WEDGE HOLOGRAPHY AS GENERALIZATION OF ADS/CFT

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Based on [arXiv:2007.06800]
collaboration with
Ibrahim Akal, Tadashi Takayanagi, Zixia Wei

Summary①

[Akal, YK, Takayanagi, Wei]

Generalization

Original Holography
$AdS_{d+1} = CFT_d$

→

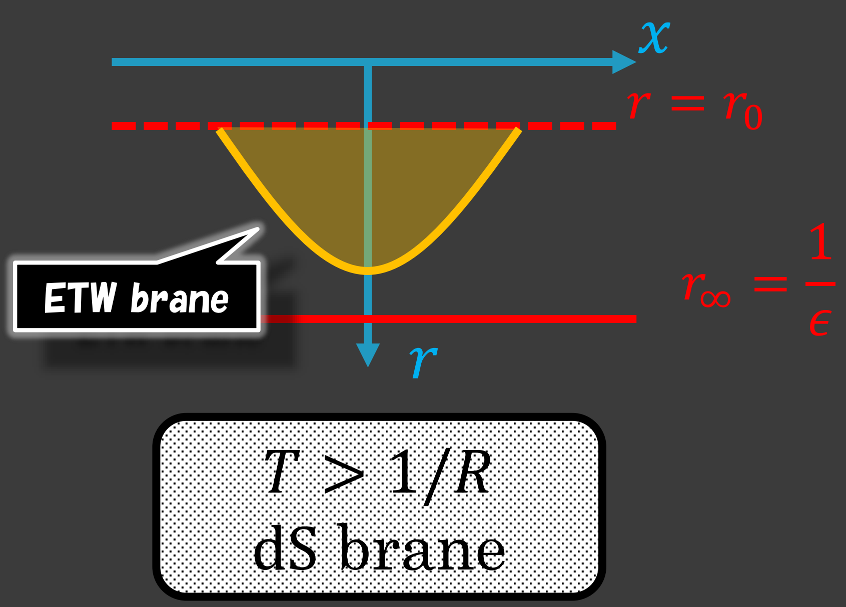
Our Wedge Holography
$AdS_{d+1} = CFT_{d-1}$

Motivation

- New laboratory
- New construction of CFT_1
- Construction is more definite than doubly holography
- Joining gravity from AdS_{d+1} and CFT_{d-1}
  ⇒radiation setup

Summary②

[Akal, YK, Takayanagi, Wei]



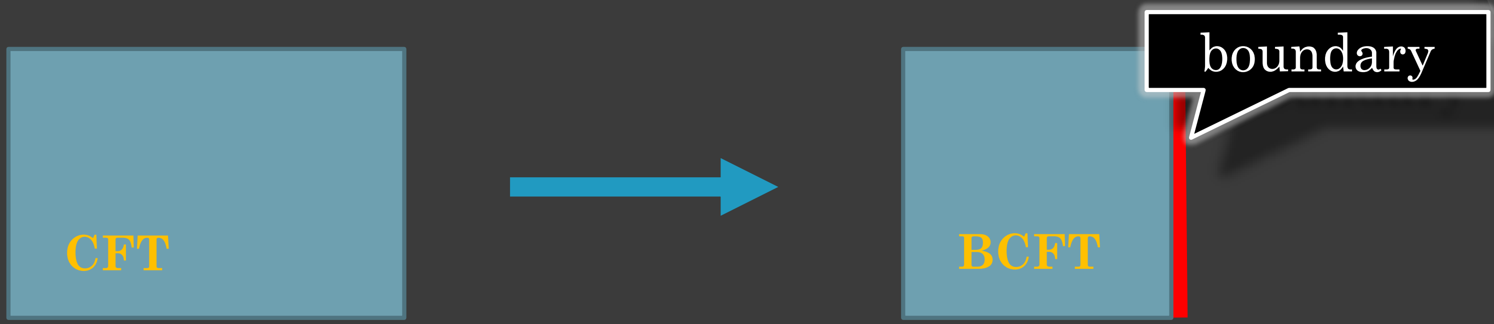
No asymptotic boundary
⇒ No CFT dual?

Motivation

- New laboratory
- dS braneworld
- Universe creation setup

Contents

Boundary CFT



Boundary breaks conformal symmetry…

Conformal Symmetry of BCFT is part of conformal symmetry preserving the bdy. position.

$$SO(2,d) \rightarrow SO(2,d-1)$$

Boundary CFT

Definite definition of $BCFT_2$ is given by [Cardy].

Difference I: **Symmetry**

Symmetry of **CFT** is

$$Vir \times \overline{Vir}$$

condition to preserve bdy.

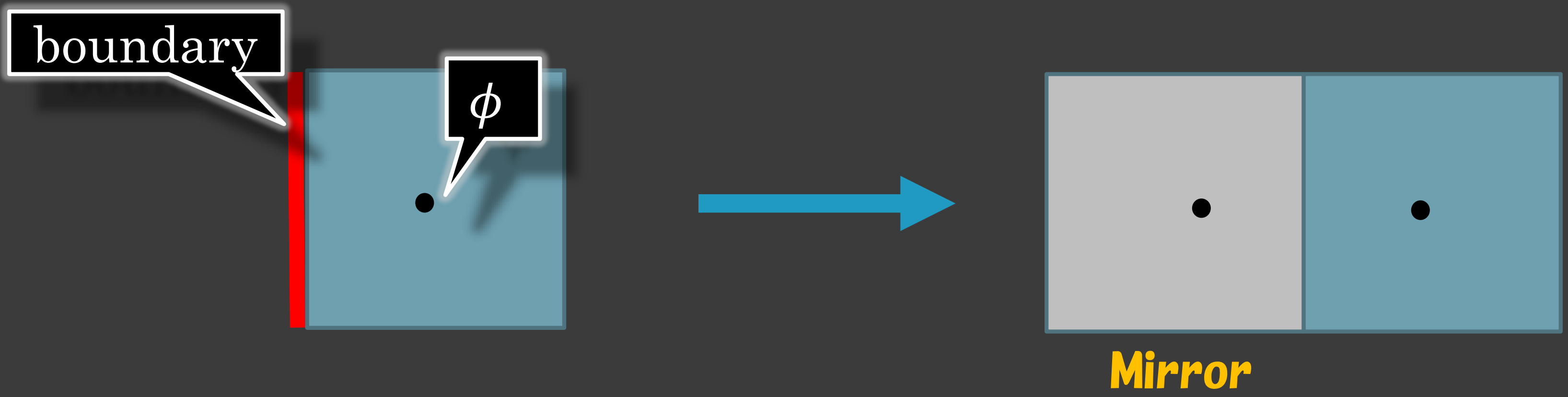
On the other hand, bdy. condition of **BCFT**,

$$T_{xy} = 0 \quad \Leftrightarrow \quad L_n - \bar{L}_{-n} = 0$$

shows the dependency of Vir and $\overline{Vir}$.
Instead, symmetry of **BCFT** is given by

$$Vir$$

Boundary CFT

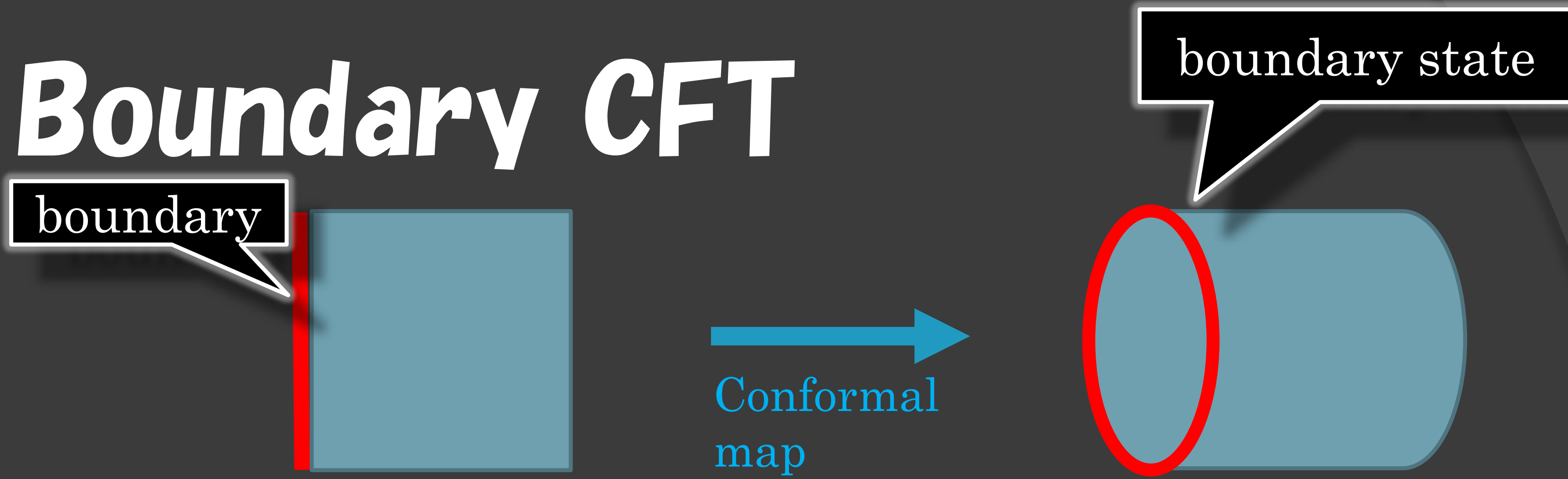


Mirror method:

$$\langle \phi(z) \rangle_{bdy} = \langle \phi(z) \phi(z^*) \rangle \propto \frac{1}{|z - z^*|^{2h_\phi}}$$

with $z = x + iy$.
The second eq. comes from standard conformal Ward id.

Boundary CFT



Difference III: **Basic contents**

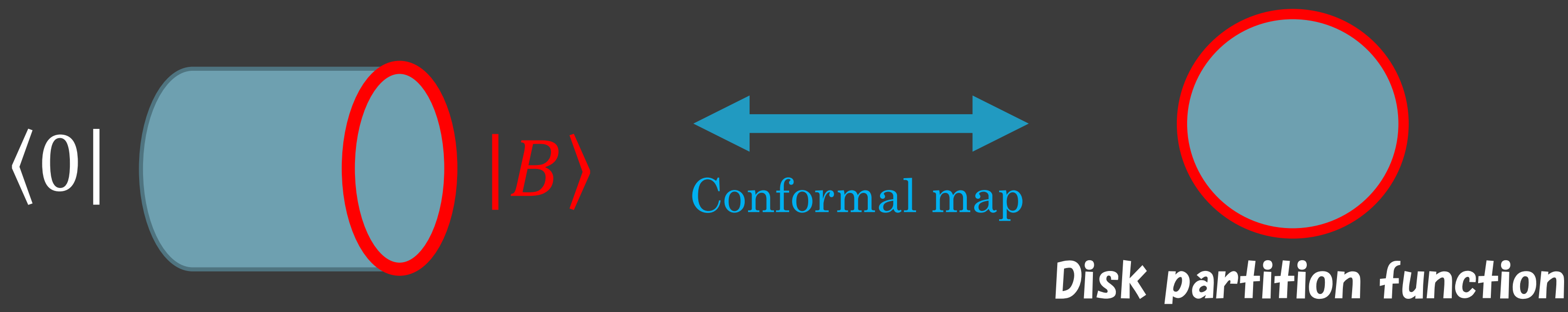
CFT is specified by central charge, primary state, OPE coefficient

BCFT has extra contents, boundary state, OPE coefficient among bulk states and boundary states

With all of them, we can explicitly calculate any correlators

Boundary entropy [Affleck, Ludwig]

DoF of boundary is measured by **boundary entropy**.

