

Lec 5

2021 / Jul 19

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Today

\* Factorization algebra  $\rightsquigarrow$  Yangian  
Koszul duality

Lec 5

\* Integrable Field Theory  $\rightsquigarrow$  4d CS + surface defect

↓  
"uplift"

\* Twistor, ASDYM, hol CS

Lec 6

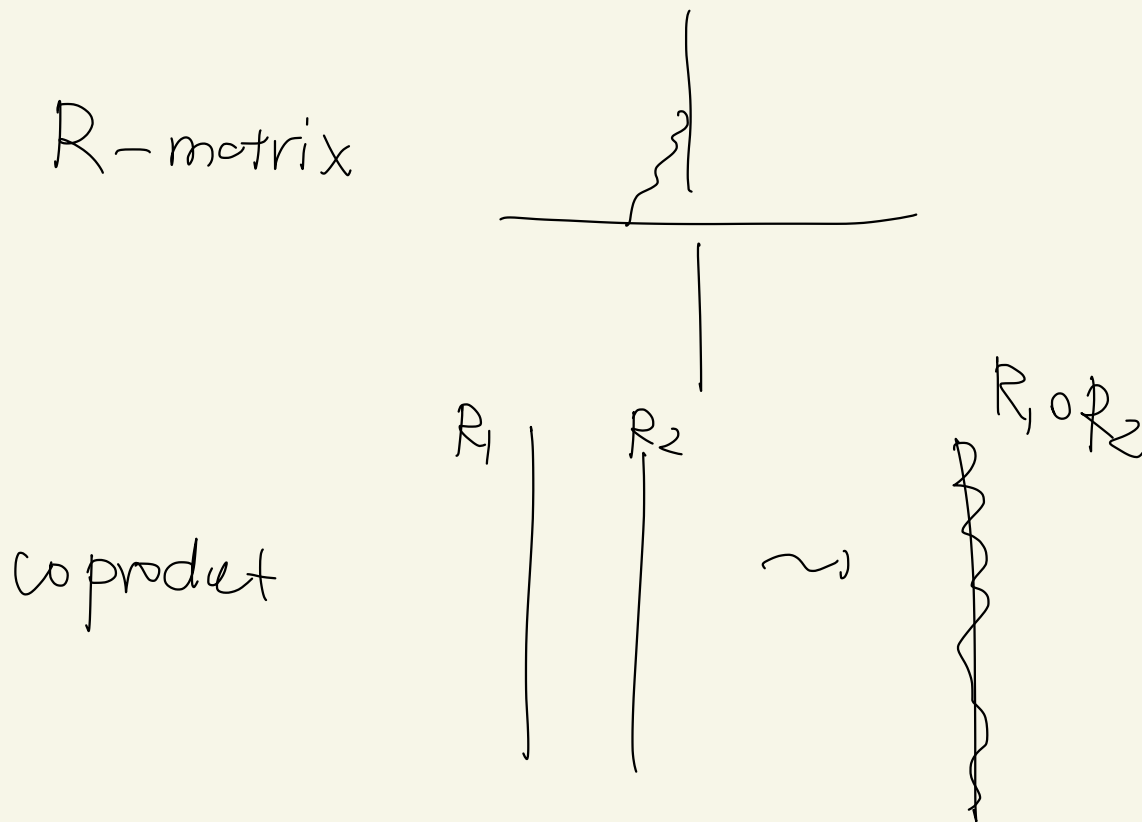
