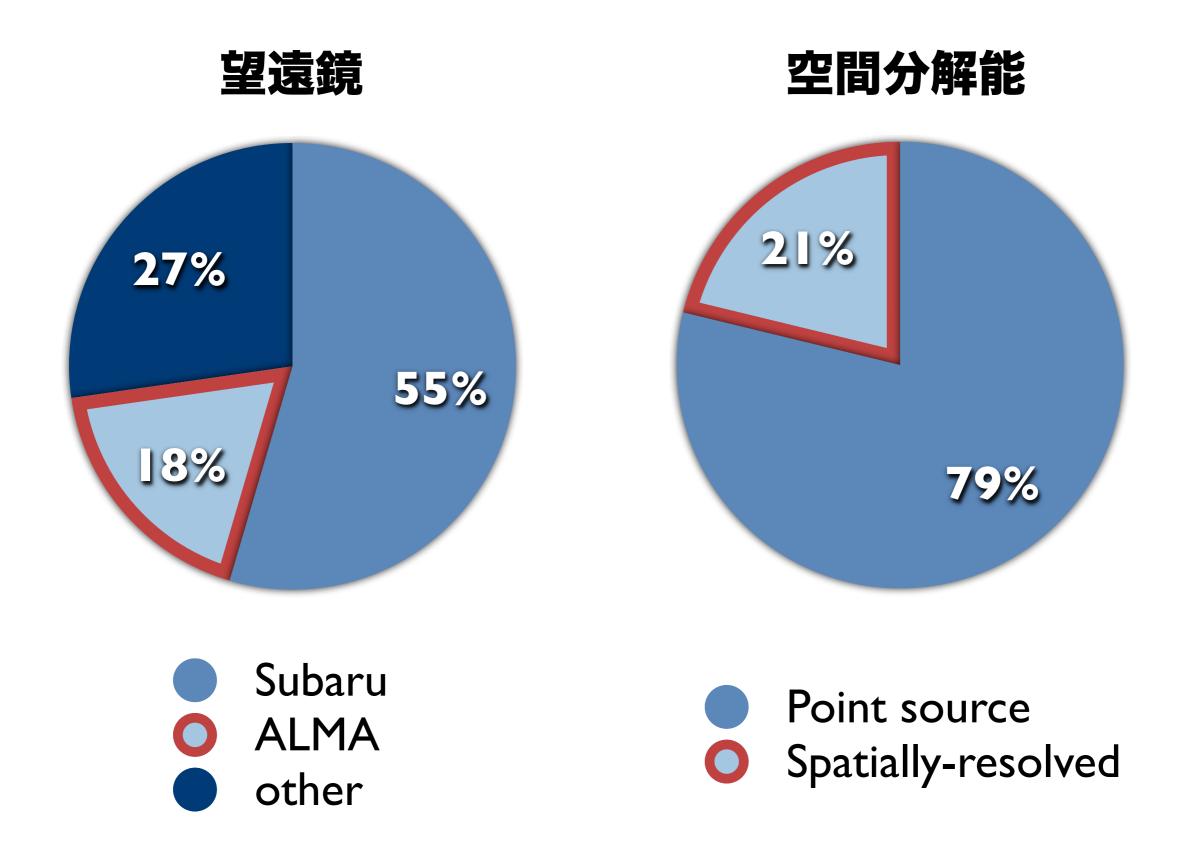
銀河進化研究会、6月7-9日、大阪大学

遠方銀河におけるガスの運動学

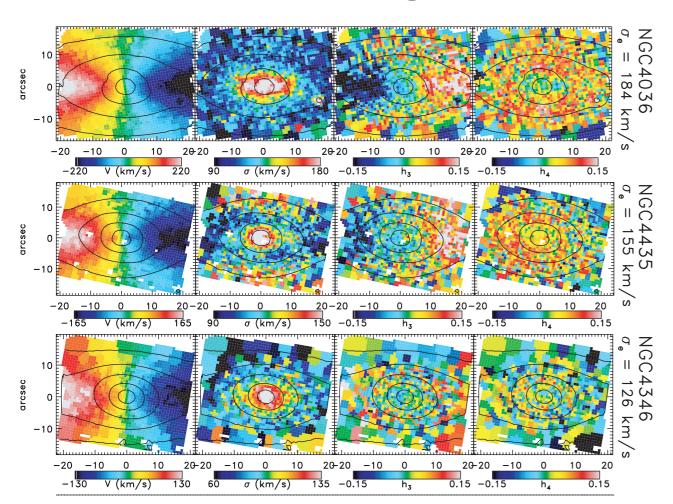
但木謙一 (国立天文台)

銀河進化研究会2017の講演内容(遠方銀河観測)



