

Is Gravity the Weakest Force?

Masahito Yamazaki
(Kavli IPMU, Tokyo)

Pacific 2019, Moorea
2019 / Sep / 1

Thanks to collaborations :

Satoshi Shirai + M Y

1904.10577 [hep-th]

Alexander Kusenko, Volodymyr Takhistov,

Masaki Yamada + M Y

1908.10930 [hep-th]



My rule today:

No complicated equations/figures

No systematic references

Loose in factors ($2 = \pi = 1$)

My rule today:

No complicated equations/figures

No systematic references

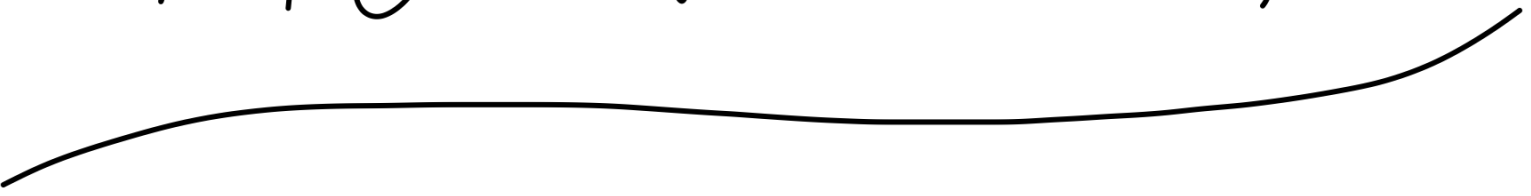
Loose in factors ($2 = \pi = 1$)

OK to be naive/speculative

Attempt for big picture/
fundamental questions

Be relaxed!

Motivation

A hand-drawn, slightly wavy underline line that starts under the 'M' and ends under the 'h' of the word 'Motivation'.

Forces in Nature

Gravity

EM

Weak

Strong

Forces in Nature

Gravity $g_{\mu\nu}$

EM A_μ $U(1)$

Weak A_μ^a $SU(2)$

Strong A_μ^a $SU(3)$

Forces in Nature

@ Gravity	$g_{\mu\nu}$	}	long-ranged
@ EM	A_μ $U(1)$		
Weak	A_μ^a $SU(2)$	}	short-ranged (NOT today)
Strong	A_μ^a $SU(3)$		

Forces in Nature

			spin 2		spin 1
@ Gravity	$g_{\mu\nu}$				
@ EM	A_μ	$U(1)$			
Weak	A_μ^a	$SU(2)$			
Strong	A_μ^a	$SU(3)$			
				long-ranged	
				short-ranged	
				(NOT today)	

Forces in Nature

		spin 2		
ⓐ Gravity	$g_{\mu\nu}$			
ⓐ EM	A_μ	$U(1)$		
				long-ranged
Weak	A_μ^a	$SU(2)$		
Strong	A_μ^a	$SU(3)$		
				short-ranged (NOT today)

??

ϕ ?