Yesterday

3d  $k=2$  line

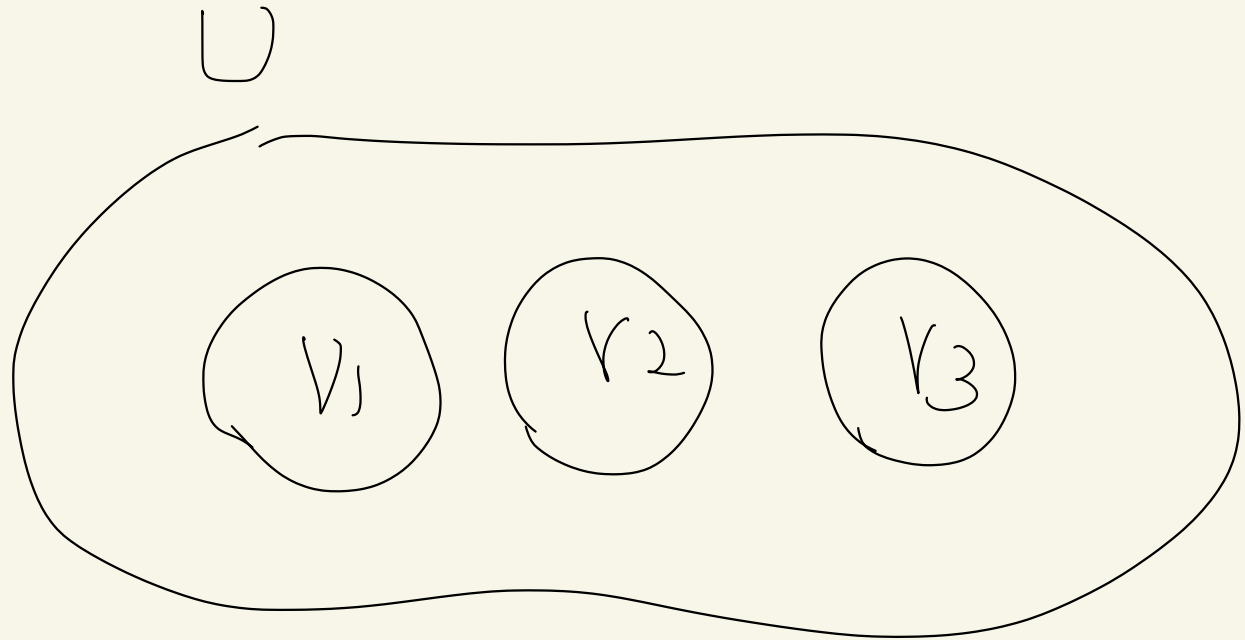
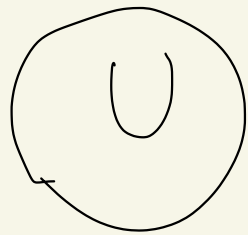
$Y_{\hbar}(\mathfrak{g})$ : Yangian  $\leftarrow$  Feynman diagram computation of  
 $\downarrow \hbar=0$  4d CS  $\hbar$   
 $U(\mathfrak{g}[[\hbar]])$  "algebra of loop operators"

g-cluster alg.