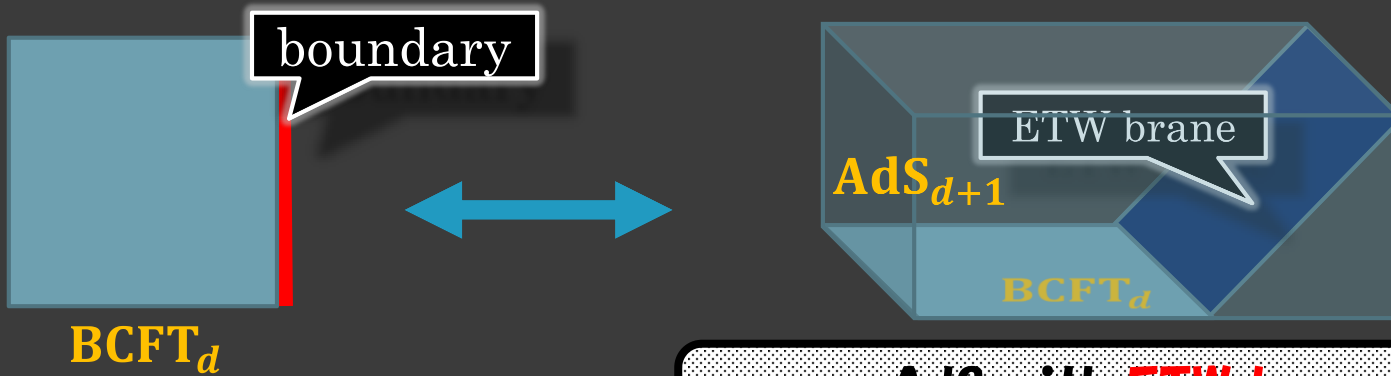


$$S_{bdy} = \log g_B, \quad g_B = \langle 0|B\rangle \simeq Z_{disk}$$

Interesting point:

- ***g*-theorem** [Friedan, Konechny]
  S_{bdy} decreases under the RG flow like central charge
  ⇒consistent with DoF measure interpretation

Gravity dual of BCFT [Fujita, Tonni, Takayanagi]



BCFT

AdS with ETW brane

$\partial(\mathrm{ETW}) = \mathrm{bdy.\,of\,CFT}$

Induced metric: $h_{\mu\nu} = g_{\mu\nu} - n_\mu n_\nu$,
Extrinsic curvature: $K_{\mu\nu} = h_\mu{}^\rho h_\nu{}^\lambda \nabla_\rho n_\lambda$

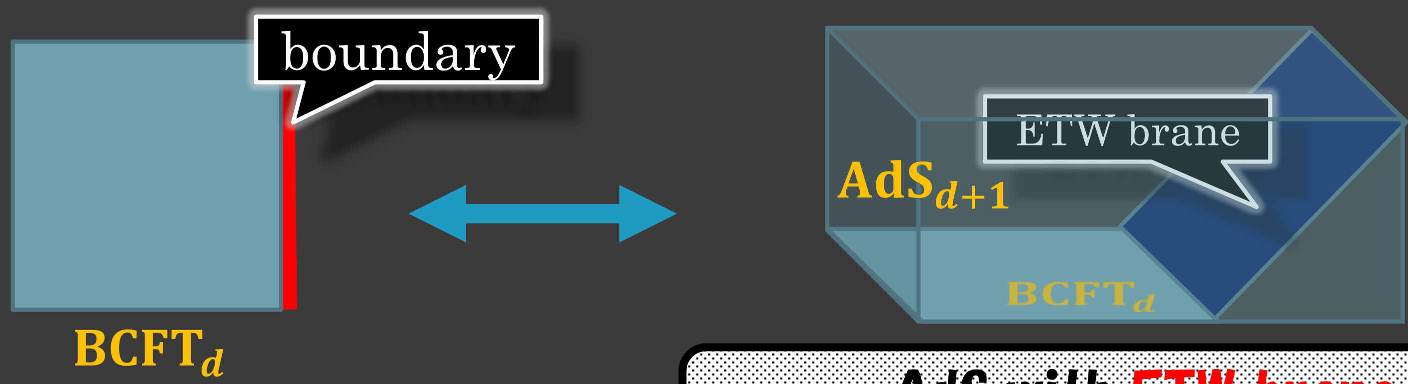
Gravity action:

$$I = -\frac{1}{16\pi G_N}\int_N \sqrt{g}(R - 2\Lambda) - \frac{1}{8\pi G_N}\int_Q \sqrt{h}(K - T)$$

Neumann b.c. is imposed on the brane (Einstein eq. of brane).

$$K_{ab} - K h_{ab} = -T h_{ab}$$

Gravity dual of BCFT [Fujita, Tonni, Takayanagi]



BCFT

AdS with ETW brane
$\partial(\mathrm{ETW}) = \mathrm{bdy.\,of\,CFT}$

Interesting point of AdS/BCFT:

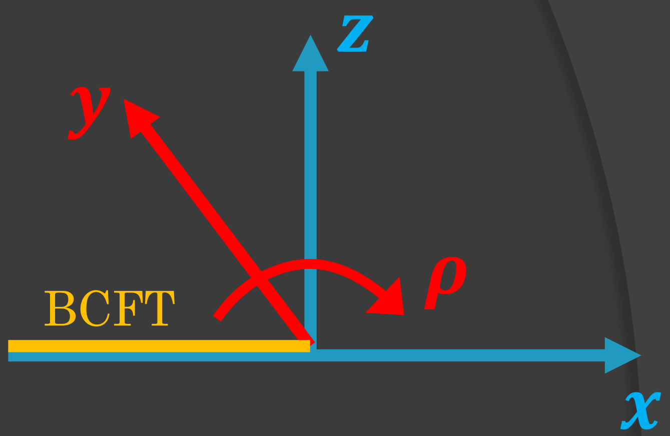
Boundary provides **dynamical** spacetime on ETW (see also [Randull-Sandrum])
This is the reason why resent progresses about Page curve utilize AdS/BCFT

Gravity dual of BCFT

Simple example:

Poincare patch:

$$ds^2 = R^2\left(\frac{dz^2 - dt^2 + dx^2 + d\vec{w}^2}{z^2}\right)$$



Let us consider a BCFT on a UHP (x < 0).

It is useful to use the coordinate $\left(x = \frac{y}{\cosh\frac{\rho}{R}}, z = \frac{y}{\cosh\frac{\rho}{R}}\right)$,

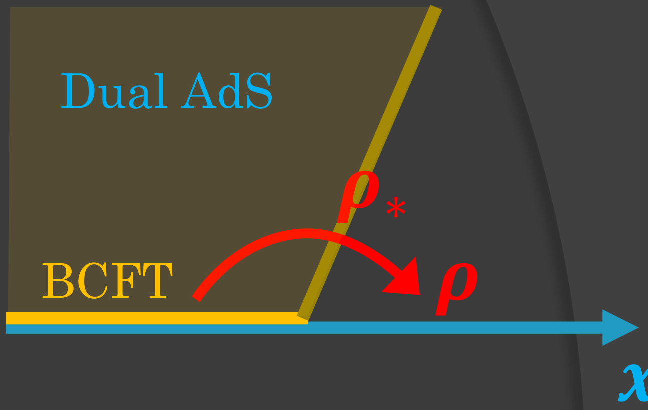
$$ds^2 = d\rho^2 + R^2\cosh^2\frac{\rho}{R}\left(\frac{dy^2 - dt^2 + d\vec{w}^2}{y^2}\right)$$

The bdy. position $\rho = \rho(y)$ is fixed by the Neumann b.c.

$$K_{ab} - Kh_{ab} = -Th_{ab}$$

Gravity dual of BCFT

$$K_{ab} - Kh_{ab} = -Th_{ab}$$



A solution to the Neumann b.c. is

$$\rho = \rho_* \qquad s.t. \quad T = \frac{d-1}{R} \tanh \frac{\rho_*}{R}$$

By an appropriate map from this UHP, the gravity dual of a disk partition function (= **boundary entropy**) is obtained.

$$S_{bdy} = -I_{disk} = \frac{\rho_*}{4G_N}$$

S_{bdy} becomes large as the size of the space is increased (i.e. ρ_* is increased)

⇒boundary becomes **classical** (explained later)

Ryu-Takayanagi formula

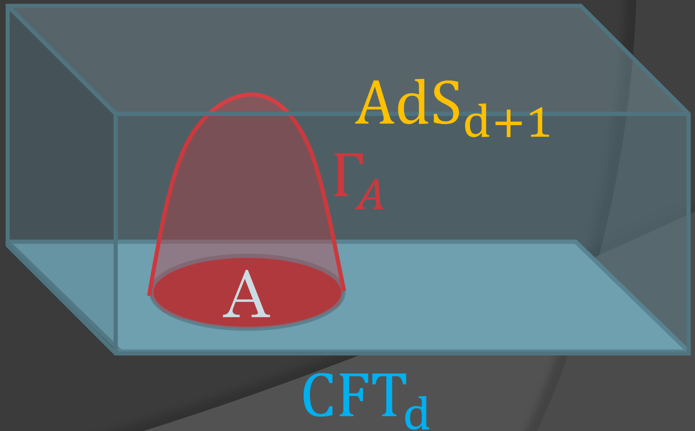
[Ryu, Takayanagi]

EE can probe how much information about $\bar{A}$ can be extracted from subregion A in CFT.

$$S_A = \mathrm{tr}\,\rho_A \log \rho_A\,, \qquad \rho_A = \mathrm{tr}_{\bar{A}}\,\rho$$

Gravity dual of EE is

$$S_A = \min_{\substack{\Gamma_A \\ \partial\Gamma_A = \partial A}} \left(\frac{\mathrm{Area}(\Gamma_A)}{4G_N} \right)$$

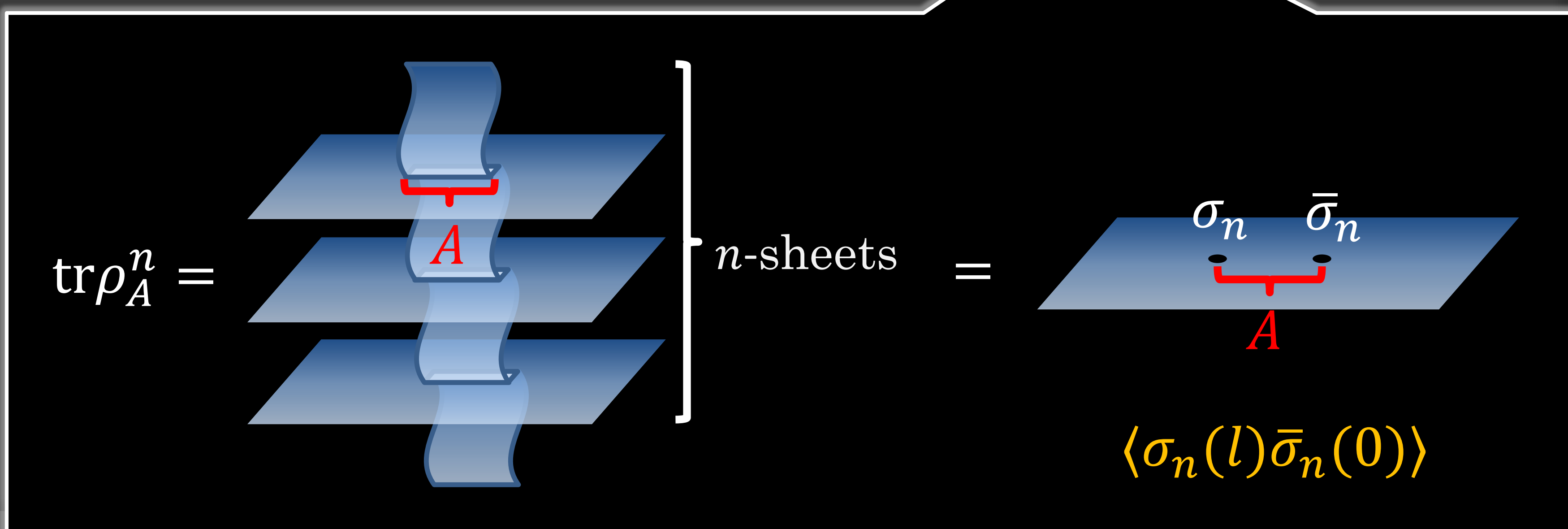


⇒ Gravity calculation is easy

Ryu-Takayanagi formula

CFT calculation is also not so hard (replica trick).