spin 0

Gravity versus EM

Microscopically

$$F_{\text{gravity}} = \frac{m^2}{8\pi M_{\text{pl}}^2} \frac{1}{r^2}$$



$$F_{\text{EM}} = \frac{e^2}{4\pi} \frac{1}{r^2}$$

Gravity versus EM

Macroscopically

$$F_{\text{gravity}} = \frac{m^2}{8\pi M_{\text{Pl}}^2} \frac{1}{r^2}$$

m : huge gravity accumulates



$$F_{\text{EM}} = \frac{e^2}{4\pi} \frac{1}{r^2}$$

$e \sim 0$
cancellation of opposite charges

Gravity versus EM

length scale

microscopic

macroscopic

F_{EM}
↓
 $F_{gravity}$

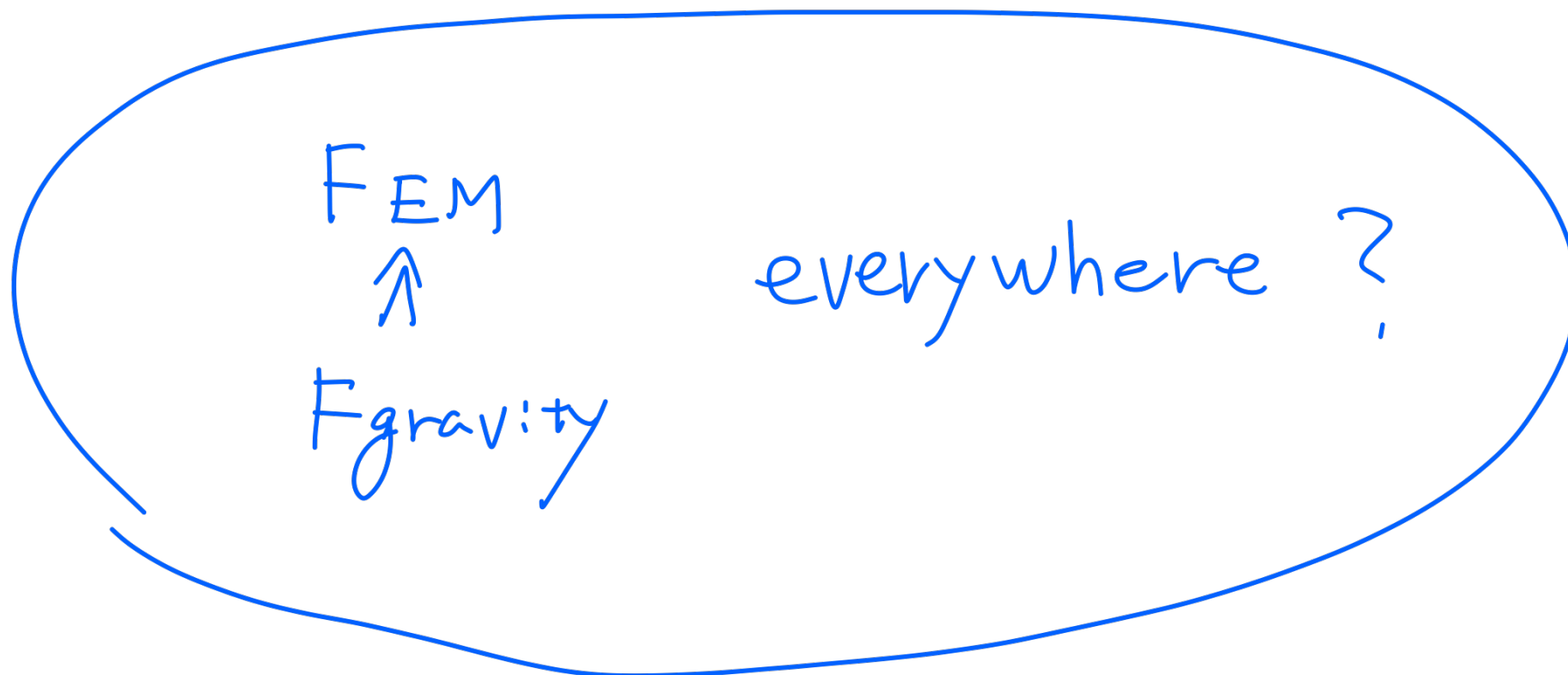
F_{EM}
↑
 $F_{gravity}$

Gravity versus EM

length scale

microscopic

macroscopic



F_{EM} small means e : gauge coupling small

eg. electron

$$e \lesssim \frac{m}{M_{pe}} \sim \frac{0.5 \text{ MeV}}{10^{18} \text{ GeV}} \sim 10^{-21}$$

F_{EM} small means e : gauge coupling small

eg. electron

$$e \lesssim \frac{m}{M_{Pl}} \sim \frac{0.5 \text{ MeV}}{10^{18} \text{ GeV}} \sim 10^{-21}$$

sounds very unnatural,

but is technically natural

$$\mathcal{L} = \int \sqrt{g} \left(\frac{1}{e^2} F_{\mu\nu} F^{\mu\nu} + R + \Lambda \right)$$

things are different if we consider
quantum gravity / strings
(swampland paradigm)



(Gange) WGC

No Global Symmetry

In the extreme limit $e \rightarrow 0$

then $U(1)$ gauge sym. $\rightarrow U(1)$ global sym.