3-mfd



(P)FA



$F(U)$   
↕

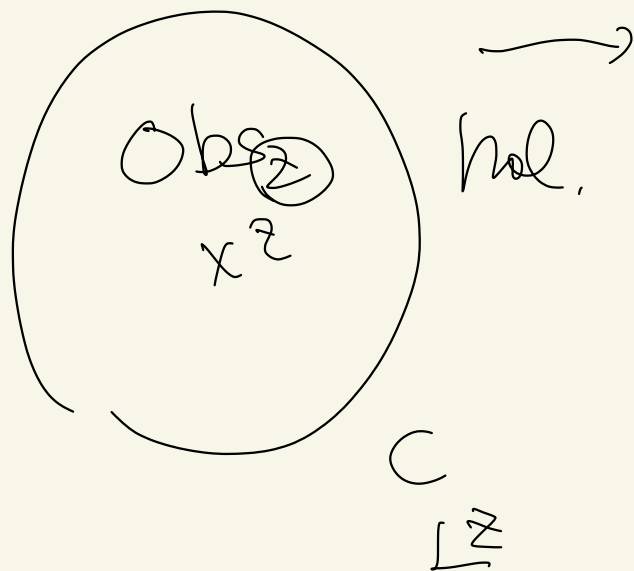
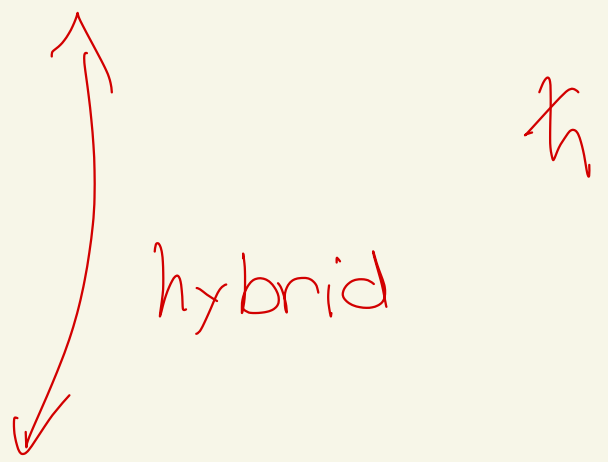
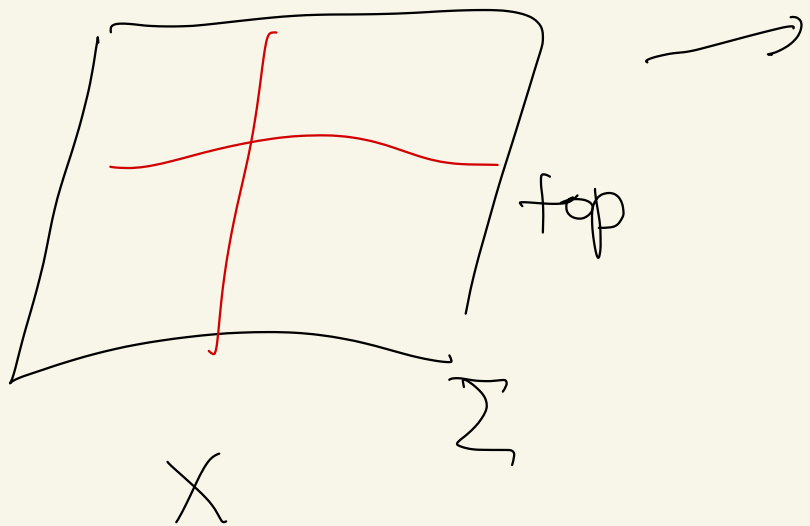
cochain complex  
operators on  $U$

$$\bigoplus_i F(v_i) \rightarrow F(U)$$

4d CS

2d

TQFT : locally const. FA / E<sub>2</sub>-alg



2d holomorphic / translationally inv.



OPE / VOA

$T_2: z \rightarrow z + k$

$Obs_0 \otimes Obs_0 \rightarrow Obs_0 \otimes \mathbb{C}((z))$

classical,  $\hbar = 0$

Obs  $\hbar=0$   
 $z=0$  = { classical alg of local op. of 4d CS at  $z=0$  } (over  $\mathbb{C}$ )

4d bulk =  $\mathbb{F}_2$ -alg =  $\mathbb{C}^* \left( \mathcal{G}[[z]] \right)_{d_{CE}}$

$\mathcal{OBRST}(\mathbb{C}) = \int_{bc} \mathbb{C}^b \mathbb{C}^c \langle \mathbb{C}, \mathbb{C} \rangle$

Koszul duality

$S \simeq d_2 \wedge A \wedge dA$

$A \rightarrow g^t Ag + g^{-1} dg$

1d defect

$z=z_0$   
 $\mathcal{U}(\mathcal{G}[[z]])$

= { classical alg of loop op located at  $z=0 \in \mathbb{C}$  }  
 $z_0$   
Hopf alg

$\mathcal{G}[[z]]$

Koszul duality: (universal) coupling between  
bulk & defect

phys

$A$ : bulk alg.  $\xrightarrow{KD}$   $(A!)$   
 $B$ : line defect alg.