$$S_A = \mathrm{tr}\,\rho_A \log \rho_A = \lim_{n\to 1}\frac{1}{1-n}\log \mathrm{tr}\rho_A^n, \qquad \rho_A = \mathrm{tr}_{\bar{A}}\,\rho$$



Ryu-Takayanagi formula

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$$\mathrm{tr}\rho_A^n = \langle \sigma_n(l)\bar{\sigma}_n(0)\rangle = \frac{1}{l^{2h_{\sigma_n}}}$$

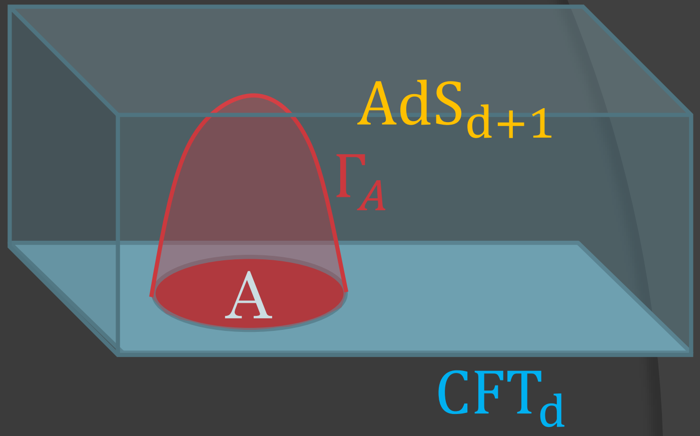
where $h_{\sigma_n} = \frac{c}{24}\left(n - \frac{1}{n}\right)$. The second eq. is just Ward id.

As a result,

$$S_A = \frac{c}{3} \log \frac{l}{\epsilon}$$

Ryu-Takayanagi formula

$$S_A = \min_{\substack{\Gamma_A \\ \partial\Gamma_A = \partial A}} \left(\frac{\text{Area}(\Gamma_A)}{4G_N} \right)$$



Usefulness

- Both gravity and CFT calculation are ***easy***
  ⇒ For this reason, we will use this to justify our new holography.
- Useful to probe global feature
  - Which is part of gravity dual to subregion A in CFT?
    Sub-region/sub-region duality (see [Suzuki, YK, Takayanagi, Umemoto])
  - c-theorem [Myers-Sinha]

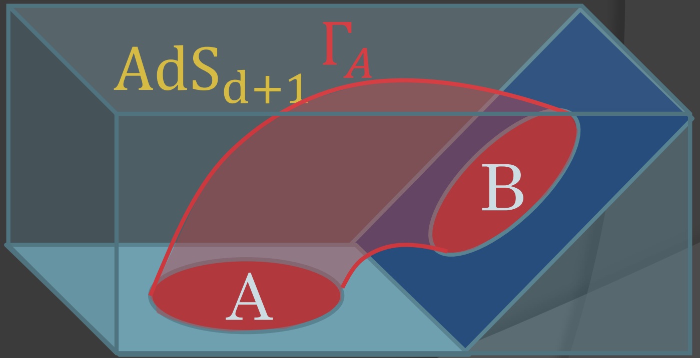
⇒ How is this generalized to AdS/BCFT?

Ryu-Takayanagi formula

Difference

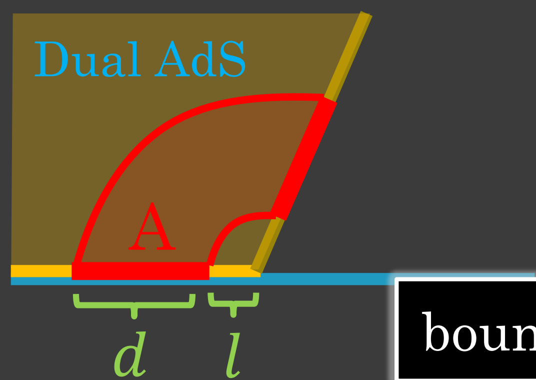
- RT surface can ***end on ETW brane*** (which satisfies homologous cond. from string perspective.)
- DoF of boundary (boundary entropy) also contributes to S_A.

$$S_A = \min_{\substack{\Gamma_A \\ \partial\Gamma_A = \partial A \cup \partial B}} \left(\min_B \frac{\text{Area}(\Gamma_A)}{4G_N} \right)$$

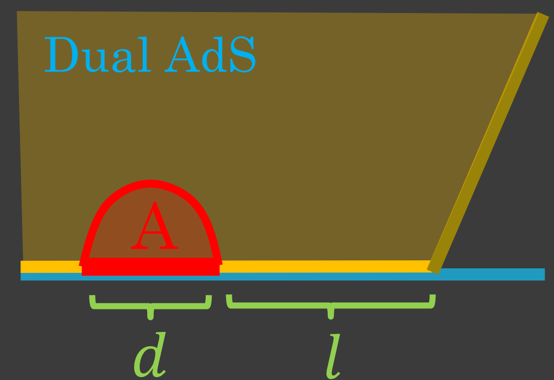


Ryu-Takayanagi formula

Simple example:



transition



$$S_A = \frac{c}{6}\log\frac{l}{\epsilon} + \frac{c}{6}\log\frac{l+d}{\epsilon} + 2\log g_B$$

boundary entropy

$$S_A = \frac{c}{3}\log\frac{d}{\epsilon}$$

This entanglement entropy can be completely reproduced by **CFT** calculation (replica trick and ***mirror method***)

Contents

Developments from AdS/BCFT

[Akal, YK, Takayanagi, Wei]

Original Holography

$d+1$ **dimensional AdS = d dimensional CFT**

↓ *Generalization*

Our Wedge Holography

$d+1$ **dimensional AdS = $d-1$ dimensional CFT**

Developments from AdS/BCFT

[Akal, YK, Takayanagi, Wei]

Original Holography

$d+1$ dimensional AdS = d dimensional CFT

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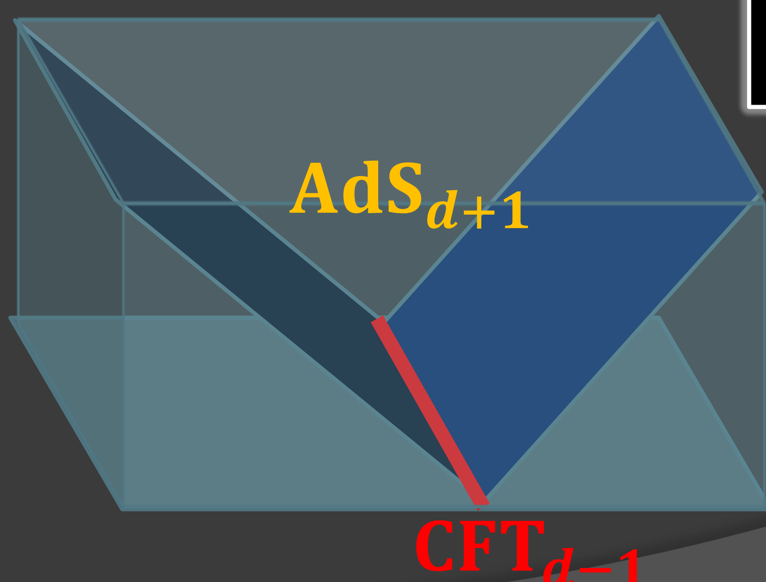
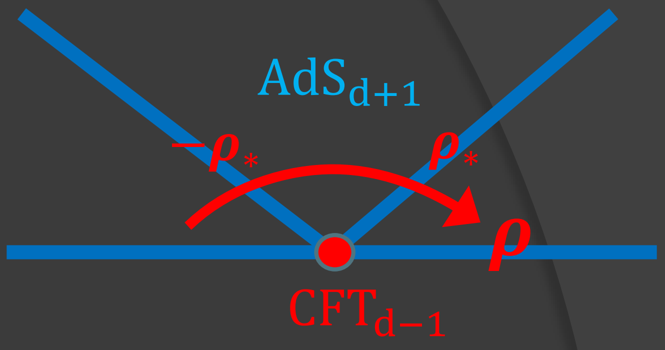
Our Wedge Holography

$d+1$ **dimensional AdS = $d-1$ dimensional CFT**



Wedge Holography



Definition of wedge:

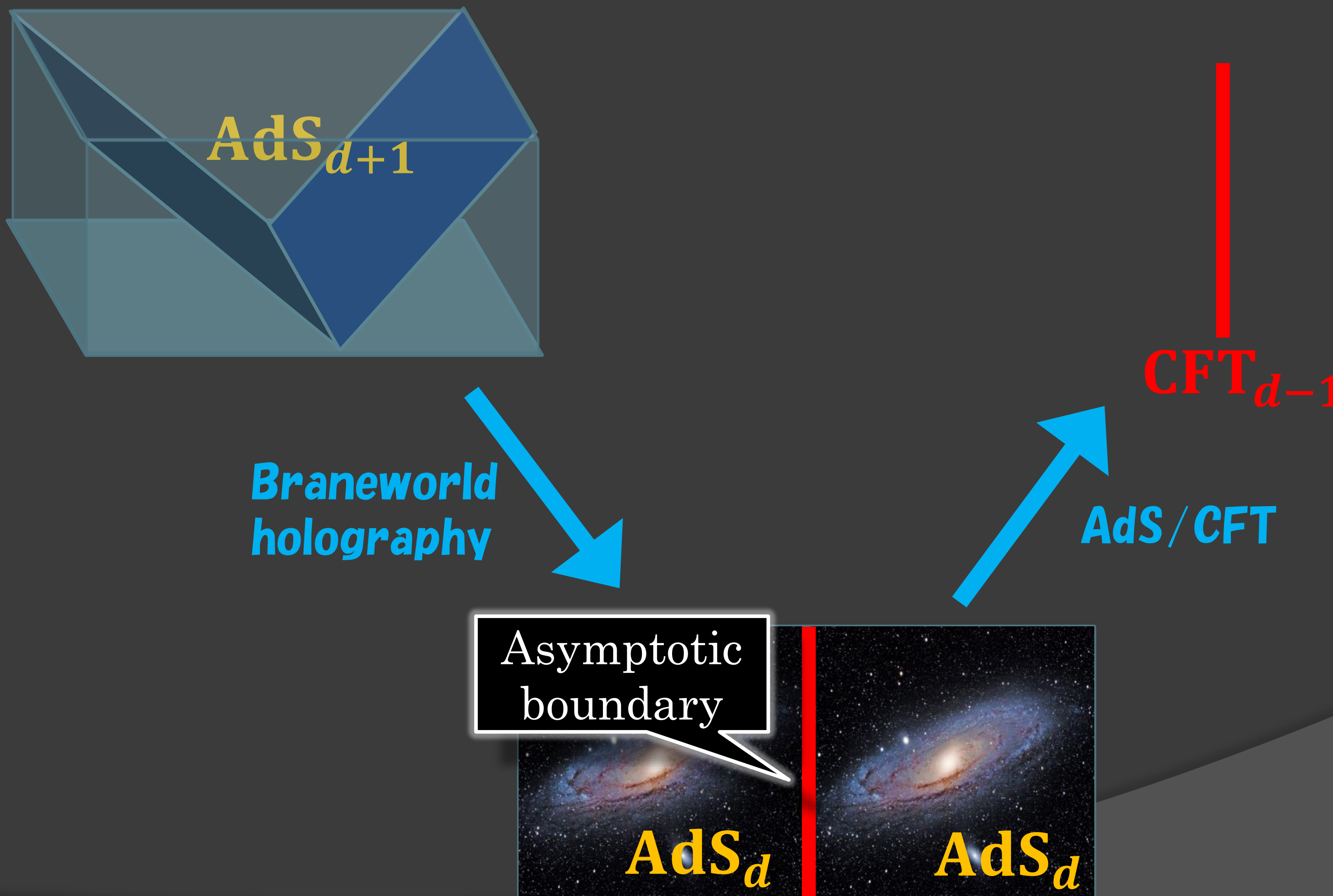
$$ds^2_{AdS_{d+1}} = d\rho^2 + R^2\cosh^2\frac{\rho}{R}\left(\frac{dy^2 - dt^2 + d\vec{w}^2}{y^2}\right)$$
$$= d\rho^2 + \cosh^2\frac{\rho}{R}\, ds^2_{AdS_d}$$

