

Development of a Spectral Model Based on Charge Transport for the Swift/BAT 32K CdZnTe Detector Array

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Abstract

The properties of 32K CdZnTe ($4 \times 4 \text{ mm}^2$ large, 2 mm thick) detectors have been studied in the pre-flight calibration of the Burst Alert Telescope (BAT) on-board the Swift Gamma-ray Burst Explorer (scheduled for launch in November 2004). In order to understand the energy response of the BAT CdZnTe array, we first quantify the mobility-lifetime ($\mu\tau$) products of carriers in individual CdZnTe detectors, which produce a position dependency in the charge induction efficiency and results in a low energy tail in the energy spectrum. Based on a new method utilizing ^{57}Co spectra obtained at different bias voltages, the $\mu\tau$ for electrons ranges from 5.0×10^{-4} to $1.0 \times 10^{-2} \text{ cm}^2\text{V}^{-1}$ while the $\mu\tau$ for holes ranges from 1.3×10^{-5} to $1.8 \times 10^{-4} \text{ cm}^2\text{V}^{-1}$. We find that this wide distribution of $\mu\tau$ products explains the large diversity in spectral shapes between CdZnTe detectors well. We also find that the variation of $\mu\tau$ products can be attributed to the difference of crystal ingots or manufacturing harness. We utilize the 32K sets of extracted $\mu\tau$ products to develop a spectral model of the detector. In combination with Monte Carlo simulations, we can construct a spectral model for any photon energy or any incident angle.

Key words: Gamma-ray astronomy, Swift, BAT, CdZnTe, mobility-lifetime products, spectral model

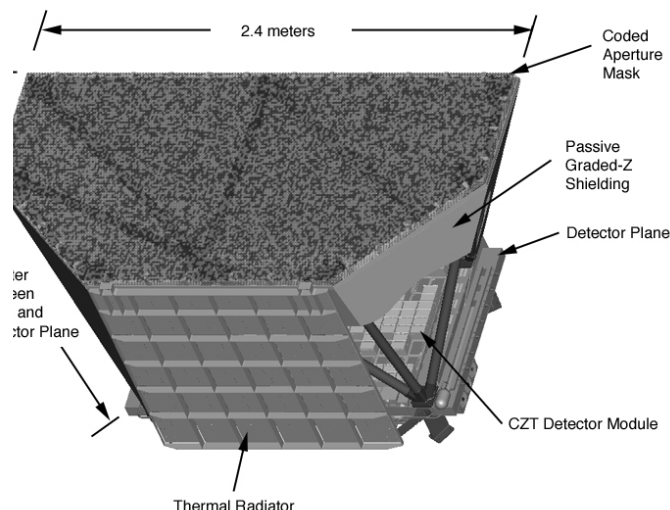


Fig. 1. The Burst Alert Telescope (BAT) on-board the Swift spacecraft. The BAT has a 3 m² D-shaped coded aperture mask with 5 mm pixels. The CdZnTe array is 5,243 cm² with 4 mm detectors. The Graded-Z Fringe Shield, which is not fully drawn to show the detector array, reduces the background due to cosmic diffuse emission. (NASA/GSFC)

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1 Introduction

Swift, scheduled for launch in November 2004, is a NASA MIDEX mission dedicated to the study of astronomical gamma-ray bursts (GRBs) [1]. It carries the Burst Alert Telescope (BAT) which will detect and localize the GRBs. The BAT is a coded aperture instrument with a 1.4 steradian field-of-view (half coded) (shown in Fig. 1). The detector plane is an assembly of 32,768 planar Cadmium Zinc Telluride (CdZnTe, CZT) detectors each 4 × 4 mm² and 2 mm thick. They form a 1.2 m × 0.6 m sensitive area, and provide imaging capability and spectroscopy [2]. In addition to detecting GRBs, BAT will also conduct a sensitive hard X-ray all-sky survey over the course of the mission.

Recently, Cadmium Telluride (CdTe) and CdZnTe semiconductor materials have attracted much attention as hard X-ray and gamma-ray detectors among the semiconductor materials with a wide band gap. Their good energy resolution and the fact that these can be fabricated into compact arrays make them very attractive in comparison with inorganic scintillation detectors coupled to either photodiodes or photo-multiplier tubes (e.g. [3]). Because of these advantages, a large CdTe gamma-ray camera is already operating in space for the International Gamma Ray Astrophysics Laboratory (INTEGRAL) mission. It is also a coded aperture instrument

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which utilizes 16,384 planar CdTe ($4 \times 4 \text{ mm}^2$, 2 mm thick) detectors to form a sensitive area of $2,621 \text{ cm}^2$. This provides high resolution spectroscopy with an energy resolution of 9% (FWHM) at 100 keV and fine imaging with an angular resolution of 12 arcmin [4] [5].

However, despite recent advances, the considerable charge loss in CdTe and CdZnTe limits their capability as high resolution spectrometers. Due to the low mobility and the short lifetime of the carriers, the electron and hole pairs generated by gamma-ray irradiation cannot be fully collected. As a result, a broad low energy tail is seen in the detector's energy spectrum. In the case of CdZnTe material, it is also known that the current High Pressure Bridgman (HPB) technique produces only polycrystals with a non-uniform distribution of charge transport properties. This can result in some degree of variation in charge transport properties between CdZnTe detectors. Therefore, it is necessary to measure these properties for each CdZnTe detector in order to understand the BAT energy response.

In this paper, we report the features of individual CdZnTe spectra. Then we evaluate the key mobility and lifetime parameters both for electrons and holes. Finally, we describe how these are used to develop the spectral model for the individual detectors and for the entire CdZnTe array.

2 CdZnTe array in the Burst Alert Telescope

The BAT detector plane ($5,243 \text{ cm}^2$) is composed of 32,768 pieces of $4 \text{ mm} \times 4 \text{ mm} \times 2 \text{ mm}$ CdZnTe detector elements. Groups of 128 detector elements are assembled into 8×16 sub-arrays. The signals from each of the 128 detectors are read out by a single analog signal processing ASIC (XA1). This 128-channel XA1 readout chip was designed and produced by IDEAS (Hovik, Norway). A Detector Module (DM) contains two such sub-arrays, and a block is made of eight DMs. There are in total 16 blocks mounted in the Detector Array Plane.

The pulse heights are digitized by an ADC unit on the XA1 and transmitted to a data handling system along with the channel ID for the event. Thus, the events are accumulated into individual spectra for each CdZnTe pixel. There is an electronic calibration pulser circuit within each DM which can sequentially inject a specified number of charge pulses of a specified level into each of the 128 channels. The XA1 has a dynamic range up to $\sim 500 \text{ keV}$ which is not perfectly linear, and we can correct the offset, gain, and nonlinearity of each of the 32K detectors with the use of this calibration pulser.