No Global Symmetry

In the extreme limit $e \rightarrow 0$

then $U(1)$ gauge sym. $\rightarrow U(1)$ global sym.

However, semiclassical GR / string
forbids exact global sym.

$\leadsto e \rightarrow 0$ not possible

["known" from long ago;
see e.g. Banks-Seiberg (10)]

(Gauge) Weak Gravity Conjecture

[Arkani-Hamed, Motl, Nicolis, Vafa ('06)]

GWGC

$\exists$ a particle s.t.

$$F_{\text{gauge}} = \frac{(eq)^2}{4\pi r^2} \geq F_{\text{gravity}} = \frac{m^2}{8\pi M_{\text{pl}}^2}$$

(q : charge, m : mass)

"gravity as the weakest force"

(Gauge) Weak Gravity Conjecture

[Arkani-Hamed, Motl, Nicolis, Vafa '06]

GWGC

$\exists$ a particle s.t.

$$m \leq \sqrt{2} e g$$

(q : charge, m : mass)

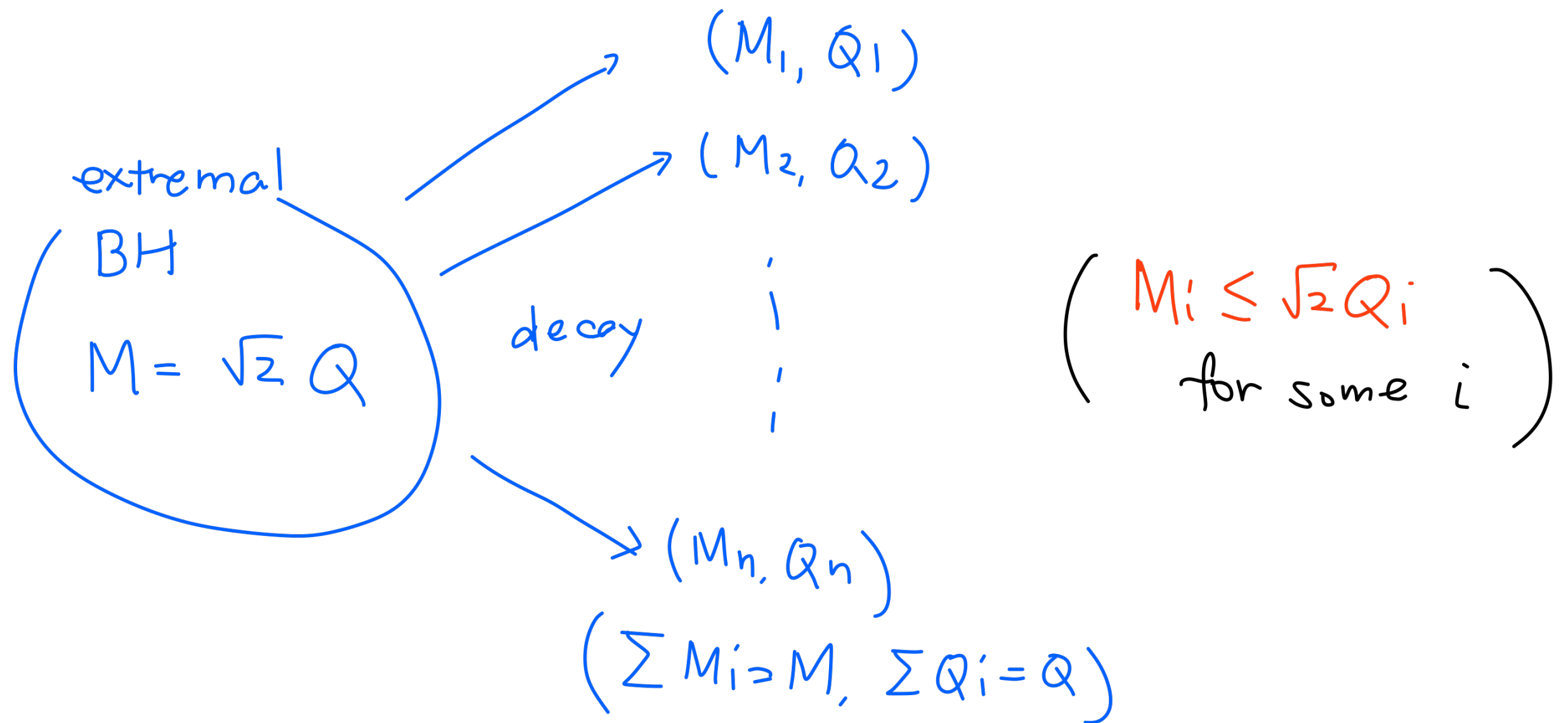
This bound

$$M \leq \sqrt{2} Q$$

is opposite from BH extremality bound

$$M \geq \sqrt{2} Q$$

Decay of extremal BH (if it happens)
requires WGC particle



Gauge-Scalar

WGC



Forces in Nature

@ Gravity

$g_{\mu\nu}$

spin 2

@ EM

A_μ

$U(1)$

spin 1

long-ranged

Weak

A_μ^a

$SU(2)$

Strong

A_μ^a