$(A \& B: \text{coupling}) \iff \text{Hom}(A!, B)$

$MC(A \otimes B)$

Maurer-Cartan

# Koszul duality

$V$ : vect sp.

(math)

$$A = T(V) \cong \bigoplus_{n \geq 0} V^n$$

$$\mathbb{R} \subset V \otimes V$$

$$\text{Sym}(V)$$

KD

$$A^! = T(V^*)$$

$$\mathbb{R}^+ \subset V^* \otimes V^*$$

$$\Lambda^*(V^*)$$

$$\mathcal{U}(\mathfrak{g}[z])$$

↓ Abelian

$$\text{Sym}(\mathfrak{g}[z])$$

KD

$$\Lambda^*(\mathfrak{g}^*[z])$$

↑ (Abelian)  $d_{CE} = 0$

$$C^*(\mathfrak{g}[z], d_{CE})$$



Quantum mechanically

$$\text{Obs}_{z=0}^{\hbar \neq 0} \leftarrow \mathbb{C}[[\hbar]]$$

bulk

•  $E_2$ -algebra (w/ augmentation)

•  $T_\lambda: \text{Obs}_{z=0}^{\hbar} \rightarrow \text{Obs}_{z=\lambda}^{\hbar}$  translation

• OPE  $\text{Obs}_0 \otimes \text{Obs}_0 \rightarrow \text{Obs}_0 \otimes \mathbb{C}((z))$

top  
transl.  
hol.

KD

$$z \rightarrow z + \lambda$$

$U(\mathfrak{g}[[z]])$

defect

$Y_{\hbar}(\mathfrak{g})$

• Hopf alg

$U(\mathfrak{g}[[z]])$

•  $T_\lambda Y_{\hbar}(\mathfrak{g}) \rightarrow Y_{\hbar}(\mathfrak{g})$

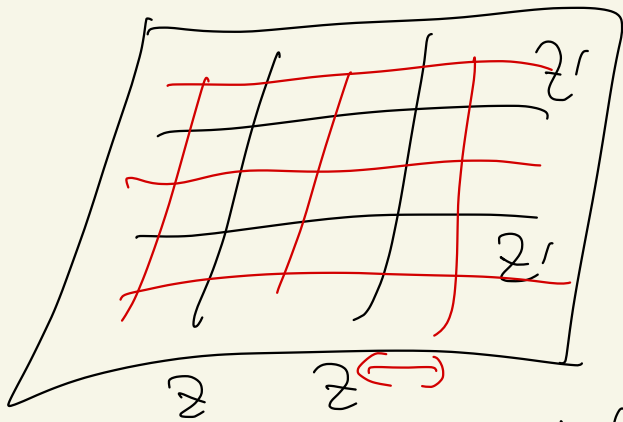
$Y_{\hbar}(\mathfrak{g}) \otimes \mathbb{C}[[\hbar]] \rightarrow \mathbb{C}[[\hbar]]$

Yangian

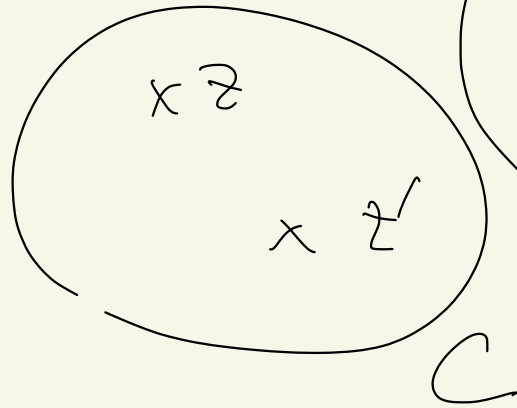
•  $R_{V,W}: T_\lambda W \otimes V \cong V \otimes T_\lambda W$