The CdZnTe array will have a nominal operating temperature of $20 \text{ }^\circ\text{C}$ and its thermal gradients (temporal and spatial) will be kept to within $\pm 1 \text{ }^\circ\text{C}$. The typical bias voltage is 200 V, with a maximum of 300 V. The CdZnTe material used in the BAT is manufactured by eV Products Inc. (Saxonburg, PA) using the HPB technique. All detectors were tested by the manufacturer to meet minimum performance specifications. In addition to mechanical and visual inspections, the detector performance specifications listed in Table 1 were required for the energy spectrum from a ^{57}Co radioactive source. A typical ^{57}Co spectrum obtained with a single CdZnTe detector is shown in Fig. 2. We can see 14.4, 122 and 136 keV lines and also a prominent low energy

Table 1

Required detector performance specifications. Peak-to-valley ratio at 14.4 keV is defined as $\text{Peak}_{14}/L_{\text{Valley}}$ where Peak_{14} = average counts from 3 channels around 14.4 keV and L_{Valley} = average counts from 3 channels in the valley on the left of the 14.4 keV peak. Peak-to-valley ratio at 122 keV is defined as $\text{Peak}_{122}/\text{Peak}_{100}$ where Peak_{122} and Peak_{100} are averages counts from 3 channels around 122 and 100 keV, respectively.

Detector performance	Requirement
Peak-to-valley ratio at 14.4 keV	> 4.5
Peak-to-valley ratio at 122 keV	> 2.5
FWHM of the 122 keV peak	< 9.5 keV
FWHM of the pulser peak	< 3.75 keV
Leakage current drawn by each detector	< 20 nA
Leakage current stability	> 3 hrs
Measured 122 keV detection efficiency	> 0.75

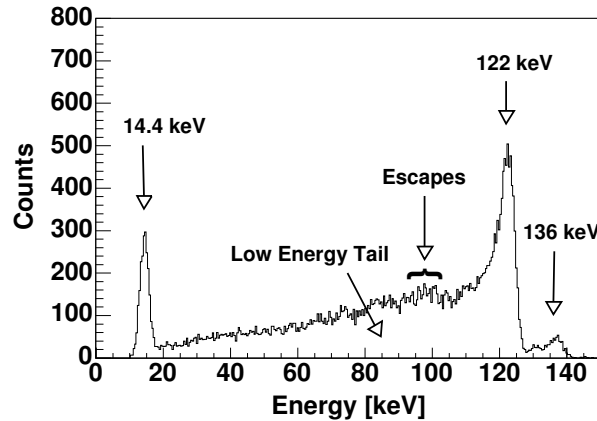


Fig. 2. A typical ^{57}Co spectrum obtained with a single CdZnTe detector, showing 14.4, 122 and 136 keV lines and also a prominent low energy tail which accompanies the 122 keV line. Small humps around 100 keV are escape peaks associated with the 122 keV main peak.

tail which accompanies the 122 keV line. Small humps around 100 keV are escape peaks associated with the 122 keV main peak, which are due to the escape of photo-fluorescence photons (23 keV for Cd and 27 keV for Te) from the detector. As described in the Introduction, the charge transport properties of the CdZnTe material were the major issue for the selection of CdZnTe detectors. Therefore, peak-to-valley ratios calculated at 122 and 14.4 keV lines were used in the screening.

3 Spectral Properties

The intensive pre-flight calibration has been performed at the NASA/Goddard Space Flight Center [6]. The calibration is divided into two phases. In the first phase, measurements were made with each block individually. This was carried out right after the blocks were constructed. In the second phase, measurements were made with the entire array inside a large clean tent,

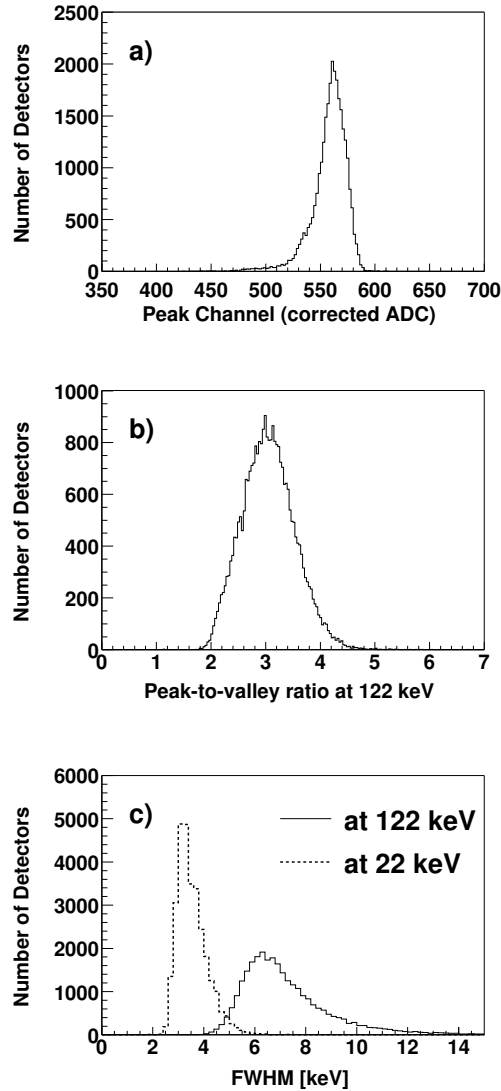


Fig. 3. Distributions of spectral parameters for 32K CdZnTe detectors, a) peak channel of 122 keV line (after corrected for electronics), b) peak-to-valley ratio at 122 keV, and c) energy resolutions (FWHM) at 122 keV and 22 keV.

after all sixteen blocks were installed on the detector array plate. In order to understand the BAT instrument's response to incident 15-150 keV gamma-rays, we acquired spectra using various radioactive sources (namely, ^{241}Am , ^{57}Co , ^{109}Cd , ^{133}Ba and ^{137}Cs) at different bias voltages, temperatures and incident angles.

