_n,\text{free}|f_1,f_2,\text{free}\rangle$$

$$+ \prod_{i=1}^{n}\int d^dy_i\, g_i(y_i)^* \frac{i}{\sqrt{Z}}(\partial_{y_i}^2+m^2) \prod_{j=1}^{2}\int d^dx_j\, f_j(x_j)\frac{i}{\sqrt{Z}}(\partial_{x_j}^2+m^2)$$

$$\langle 0|T\,\mathcal{O}(y_1)\cdots\mathcal{O}(y_n)\mathcal{O}(x_1)\mathcal{O}(x_2)|0\rangle$$

S-matrix is described by correlation functions.

This is the LSZ reduction formula.

Lehmann, Symanzik, Zimmermann

There is also a formula for $\langle g_1,\cdots,g_n,\text{free}|S|f_1,\cdots,f_m,\text{free}\rangle$

which is a bit more complicated.

It simplifies if $\langle g_a|f_b\rangle=0 \quad \forall a,b$:

$$\langle g_1,\cdots,g_n,\text{free}|S|f_1,\cdots,f_m,\text{free}\rangle$$

$$= \prod_{i=1}^{n}\int d^dy_i\, g_i(y_i)^* \frac{i}{\sqrt{Z}}(\partial_{y_i}^2+m^2) \prod_{j=1}^{m}\int d^dx_j\, f_j(x_j)\frac{i}{\sqrt{Z}}(\partial_{x_j}^2+m^2)$$

$$\langle 0|T\,\mathcal{O}(y_1)\cdots\mathcal{O}(y_n)\mathcal{O}(x_1)\cdots\mathcal{O}(x_m)|0\rangle.$$

In Fourier modes :

$$\langle 0 | T\, \mathcal{O}(x_1) \cdots \mathcal{O}(x_s) | 0 \rangle = \int \prod_{i=1}^{s} \frac{d^d p_i}{(2\pi)^d}\, e^{-i p_i x_i}\, G(p_1, \cdots, p_s)$$

$$\langle g_1, \cdots, g_n, \text{free} | (S-1) | f_1, f_2, \text{free} \rangle$$

$$= \int \prod_{i=1}^{n} d^d y_i \frac{d^d l_i}{(2\pi)^d}\, g_i(y_i)^* e^{-i l_i y_i} \frac{i}{\sqrt{Z}} (-l_i^2 + m^2)$$

$$\prod_{j=1}^{2} d^d x_j \frac{d^d k_j}{(2\pi)^d}\, f_j(x_j)\, e^{-i k_j x_j} \frac{i}{\sqrt{Z}} (-k_j^2 + m^2)$$

$$G(l_1, \cdots, l_n, k_1, k_2)$$

$$\left[\begin{aligned} & f(x) = \int \frac{d^{d-1} \mathbb{p}}{(2\pi)^{d-1} 2\omega_{\mathbb{p}}}\, e^{-i p_{\mathbb{p}} x}\, \tilde{f}(\mathbb{p}) \\ & \int d^d x \frac{d^d k}{(2\pi)^d} f(x)\, e^{-ikx} F(k) = \int \frac{d^{d-1} \mathbb{p}}{(2\pi)^{d-1} 2\omega_{\mathbb{p}}} \tilde{f}(\mathbb{p})\, F(k) \Big|_{k = -p_{\mathbb{p}}} \end{aligned} \right.$$

$$= \int \prod_{i=1}^{n} \frac{d^{d-1} \mathbb{q}_i}{(2\pi)^{d-1} 2\omega_{\mathbb{q}_i}} \tilde{g}_i(\mathbb{q}_i)^* \prod_{j=1}^{2} \frac{d^{d-1} \mathbb{p}_j}{(2\pi)^{d-1} 2\omega_{\mathbb{p}_j}} \tilde{f}_j(\mathbb{p}_j)$$