Integrable Field

Theories

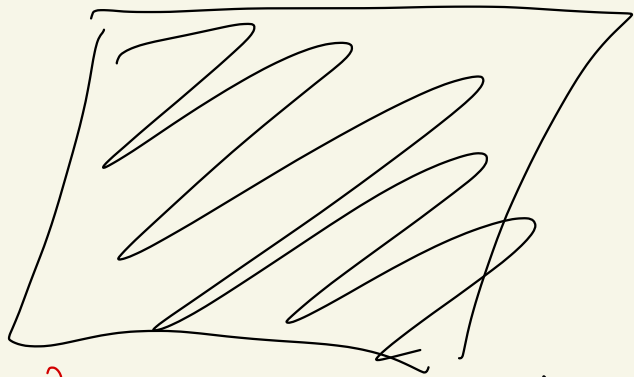


1d defect



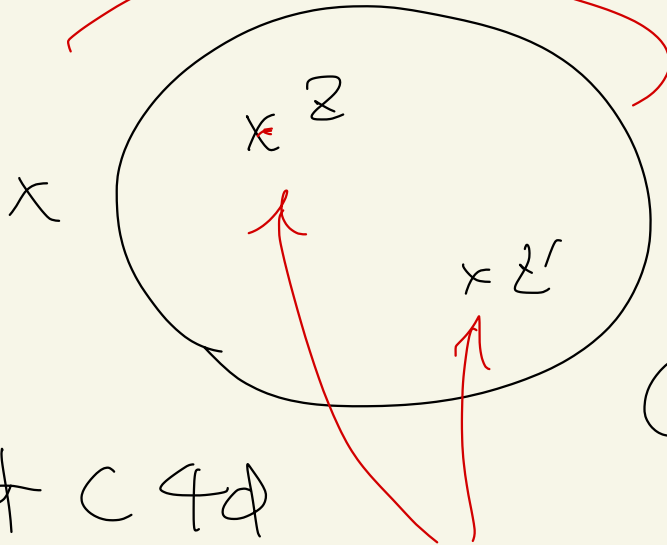
disorder  $\leftarrow$  A b.c.

order  $\leftarrow$  couple to defect theory



$\mathbb{R}^2$

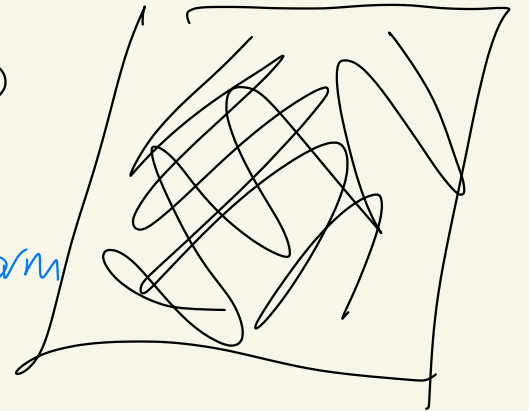
2d defect  $\subset$  4d  
"surface defect"



b.c. for gauge field A

"integrate out C"