3.1 Properties of Individual CdZnTe Spectra

Even after correction for electronics gain, offset and nonlinearity, we still observe a significant variation in peak channels as can be seen in Fig. 3 (a), which shows a distribution of the peak channels of the 122 keV line extracted from 32K spectra. The data were obtained during the

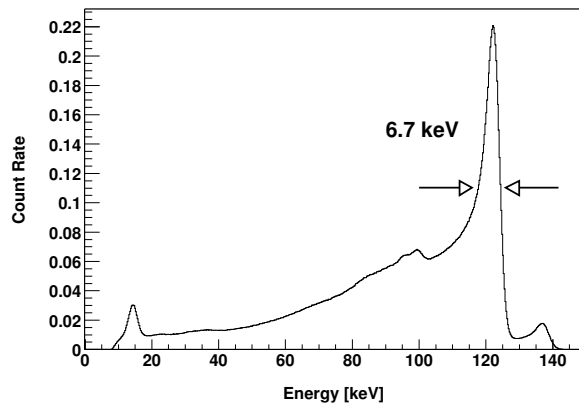


Fig. 4. Summed spectrum of 32K CdZnTe detectors (the entire BAT array) for ^{57}Co . The 122 and 136 keV lines are clearly seen. The energy resolution is 6.7 keV (FWHM) at 122 keV.

block calibration with a ^{57}Co source placed 1 m above the center of the block detector plane. The detectors were operated with a bias voltage of 200 V at a temperature of 20 °C.

The peak channel varies more than 10% from detector to detector. We consider this to be a result of differences in charge induction efficiency between detectors. The peak-to-valley ratio is another parameter which characterizes the charge transport properties. We define it as the ratio of counts at 122 keV to counts at 100 keV where the counts at these energies are averages over 5 neighboring channels. As shown in Fig. 3 (b), the peak-to-valley ratio ranges from 2 to 4.

The energy resolutions measured at 122 keV are shown in Fig. 3 (c) along with those at 22 keV derived from ^{109}Cd spectra. The average value of FWHM at 122 keV is much greater than at 22 keV, even after taking into account the expected increase due to Fano noise. This is because the 122 keV peak has an asymmetric shape due to the low energy tail, which is prominent only for higher energy lines. The wider distribution of FWHM at 122 keV is due to the variation in the shape of low energy tail.

3.2 Summed Spectrum for the Entire CdZnTe Array

The variation of peak channel shown in §3.1 indicates that it is not sufficient to correct only for the readout electronics when we construct a summed spectrum for the entire CdZnTe array. The shift of photo-peak due to the charge transport properties in CdZnTe materials should also be corrected. For this, we simply use the actual peak channels of the 22.1, 59.5, 88.0 and 122.1 keV lines. The peak channels are fitted well with a quadratic function of energy, which is used to convert the ADC channel to energy for each 32K CdZnTe pixel.

To obtain the ^{57}Co summed spectrum, we acquire data at the level of the array calibration, in which the detector array is located 1 m below the coded aperture mask. The ^{57}Co source is placed at 3 m above the center of the detector plane, or 2 m above the center of the coded

aperture mask. Therefore, the detector signals are modulated by the mask pattern. To obtain the summed spectrum, we also applied the “mask weighting” method, which utilizes the mask modulation (i.e., how much the detectors are shadowed) as summation weights [7]. This enables us to eliminate flat non-modulated components such as scattering. The summed ^{57}Co spectrum obtained with this method is shown in Fig. 4. The 122 and 136 keV lines are clearly seen. The energy resolution is 6.7 keV (FWHM) at 122 keV.

4 Charge Transport Properties in Individual CdZnTe Detectors

4.1 Fitting Procedure to Extract $\mu\tau$ products

The mobility and lifetime for electrons and for holes are key parameters when we model the spectral response of a CdZnTe detector, because once we know these values, the charge induction efficiency can be calculated at each depth in the detector and applied to any incident photon energy. Here we deduce $(\mu\tau)_e$ and $(\mu\tau)_h$ for each detector by following the concept described in Sato et al.[8]. We first derive a spectral fitting function to be applied to the 122 keV peak in the ^{57}Co spectrum. We assume the ideal situation in which the electric field (E) throughout the detector volume is constant. Then the charge induction efficiency is given by the Hecht equation [9]

$$\eta(x) = \frac{(\mu\tau)_e E}{D} \left[1 - \exp\left(-\frac{D-x}{(\mu\tau)_e E}\right) \right] + \frac{(\mu\tau)_h E}{D} \left[1 - \exp\left(-\frac{x}{(\mu\tau)_h E}\right) \right], \quad (1)$$

where D is the detector thickness and $(\mu\tau)_e$ and $(\mu\tau)_h$ are mobility-lifetime products for electrons and holes, respectively. It is described as a function of the interaction depth (x) from the cathode surface for a specific set of $\mu\tau$ products.

The pulse height from CdZnTe depends on the depth where the interaction takes place. Therefore, when we construct a spectral model, we need to know the “depth distribution”, which indicates how many photons deposit energy at different depths in a detector. For this, we perform a Monte Carlo simulation, which can also incorporate the photo-fluorescence (23 keV for Cd and 27 keV for Te) properly. The obtained distribution is shown in Fig. 5 as a function of the depth from the cathode surface, $\frac{dP(x)}{dx}$. After multiplying it by a certain small thickness dx , it gives the probability that a photon will deposit all of its energy (producing the main photo-peak) within dx at each depth x for an unit photon flux (photons/cm²) incident normal to the top of the detector. In addition to the 122 keV line, we obtain the distribution for the 136 keV line because a considerable amount of its low energy tail impinges on the energy region of 122 keV line. We can see the decrease of full energy deposits at both ends of a detector, which is due to the escape of photo-fluorescence photons from the detector surface.

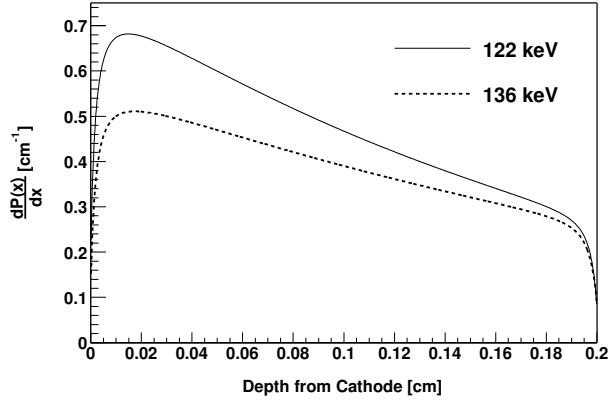


Fig. 5. The “depth distribution” for 122 and 136 keV lines, which gives the differential probability of depositing all of the energy at various depths in a detector for an unit photon flux (photons/cm²). The decreases at both ends of a detector are due to escapes from the cathode and anode surfaces.