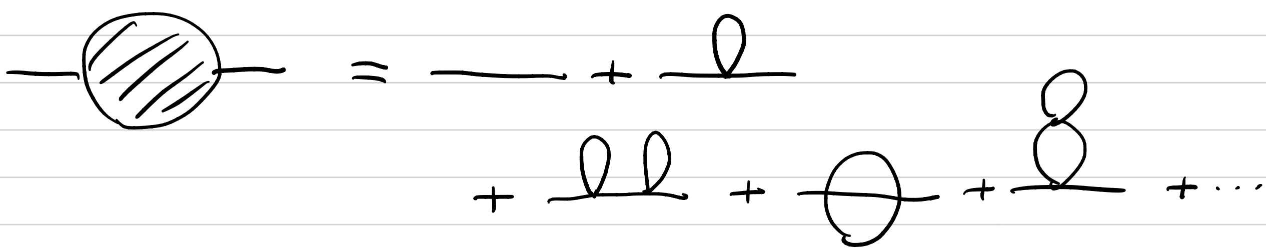
$$\prod_{i=1}^{n} \frac{-i(l_i^2 - m^2)}{\sqrt{Z}} \prod_{j=1}^{2} \frac{-i(k_j^2 - m^2)}{\sqrt{Z}}\, G(l_1, \cdots, l_n, k_1, k_2) \Bigg|_{\substack{l_i = p_{\mathbb{q}_i} \\ k_j = -p_{\mathbb{p}_j}}}$$

S-matrix in perturbation theory

— Diagrammatic expressions

ϕ : an elementary field s.t. $\langle \phi(x) \rangle = 0$

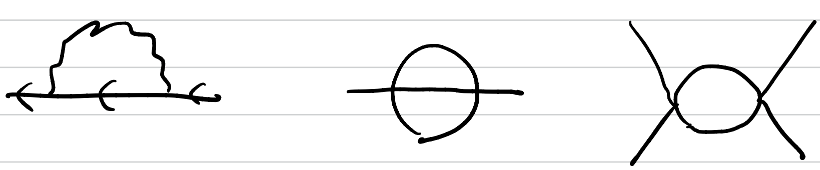
Its two point function $\langle \phi(x)\phi(y) \rangle$ is a sum of connected diagrams



1PI (one particle irreducible) diagram

= a connected diagram which is still connected if an internal line is cut.

e.g.

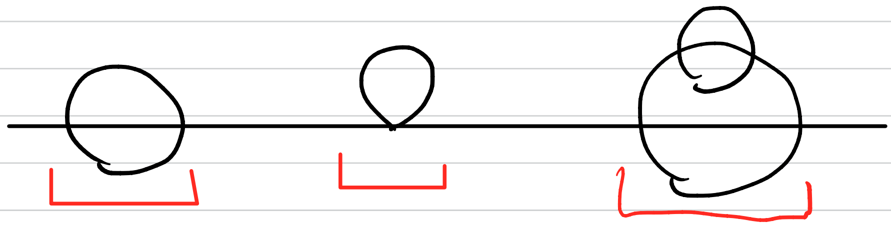


1PI

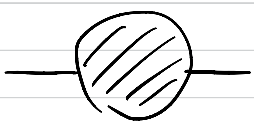
not 1PI

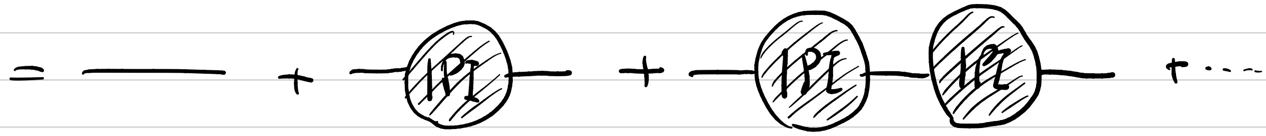
Any connected diagram with two external lines can be decomposed into 1PI blocks:

e.g.



three 1PI blocks





no 1PI　　one 1PI　　two 1PI's

$$= \frac{i}{p^2 - m_0^2} + \frac{i}{p^2 - m_0^2}\left(-iM^2(p^2)\right)\frac{i}{p^2 - m_0^2}$$

$+i0$ omitted

$$+ \frac{i}{p^2 - m_0^2}\left(-iM^2(p^2)\right)\frac{i}{p^2 - m_0^2}\left(-iM^2(p^2)\right)\frac{i}{p^2 - m_0^2} + \cdots$$

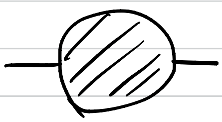
$$= \frac{i}{p^2 - m_0^2}\left(1 + \frac{M^2(p^2)}{p^2 - m_0^2} + \left(\frac{M^2(p^2)}{p^2 - m_0^2}\right)^2 + \cdots\right)$$

$$= 1\Big/\left(1 - \frac{M^2(p^2)}{p^2 - m_0^2}\right)$$

$$= \frac{i}{p^2 - m_0^2 - M^2(p^2)}$$

A zero of $p^2 - m_0^2 - M^2(p^2)$ is mass^2 of a particle.