$SU(3)$

short-ranged

(NOT today)

??

ϕ ?



spin 0

Scalar ?

$F_{\text{grav.}}$, F_{gauge} , F_{scalar}



however, typically not "long-range"

Since scalar ϕ is massive

$F_{\text{grav.}}$, F_{gauge} , F_{scalar}



however, typically not "long-range"

since scalar ϕ is massive

a) SUSY

b) very light ($m \lesssim H^{-1}$)

c) "IR/UV"

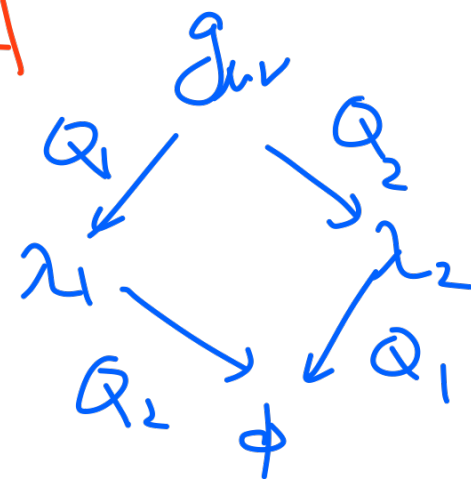
} IR

Let's assume

a) SUSY

for the moment

BPS extremal BH in 4d $N=2$ SUGRA



BH extremality bound

[Ceresole, d'Auria, Ferrara '95]

$$\begin{aligned}
 \underbrace{Q^2}_{\text{gauge}} &\geq |Z|^2 + g^{i\bar{j}} D_i Z D_{\bar{j}} \bar{Z} \\
 &= \underbrace{M^2}_{\text{gravity}} + \underbrace{g^{i\bar{j}} (\partial_i M) (\partial_{\bar{j}} M)}_{\text{scalar}}
 \end{aligned}$$

$$|Z| = M$$

Scalar Force

$$\mathcal{L} \supset m(\phi) \bar{\chi} \chi$$

Scalar Force

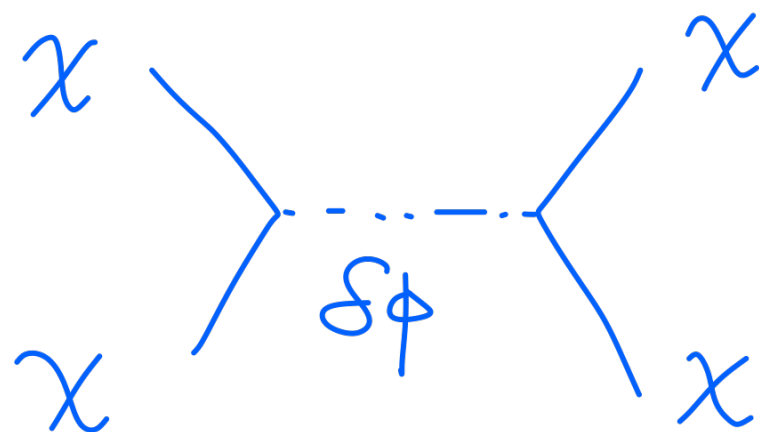
$$\mathcal{L} \supset m(\phi) \bar{\chi} \chi = m(\langle \phi \rangle) \bar{\chi} \chi + \boxed{(\partial_\phi m) \delta \phi \bar{\chi} \chi} + \dots$$