The detector response for monochromatic gamma-rays is generated in the following way. We first divide a detector into a number of slices. The relative pulse height is calculated from the Hecht equation at each slice with respect to a certain channel, called CH_{\max} . We fill the model histogram at the calculated pulse height channel with the interaction probability given by the “depth distribution”. This operation is iterated for all the slices from the cathode side to the anode side, which is schematically shown as follows:

$$\begin{aligned} & \sum_{x=\text{cathode}}^{\text{anode}} \text{Fill} (CH_{\max} \times \eta(x), \frac{dP(x)}{dx} \times dx) \\ \implies & f (E_{\gamma}, (\mu\tau)_e, (\mu\tau)_h, E, \sigma, CH_{\max}), \end{aligned} \quad (2)$$

where E_{γ} is the energy of incident gamma-rays and σ is the resolution excluding the tail structure. The resultant spectrum gives the peak shape and is determined by the 6 parameters listed in the above equation. While the product of $(\mu\tau)_e \times E$ mainly determines the peak channel, the product of $(\mu\tau)_h \times E$ determines the amount of the low energy tail. Since the parameter CH_{\max} , the channel where Hecht equation becomes unity, is unknown beforehand and cannot be determined from a single spectrum, we acquire ⁵⁷Co spectra at three different bias voltages (100, 150, 200 V) to utilize the peak shifts due to the changes in the strength of electric field (E). Then, we simultaneously fit the three spectra around the 122 keV main peaks with free parameters of a single set of $\mu\tau$ for electrons and $\mu\tau$ for holes, CH_{\max} , and also 3 sets of σ and normalization (i.e. 9 free parameters in total). This therefore allows us to extract the $\mu\tau$ products from the individual detectors. It is necessary to use the data corrected for electronics gain, offset and nonlinearity in the way described in §3.1. To take into account the 136 keV line component with its low energy tail extending to the energy region of the 122 keV line, the response for the 136 keV line is also calculated using the same parameters for 122 keV except for E_{γ} , and then added to the models.

We apply the model fitting to a single CdZnTe detector. In the model calculation, the number of slices is chosen to be 1000 to calculate the “depth distribution” precisely to the thickness of

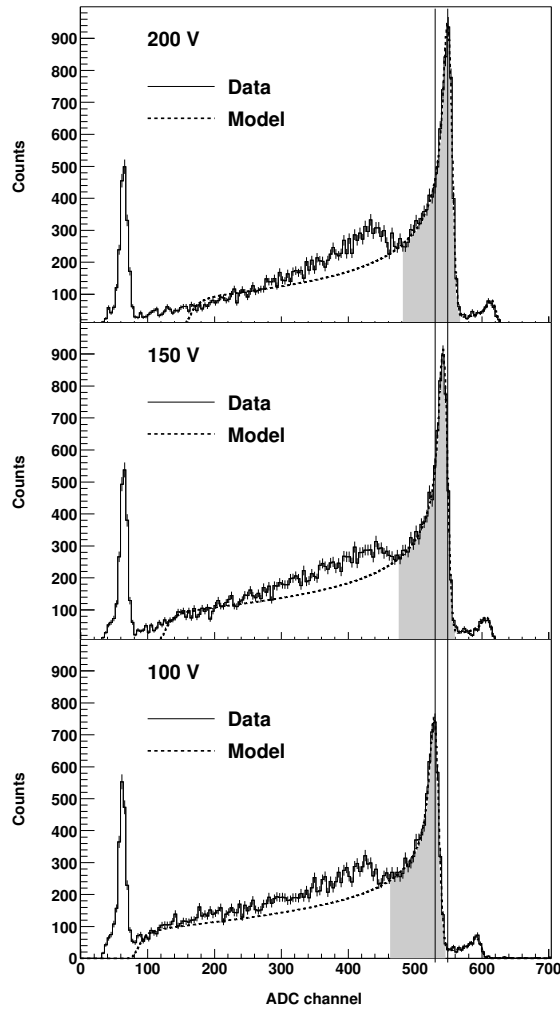


Fig. 6. The ^{57}Co spectra taken at three different bias voltages. Two guide lines are drawn to show how much the 122 keV peak shifts with the bias voltage. The shadowed regions are fitted with the models simultaneously. The $\mu\tau$ for electrons and $\mu\tau$ for holes are determined as $2.57 \pm 0.03 \times 10^{-3}$ and $5.80 \pm 0.33 \times 10^{-5} \text{ cm}^2\text{V}^{-1}$, respectively, with $\chi^2/\nu = 241/244$.

$2 \mu\text{m}$. The region used in the fit is selected above 107 keV to eliminate the escape peaks and to minimize the effect of the scattered component. The region is also restricted to the bins which contain over 10 counts for statistics. This usually excludes the 136 keV peak (but includes its tail component lying under the 122 keV component). For the particular detector shown in Fig. 6, $\mu\tau$ for electrons and $\mu\tau$ for holes are determined as $2.57 \pm 0.03 \times 10^{-3}$ and $5.80 \pm 0.33 \times 10^{-5} \text{ cm}^2\text{V}^{-1}$, respectively, with $\chi^2/\nu = 241/244$. As shown in Fig. 6, it is clear that the shift of the peak channel and the change of the low energy tail with respect to the change of bias voltage can be reproduced by our model.

We implemented our fitting procedure using the ROOT analysis tool kit[10]. The χ^2 fitting is performed by the MINUIT package in ROOT. It takes ~ 2 minutes to fit a set of spectra for each detector with an AMD Athlon(TM) XP 3000+ machine.

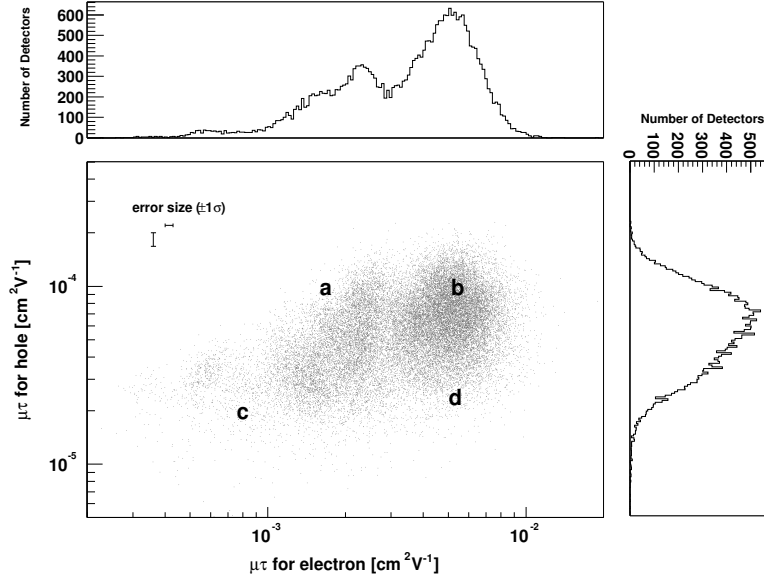


Fig. 7. A cross-plot of $\mu\tau$ for electrons and $\mu\tau$ for holes, which were determined by the spectral fitting method. The two parameters range over more than 1 order of magnitude, showing a slight correlation. The projections are also shown in the top and right panels for $\mu\tau$ for electrons and $\mu\tau$ for holes, respectively.