e.g. if $M^2(p^2) = \alpha p^2 + \beta$



$$= \frac{i}{(1-\alpha)p^2 - m_0^2 - \beta} = \frac{(1-\alpha)^{-1} i}{p^2 - (1-\alpha)^{-1}(m_0^2+\beta)}$$

$$\Rightarrow \quad m^2 = (1-\alpha)^{-1}(m_0^2+\beta), \quad Z = (1-\alpha)^{-1}.$$

In general, if $p^2 - m_0^2 - M^2(p^2)$ has zeros at $p^2 = m_1^2, m_2^2, \cdots$

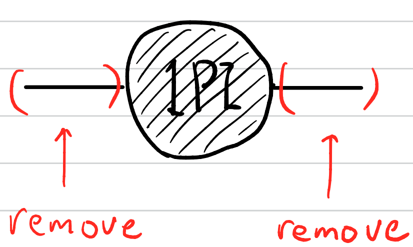
then, these are mass^2 of particles and

$Z_i = |\langle 0, i | \phi(0) | 0 \rangle|^2$ for the i-th particle is given by

$$Z_i^{-1} = \frac{d}{dp^2}\left(p^2 - M^2(p^2)\right)\Big|_{p^2 = m_i^2} = 1 - \frac{dM^2}{dp^2}(m_i^2).$$

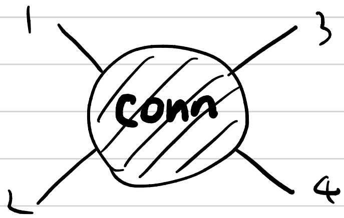
Thus the spectrum of particles & their coupling to ϕ