近年の運動学的研究におけるマイルストーン

ATLAS3D



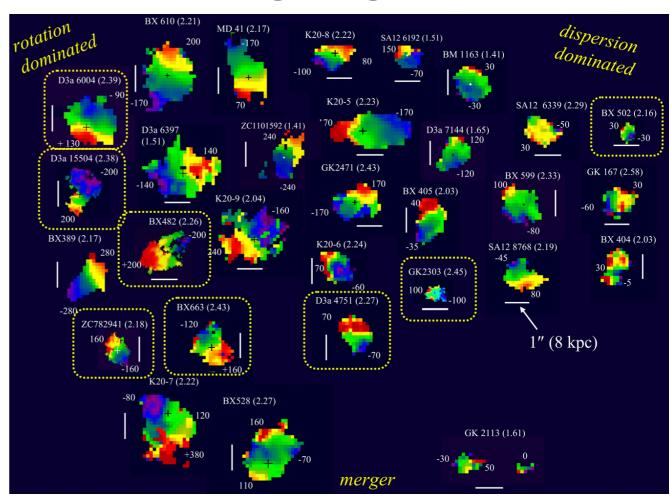
Cappellari+11

近傍早期型銀河



slow rotator fast rotator

SINS



Foerster Schreiber+09

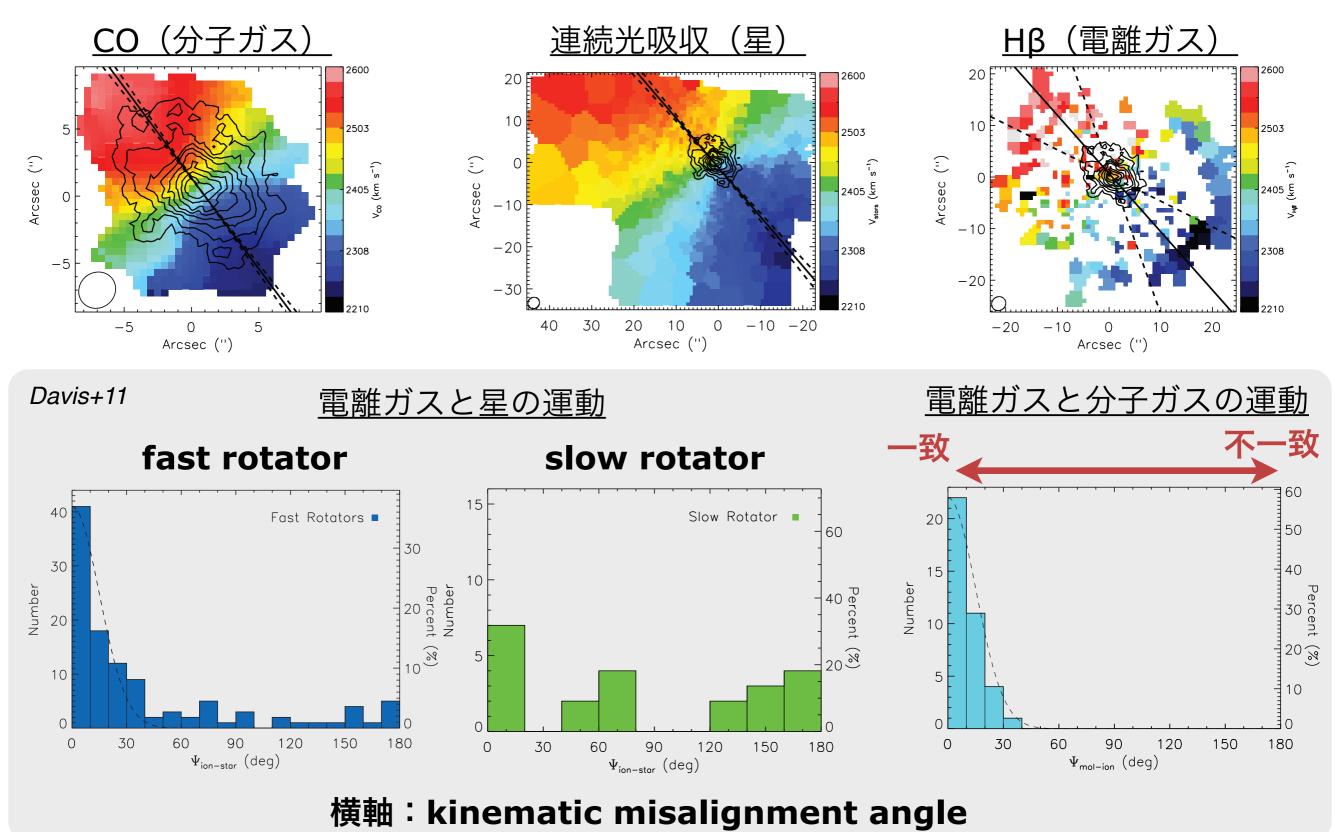
遠方星形成銀河



merger -> 回転円盤銀河

運動学のトレーサー

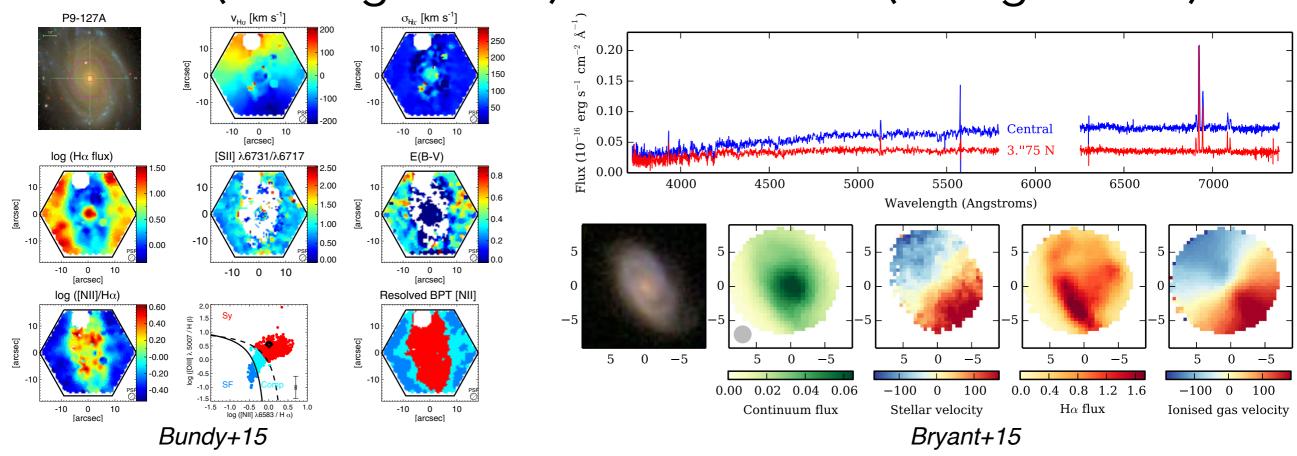
ATLAS^{3D}銀河の例



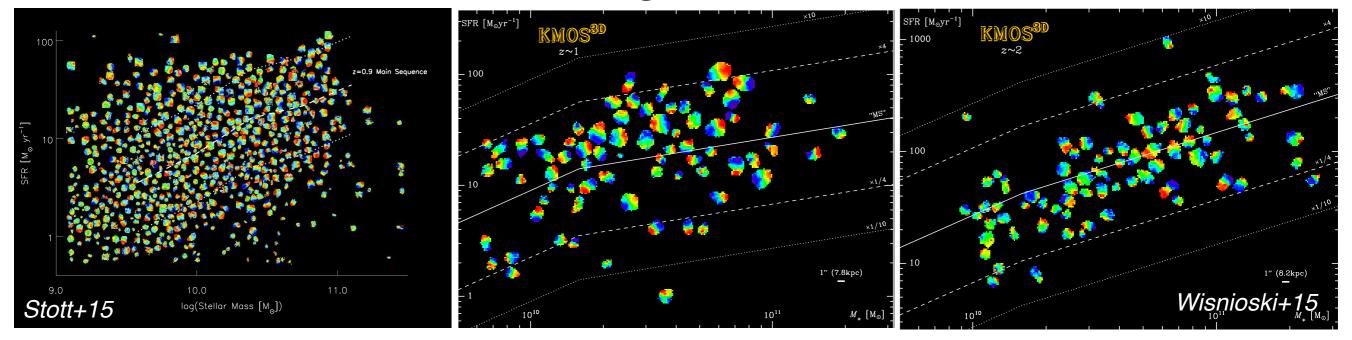
最近の大規模探査(電離ガス・星)

MaNGA (~10k galaxies)

SAMI (~3k galaxies)

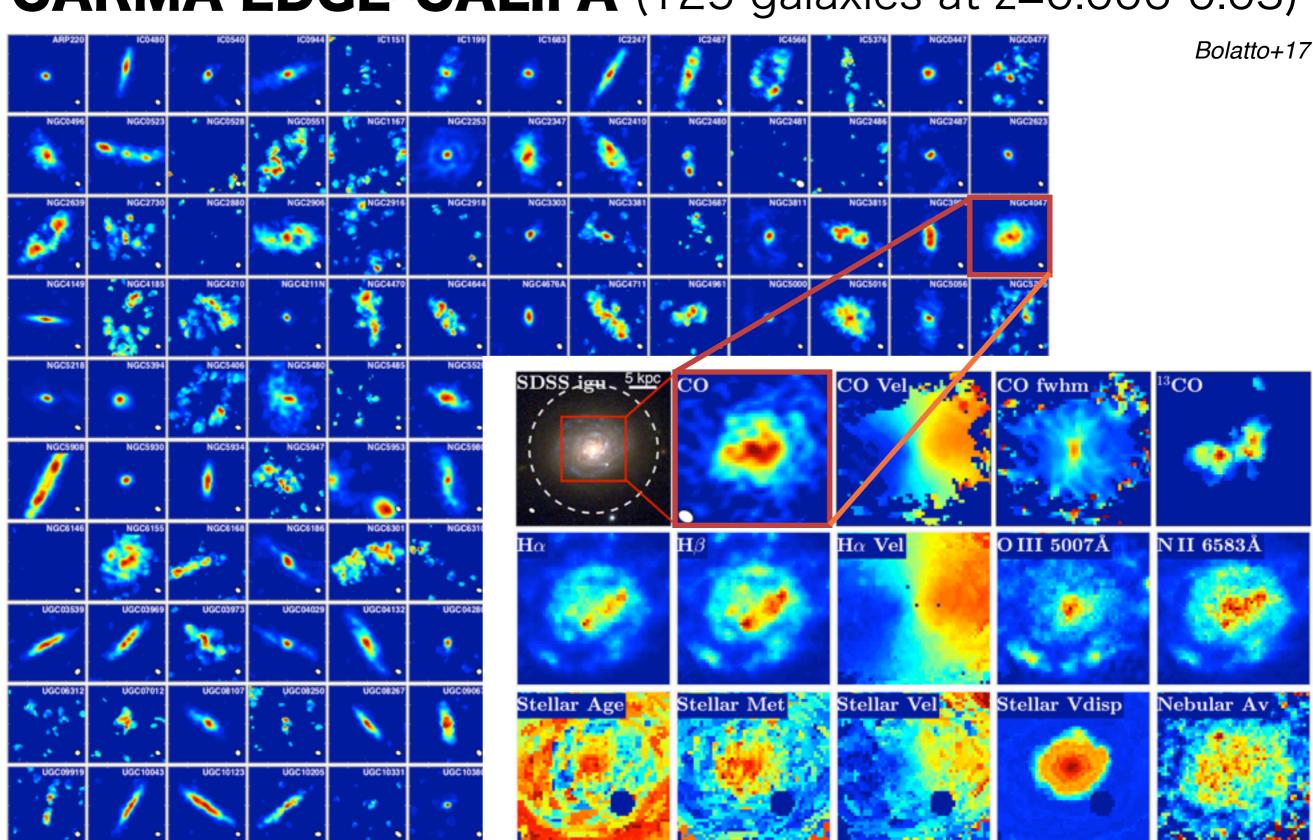


KMOS (~1k galaxies at z>0.5)