4.2 Finding $\mu\tau$ products for the 32K CdZnTe detectors

As indicated in §3.1, the diversity in spectral shape is produced by the variation in the charge transport properties between CdZnTe detectors. Here, we evaluate such variations in terms of $\mu\tau$ products by applying the spectral fitting method to all the 32K CdZnTe detectors.

In the individual block calibration, we acquired ^{57}Co spectra at the three different bias voltages (100, 150, 200 V). The temperatures were maintained at 20 °C, though, because of the limited calibration time, two of the blocks had to be operated at 10 °C. In order to minimize the scattered component in the spectra, the measurements were performed in such a way that the radioactive source was mounted inside a box made of lead with a hole cut in it large enough to irradiate all 2048 CdZnTe detectors in the block. A large shielding plate was attached to the box to prevent photons scattering off the room’s ceiling. Then, we applied the fitting method to 122 keV peaks for the 32K CdZnTe detectors. We successfully extracted 29,347 ($\sim 90\%$) sets of $\mu\tau$ products as shown in Fig. 7. The two parameters exhibit remarkably wide distributions, ranging over more than 1 order of magnitude, showing a slight correlation. The $\mu\tau$ for electrons ranges from 5.0×10^{-4} to 1.0×10^{-2} cm^2V^{-1} , and $\mu\tau$ for holes ranges from 1.3×10^{-5} to 1.8×10^{-4} cm^2V^{-1} . Typical errors are 3.5 % (1σ) and 8.7 % (1σ) for the $\mu\tau$ for electrons and $\mu\tau$ for holes, respectively.

Examples of the extreme differences in the spectra due to the different $\mu\tau$ products are shown in Fig. 8, where the symbols “a” to “d” correspond to the location in Fig. 7. Though the selected detectors are exposed to the ^{57}Co source for the same duration, there are large differences in spectral shapes. Detectors (a) and (b) have larger values of the $\mu\tau$ for holes than detectors (c) and (d), and their spectra show higher 122 keV peaks than those of detectors (c) and (d). In

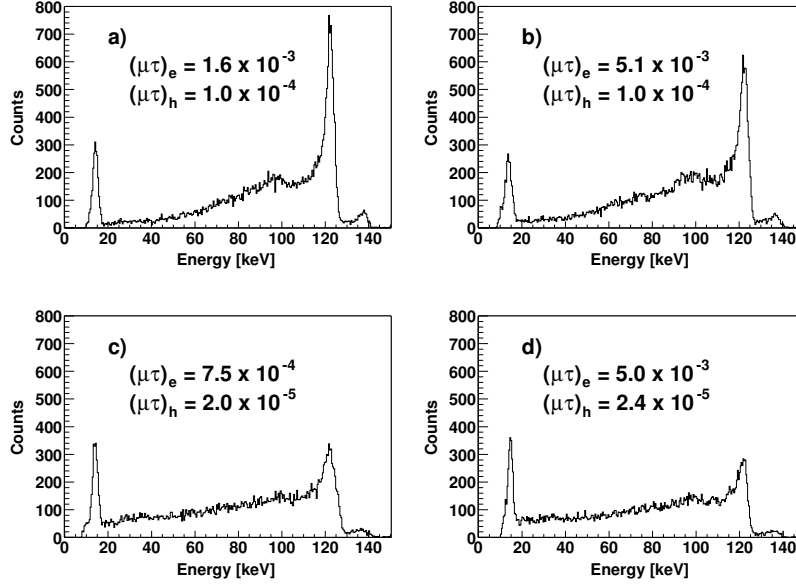


Fig. 8. Spectra obtained with single CdZnTe detectors with various sets of $\mu\tau$ products, corresponding to the symbols in the $\mu\tau$ products cross-plot (Fig.7). The detectors are exposed to a ^{57}Co source for the same duration.

contrast, there are more counts in the lower energy tail in the spectra (c) and (d). The more interesting feature is that spectrum (a) with a smaller value of $\mu\tau$ for electrons exhibits a higher 122 keV peak than spectrum (b). Likewise, spectrum (c) has a higher peak than spectrum (d). To investigate this relation between the $\mu\tau$ products and the spectral shapes, we use the measured peak-to-valley ratios at 122 keV, which are calculated in section 3.1. We first subdivide the $\mu\tau$ products cross-plot into grids with equal log-scale intervals. Then, the peak-to-valley ratios are averaged over the detectors in each grid and shown in grey-scale in Fig. 9. First, it is evident that the peak-to-valley ratio increases as the $\mu\tau$ for holes becomes higher. Second, the detectors with largest values of peak-to-valley ratio are found in top and left edge of the $\mu\tau$ distributions. This trend can be recognized as a result of the closer values between $\mu\tau$ for electrons and $\mu\tau$ for holes. In such a detector, there will be a smaller difference in charge induction efficiency throughout the detector, and as a result, a higher 122 keV photo-peak can be formed in the spectrum.

Furthermore, we find that the extracted $\mu\tau$ products are somewhat correlated to the DM or block location [11], which turns out to be strongly correlated to the detector boule number. In Fig. 10, the $\mu\tau$ products for $\sim 8,000$ detectors are shown against detector number sorted by ingot ID. It is realized that the diversities of $\mu\tau$ products can be attributed to the difference of crystal ingots or manufacturing harness, though there remains a large variation in the $\mu\tau$ for holes even in the same ingots.

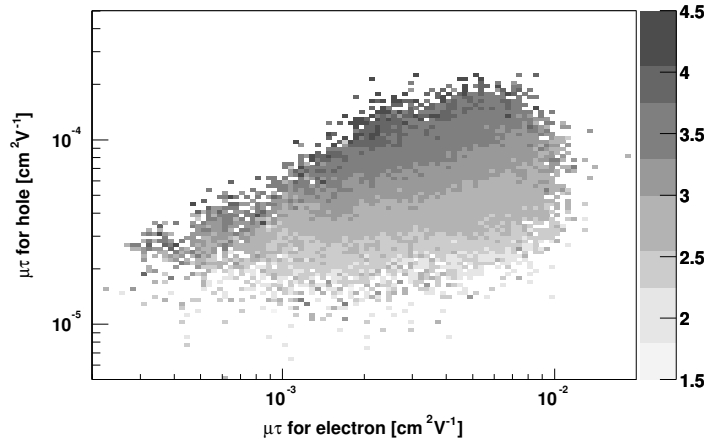


Fig. 9. The $\mu\tau$ products cross-plot divided in grids with equal intervals in log-scale, showing the measured peak-to-valley ratio averaged over the detectors in each grid in grey-scale. As the $\mu\tau$ hole becomes higher, the peak-to-valley ratio increases. The detectors with largest values of peak-to-valley ratio are found in top and left edge of the $\mu\tau$ distributions.