can be found by the sum of 1PI diagrams for $\langle \phi(x) \phi(y) \rangle$

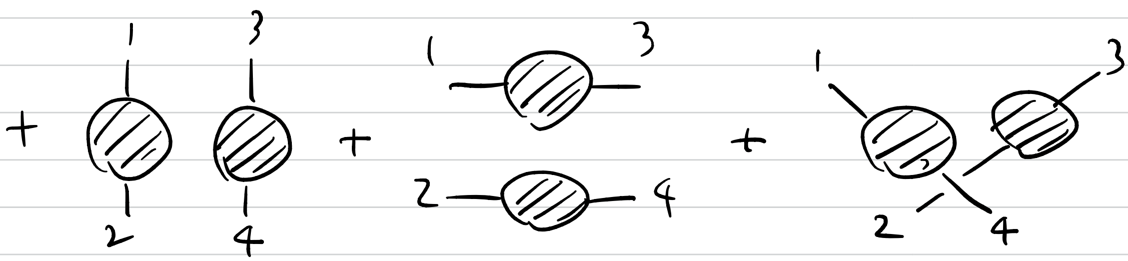


$$= -i M^2(p^2)$$

Four point function $\langle \phi(x_1)\phi(x_2)\phi(x_3)\phi(x_4) \rangle$

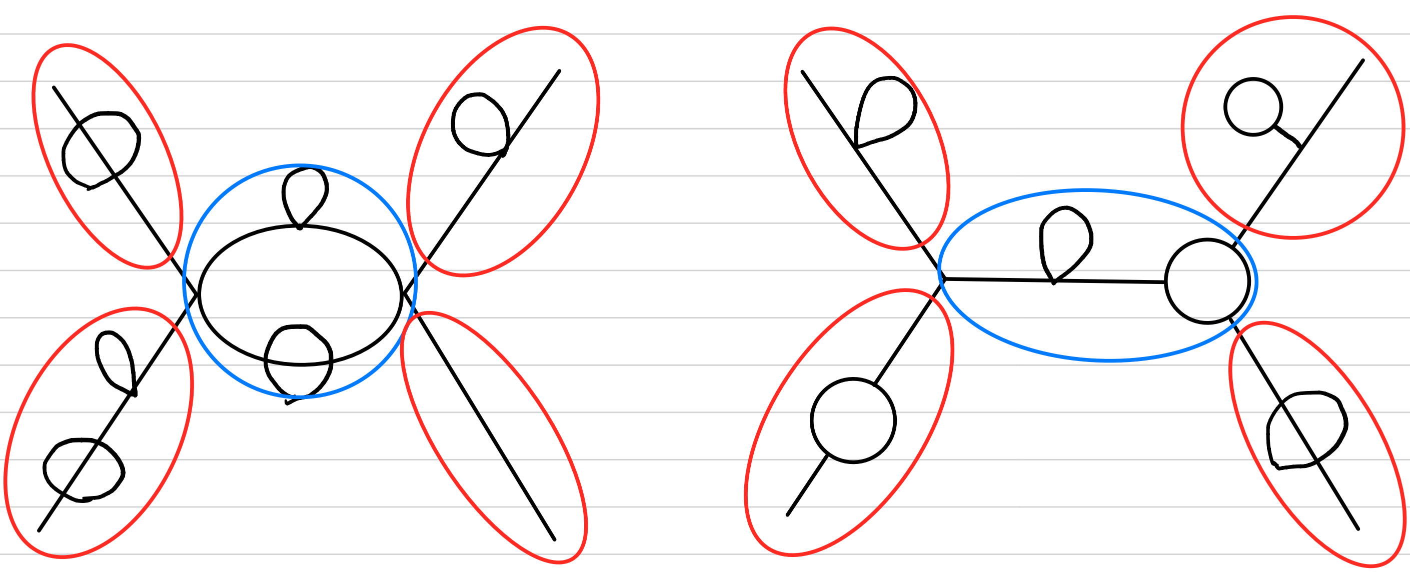
$=$ 

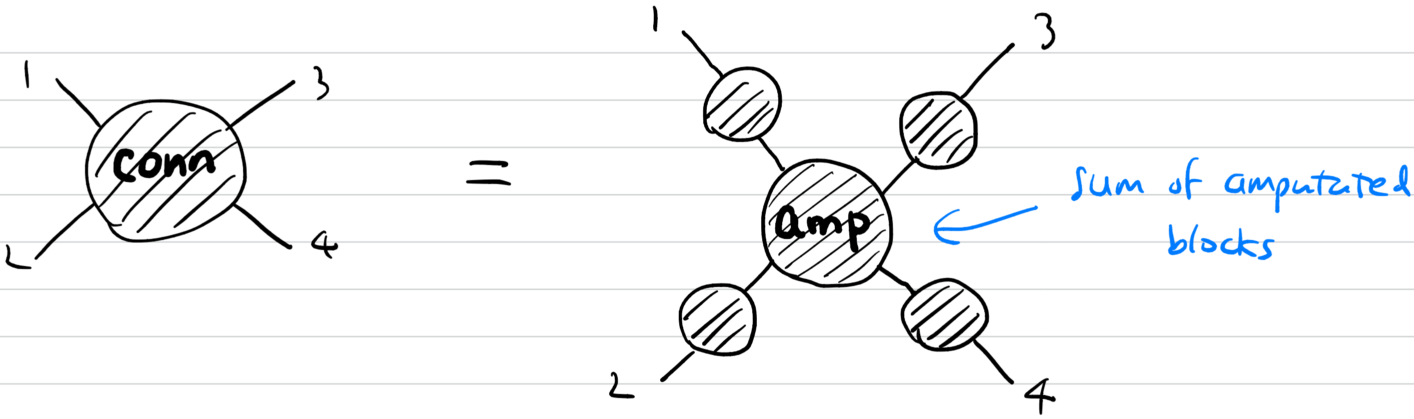
← connected

$+$ 

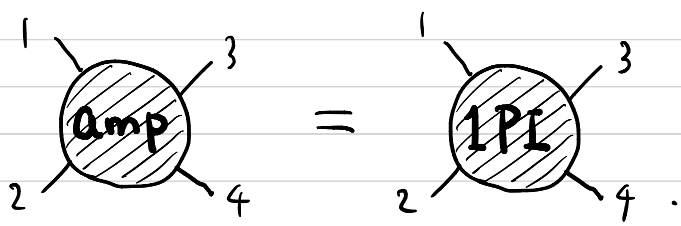
Any connected diagram with n external lines can be decomposed into n external parts and one amputated block

$n=4$ examples :