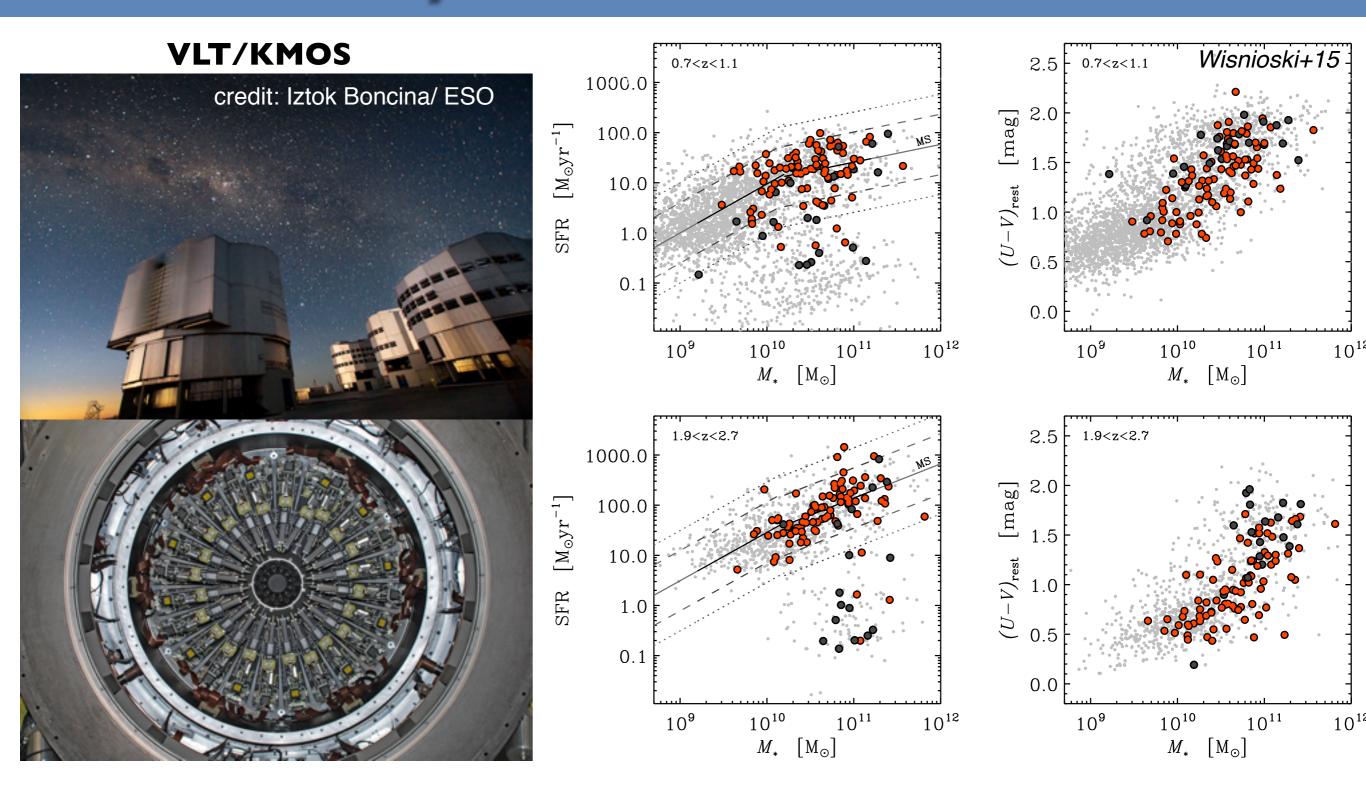


最近の大規模探査(分子ガス)

CARMA EDGE-CALIFA (125 galaxies at z=0.006-0.03)



KMOS^{3D} survey



Ha IFU survey in an unbiased sample of >600 galaxies at 0.7<z<2.7

IFUデータの解析



パイプライン

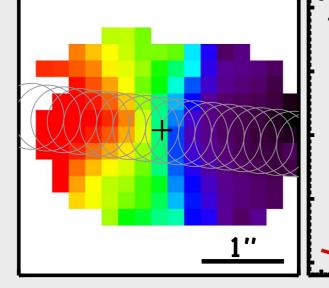
3D CUBE

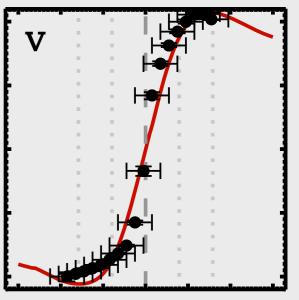
H - - - - | Cos3_04796 - - |

ガウシアン フィッティング Vcir, σ_0 , Mdyn, (Mhalo, λ)

ディスクモデリング

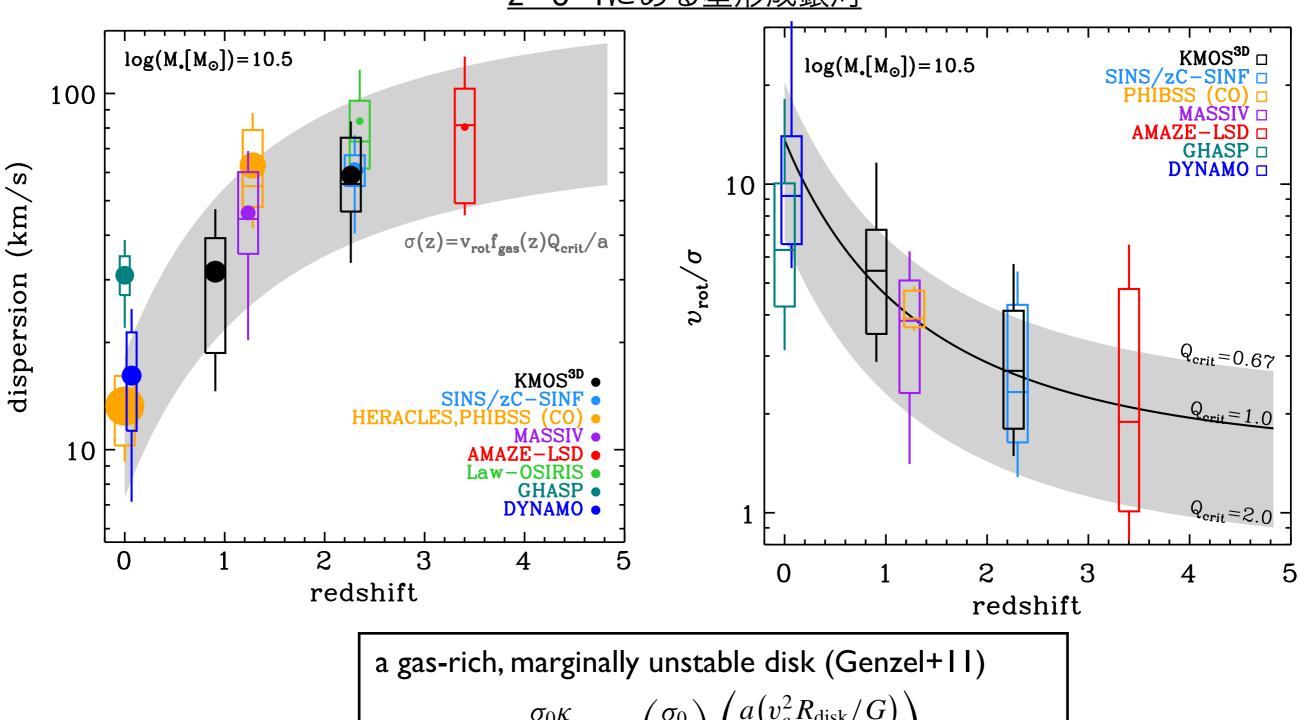
velocity map





σ、v/σの赤方偏移進化

z=0-4にある星形成銀河



 $Q_{\rm gas} = \frac{\sigma_0 \kappa}{\pi G \Sigma_{\rm gas}} = \left(\frac{\sigma_0}{v_c}\right) \left(\frac{a \left(v_c^2 R_{\rm disk}/G\right)}{\pi R_{\rm disk}^2 \Sigma_{\rm gas}}\right)$ $= \left(\frac{\sigma_0}{v_c}\right) \left(\frac{a M_{\rm tot}}{M_{\rm gas}}\right) = \left(\frac{\sigma_0}{v_c}\right) \left(\frac{a}{f_{\rm gas}}\right).$