Restricted to **wedge** subspace $-\rho_* \leq \rho \leq \rho_*$

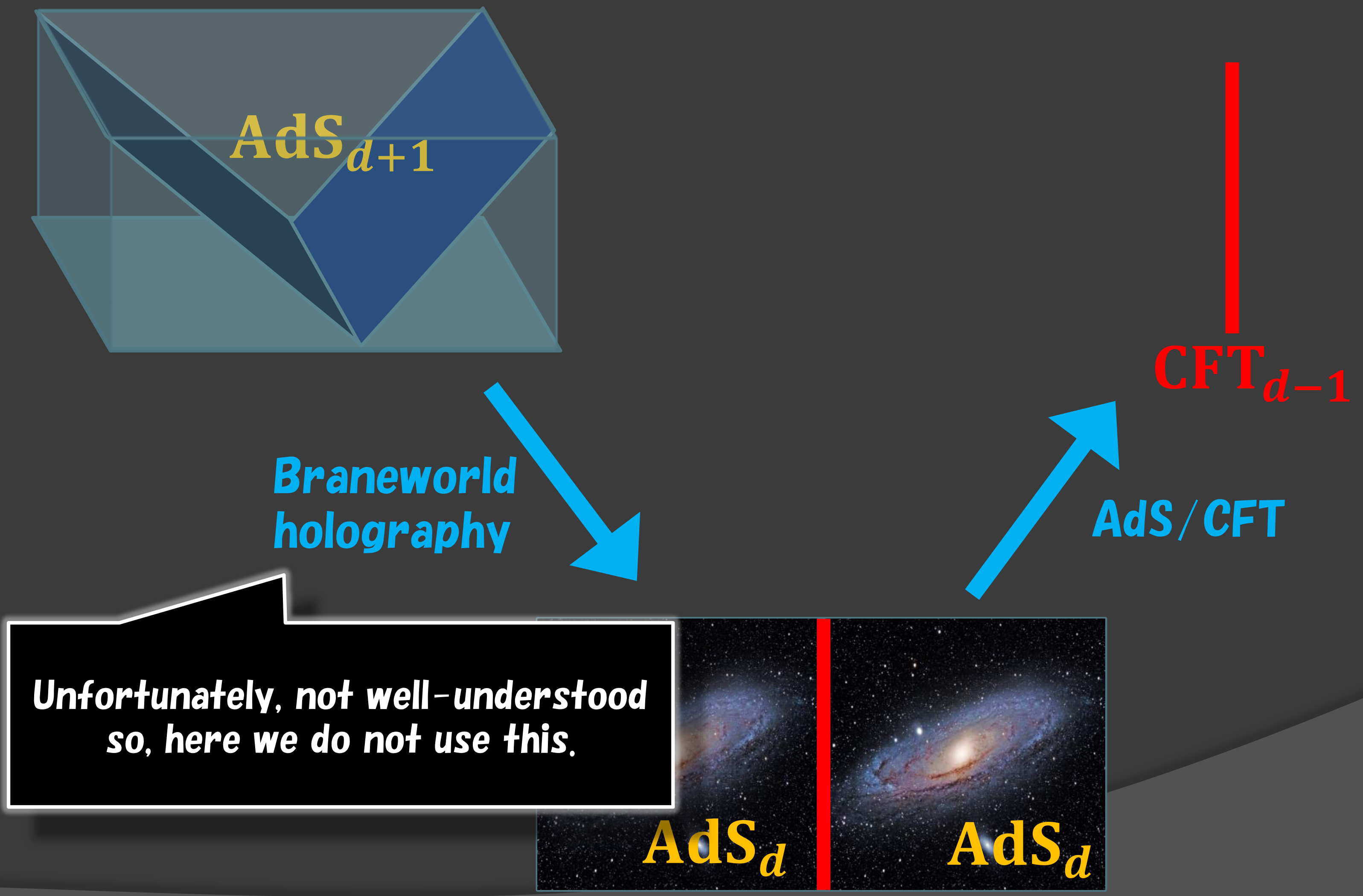
⇒asymptotic boundary is $d-1$ dimensional theory

Our proposal: This is CFT$_{d-1}$

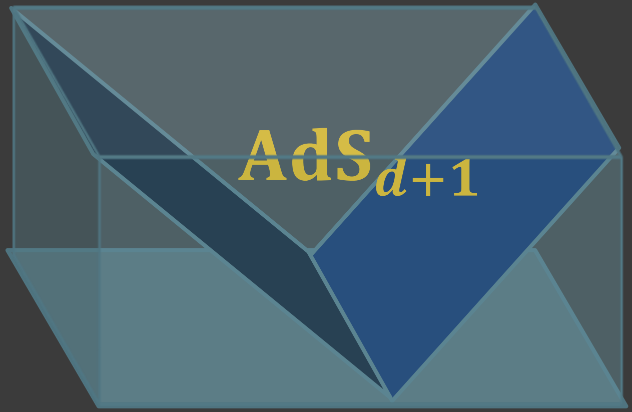
Naive derivation?

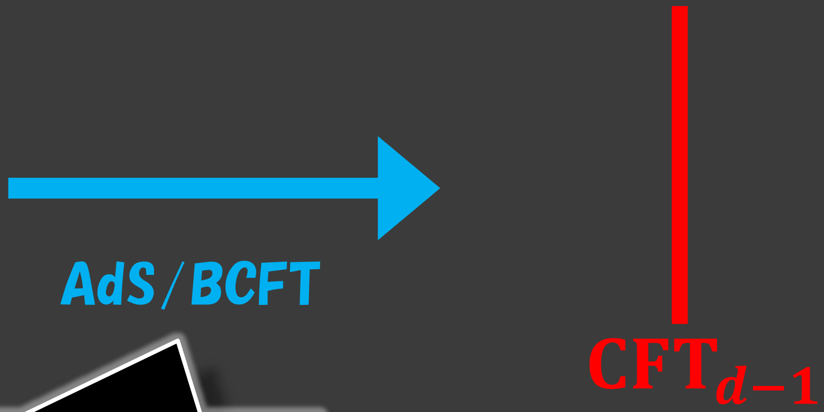


Naive derivation?



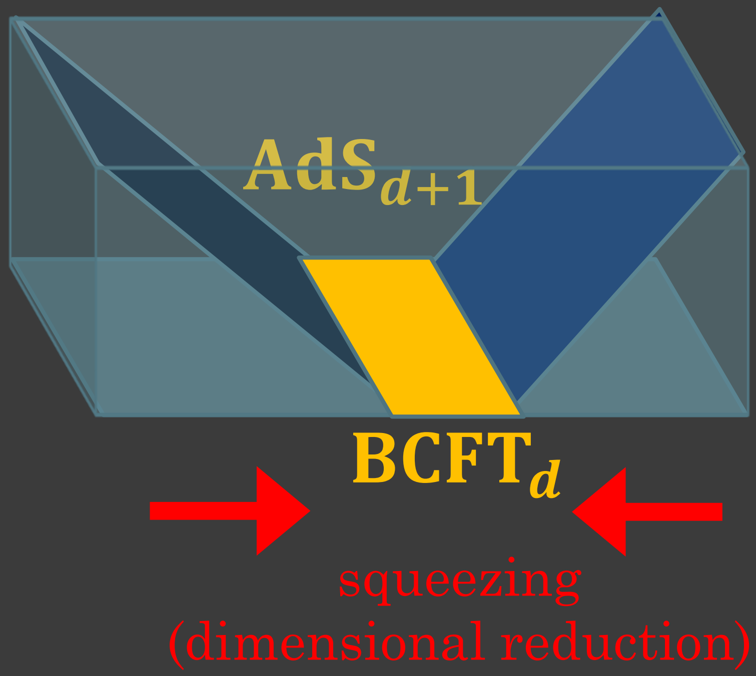
Naive derivation?





Instead, we employ AdS/BCFT,
clear understanding,
some justifications

Derivation from AdS/BCFT



Definition: **Limit of BCFT** (very clear!)

Derivation from AdS/BCFT

Definition: Limit of BCFT (very clear!)

⇒But is this really CFT?

Justifications:

- Brane world holography + AdS/CFT (at least large ρ_*)
- Boundary of BCFT_d can be thought of as CFT_{d-1} , because OPE of boundary state also satisfies axiom of CFT. In this sense, our CFT_{d-1} can be interpreted as two boundaries interacting with each other through bulk.
- ***From now on, we see matching of gravity calculation and CFT calculation for some physical quantities.***

Free energy on Sphere

One check of holography principle can be done by seeing ***free energy***.

The on-shell action of $\mathrm{AdS_{d+1}}$ (without boundaries) has the following form (see [Henningson, Skendteries]),

$$I = \#\frac{1}{\epsilon^{d}} + \#\frac{1}{\epsilon^{d-2}} + \cdots + \#\frac{1}{\epsilon^{2}} - \#\log\epsilon + O(1) \qquad (\text{if } d = odd)$$

$$I = \#\frac{1}{\epsilon^{d}} + \#\frac{1}{\epsilon^{d-2}} + \cdots + \#\frac{1}{\epsilon} + O(1) \qquad (\text{if } d = even)$$

In fact, $\mathrm{CFT_d}$ has the same form.

$$I_{AdS_{d+1}} = I_{CFT_d}$$

This is one consistency check of the $\mathrm{AdS_{d+1}}$ /$\mathrm{CFT_d}$ correspondence.

Free energy on Sphere

Let us consider $\mathrm{AdS}_{d+1}/\mathrm{CFT}_{d-1}$.

Of course, without the wedge, we can conclude

$$I_{AdS_{d+1}} \neq I_{CFT_{d-1}}$$

Therefore, the contribution from the **wedge** should have some non-trivial effects. We will see this.

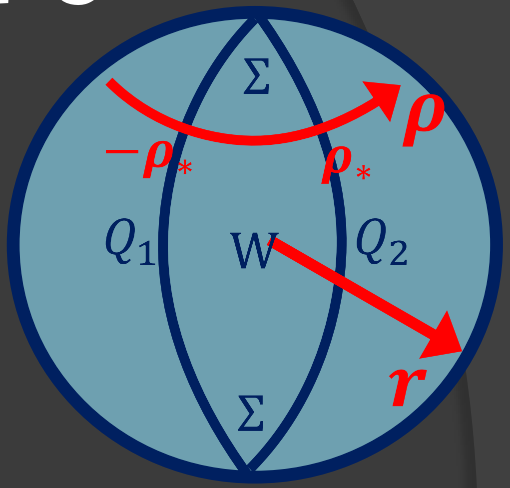
For simplicity, let us focus on $\mathrm{AdS}_4/\mathrm{CFT}_2$.

We know

$$I_{CFT_2} = \# \frac{1}{\epsilon^2} + \frac{c}{6}\chi(\Sigma)\log\epsilon + O(1)$$

Let's calculate the gravity action of wedged AdS_4 directly!