$\omega$   
C'  $\xrightarrow{\text{rel. form}}$



effective 2d integrable FT

$\mathbb{R}^2$

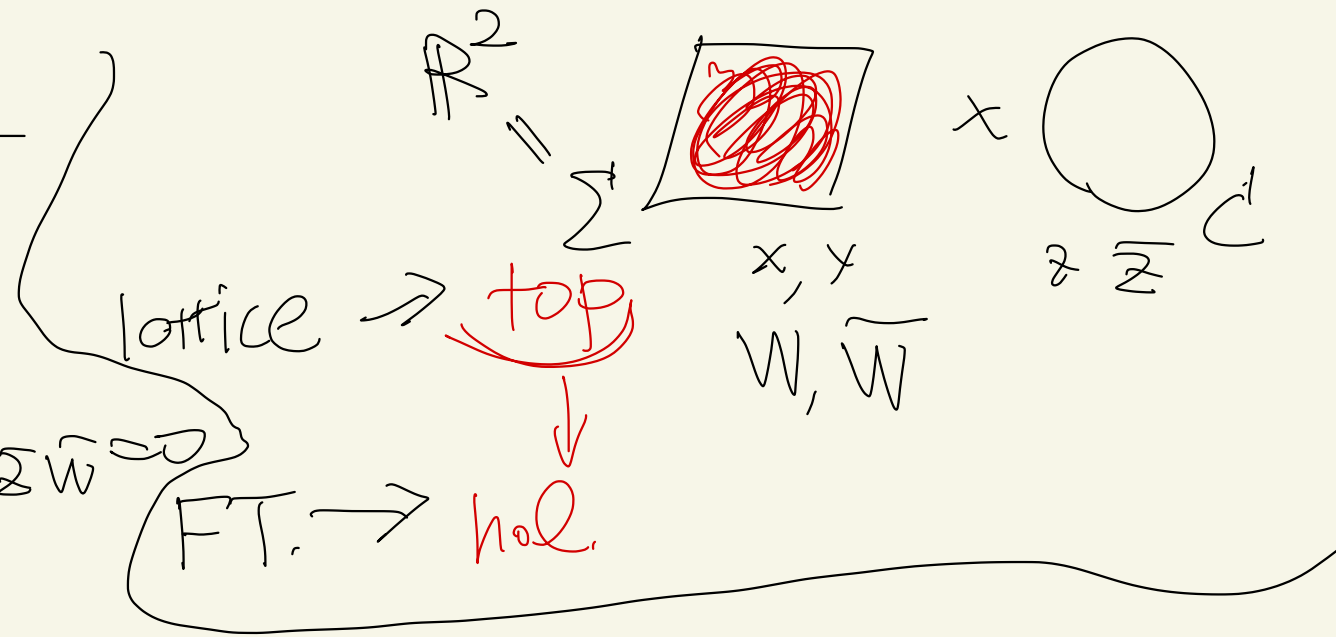
Why integrable?

4D CS



e.o.m.  $F_{w\bar{w}} = F_{\bar{z}w} = F_{z\bar{w}} = 0$

↓ ( $A_{\bar{z}} = 0$  gauge)



\*  $\partial_{\bar{z}} A_w = \partial_{\bar{z}} A_{\bar{w}} = 0$   $\leftarrow$   $A_w, A_{\bar{w}} : z$ -hol.

@  $\partial_w A_{\bar{w}} - \partial_{\bar{w}} A_w + [A_w, A_{\bar{w}}] = F_{w\bar{w}} = 0$   $\leftarrow$   $w\bar{w}$ -direction flat

↓

\*  $\mathcal{L}(z) \equiv A_w(z) dw + A_{\bar{w}}(z) d\bar{w} : 1$ -form on  $\mathbb{R}^2 = \sum_{w, \bar{w}}$

@  $d\mathcal{L}(z) + \mathcal{L}(z) \wedge \mathcal{L}(z) = 0 : Lax$  connection  
"classical integrable"

$$\partial_t \left( \text{Tr}_R P \exp \int_{\mathbb{R}} \mathcal{L}(z) \right) = 0$$