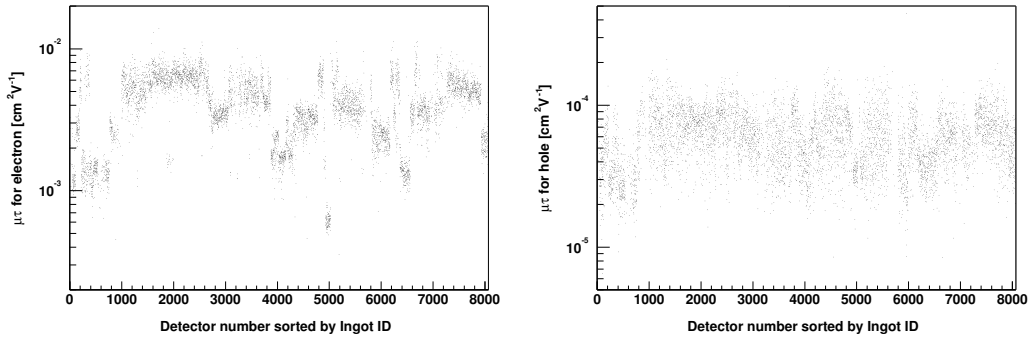


Fig. 10. The extracted $\mu\tau$ products against detector number sorted by ingot ID, indicating a strong correlation between the $\mu\tau$ products and the crystal ingots or manufacturing harness.

4.3 Temperature Dependence of $\mu\tau$ products

The carrier mobilities are intrinsically dependent on temperature, limited by phonon scattering at higher temperatures and by impurity scattering at lower temperatures. This dependence in HPB CdZnTe has been reported in several papers [12]. However, since it is highly dependent on material qualities, it is not clear how the carriers in BAT CdZnTe detectors behave with temperature. Thus, for one specific block (2,048 CdZnTe detectors), we take the same fitting procedure to extract $\mu\tau$ products from ^{57}Co spectra acquired at lower temperatures of 0 and 10 °C in addition to the data at 20 °C.

The distributions of extracted $\mu\tau$ products are shown in Fig. 11. There are trends in both distributions with the $\mu\tau$ product shifting slightly downward with a decrease in temperature, but the overall shape is not affected very much. For individual detectors, the relative values of $\mu\tau$ products are calculated with respect to the values at 20 °C. Since we observe a large variation

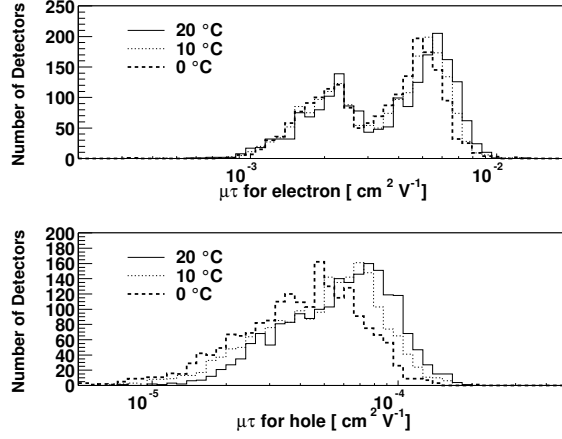


Fig. 11. The distributions of $\mu\tau$ for electrons (top) and $\mu\tau$ for holes (bottom) at three different temperatures (20, 10, 0 °C), showing trends for both distributions to shift slightly downward with decrease of temperature, but not to change the overall shape very much.

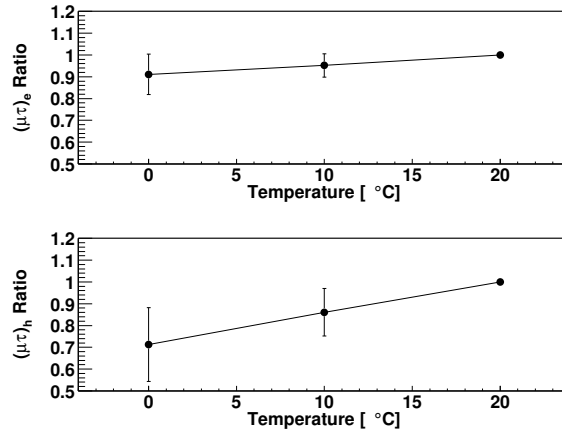


Fig. 12. For individual detectors, the relative values of $\mu\tau$ products are calculated with respect to the values at 20 °C. The average values and the RMS variations between detectors are shown for the $\mu\tau$ for electrons (top) and $\mu\tau$ for holes (bottom).

from detector to detector in the way that $\mu\tau$ products change from 20 to 10, and 0 °C, we show the averaged values and the RMS variations between detectors in Fig. 12 after subtraction of the deviation due to $\mu\tau$ extraction errors. We find that the $\mu\tau$ for electrons decreases by $\sim 9\%$ while the $\mu\tau$ for holes decrease by $\sim 29\%$ from 20 to 0 °C on average.

5 Development of the Spectral Model

Our final goal is to develop an energy and angular response for the entire CdZnTe array from 15 to 150 keV and from 0 to 60 ° for a large FOV of the BAT instrument. Using the model calculations described in the previous section, it is possible to calculate the pulse heights for any incident photon energy and to obtain the “depth distribution” for any incident angle with

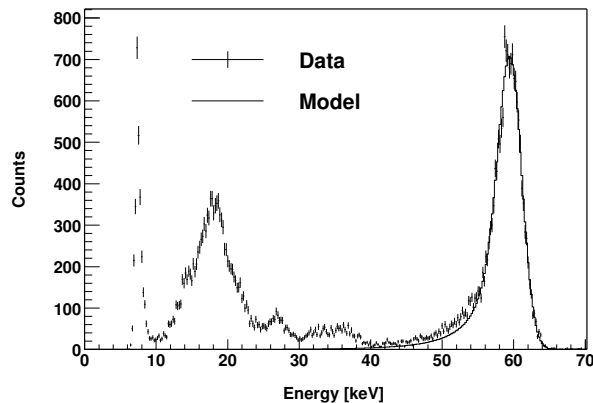


Fig. 13. An ^{241}Am spectrum obtained with a single CdZnTe detector, which is fitted with the model function calculated for 59.5 keV using the values of μT products derived from the spectral fitting of 122 keV line.

the use of Monte Carlo simulations.