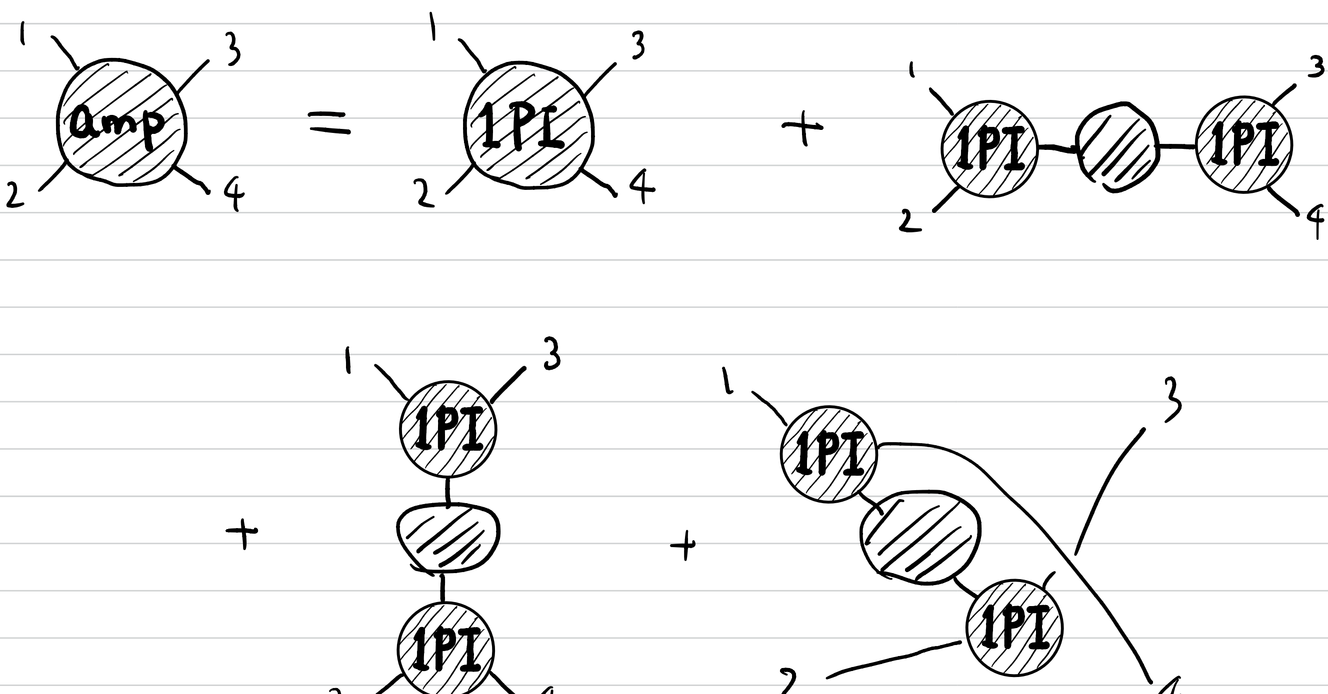




In ϕ^4 theory (or any theory with $\phi \to -\phi$ symmetry)



In general e.g. $U(\phi) = \frac{\lambda_3}{3!}\phi^3 + \frac{\lambda_4}{4!}\phi^4 + \ldots$ with $\lambda_3 \neq 0$,

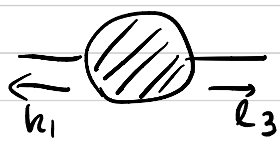


$$\langle g_3, g_4, \text{free} \,|\, (S-1) \,|\, f_1, f_2, \text{free} \rangle$$

$$= \int \prod_{i=3}^{4} \frac{d^{d-1}\mathbf{p}_i}{(2\pi)^{d-1} 2\omega^i_{\mathbf{p}_i}} \tilde{g}_i(\mathbf{p}_i)^* \prod_{j=1}^{2} \frac{d^{d-1}\mathbf{p}_j}{(2\pi)^{d-1} 2\omega^j_{\mathbf{p}_j}} \tilde{f}_j(\mathbf{p}_j)$$

$$\prod_{i=3}^{4} \frac{-i(l_i^2 - m_i^2)}{\sqrt{Z_i}} \prod_{j=1}^{2} \frac{-i(k_j^2 - m_j^2)}{\sqrt{Z_j}} G(l_3, l_4, k_1, k_2) \Bigg|_{\substack{k_j = -p^j_{\mathbf{p}_j} \\ l_i = p^i_{\mathbf{p}_i}}}$$

where $\omega^j_{\mathbf{p}} = \sqrt{\mathbf{p}^2 + m_j^2}$, $p^j_{\mathbf{p}} = (\omega^j_{\mathbf{p}}, \mathbf{p})$.