Free energy on Sphere



Metric of Euclidian AdS_{d+1}:

$$ds^2 = d\rho^2 + R^2 \cosh^2\frac{\rho}{R}\left(d\eta^2 + \sinh^2\eta\, d\Omega_{d-1}^2\right)$$
$$= dr^2 + R^2 \sinh^2\frac{r}{R}\left(d\theta^2 + \cos^2\theta\, d\Omega_{d-1}^2\right)$$

On-shell action gravity action on this wedge geometry,

$$I = -\frac{1}{16\pi G_N}\int_W \sqrt{g}(R - 2\Lambda) - \frac{1}{8\pi G_N}\int_{Q_1\cup Q_2}\sqrt{h}(K - T) - \frac{1}{8\pi G_N}\int_\Sigma \sqrt{h}\, K$$

For the wedged AdS_4, we have

$$I_{WAdS_4} = -\frac{R^2}{2G_N\epsilon^2}\sinh\frac{\rho_*}{R} + \frac{R^2}{G_N}\sinh\frac{\rho_*}{R}\log\epsilon + O(1)$$

Free energy on Sphere

$$I_{WAdS_4} = -\frac{R^2}{2G_N\epsilon^2}\sinh\frac{\rho_*}{R} + \frac{R^2}{G_N}\sinh\frac{\rho_*}{R}\log\epsilon + O(1)$$

The form for CFT_2 is

$$I_{CFT_2} = \#\frac{1}{\epsilon^2} + \frac{\mathrm{c}}{6}\chi(\Sigma)\log\epsilon + O(1)$$

The ϵ dependence perfectly matches !!

The central charge (i.e. degrees of freedom) of the wedged AdS_4 is ($\chi(\mathrm{S}^2) = 2$)

$$c = \frac{3R^2}{G_N}\sinh\frac{\rho_*}{R}$$

Next, we will check this result in an independent way.

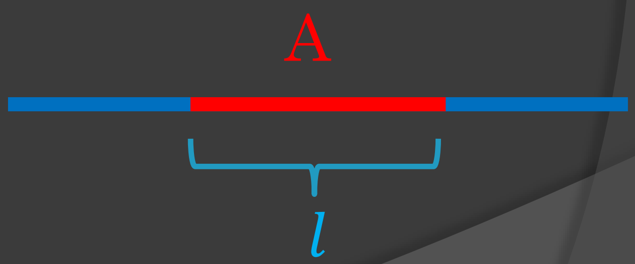
Entanglement entropy

Entanglement entropy can be evaluated from both **CFT** and **gravity** side. Therefore, it is very useful to check our new holographic principle.

For simplicity, let us focus on AdS_4/CFT_2 and consider a single interval A with size l.

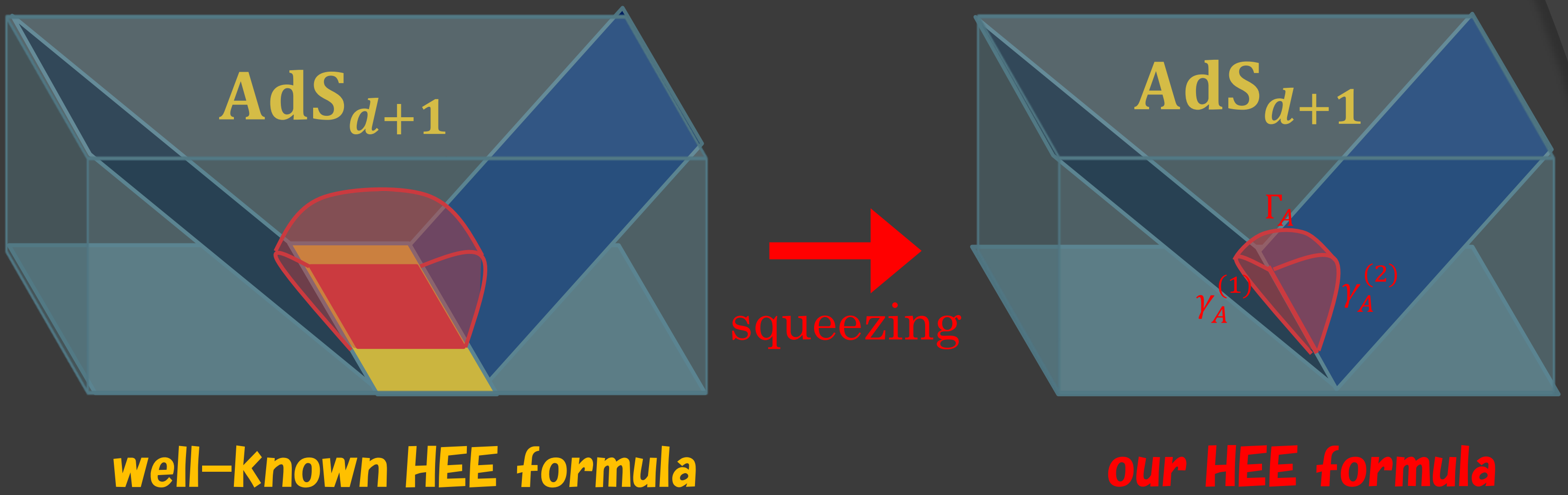
In this case, the entanglement entropy in CFT_2 is known as

$$S_A = \frac{c}{3}\log\frac{l}{\epsilon}$$



Let us compare this with gravity calculation.

Holographic EE



The holographic entanglement entropy for our holography is naturally defined by recalling its definition,

$$S_A = \min_{\substack{\gamma_A^{(1)},\gamma_A^{(2)} \\ \partial\gamma_A^{(1,2)}=\partial A}} \left(\min_{\substack{\Gamma_A \\ \partial\Gamma_A=\gamma_A^{(1)}\cup\gamma_A^{(2)}}} \left(\frac{\text{Area}(\Gamma_A)}{4G_N} \right) \right)$$

Holographic EE

This formula in 4d case leads to

$$S_A = \frac{R^2}{G_N} \sinh \frac{\rho_*}{R} \log \frac{l}{\epsilon}$$

This form matches with the CFT_2 result.

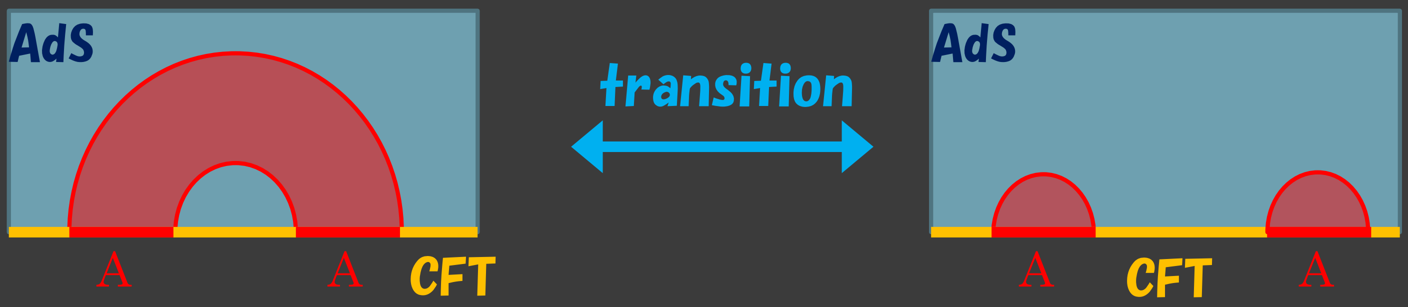
We can extract the **central charge** from this result as

$$c = \frac{3R^2}{G_N} \sinh \frac{\rho_*}{R}$$

This is completely same as the previous result!

Note: we can see this matching for more general cases (higher dimension and intervals).

Holographic EE



Double interval case:

We should find ***transition***, like the standard HEE.

> **Implication:**
> ***CFT dual to Wedge geometry is similar to holographic CFT (i.e. sparse CFT)***

Finding explicit constraints on the CFT data is interesting future work!

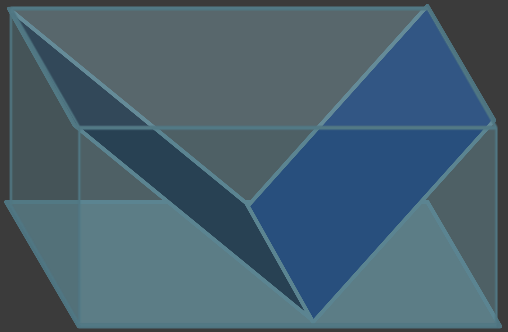
Comments

- CFT_{d-1} ***dual to*** AdS_{d+1}
  - $c = \frac{3R^2}{G_N} \sinh \frac{\rho_*}{R}$
  - Sparseness (not the same as HKL sparseness)
  - Large DoF of AdS_{d+1} leads to infinite tower of primary operators, which are interpreted as KK modes in brane world.

    $\Rightarrow$NOT holographic CFT
  - Zero size limit seems to be singular, but from gravity perspective, quantity in CFT_{d-1} are well-defined as wee saw.

Comments

- **Braneworld holography**



Wedged AdS_{d+1}

Joining gravity

Accessible

CFT_{d-1}

Comments

- **Braneworld holography**

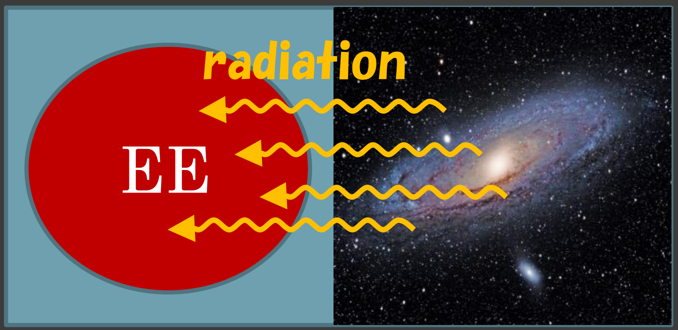
Joining gravity

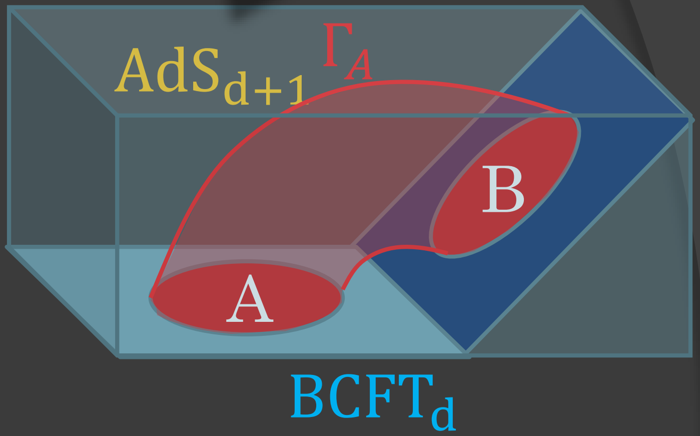
From compactification,
$$\frac{1}{G_N^{(d)}} = \frac{R}{G_N^{(d+1)}} \sinh\frac{\rho_*}{R}$$

- $$c_{AdS_3} \times 2 = \frac{3R}{2G_N^{(3)}} \times 2 = \frac{3R^2}{G_N^{(4)}} \sinh\frac{\rho_*}{R}$$

  is consistent with the central charge derived from EE.