$\uparrow$
 $\phi = \langle \phi \rangle + \delta \phi$

Scalar Force

$$\mathcal{L} \supset m(\phi) \bar{\chi} \chi = m(\langle \phi \rangle) \bar{\chi} \chi + \boxed{(\partial_\phi m) \delta\phi \bar{\chi} \chi} + \dots$$

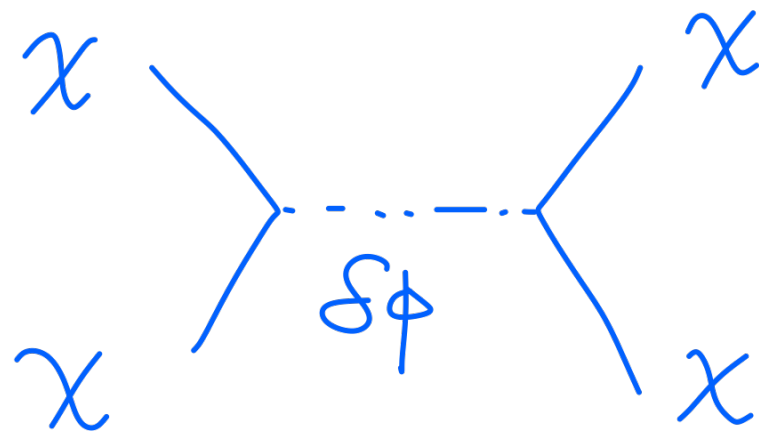
$\phi = \langle \phi \rangle + \delta\phi$



Scalar Force

$$\mathcal{L} \supset m(\phi) \bar{\chi} \chi = m(\langle \phi \rangle) \bar{\chi} \chi + \boxed{(\partial_\phi m) \delta\phi \bar{\chi} \chi} + \dots$$

$\phi = \langle \phi \rangle + \delta\phi$



$$F \sim \frac{(\partial_\phi m)^2}{r^2}$$

" $\partial_\phi m$: scalar charge"

GSGC 1 (repulsive force conjecture)

$$F_{\text{gauge}} \geq F_{\text{gravity}} + F_{\text{scalar}}$$

GSGC 2

[Palti ('17), ...]

$$\frac{Q}{M} \geq \left(\frac{Q}{M} \right)_{\text{extremal BH}}$$

* GSGC 1 $\neq$ GSGC 2 but closely related.

[Heidenreich, Reece, Rudelius ('19)]

Scalar WGC

In many situations, Coulomb repulsions
const energy and charges tends to neutralize

~~F_{gauge}~~

F_{gravity}

F_{scalar}



competition?

We can again start w/ a) SUSY

Previously for extremal BH

$$Q^2 = M^2 + g^{i\bar{j}} D_i M D_{\bar{j}} \bar{M}$$

There is a similar relation

[Ceresole, d'Auria, Ferrara '95]

$$\left(Q \Big|_{N_{IJ} \rightarrow F_{IJ}} \right)^2 = M^2 - g^{i\bar{j}} D_i M D_{\bar{j}} \bar{M}$$

SWGC 1 [Palti ('17, 19)]

$\exists$ particle

$$\frac{\overset{\parallel}{F_{\text{scalar}}}}{\underset{4\pi r^2}{|2\phi m|^2}} \geq \frac{\overset{\parallel}{F_{\text{gravity}}}}{\underset{8\pi M_{\text{pl}} r^2}{m^2}}$$

i.e.,

$$|2\phi m| \geq \frac{m}{M_{\text{pl}}}$$

SWGC 1' [Palti ('17, 19)]