5.1 Energy and Angular Response for a Single CdZnTe Spectrum

5.1.1 Response for the 59.5 keV Line

We first calculate a model function for 59.5 keV to be applied to an ^{241}Am spectrum obtained with a specific CdZnTe detector as shown in Fig. 13. In the 6 parameters of the model function (eq. 2), the μT products are set to the values determined through the spectral fitting of the 122 keV line while the σ and CH_{max} are allowed to vary to adjust the resolution and the gain. We can see the model reproduces the asymmetric shape of photo-peak very well from ~ 55 to ~ 64 keV. However, there remains an excess over the tail structure below 55 keV, which is observed in most of the detectors. After a more detailed study, we determined that this is most likely due to the combination of Compton scattering from a plastic package surrounding the radioactive source and an illumination onto the side surface of a detector, which impinges on deep regions below the cathode surface where the charge induction efficiency is relatively low. Another possibility is a non-uniformity in the μT products within a detector. Suzuki et al. [13] have reported that there exist variations in μT products on scales less than 1 mm. Here it should also be noted that the lower part of the tail component is somewhat uncertain because the μT products are extracted with the fitting of the restricted energy region around the main peaks (see Fig. 6).

5.1.2 Response for Low Energy Lines

The response for lower energy lines are investigated with a ^{109}Cd spectrum. We calculate model functions for 22 and 25 keV lines to be almost Gaussians because most of the interactions occur in shallow regions below the cathode surface. However, we observe a low energy tail component

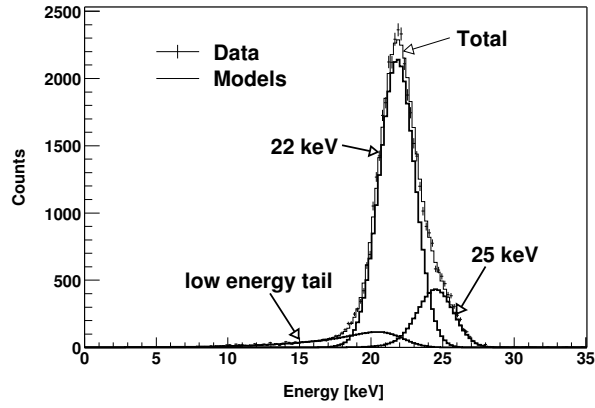


Fig. 14. A ^{109}Cd spectrum obtained with a single CdZnTe detector, which is fitted well with the model composed of 22, 25 keV and an additional low energy tail component.

below 20 keV as shown in Fig. 14. To evaluate the low energy tail, we introduce an empirical formula:

$$f(x) = \exp\left(\frac{x - x_{\text{peak}}}{\lambda}\right) \quad \text{if } x < x_{\text{peak}}, \quad (3)$$

where x and x_{peak} are the energy and peak channel, and λ is a parameter to determine how the tail decreases with separation from the main peak. The component is added to the model after smearing with the energy resolution, which is fitted well with the actual data (Fig. 14). It is partly explained with side illumination onto the side surface of the detector and with Compton scattering from the plastic tablet of the radioactive source. However, if the incident photon energy is higher than the K-edges of Cd and Te (26.7 and 31.8 keV), such as the 31 and 35 keV lines from a ^{133}Ba source, the situation becomes more complicated because escape photons from adjacent detectors contaminate the energy region below the main peaks.

5.1.3 Angular Response

The “depth distribution” depends on the incident angle. Fig. 15 (*left*) demonstrates how the distribution changes from an incident angle of 0 to 30 and 60°. The incident angle affects the energy response in two main ways: one is a decrease of the effective area, and the other is a decrease of the interaction depth. This changes the shape and amount of the low energy tail as presented in Fig. 15 (*right*).

5.1.4 Modeling the Escapes

Another important feature in CdZnTe spectra is “escape peaks”, which occur as a result of photo-fluorescence. The prominent peaks are of Cd-K and Te-K escapes, which appear at 23 and 27 keV below main photo-peak, respectively. To model the escape peaks, we count the events where the energy depositions are smaller by 23 or 27 keV than the full incident photon

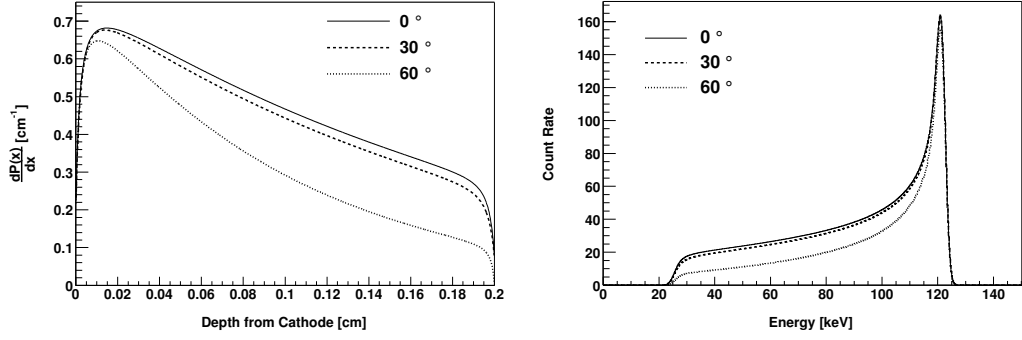


Fig. 15. The “depth distributions” for a stream of 122 keV photons incident on a detector at three different angles (0, 30 and 60 °) (*left*), the difference of which results in change of amount of tail and total area of the model spectrum (*right*).