Disk modeling in KMOS^{3D}

in Wuyts+16 paper, we derive M_{dyn} and σ_0 by contrasting the data with convolved disk models

data

LINEFIT code (Davies+09, Forster Schreiber+09)

- I. extracting spectra in each spatial pixel or aperture
- 2. Gaussian fitting to derive rotation velocity and velocity dispersion

contrast

disk model

DYSMAL code (Davies+11, Cresci+09)

characterizing morphology

r_e: effective radius

n: Sersic index

h: scale height

i: inclination

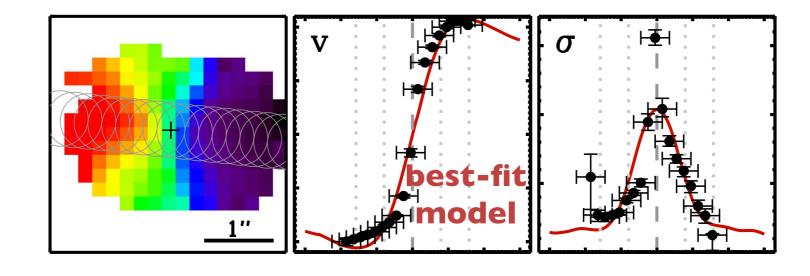
PA: position angle

from HST 0.2"resolution map

characterizing kinematics

v: rotation velocity (M_{dyn})

σ: velocity dispersion



high-z disks are highly turbulent

pressure support (asymmetric drift)

hydrostatic equation

Burkert+10

$$\frac{v_{\mathrm{rot}}^2}{r} = f_g(r) + \frac{1}{\rho} \frac{dp}{dr},$$
 (1) gravitational force pressure gradient (thermal pressureは無視できる)

$$v_0^2 \equiv f_g \times r$$
 (zero-pressure)

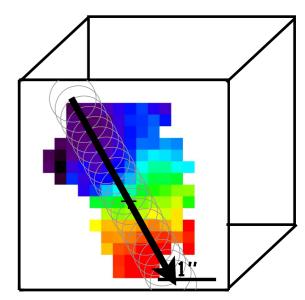
$$v_{\text{rot}}^2 = v_0^2 + \frac{r}{\rho} \frac{dp}{dr} = v_0^2 + \frac{1}{\rho} \frac{d}{dlnr} (\rho \sigma^2).$$
 (2)

exponential disk (see Burkert+10 and Binney&Tremaine 08)

circular velocity
$$v_{\text{rot}}^2 = v_0^2 - 2\sigma^2 \left(\frac{r}{r_d}\right). \tag{11}$$
 observed pressure support

σが大きいほど見かけの回転速度が遅くなる

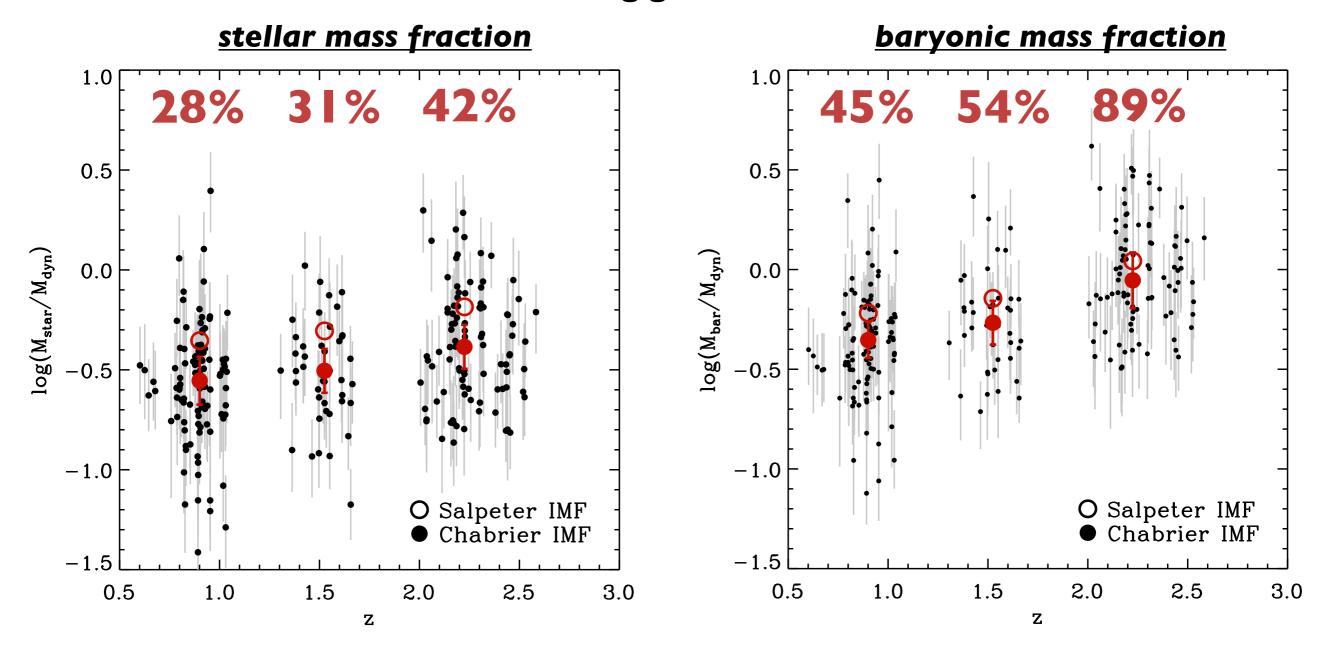
dynamical mass



Dynamical mass

Wuyts+16

240 star-forming galaxies at 0.6<z<2.6



遠方銀河の内側(R<R_e)では、 ほとんどの質量はバリオン(星+ガス)が担っている

Exponential Disk within an NFW Dark Matter Halo (Mo+98)

Burkert+16

DM halo parameter

 $(M_{DM}, R_{virial}, V_{virial}, \lambda)$

disk parameter

(Mdisk, Rdisk, Vdisk, O0)

観測可

NFW Dark Matter Halo (Navarro+97)

$$\rho_{DM}(r) = \frac{4\rho_c}{(r/r_s)(1 + r/r_s)^2}$$

$$v_{DM}^{2}(r) = V_{200}^{2} \left(\frac{r_{200}}{r}\right) \frac{\ln(1 + r/r_{s}) - (r/r_{s})/(1 + r/r_{s})}{\ln(1 + c) - c/(1 + c)}.$$

Exponential thin disk (Freeman+70)

$$\Sigma(r) = \Sigma_0 \times \exp\left(-\frac{r}{r_d}\right)$$

$$v_{\text{disk}}^2(r) = 4\pi G \Sigma_0 r_d y^2 [I_0(y) K_0(y) - I_1(y) K_1(y)]$$

$$v_{\rm circ}^2 = v_{\rm disk}^2 + v_{\rm DM}^2$$
, 観測された回転速度において、 diskで説明できない分は DMからの寄与と考える