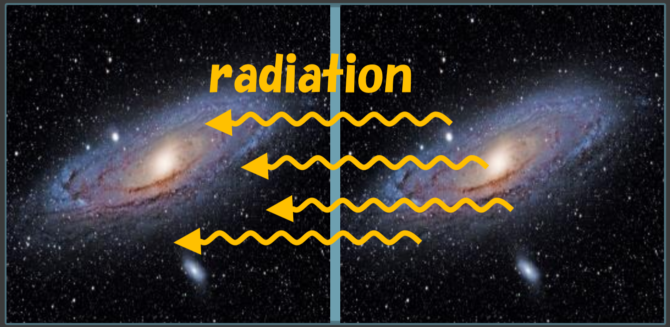
Comments





double holography
[Almheri, Mahajan, Maldacena, Zhao]

Accessible

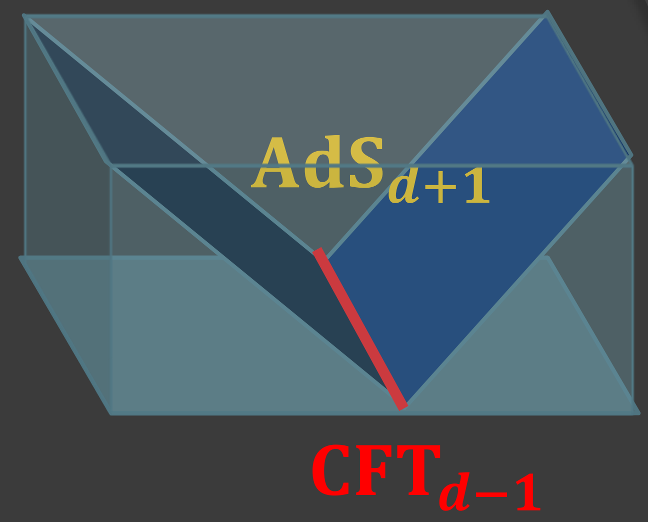


More realistic setup
[Geng, Karch, Paradavila, Raju, Randall, Riojas, Shashi]

Summary①

Wedge Holography

$\mathrm{AdS}_{d+1} = \mathrm{CFT}_{d-1}$



Justifications

- Brane world picture:

  Brane world holography + AdS/CFT (at large ρ_*)
- BCFT picture:

  CFTs on two boundaries with bulk interaction
- Matching of physical quantities

  explicit calculation of DoF and EE from both sides

Summary①

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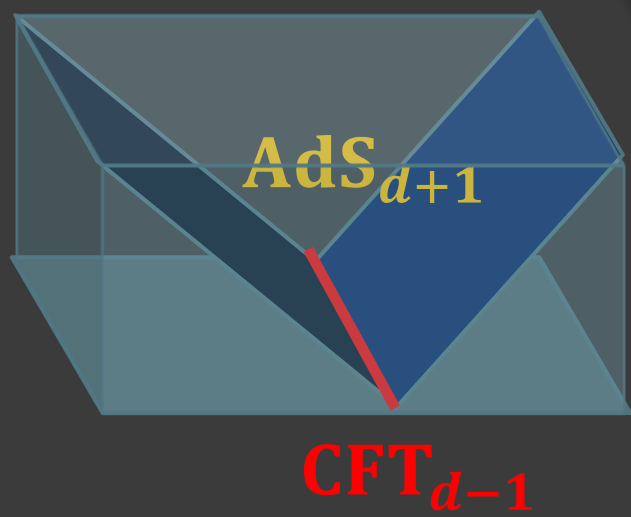
$\mathrm{AdS}_{d+1} = \mathrm{CFT}_{d-1}$



Future direction

- Employing this new laboratory
- Identifying universal properies of CFT_{d-1}
- Joining gravity from AdS_{d+1} and CFT_{d-1}
  ⇒more general double holography setup
- More justification (e.g. realization via string theory)

Appendix
(Another wedge holography)

W-holography at finite temp.

Let us consider more general case.

BTZ is simplest one other than Poincare AdS.

Metric of BTZ:

$$ds^2 = -(r^2 - r_0^2)dt^2 + R^2 \frac{dr^2}{r^2 - r_0^2} + r^2 dx^2$$

Coordinate transformation

$$t' \pm x' = \pm e^{\frac{r_0}{R}(x \pm t)} \sqrt{1 - \frac{r_0^2}{r^2}}, \quad z = \frac{r_0}{r} e^{\frac{r_0}{R}x}$$

Poincare:

$$ds^2 = R^2 \left(\frac{dz^2 - dt'^2 + dx'^2}{z^2} \right)$$

W-holography at finite temp.

General solution to Neumann bdy. in Poincare metric

$$(z-\alpha)^2 + (x'-p)^2 - (t'-q)^2 = \beta^2$$

Its tension is $T = \frac{\alpha}{\beta R}$

Coordinate transformation

$RT < 1$: AdS_2 brane

$$r(x) = \frac{r_0 T R}{\sqrt{1 - T^2 R^2}\sinh\frac{r_0 x}{R}}$$

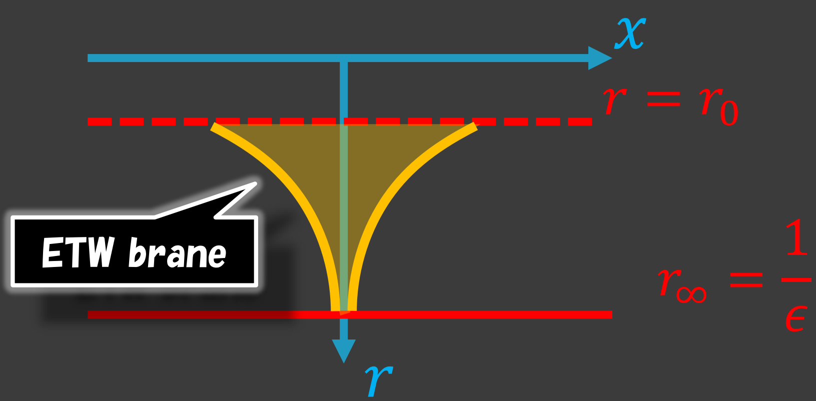
$RT = 1$: R_2 brane

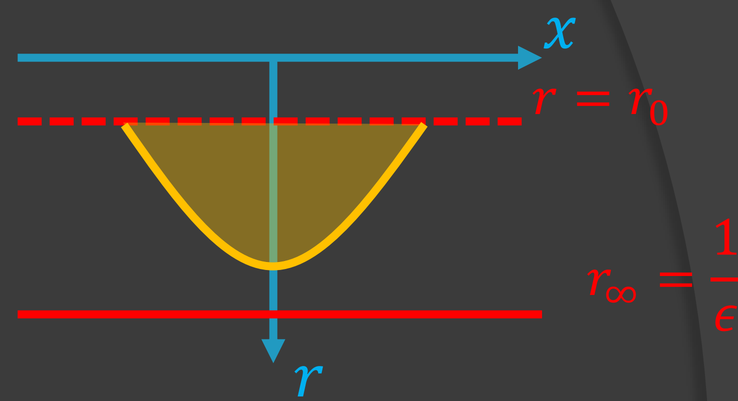
$$r(x) = 2 r_0 e^{\frac{r_0 x}{R}}$$

$RT > 1$: dS_2 brane

$$r(x) = \frac{r_0 T R}{\sqrt{T^2 R^2 - 1}\cosh\frac{r_0 x}{R}}$$

W-holography at finite temp.





$T < 1/R$
AdS_2 brane

$T > 1/R$
dS_2 brane

AdS brane case:

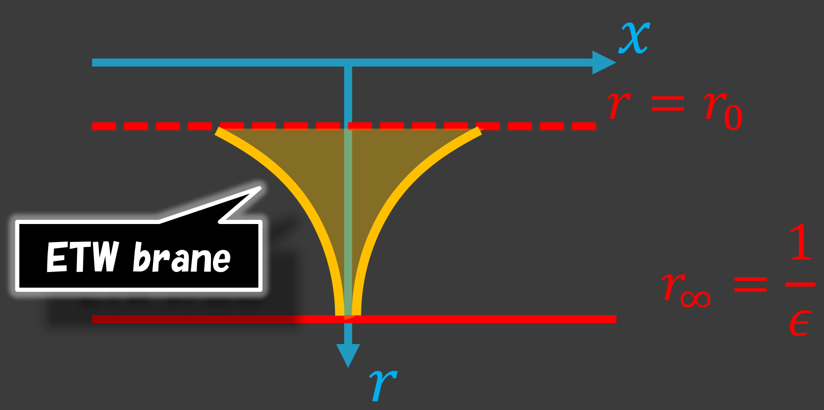
$$I_W = -\frac{\beta \sinh \rho_*}{4\pi G_N \epsilon} - \frac{\rho_*}{2G_N}$$

$= 2 \times S_{bdy}$

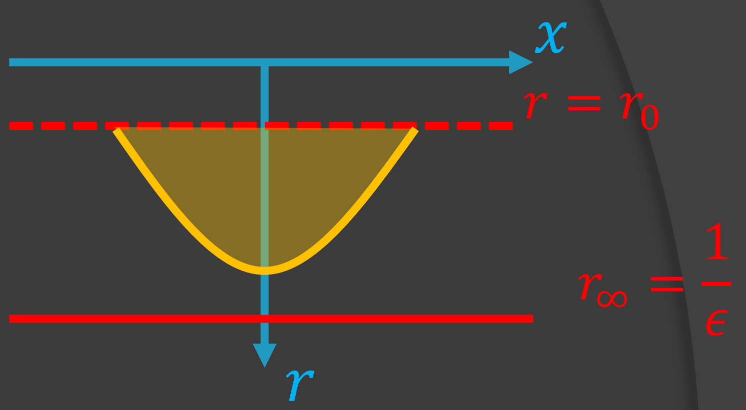
Gravity action is consistent with anomaly of CFT_1.

DoF of CFT_1 is equal to DoF of two AdS branes.

W-holography at finite temp.



$T < 1/R$
AdS_2 brane



$T > 1/R$
dS_2 brane

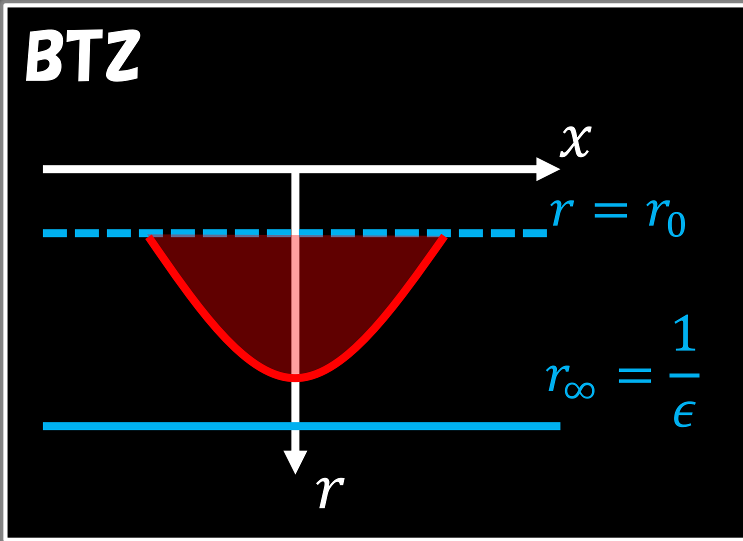
dS brane case?

No asymptotic boundary ⇒ No CFT dual ?

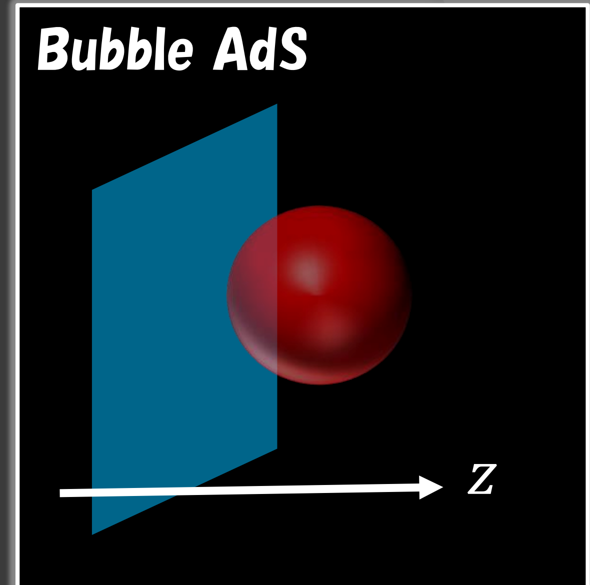
Contents

- Introduction
- Basics of AdS/BCFT
  - Quick lesson of BCFT
  - Gravity dual
- **Wedge holography**
  - Codimension two holography
  - **Another wedge holography**
- Summary

Imaginary BCFT



back to (Euclidean) Poincare

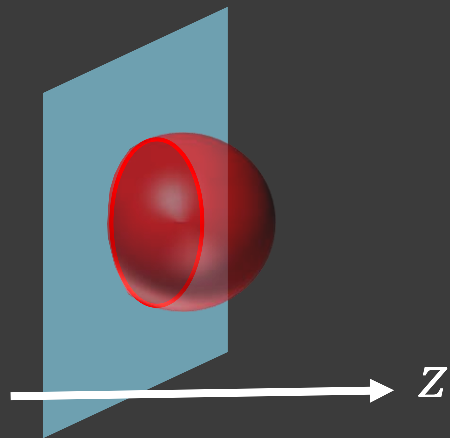


Profile of **bubble** brane is

$$(z-\alpha)^2 + x^2 + t^2 = \beta^2 \qquad \text{with} \qquad |\alpha| > |\beta|$$

Bubble AdS had not been considered in AdS/BCFT context.
Here, we give the CFT dual of this bubble AdS.