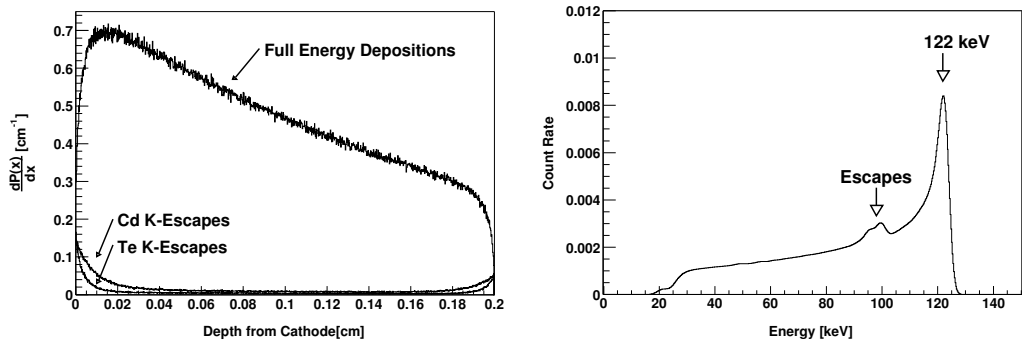


Fig. 16. The “depth distributions” for K-escapes for a stream of 122 keV photons (*left*), which enables us to calculate a model spectrum including escape peaks (*right*).

energy and derive the “depth distributions” for escape peaks, which are shown in Fig. 16 (*left*). Then, by applying the charge induction efficiency such as with the main photo-peak, we can obtain the model of escape peaks as shown in Fig. 16 (*right*).

5.2 Composite Model for the Entire CdZnTe Array

To develop the spectral model for the entire CdZnTe array, it is necessary to take into account the large range of $\mu\tau$ products. For this, we first divide the 32K detectors into 9-by-7 groups with equal log-scale intervals. The percentage of detectors included in each grid is shown in Fig. 17. Then, we construct spectral models with $\mu\tau$ products at the center of each grid. The models are calculated for the most populated groups with the criteria of $> 0.1\%$. Thus, there are 35 spectral models based on the model functions developed for single CdZnTe detectors. Finally, we sum them, using the fraction of detectors as a weight. This gives us a composite spectral model which represents the energy response for the entire CdZnTe array.

The resultant spectral model is shown in Fig. 18 with an actual calibration spectrum resulting from a ^{57}Co source positioned slightly off-axis. Although there still remains the discrepancy in

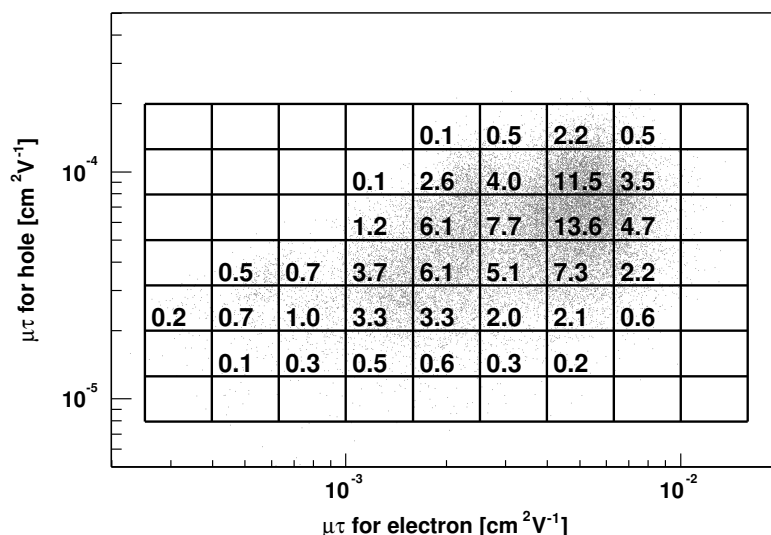


Fig. 17. The $\mu\tau$ products cross-plot divided into 9-by-7 groups with equal intervals in log-scale, showing the numbers of detectors included in each grid as percentages of the total.

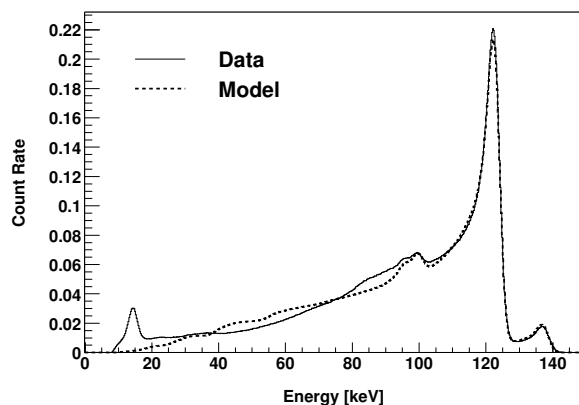


Fig. 18. A ^{57}Co spectrum summed over the entire CdZnTe array, which is fitted with the composite spectral model for 122 and 136 keV lines.

the lower part of the tail, the model represents the overall spectral shape very well. To demonstrate how the model can reproduce the summed spectrum over a wide energy range, we construct a spectral model for ^{133}Ba spectrum, which includes 31, 35, 53 and 81 keV lines. The result is shown in Fig. 19, which looks almost perfect. The model reproduces the tail component given by the $\mu\tau$ products of electrons and holes, the tail component for low energy lines, and also escape peaks. Note that the values of the $\mu\tau$ products of electrons and holes are fixed at those values obtained in §4.

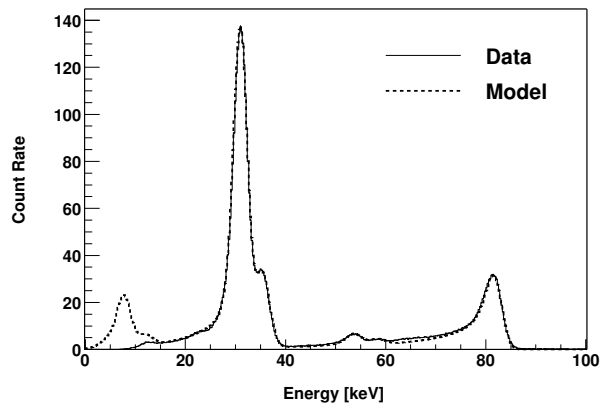


Fig. 19. A ^{133}Ba summed spectrum fitted with the composite spectral model, showing a good agreement in the wide energy range from 15 to 85 keV. lines.

6 Conclusion

Based on the fitting method applied to ^{57}Co spectra, we have succeeded to extract $\mu\tau$ products both for electrons and holes from 32K individual CdZnTe detectors. The reason for this procedure is to calculate a spectral model following the actual physical processes from the photon-detector interaction to the charge transport. Once the $\mu\tau$ products are obtained, it allows us to predict the energy response for any photon energy or any incident angle. The actual calibration data is reproduced well by the model calculated with the Hecht equation and the interaction depth distributions for each energy. However, since various structures can be produced by scattered components and other geometrical effects, it is also necessary to estimate carefully such contributions with the use of Monte Carlo simulations in each case. Another important point is to deal with the large distributions of 32K sets of $\mu\tau$ products. By grouping the detectors, we have successfully developed the spectral model for the entire CdZnTe array, which is to be included in the software of the Detector Response Matrix (DRM) generator for the BAT instrument.

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