DMからの寄与と考える

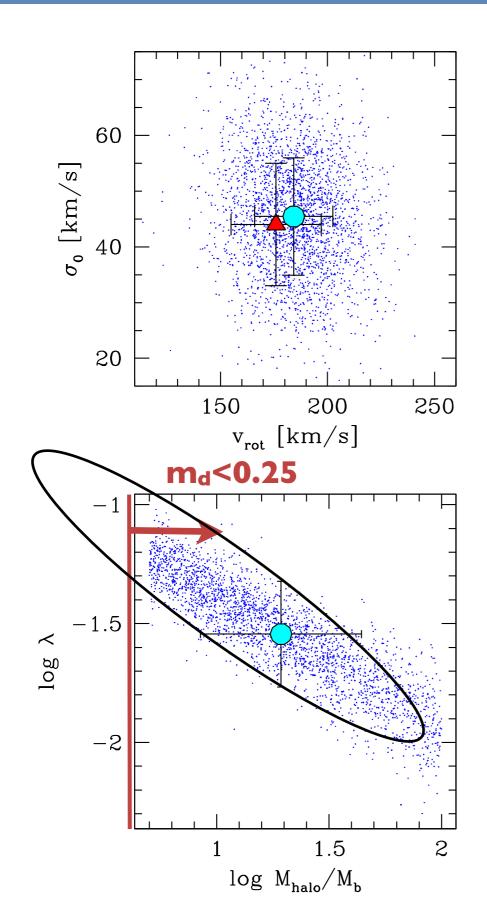
原理的には観測データからhalo massとhalo spin parameterが得られる

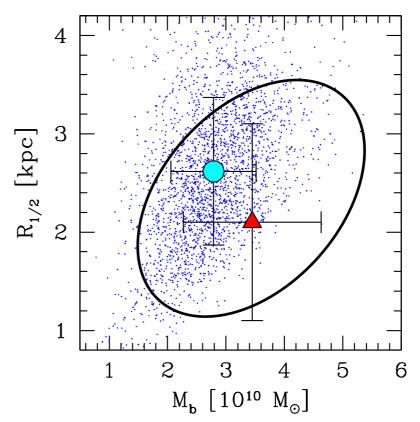
halo mass

halo spin parameter

$$M_{\rm DM} = \frac{v_{\rm virial}^2 R_{\rm virial}}{G} = \frac{v_{\rm virial}^3}{10 GH(z)}$$
 $\lambda = \sqrt{2} \times \left(\frac{R_d}{R_{\rm virial}}\right) \times \left(\frac{j_d}{j_{\rm DM}}\right)^{-1}$

Monte-Carlo search of halo parameter





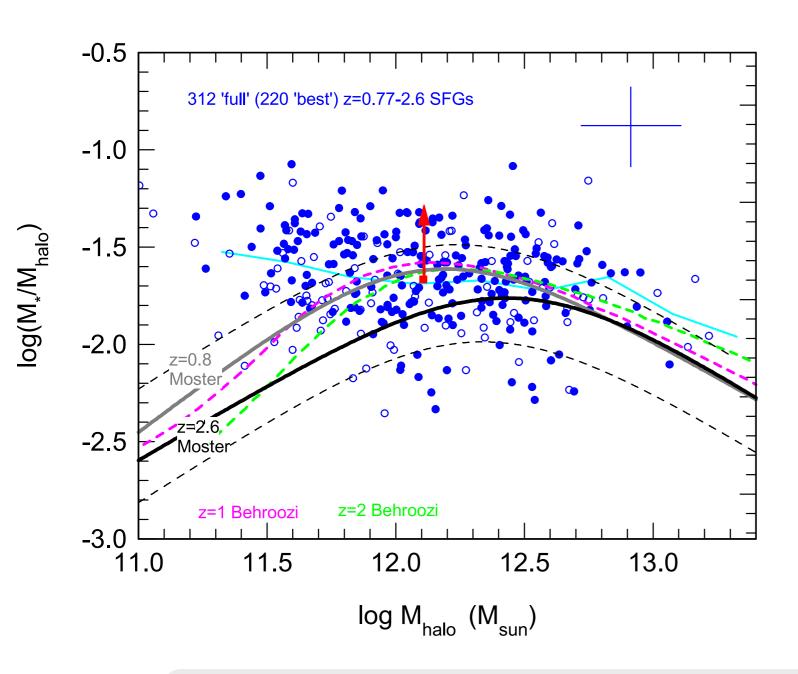
Burkert+16

red: observed properties and uncertainties cyan: mean of $\frac{\text{converged}}{\text{m}_{\text{d}} < 0.25}$

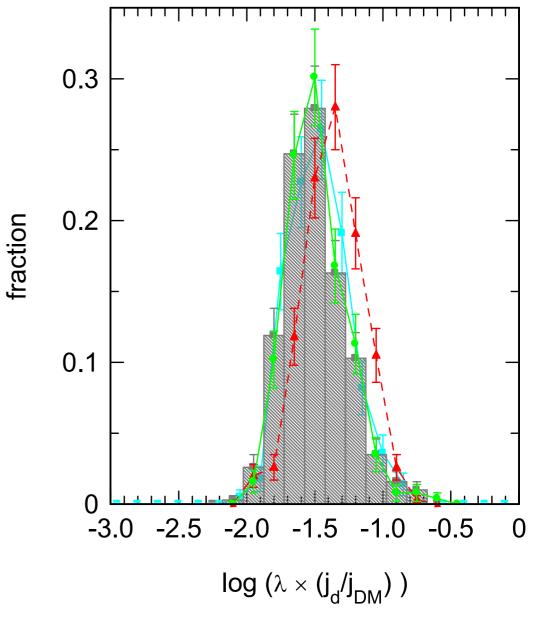
(cosmic baryonic fraction, 0.17)

Monte-Carlo search of halo parameter

Burkert+16



N=220 'best' MC-NFW no adiabatic contraction
N=312 'full' MC-NFW no adiabatic contraction
N=304 'full', MC-NFW. with adiabatic contraction
N=256 'full' MC-NFW, no ad.contr., with n_{Sersic}-correction



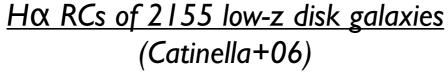
halo mass

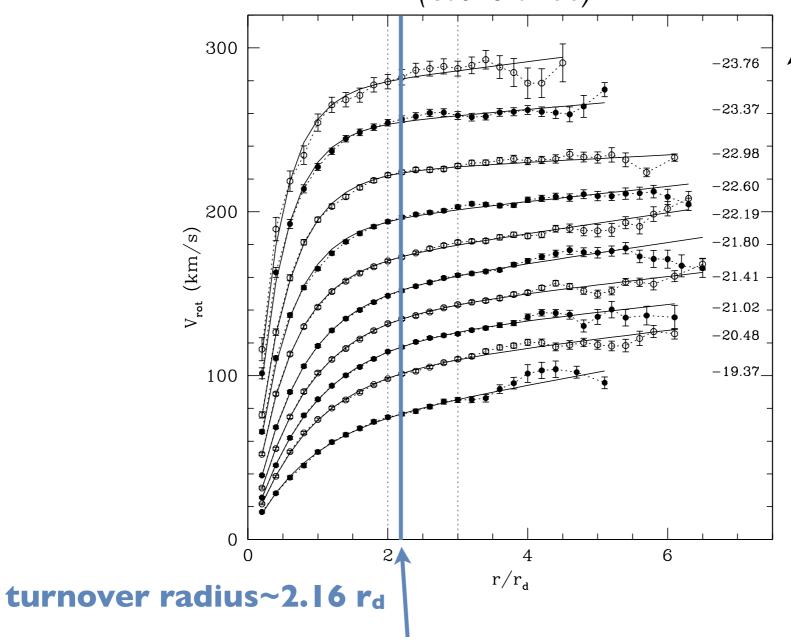
$$M_{\rm DM} = \frac{v_{\rm virial}^2 R_{\rm virial}}{G} = \frac{v_{\rm virial}^3}{10GH(z)}$$

halo spin parameter

$$\lambda = \sqrt{2} \times \left(\frac{R_d}{R_{\text{virial}}}\right) \times \left(\frac{j_d}{j_{\text{DM}}}\right)^{-1}$$