Imaginary BCFT

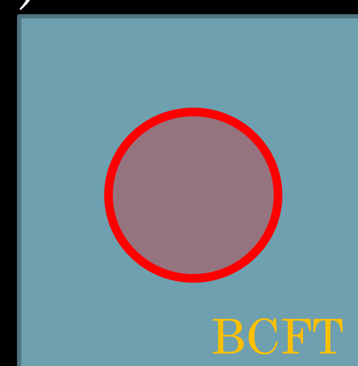


$T < 1/R$: AdS_2 brane

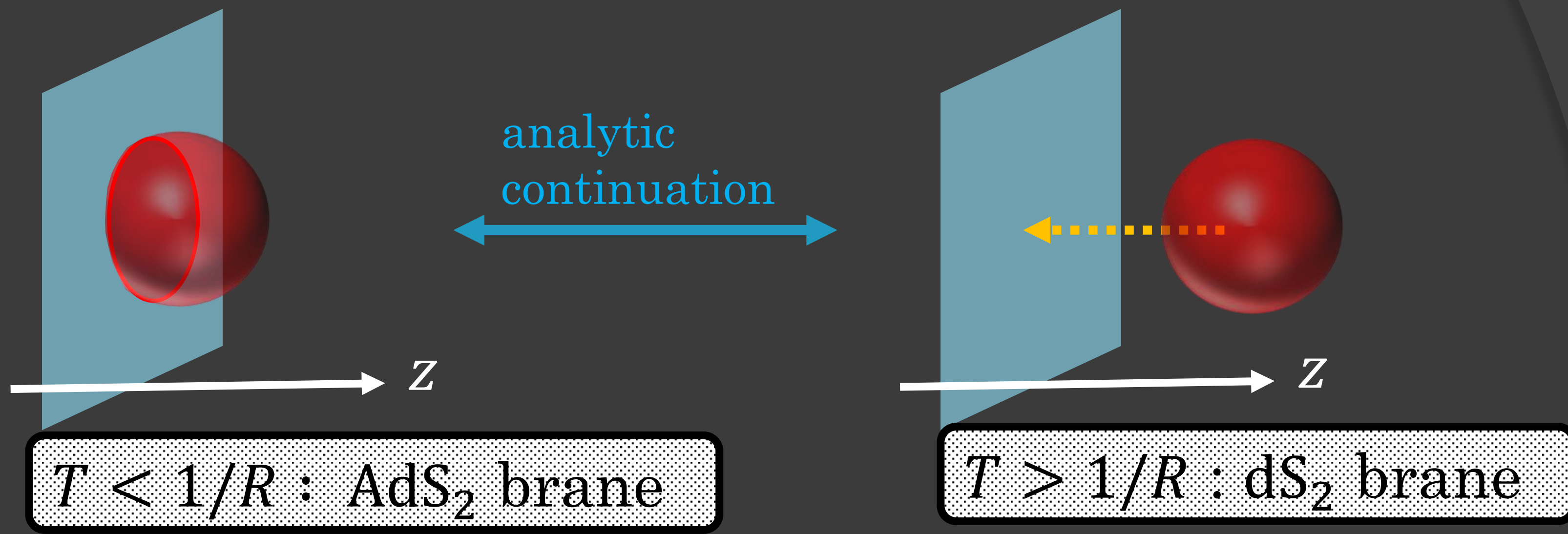
This geometry has BCFT dual (**Known!**).
Its bdy. is circle with radius

$$r^2 = \beta^2 - \alpha^2$$

Note that $T < 1/R \iff \beta^2 > \alpha^2$



Imaginary BCFT



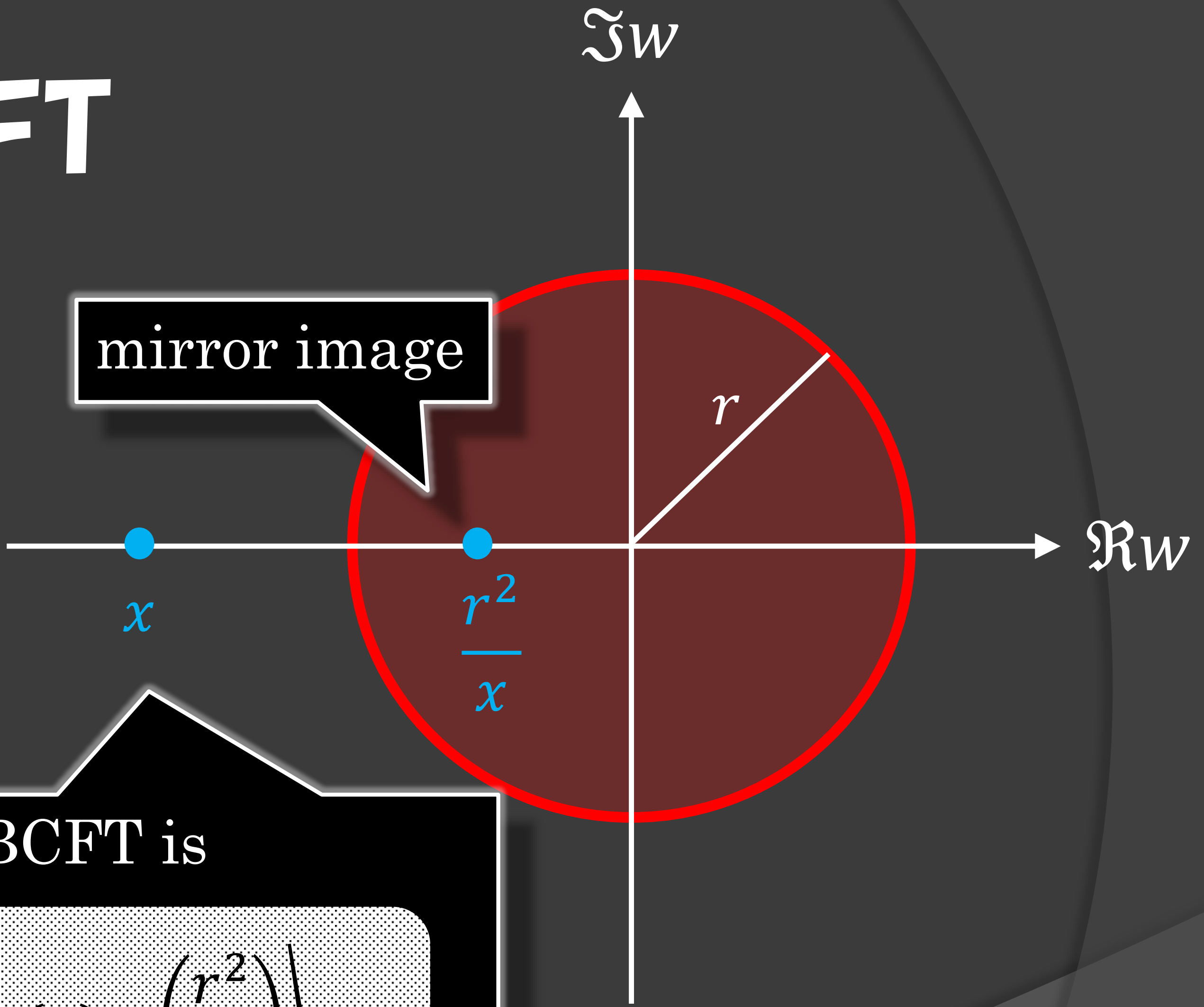
$T > 1/R$ implies

$$r^2 = \beta^2 - \alpha^2 < 0$$

Imaginary radius is problematic, but in gravity picture, this analytic continuation is not so problematic.
We formally define **imaginary BCFT**.

EE in BCFT

Let us recall that mirror method can completely fix 1-pt.correlator.



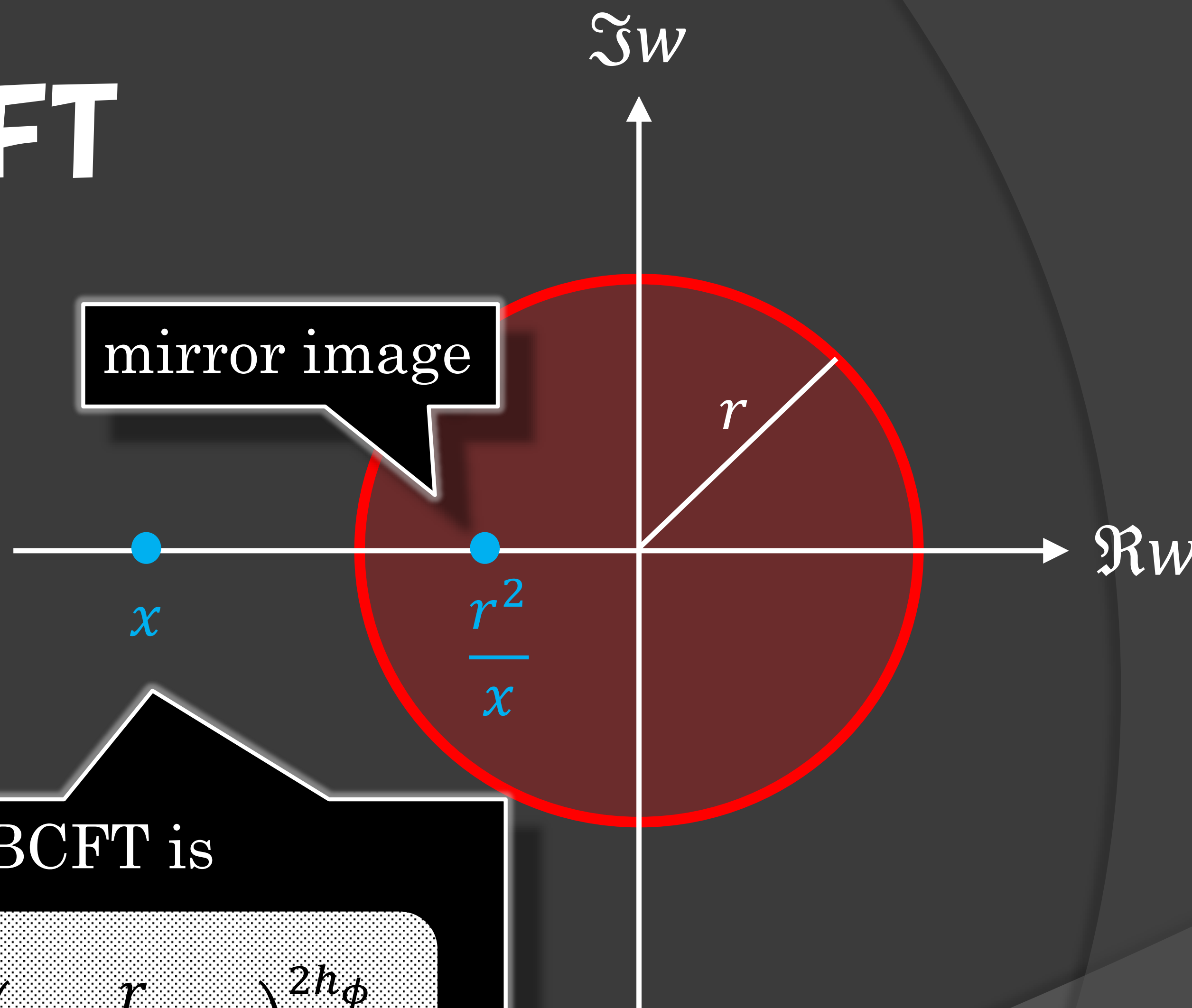
1-pt. correlator in BCFT is

$$\langle \phi(x) \rangle_{disk} \simeq \left\langle \phi(x)\phi\left(\frac{r^2}{x}\right)\right\rangle$$

up to conformal prefactor.