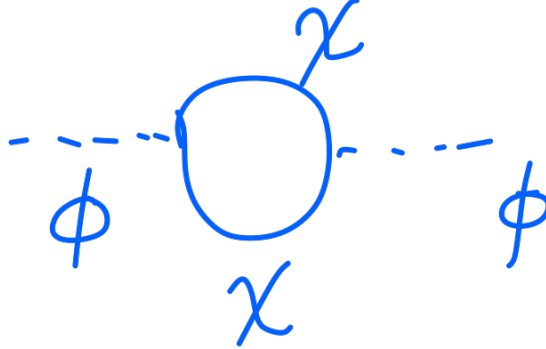
$\forall$ (GWGC) particle

$$F_{\text{scalar}} \geq F_{\text{gravity}}$$

Let's also go beyond a): SUSY

Let's also go beyond a): SUSY

Hard to expect exactly massless scalar

$$\left(\mathcal{L} \supset \phi \bar{\chi} \chi \text{ generates mass} \right)$$


The diagram shows a central circle representing a scalar loop. Four external lines are connected to the circle: a solid line at the top labeled χ , a solid line at the bottom labeled χ , a dashed line on the left labeled ϕ , and a dashed line on the right labeled ϕ .

but can expect very light scalar ϕ

(e.g. ALP, quintessence)

cf. dS conjecture
[Obied-Ooguri-Spodzyneiko]
- Vafa (18)]

b): very light $m \lesssim H^{-1}$

[cf. Shiraishi-Y (19)]

SWGC 1'

$\forall$ (GWGC) particle

$$F_{\text{scalar}} \geq F_{\text{gravity}}$$

This version is too strong

in tension w/ fifth force searches

[Shirai-Y (19)]

($\times$ Loophole by screening mechanism
e.g. chameleon)

different version?

Previously we considered

$$\mathcal{L} \supset \underbrace{m(\phi)}_{\substack{\phi: \text{very light} \\ \text{long-range forces}}} \underbrace{\bar{\chi}\chi}_{\text{general}}$$

What if we use *self-interaction*?

$$\begin{aligned}\mathcal{L} &\supset m(\phi)\phi^2 \\ &= m\phi^2 - A\phi^3 + \lambda\phi^4 + \dots\end{aligned}$$

SWGC 1

$$(2\phi m)^2 \geq \frac{m^2}{M_{Pl}^2}$$



$$m = V'' = V_{\phi\phi}$$

$$(V''')^2 \geq \frac{(V'')^2}{M_{Pl}^2}$$

SWGC 1

$$(2\phi m)^2 \geq \frac{m^2}{M_{Pl}^2}$$



$$m = V'' = V_{\phi\phi}$$

$$(V''')^2 \geq \frac{(V'')^2}{M_{Pl}^2}$$

SWGC 2

[Gonzalo, Ibanez '19]

$$2(V''')^2 - V'' V''' \geq \frac{(V'')^2}{M_{Pl}^2}$$

SWGC 2

[Gonzalo, Ibanez '19]