Outer rotation curves





brighter in I-band magnitude

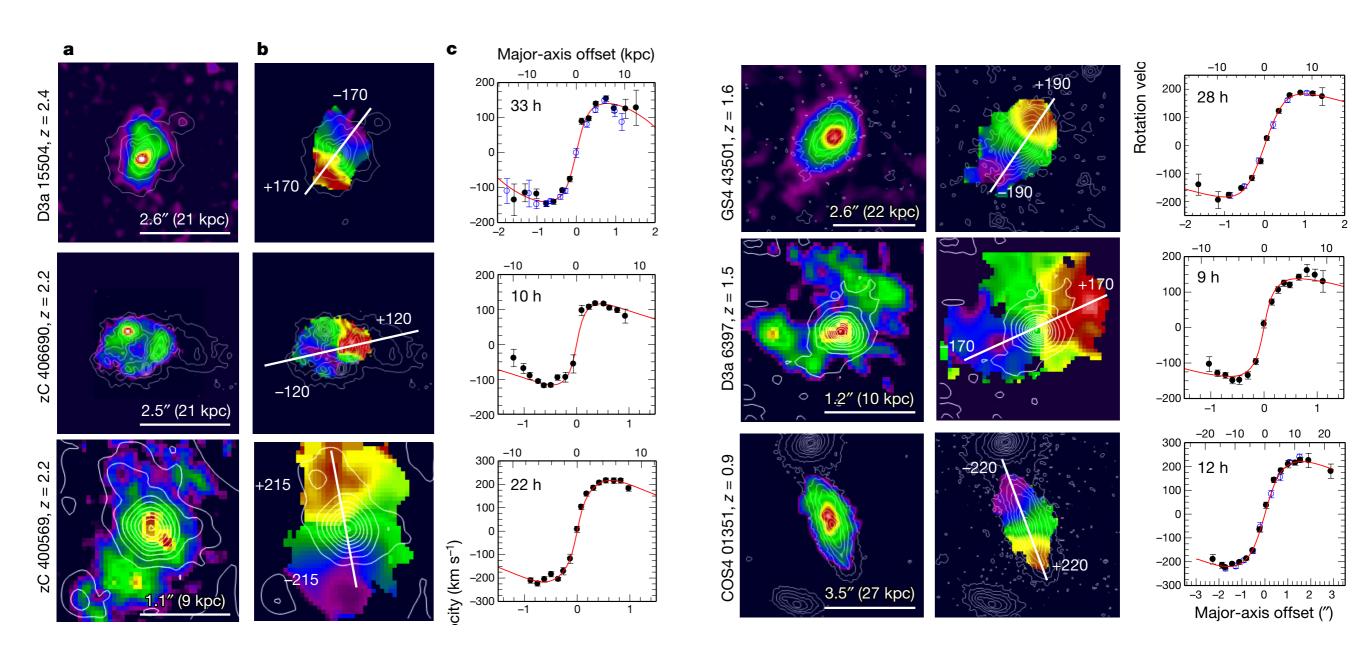
Exponential thin disk (Freeman+70)

$$\Sigma(r) = \Sigma_0 \times \exp\left(-\frac{r}{r_d}\right)$$

$$v_{\text{disk}}^2(r) = 4\pi G \Sigma_0 r_d y^2 [I_0(y) K_0(y) - I_1(y) K_1(y)]$$

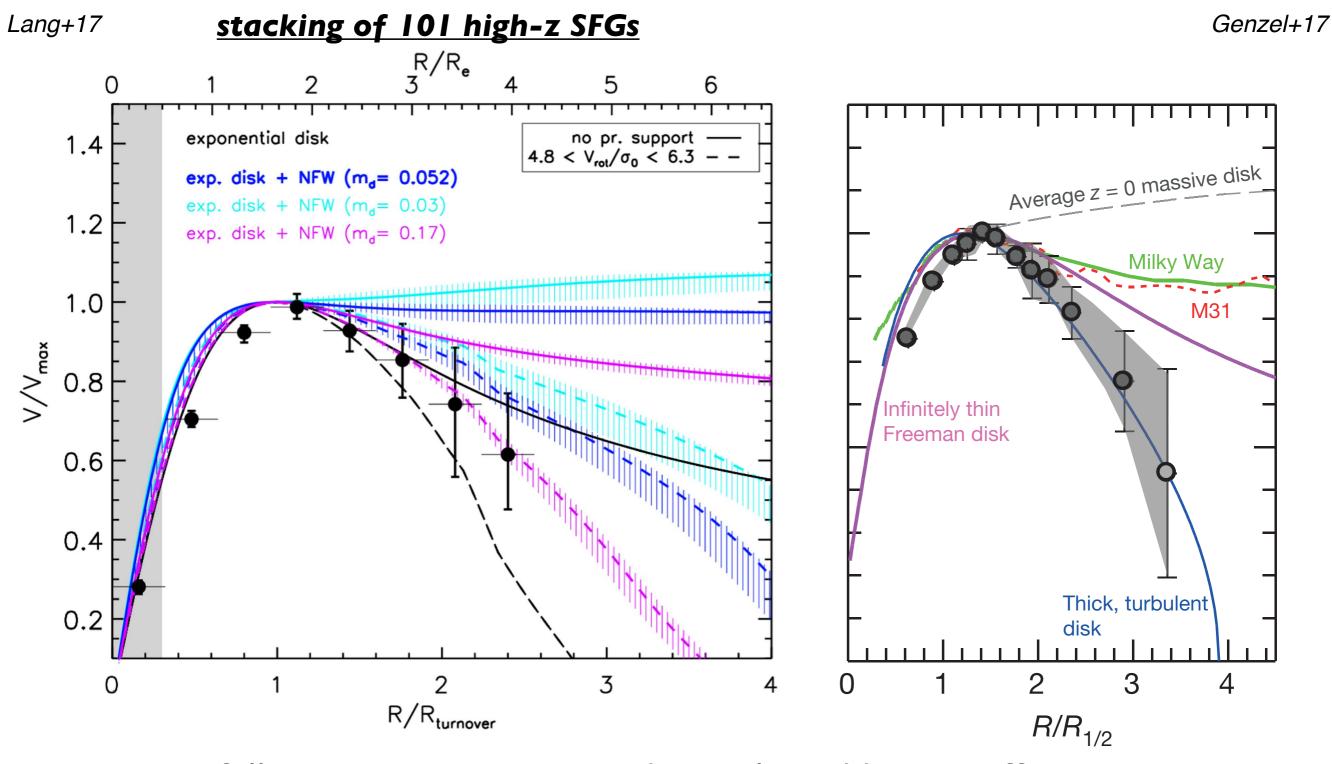
Outer rotation curves

Genzel+17



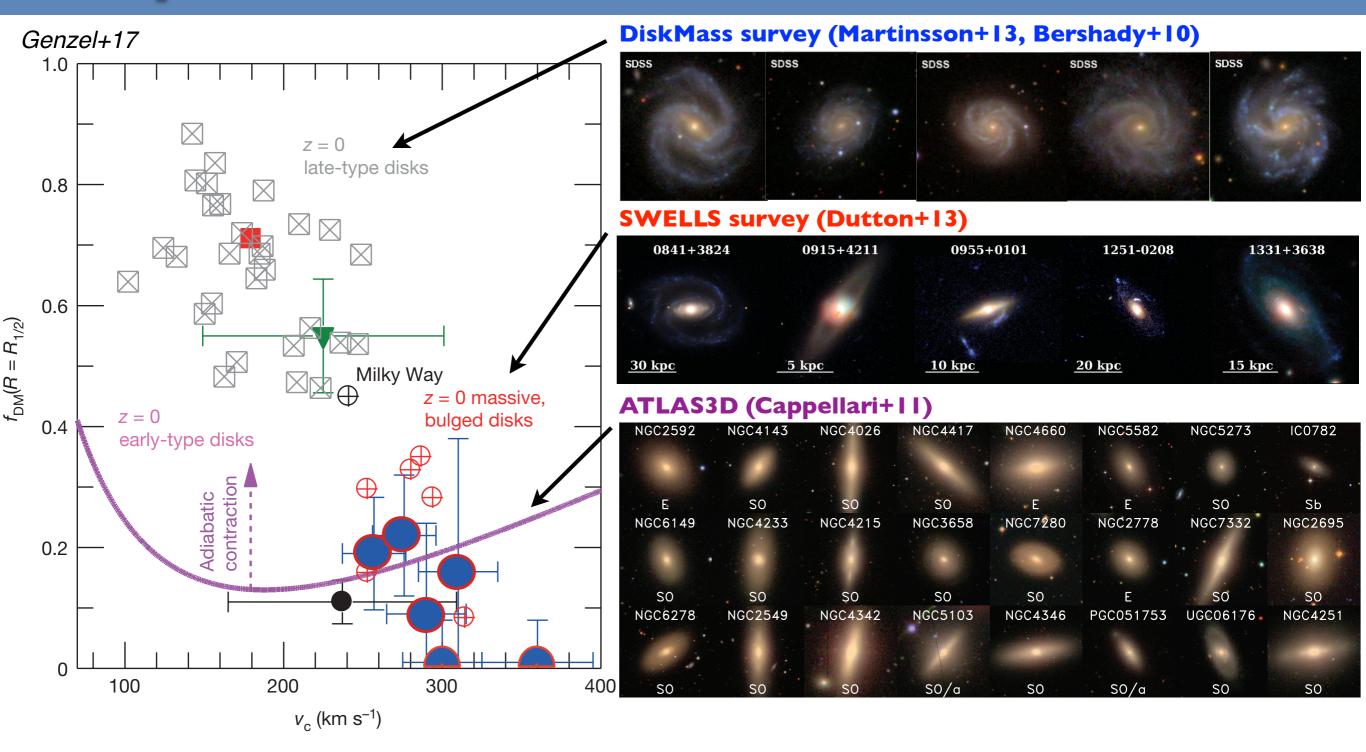
- outer rotation curves drop beyond the turnover radius are they representative of high-z SF galaxies?