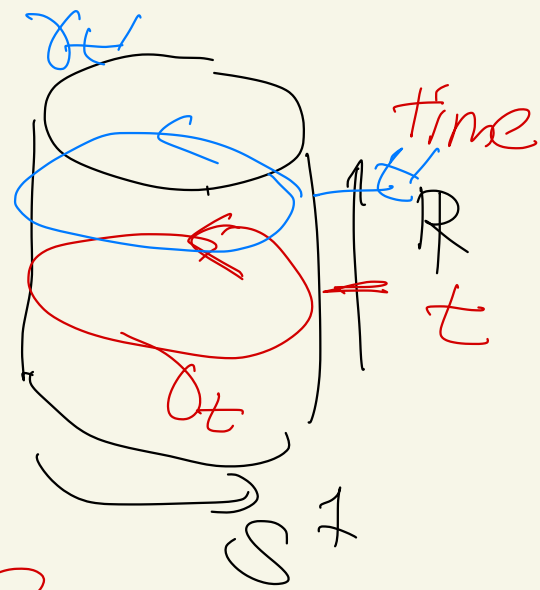
Where

$$\mathcal{L}(z) = \sum_{n \geq 0} Q_n / z^n$$

$Y_n(\sigma)$   $\leftarrow$  level  $n$

$\infty$ -many charges conserved

$\mathbb{R}^2 \rightarrow$



$$\partial_t Q_n = 0$$

Yangian non-local charge of

integrable lattice model (top. along  $\Sigma = \mathbb{R}^2$ )

$\omega$ : hol. 1-form  $\leftarrow$  no zero.



$C = \mathbb{C}, \mathbb{C}^x, \mathbb{E}$

$\begin{pmatrix} A & \omega \\ A & \bar{\omega} \end{pmatrix}$

integrable field theory (hol. along  $\Sigma$ )

$\omega$ : hol. 1-form  $\leftarrow$  zeros allowed



$C$ : higher genus ok.

$$S \Rightarrow \int \omega \wedge \text{Tr}(A \wedge dA)$$

Diagram illustrating the integral  $S \Rightarrow \int \omega \wedge \text{Tr}(A \wedge dA)$ . The integrand is enclosed in a blue box. Red arrows point from  $\omega$  to  $\infty$  and  $0$ . Blue arrows point from  $\text{Tr}(A \wedge dA)$  to  $0$  and  $\infty$ .

$\square \phi \Rightarrow$

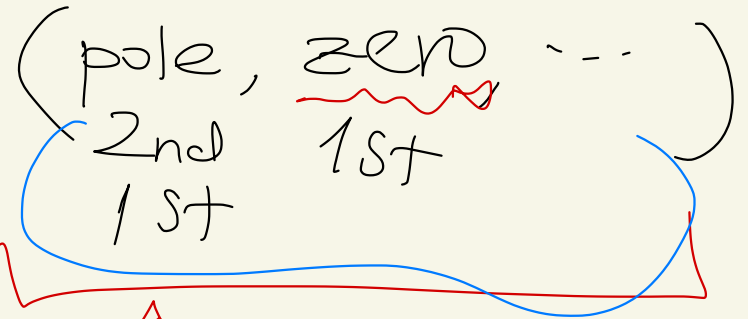
general data

- spectral curve  $C$
- hol. one-form  $\omega$  on  $C$
- b.c. for  $A$   
consistent  $\omega$

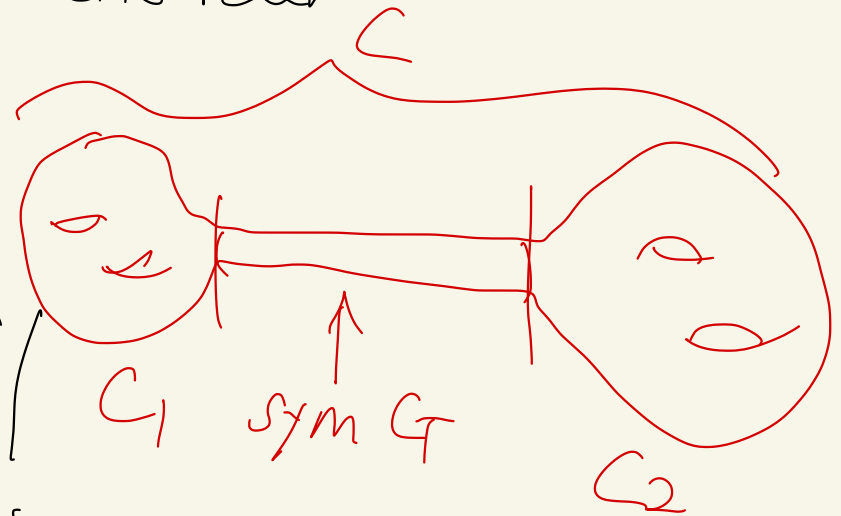
[ order surface defect  
(coupling to 2d FT) ]