$$2(V''')^2 - V'' V'''' \geq \frac{(V'')^2}{M_{\text{pl}}^2}$$

if $V'' > 0$

$$2 \frac{(V''')^2}{V''} - V'''' \geq \frac{V''}{M_{\text{pl}}^2}$$

gravity

repulsion

attraction

SWGC 2

$$2(V''')^2 - V'' V'''' \geq \frac{(V'')^2}{M_{pl}^2}$$

if $V'' > 0$

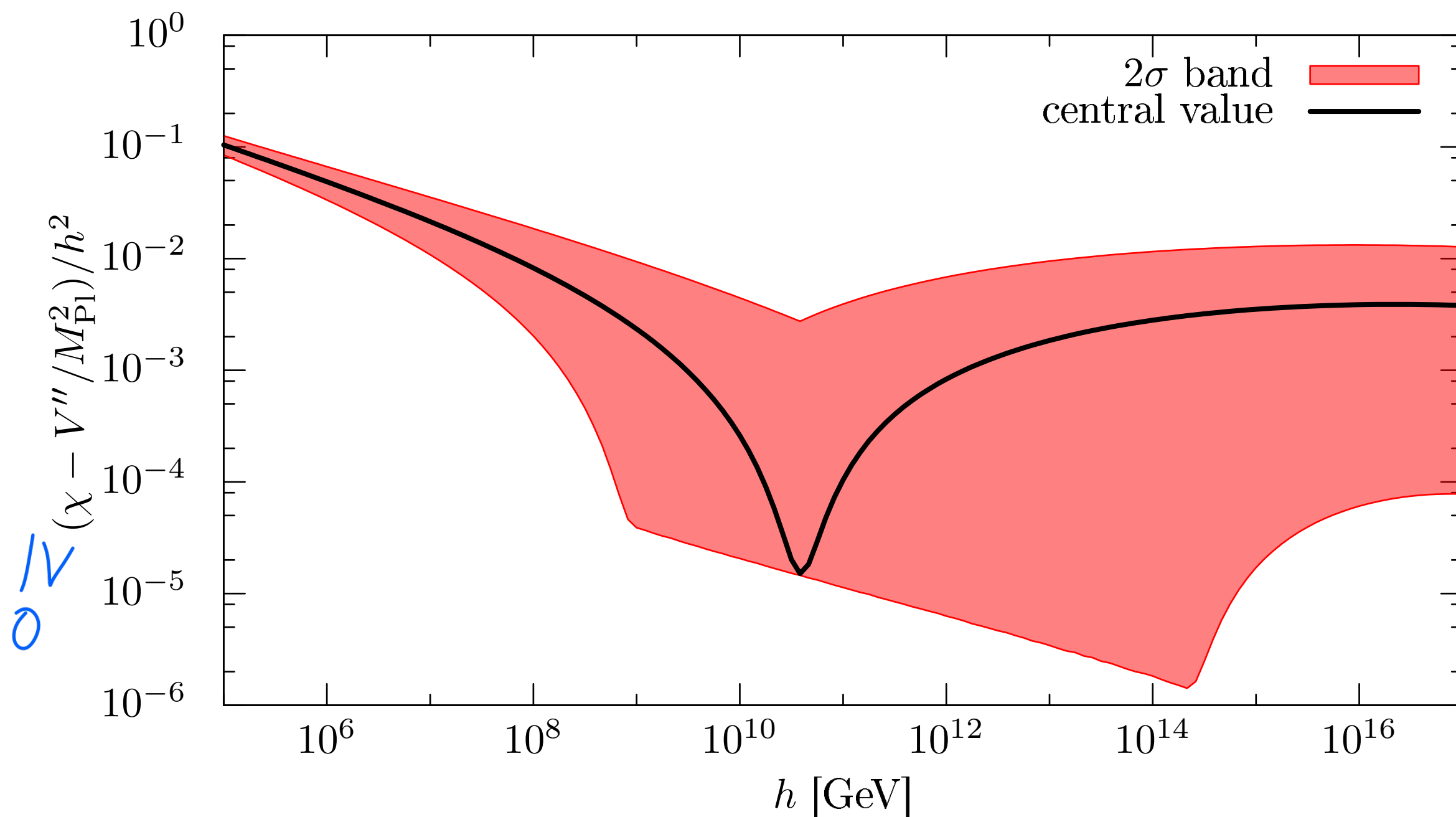
$$2 \frac{(V''')^2}{V''} - V'''' \geq \frac{(V'')^2}{M_{pl}^2}$$

i.e. schematically

$$F_{\text{scalar}}^{\text{net}} = F_{\text{scalar}}^{\text{att.}} - F_{\text{scalar}}^{\text{rep.}} \geq F_{\text{grav.}}$$

SWGc 2

e.g. Higgs [Shirai - Y ('19)]



SWGC 2

$$2(V''')^2 - V'' V'''' \geq \frac{(V'')^2}{M_{\text{Pl}}^2}$$

- This does NOT require a
very light scalar
(in general no long-range scalar force)
- hence c) "UV/IR" statement

SWGC 2

$$2(V''')^2 - V'' V'''' \geq \frac{(V'')^2}{M_{Pl}^2}$$

↓ $M_{Pl} \rightarrow \infty$

$$2(V''')^2 - V'' V'''' \geq 0$$

still non-trivial constraint

VERY different from other

swampland conjectures

(Not a swampland conjecture?)