Outer rotation curves



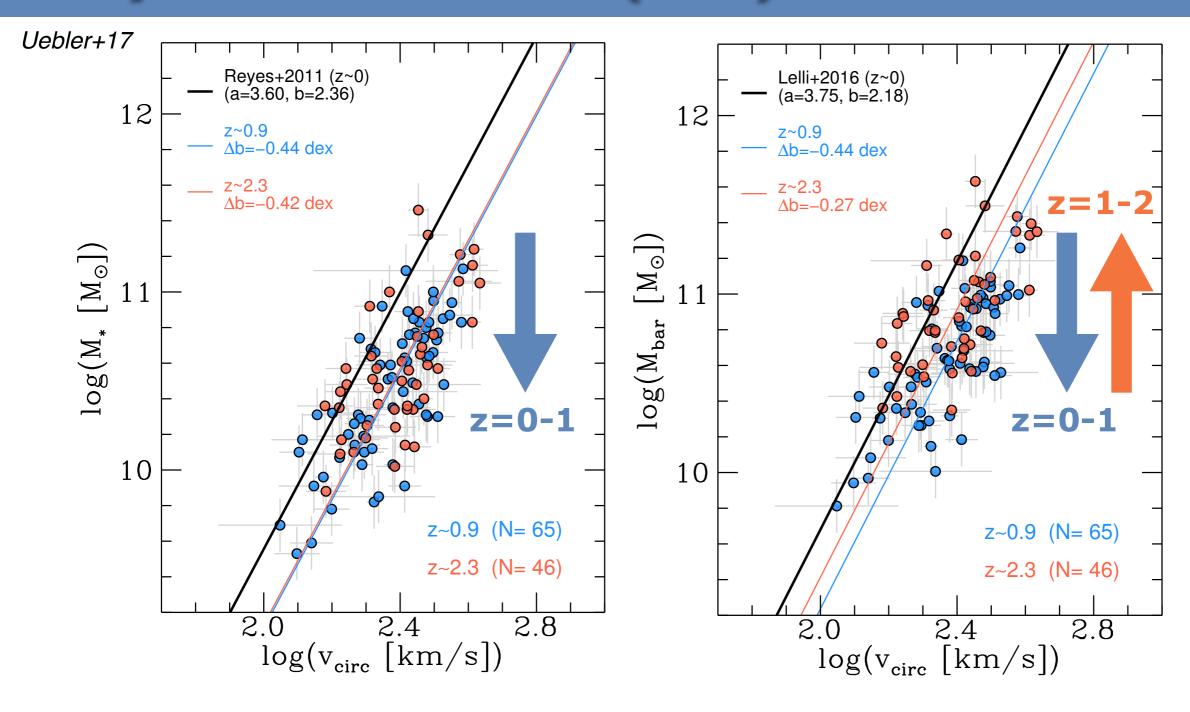
- falling rotation curves can be explained by two affects
- 1. high baryonic fraction
- 2. pressure support

Comparisons with low-redshift results



- passive galaxies are probably the descendants of massive star-forming galaxies
- the low dark-matter fractions may be preserved in the properties of the local passive population

Tully-Fisher relation (TFR)



- at fixed velocity, higher baryonic mass and similar stellar mass at $z\sim2.3$ as compared to $z\sim0.9$
- zero point offset is not monotonic

Toy model of baryonic TFR

I. DM Halo (Mo+98)

$$M_{h} = \frac{V_{h}^{3}}{10G \cdot H(z)} ; \quad R_{h} = \frac{V_{h}}{10H(z)}$$

$$M_{bar} = m_{f} \cdot M_{h} ; \quad R_{bar} = \underline{r_{f}} \cdot R_{h}$$

$$\text{constant}$$

$$M_{bar}/m_{d} = C_{1} \times R_{d}^{3}H(z)^{2}$$

$$(3)$$

2. Exponential disk (Freeman+70)

$$v_{\text{circ}}(r) = \sqrt{v_{\text{bar}}^2(r) + v_{\text{DM}}^2(r)}.$$
 (D4)
 $\int f_{\text{DM}}(r) = v_{\text{DM}}^2(r)/v_{\text{circ}}^2(r).$

$$v_{bar}^2 = (1 - f_{DM}) v_{cir}^2$$

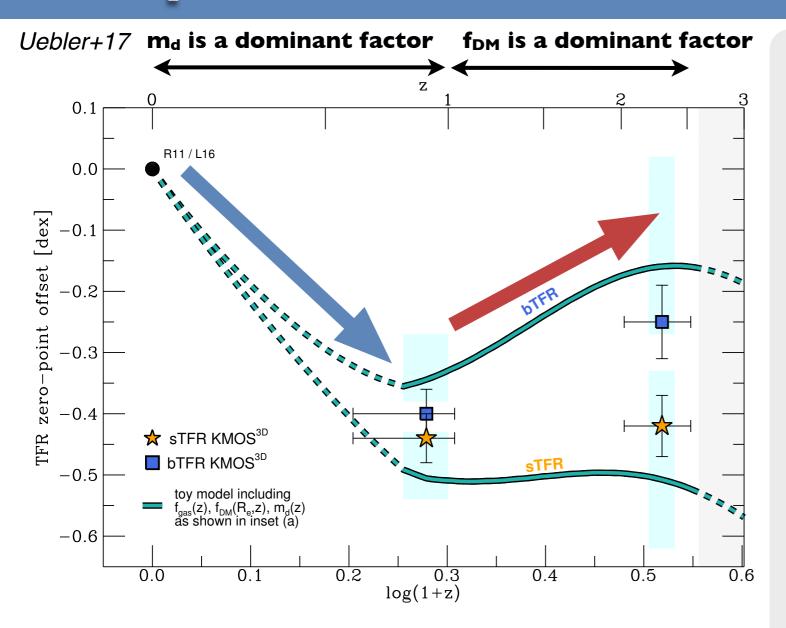
$$v_{\text{bar}}^{2}(r) = 4\pi G \Sigma_{0} R_{d} y^{2} [I_{0}(y) K_{0}(y) - I_{1}(y) K_{1}(y)], \quad (D5)$$

$$\propto M_{bar}/R_{d}$$

$$M_{bar}/R_d=C_2\times(1-f_{DM})v_{cir}^2$$

$$M_{\text{bar}} = \frac{v_{\text{circ}}^3(R_e)}{H(z)} \left[\frac{[1 - f_{\text{DM}}(R_e, z)]^{3/2}}{m_d^{1/2}(z)} \cdot C \right]$$
(4)