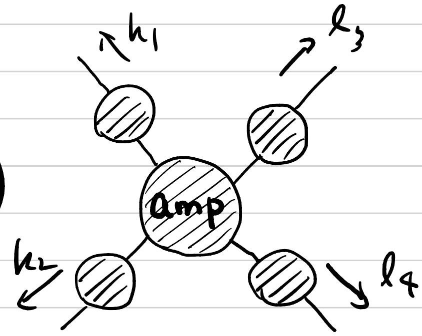
$(l_3^2 - m_3^2)(k_1^2 - m_1^2)$ 

$$\propto \delta(k_1 + l_3)(l_3^2 - m_3^2)(k_1^2 - m_1^2)\left(\sum_i \frac{i Z_i}{k_1^2 - m_i^2} + \cdots \right)$$

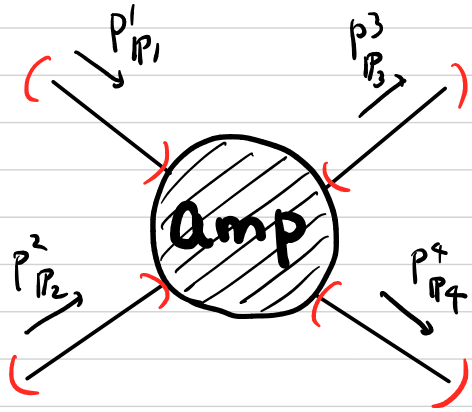
$$\longrightarrow 0 \quad \text{as} \quad \begin{matrix} k_1 \to -p^1_{\mathbf{p}_1} \\ l_3 \to p^3_{\mathbf{p}_3} \end{matrix} \quad \left(\Rightarrow \begin{matrix} k_1^2 \to m_1^2 \\ l_3^2 \to m_3^2 \end{matrix} \right)$$

Thus, the contribution of disconnected diagrams vanishes.

$$\prod_{l=3}^{4} (l_i^2 - m_i^2) \prod_{j=1}^{2} (k_j^2 - m_j^2)$$



$$\xrightarrow[\;k_j \to -p^j_{\mathbb{p}_j}\;]{\;l_i \to p^i_{\mathbb{p}_i}\;} \prod_{i=1}^{4} i Z_i \times$$

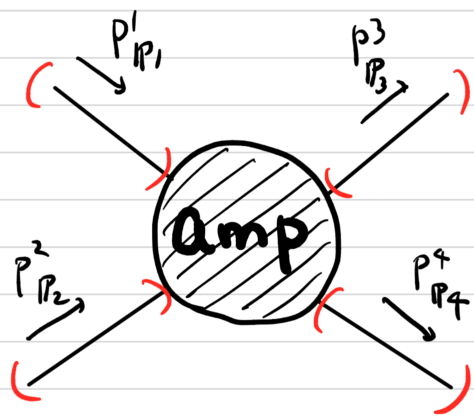


Thus

$$\langle g_3, g_4, \text{free} | (S-1) | f_1, f_2, \text{free} \rangle$$

$$= \int \prod_{i=3}^{4} \frac{d^{d-1}\mathbb{p}_i}{(2\pi)^{d-1} 2\omega^i_{\mathbb{p}_i}} \tilde{g}_i(\mathbb{p}_i)^* \prod_{j=1}^{2} \frac{d^{d-1}\mathbb{p}}{(2\pi)^{d-1} 2\omega^j_{\mathbb{p}_j}} \tilde{f}_j(\mathbb{p}_j)$$

$$\sqrt{Z_1}\sqrt{Z_2}\sqrt{Z_3}\sqrt{Z_4}$$



1PI effective action

Consider a theory of variables $\phi = (\phi_1, \cdots, \phi_N)$

measure $d\phi$ and action $S_E(\phi)$ (omit "E" below).

$$e^{-W(J)} = \int d\phi \, e^{-S(\phi) + J\cdot\phi}$$

Decompose $S(\phi) - J\cdot\phi = \underbrace{\frac{1}{2}\sum_{i,j} \phi_i A_{ij} \phi_j}_{\text{free part}} + \underbrace{\text{else}}_{\text{interaction part}}$

and evaluate $W(J)$ perturbatively.