2d  
integrable field theory

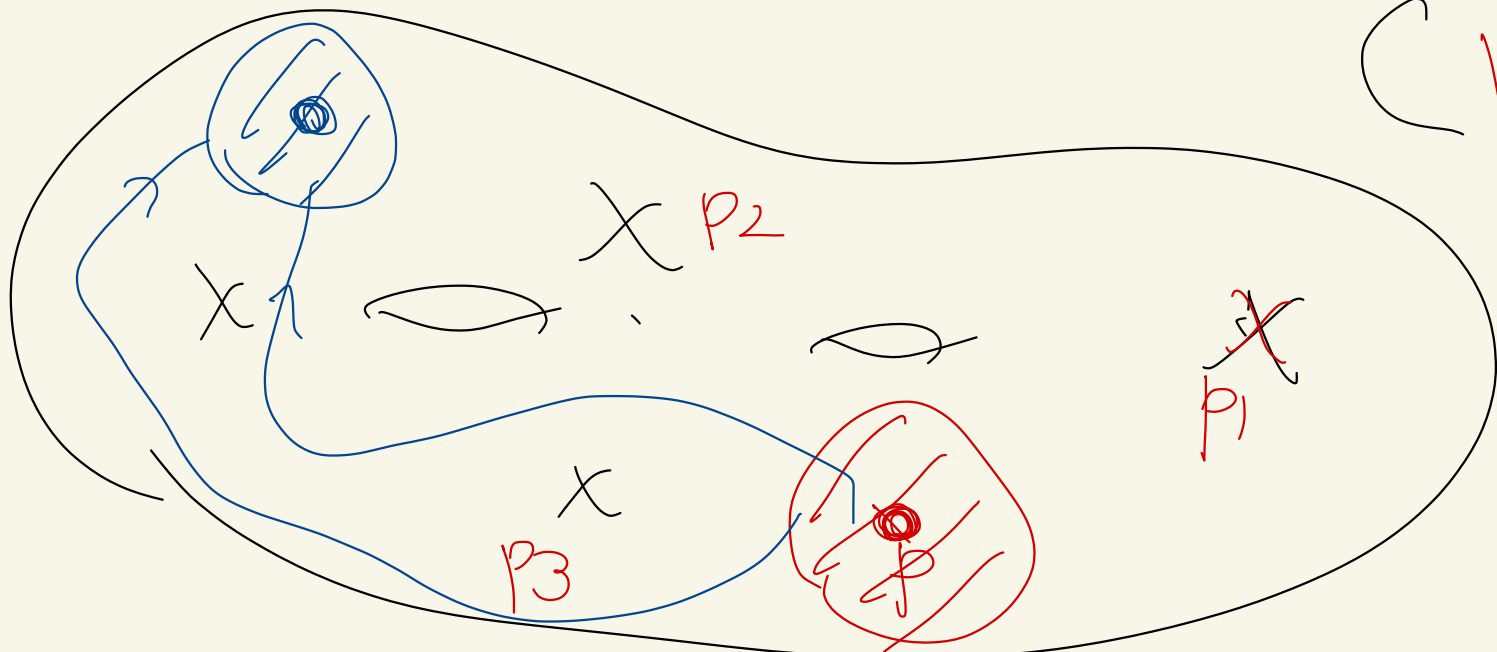
translation  $g > 1$  <sup>surface</sup>  
(higher genus  $0 \neq k$ )



↑ surface defect  
disorder



$$T_C \cong (T_{C_1} \times T_{C_2}) // G$$



$\mathbb{C}$   
 $\{p_i\}$

$$W_0(z) \Big|_{z=p} (z-p)^n Q_n$$

$\downarrow$   
 conserved charge