Interpretation of TFR zero-point offset

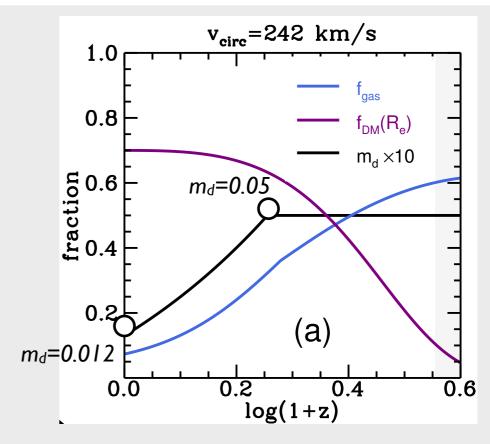


Baryonic TF relation zero-point offset term ->

$$M_{\text{bar}} = \frac{v_{\text{circ}}^3(R_e)}{H(z)} \begin{bmatrix} [1 - f_{\text{DM}}(R_e, z)]^{3/2} \\ \frac{m_d^{1/2}(z)}{} \end{bmatrix} C$$
 (4)

Stellar TF relation

$$M_* = \frac{v_{\text{circ}}^3(R_e)}{H(z)} \left[\frac{[1 - f_{\text{DM}}(R_e, z)]^{3/2} [1 - f_{\text{gas}}(z)]}{m_d^{1/2}(z)} \right] C', (5)$$



1. gas fraction (Mgas/Mstar)

$$\log\left(\frac{M_{\rm gas, mol}}{M_*}\right) \approx 0.12 - 3.62 \cdot \left[\log(1+z) - 0.66\right]^2$$

$$-0.33 \cdot \left[\log(M_* [M_{\odot}]) - 10.7\right].$$

$$Tacconi+17$$
(D8)

2. disk mass fraction (M_{bar}/M_{DM})

 m_d =0.012 at z=0: Moster+13 m_d =0.05 at z=0.8<z<2.6 Burkert+16

3. dark matter fraction $(M_{DM}/M_{bar}$ at $R_e)$

$$f_{\rm DM}(R_e) = 0.7 \cdot \exp[-(0.5 \cdot z)^{2.5}]$$

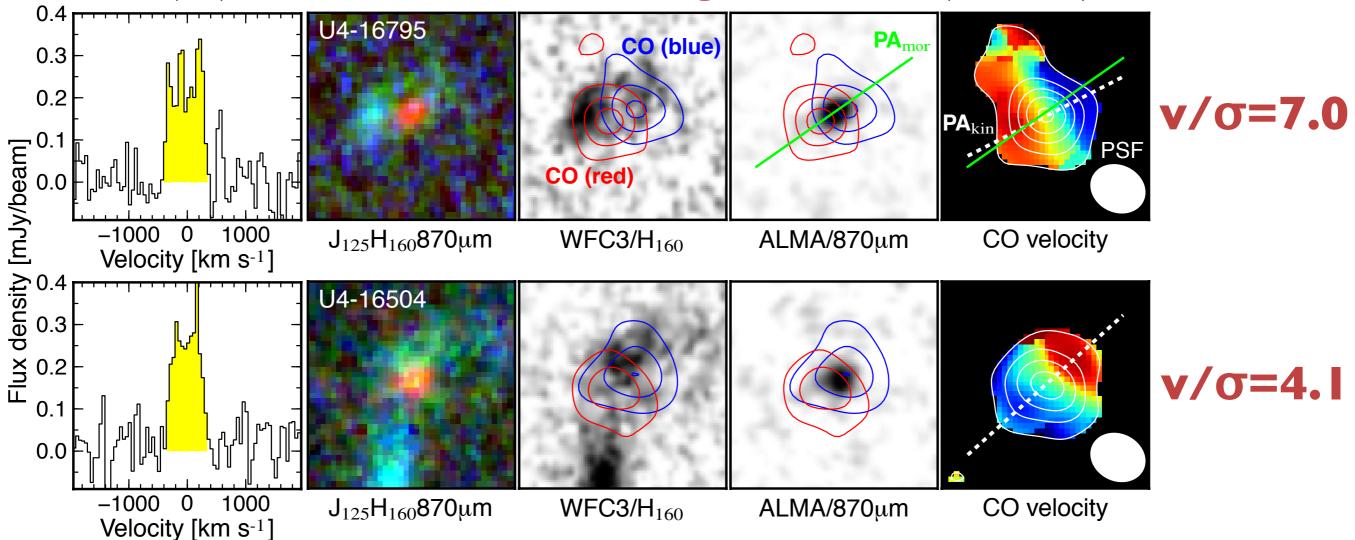
z=0: Courteau&Dutton 2015

z=0.9,2.3: Wuyts+16

CO observations with ALMA

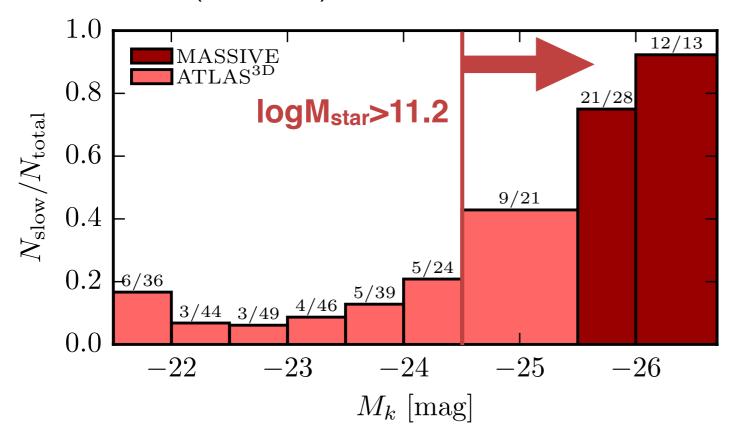


CO(3-2) observations in the most massive galaxies at z=2.5 (Tadaki+17)

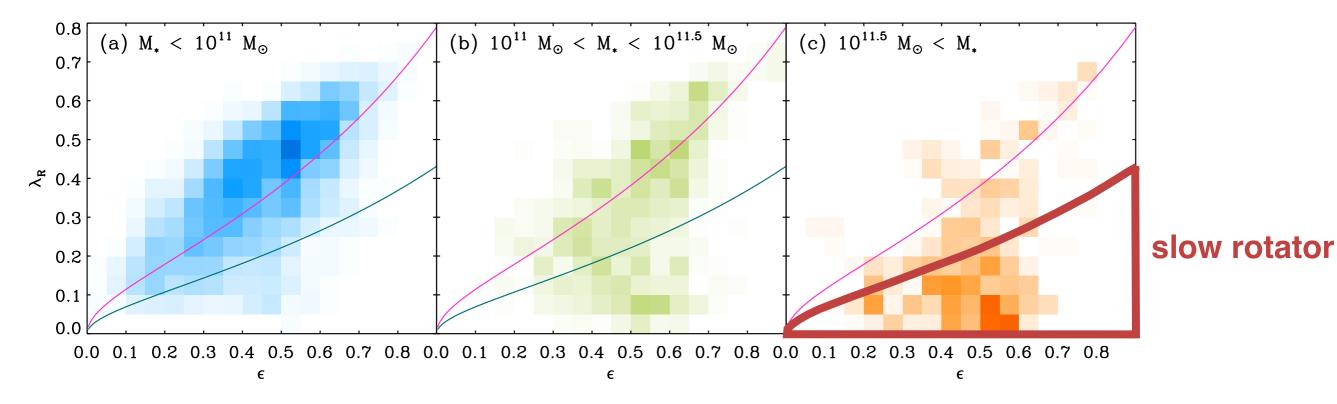


z=2にある重い銀河はslow rotatorになる?

Observations (Veale+16)

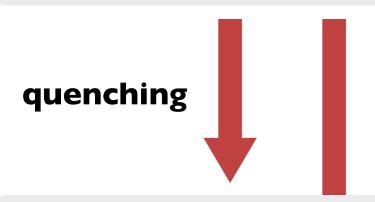


Simulations (Penoyer+17)



何が銀河の回転を止めたのか?

the most massive star-forming galaxies at high-z $V/\sigma>>1$



quiescent galaxies at high-z $V/\sigma\sim1$

size evolution

slow rotators at z=0 $V/\sigma << 1$