※ Everything below is perturbation series but we omit writing "pert" each time.

e.g. $W(J) = W_{pert}(J)$ is the sum of connected diagram.

$$-\frac{\partial}{\partial J_i} W(J) = \frac{\int d\phi \, e^{-S(\phi)+J\cdot\phi} \phi_i}{\int d\phi \, e^{-S(\phi)+J\cdot\phi}} =: \langle \phi_i \rangle_J$$

Solve $\langle \phi_i \rangle_J \overset{!}{=} \phi_i$ $\quad i = 1, \cdots, N$ for J, write the solution

as $J = J(\phi)$ and put

Unique in perturbation theory

$$\Gamma(\phi) := W(J(\phi)) + J(\phi)\cdot\phi$$

$$\frac{\partial \Gamma(\phi)}{\partial \phi_i} = \cancel{\frac{\partial J_j(\phi)}{\partial \phi_i} \cdot \underbrace{\frac{\partial W}{\partial J_j}(J(\phi))}_{\phi_j}} + \cancel{\frac{\partial J_j(\phi)}{\partial \phi_i} \cdot \phi_j} + J_i(\phi) = J_i(\phi) \,,$$

$i = 1, \cdots, N.$

Thus,

$$\phi_i^* := \langle \phi_i \rangle_{J=0} \quad \Rightarrow \quad J(\phi^*) = 0 \quad \therefore \ \frac{\partial \Gamma}{\partial \phi_i}(\phi^*) = 0.$$

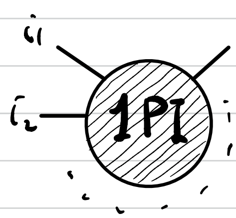
VEV of ϕ at $J=0$ is a critical point of $\Gamma(\phi)$.

Properties of $\Gamma(\phi)$

① It is a generating series of 1PI vertices

$$\Gamma(\phi) = \frac{1}{2} \log \det(A/2\pi) + \frac{1}{2} \sum_{i,j} \phi_i A_{ij} \phi_j - \sum_{n=0}^{\infty} \frac{1}{n!} \sum_{i_1 \cdots i_n} \lambda_{1PI}^{i_1 \cdots i_n} \phi_{i_1} \cdots \phi_{i_n}$$

where $\lambda_{1PI}^{i_1 \cdots i_n}$ is the 1PI vertex defined by

$$\langle \phi_{i_1} \cdots \phi_{i_n} \rangle_{1PI} =$$  $$= \sum_{j_1 \cdots j_n} \overbrace{\phi_{i_1} \phi_{j_1}} \cdots \overbrace{\phi_{i_n} \phi_{j_n}} \lambda_{1PI}^{j_1 \cdots j_n}$$

For this reason, $\Gamma(\phi)$ is called 1PI effective action.

② $\Gamma(\varphi) = \frac{1}{2}\log\det(A/2\pi)$ — The sum of 1PI vacuum diagrams of $\mathcal{T}(\varphi)$, the theory with background φ:

$$\begin{cases} \text{variables} & \xi = (\xi_1, \cdots, \xi_N) \\ \text{measure} & d_\varphi \xi = d(\varphi + \xi) \\ \text{action} & S_\varphi(\xi) = S(\varphi + \xi) \end{cases}$$

$$\int d_\varphi \xi \; e^{-S_\varphi(\xi)} = \sqrt{\frac{(2\pi)^N}{\det A}} \; e^{\text{connected vacuum diagrams}}$$

$$= e^{-\Gamma(\varphi) + \text{non-1PI conn. vac. diagrams}}$$

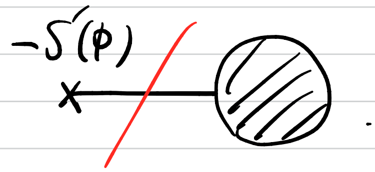
Here we take

$$S_\varphi(\xi) = \underbrace{\frac{1}{2}\xi_i A_{ij}\xi_j}_{\text{free part}} + \underbrace{\text{else}}_{\text{interaction part}}$$

③ This holds for any decomposition of $S_\varphi(\xi)$ into free + interaction. In particular, for the expansion in powers of ξ, we can take the ξ-quadratic part $\frac{1}{2}\sum_{ij}\xi_i\xi_j\partial_i\partial_j S(\varphi)$ as the free part.

$$S_\phi(\xi) = S(\phi) + S'(\phi)\,\xi + \underbrace{\frac{1}{2} S''(\phi)\,\xi^2}_{\text{free part}} + \underbrace{\frac{1}{3!} S'''(\phi)\,\xi^3 + \cdots}_{\text{interaction}}$$

- $S(\phi)$ is outside the ξ integral.
- Any diagram involving the vertex $-S'(\phi)\cdot\xi$ is not 1PI:



Thus, we can take only the cubic or higher powers in ξ as the interaction part to produce vertices.