Currently little support for SWGC2,
likely not true as a general statement?
Contrary to the claims of [Gonzalo (19)
Ibanez]

However, the main point

$$F_{\text{scalar}} \gtrsim F_{\text{gravity}}$$

could be true in many cases?

String moduli

In *String Theory*, "generically"
 $\mathcal{O}(100)$ scalar moduli

Often assumed: gravitational interaction
(cf. moduli problem)

String moduli

In *String Theory*, "generically"

$\mathcal{O}(100)$ scalar moduli

Often assumed: gravitational interaction
(cf. moduli problem)

Scenario changed if $F_{\text{scalar}} \gtrsim F_{\text{grav.}}$

If $F_{\text{scalar}} \geq F_{\text{grav.}}$

for $V(\phi) = \frac{1}{2} m^2 \phi^2 - \frac{A}{3} \phi^3 + \frac{\lambda}{4} \phi^4$

The diagram illustrates the relationship between the potential $V(\phi)$ and its derivatives. An arrow labeled V'' points to the m^2 term in the quadratic part of the potential. Two arrows labeled V''' point to the A and λ terms in the cubic and quartic parts of the potential, respectively.

If $F_{\text{scalar}} \geq F_{\text{grav.}}$

for $V(\phi) = \frac{1}{2} m^2 \phi^2 - \frac{A}{3} \phi^3 + \frac{\lambda}{4} \phi^4$

V'' points to m^2 , V''' points to A , and V'''' points to λ .

SWGC $\rightarrow$ $\underbrace{A^2}_{F_{\text{att. scalar}}} - \underbrace{\lambda m^2}_{F_{\text{rep. scalar}}} \gtrsim \underbrace{\left(\frac{m^2}{M_{\text{pl}}}\right)^2}_{F_{\text{grav.}}}$

If $F_{\text{scalar}} \geq F_{\text{grav.}}$

for $V(\phi) = \frac{1}{2} \underbrace{m^2}_{V''} \phi^2 - \frac{A}{3} \phi^3 + \frac{\lambda}{4} \phi^4$

$\swarrow V''' \quad \searrow V''''$

SWGC $\rightarrow \underbrace{A^2}_{F_{\text{att. scalar}}} - \underbrace{\lambda m^2}_{F_{\text{rep. scalar}}} \gtrsim \underbrace{\left(\frac{m^2}{M_{\text{Pl}}}\right)^2}_{F_{\text{grav.}}}$

λ is small

gravity small

$$A \gtrsim \frac{m^2}{M_{\text{Pl}}}$$

$$A^2 \gtrsim \lambda m^2$$

Q-ball formation
(Floquet exponent $\gtrsim H$)

oscillon formation

[Kusenko - Takistov - Yamada - Y]
(19)

If $F_{\text{scalar}} \geq F_{\text{grav.}}$

Scalar fields fragment into **localized lumps**

after inflation

[Kusenko-Takhistov-Yamada-Y ('19)]

Q-ball / oscillon [Kusenko-Shaposhnikov ('97), ...]

PBH [Alex's talk]

⋮

as DM?

Summary

GWGC

$$F_{\text{gauge}} \geq F_{\text{gravity}}$$

Swampland

QG/string?

Summary

GWGC

$$F_{\text{gauge}} \geq F_{\text{gravity}}$$

Swampland
QG/string?

GSWGC

$$F_{\text{gauge}} \geq F_{\text{scalar}} + F_{\text{gravity}}$$

SWG C

$$F_{\text{scalar}} \geq F_{\text{gravity}}$$

most controversial but most powerful

Q: Is gravity the
weakest force?

Q: If so, in what sense?
why?