With this understanding,

$$e^{-\Gamma(\phi)} = e^{-S(\phi)} \cdot \sqrt{\frac{(2\pi)^n}{\det S''(\phi)}} \cdot \exp\left(\text{1PI vacuum diagrams}\right).$$

That is,

$$\Gamma(\phi) = S(\phi) + \underbrace{\frac{1}{2}\log\det\left(\frac{S''(\phi)}{2\pi}\right)}_{\frac{1}{2}\operatorname{tr}\log\left(\frac{S''(\phi)}{2\pi}\right)} - \text{1PI vacuum diagrams}.$$

Consequence of ② :

recover $\hbar$ $\quad \int d_\phi \xi \, e^{-\frac{1}{\hbar} S_\phi(\xi)} = e^{-\frac{1}{\hbar}\Gamma(\phi,\hbar) + \text{others}}$

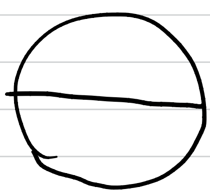
$\rightsquigarrow$ propagator $\propto \hbar$, vertex $\propto \hbar^{-1}$

A 1PI vacuum diagram with # propagator $= P$

vertices $= V$

$\propto \hbar^{P-V} = \hbar^{L-1}$

where $L = P - V + 1$ is # loops

e.g.  $\quad P = 3$, $V = 2$ $\quad L = 3 - 2 + 1 = 2$

Thus, $\Gamma(\phi,\hbar) = \sum_{L=0}^{\infty} \hbar^L \Gamma_L(\phi)$

$\Rightarrow$ $-\Gamma_L(\phi) =$ the sum of 1PI vacuum diagrams

with # loops $= L$

($\log \det(A/2\pi\hbar)$ is included in $L = 1$)

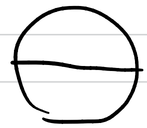
$\therefore$ $\hbar$-expansion $=$ loop expansion.

Remark

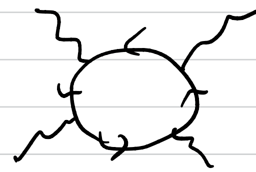
E = # external lines

P = # propagators

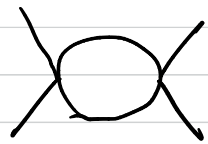
V = # vertices



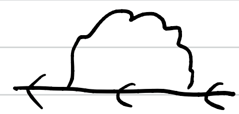
$E=0$
$P=3$
$V=2$



$E=4$
$P=8$
$V=4$



$E=4$
$P=6$
$V=2$



$E=2$
$P=4$
$V=2$

Then # internal lines $I = P - E$

and # loops $L = I - V + 1 = P - E - V + 1$ if connected

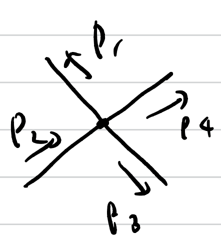
Contribution to partition/correlation function:

$$\int \prod_{v\in V} d^d y_v \int \prod_{e\in E} d^d p_e \, e^{-ip_e(x_e - y_{v(e)})} \int \prod_{i\in I} d^d p_i \, e^{-ip_i(y_{t(i)} - y_{s(i)})} F(p)$$

$$\int d^d y_v \, e^{i\sum_{\ell\in V} \varepsilon_\ell p_\ell y_v} = (2\pi)^d \delta\Big(\sum_{\ell\in V} \varepsilon_\ell p_\ell\Big)$$

sum over lines connected to v

$$\varepsilon_\ell = \begin{cases} +1 & \text{if } \ell \text{ goes out of } v \\ -1 & \text{if } \ell \text{ comes in to } v \end{cases}$$



$\rightsquigarrow \; (2\pi)^d \delta(p_1 - p_2 + p_3 + p_4)$

$$= \int \prod_{e\in E} d^d p_e \, e^{-i p_e x_e} \int \prod_{i \in I} d^d p_i \, \underbrace{\prod_{v\in V} (2\pi)^d \delta\Big(\sum_{e\in v} \varepsilon_e p_e\Big)}_{(2\pi)^d \delta\left(\sum_{e\in E} p_e\right) \prod_{v=1}^{V-1} (2\pi)^d \left(\sum_{e\in v} \varepsilon_e p_e\right)} F(p)$$

Overall momentum conservation

$\therefore$ net # of momentum integrals

$$= I - (V-1) = L.$$