EE in BCFT

Let us recall that mirror method can completely fix 1-pt.correlator.



1-pt. correlator in BCFT is

$$\langle\phi(x)\rangle_{disk} = g_B\left(\frac{r}{|w|^2 - r^2}\right)^{2h_\phi}$$

EE in BCFT

For example, EE in the right setup can be evaluated by

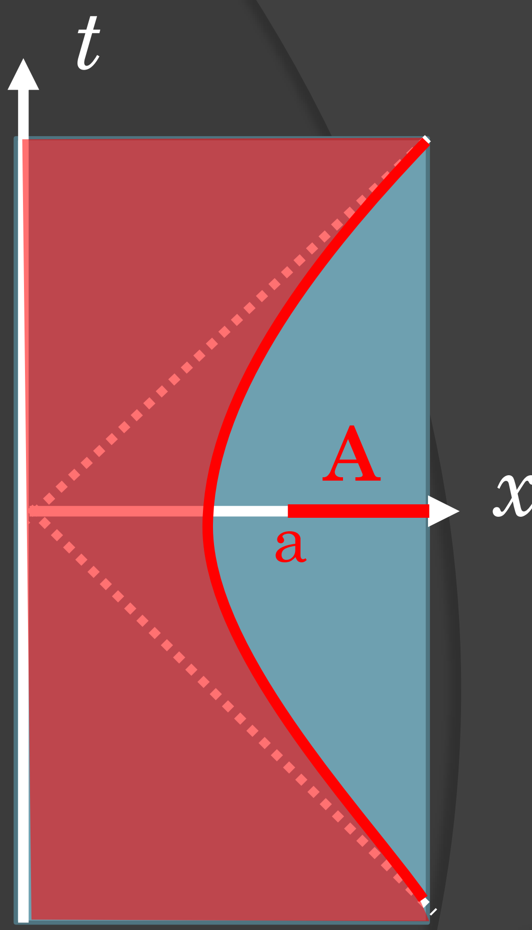
$$\langle \sigma_n(\tau + ia) \rangle_{disk} = g_B \left(\frac{r}{\tau^2 + a^2 - r^2} \right)^{2h_{\sigma_n}}$$

where σ_n is twist op. with $h_{\sigma_n} = \frac{c}{24}\left(n - \frac{1}{n}\right)$

After analytic continuation $\tau \to it$,

$$S_A(t) = \lim_{n\to 1} \frac{1}{1-n} \log\langle \sigma_n(\tau + ia) \rangle_{disk}$$
$$= \frac{c}{6} \log \frac{a^2 - t^2 - r^2}{\epsilon r} + S_{bdy}$$

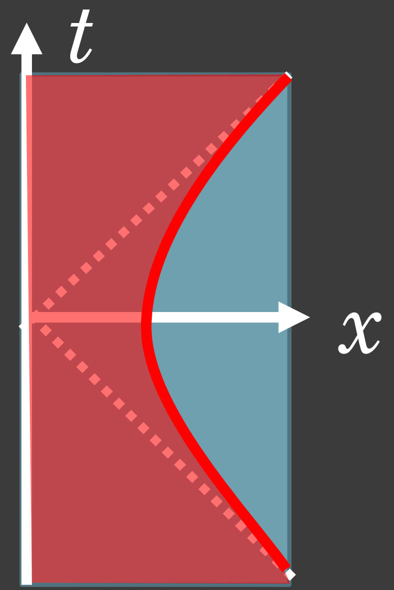
where $S_{bdy} = \log g_B$ is boundary entropy



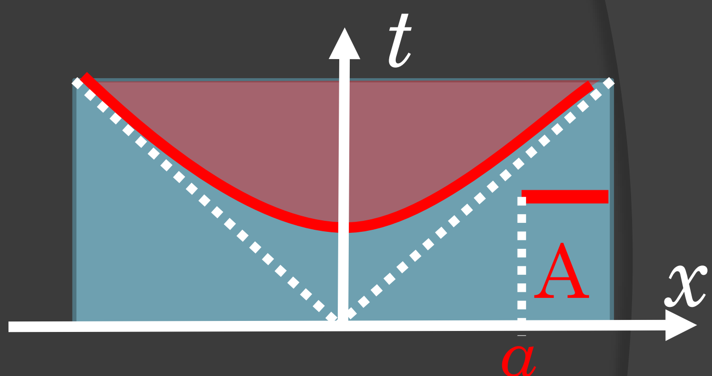
Lorentzian ver. of Euclidean disk

EE in Im-BCFT

Let us consider imaginary BCFT in physical setup



Analytic continuation
to imaginary radius

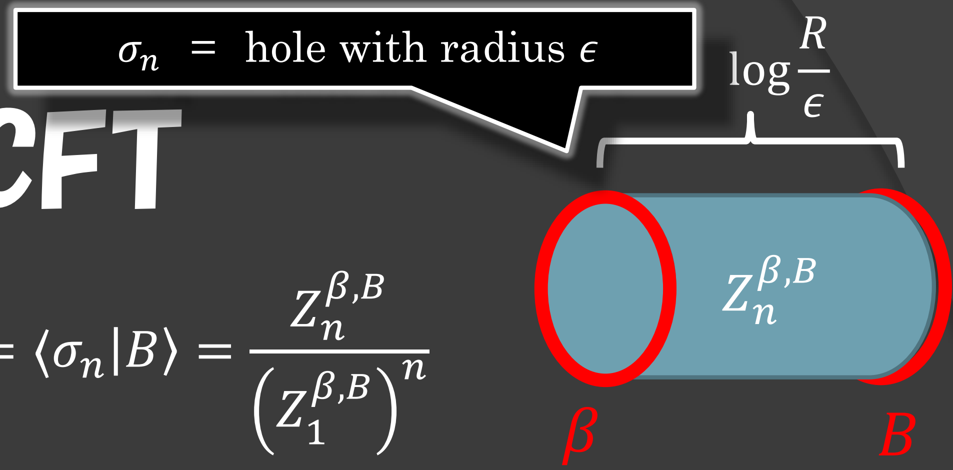


creation of universe setup

Naïve analytic continuation is imaginary.
To be more careful.

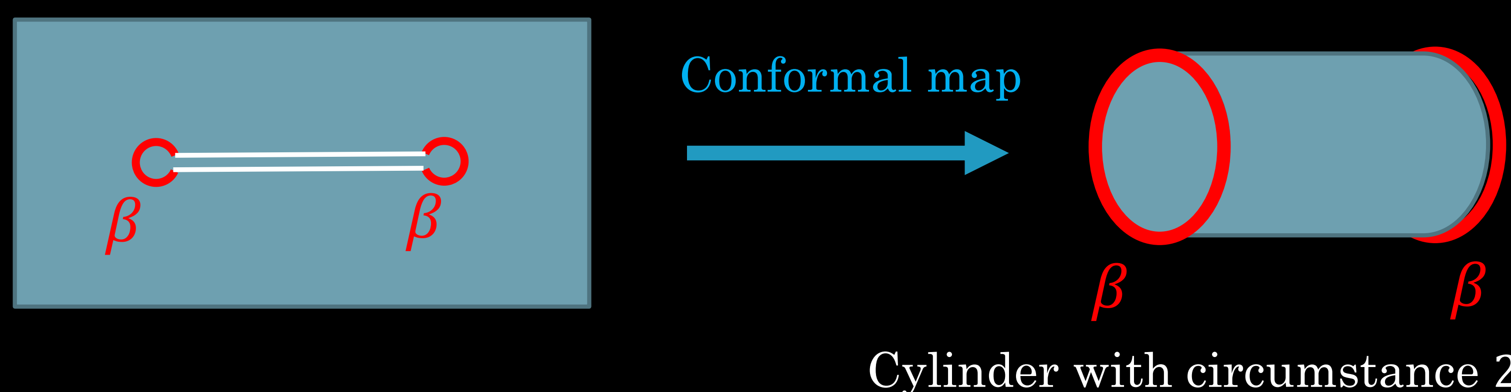
$$S_A(t) = \frac{c}{6}\log\frac{a^2-t^2+r^2}{i\epsilon r} + S_{bdy}$$

EE in lm-BCFT



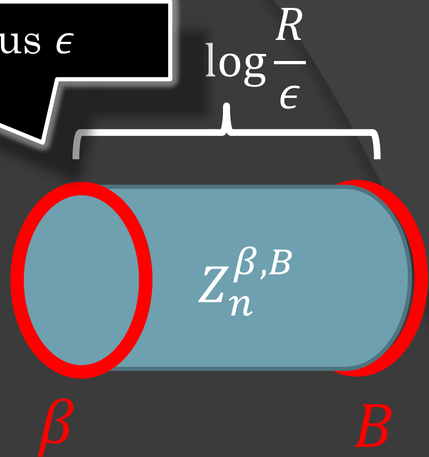
$$g_B = \langle \sigma_n | B \rangle = \frac{Z_n^{\beta,B}}{\left(Z_1^{\beta,B}\right)^n}$$

where β is boundary state related to σ_n (which can be absorbed in regularization ϵ)



EE in Im-BCFT

σ_n = hole with radius ϵ



$$g_B = \langle \sigma_n | B \rangle = \frac{Z_n^{\beta,B}}{\left(Z_1^{\beta,B}\right)^n}$$

where β is boundary state related to σ_n (which can be absorbed in regularization ϵ)

$$Z_n^{\beta,B} = \langle \beta | e^{-\frac{1}{2\pi n}\log\frac{R}{\epsilon}H} | B \rangle \xrightarrow[\epsilon\to 0]{} e^{-\frac{1}{2\pi n}\log\frac{R}{\epsilon}E_0} \langle \beta | 0 \rangle \langle 0 | B \rangle$$

R is formal radius $R = 1$ ($\rightarrow i$ later). As a result,

$$\langle \sigma_n | B \rangle = \left(\frac{R}{\epsilon}\right)^{-2h_{\sigma_n}} \left(\langle \beta | 0 \rangle \langle 0 | B \rangle\right)^{1-n}$$

Analytic continuation ($R \rightarrow iR$) leads to

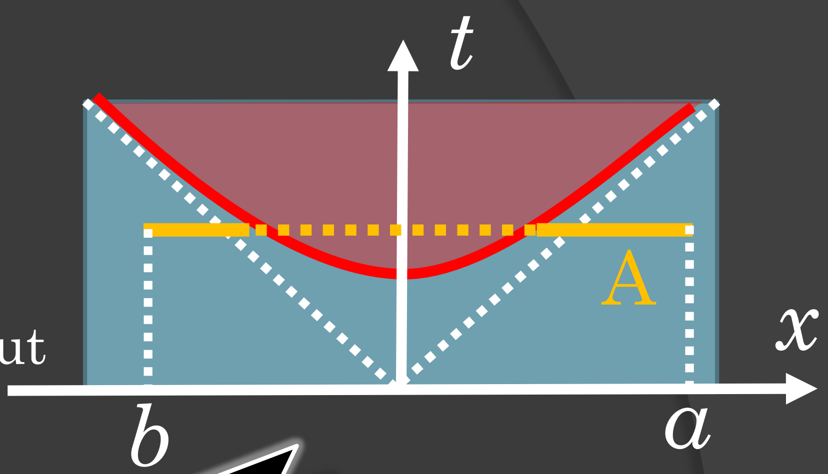
important point

$$\langle \sigma_n | Im\text{-}B \rangle = i^{-2h_{\sigma_n}} \langle \sigma_n | B \rangle$$

cancel imaginary part

$$S_A(t) = \frac{c}{6}\log\frac{a^2 - t^2 + r^2}{\epsilon r} + S_{bdy}$$

EE in Im-BCFT



CFT calculation is non-trivial but gravity calculation is trivial (just a geodesic in a cut AdS).

⇒**Consistency check** from gravity calculation is possible.

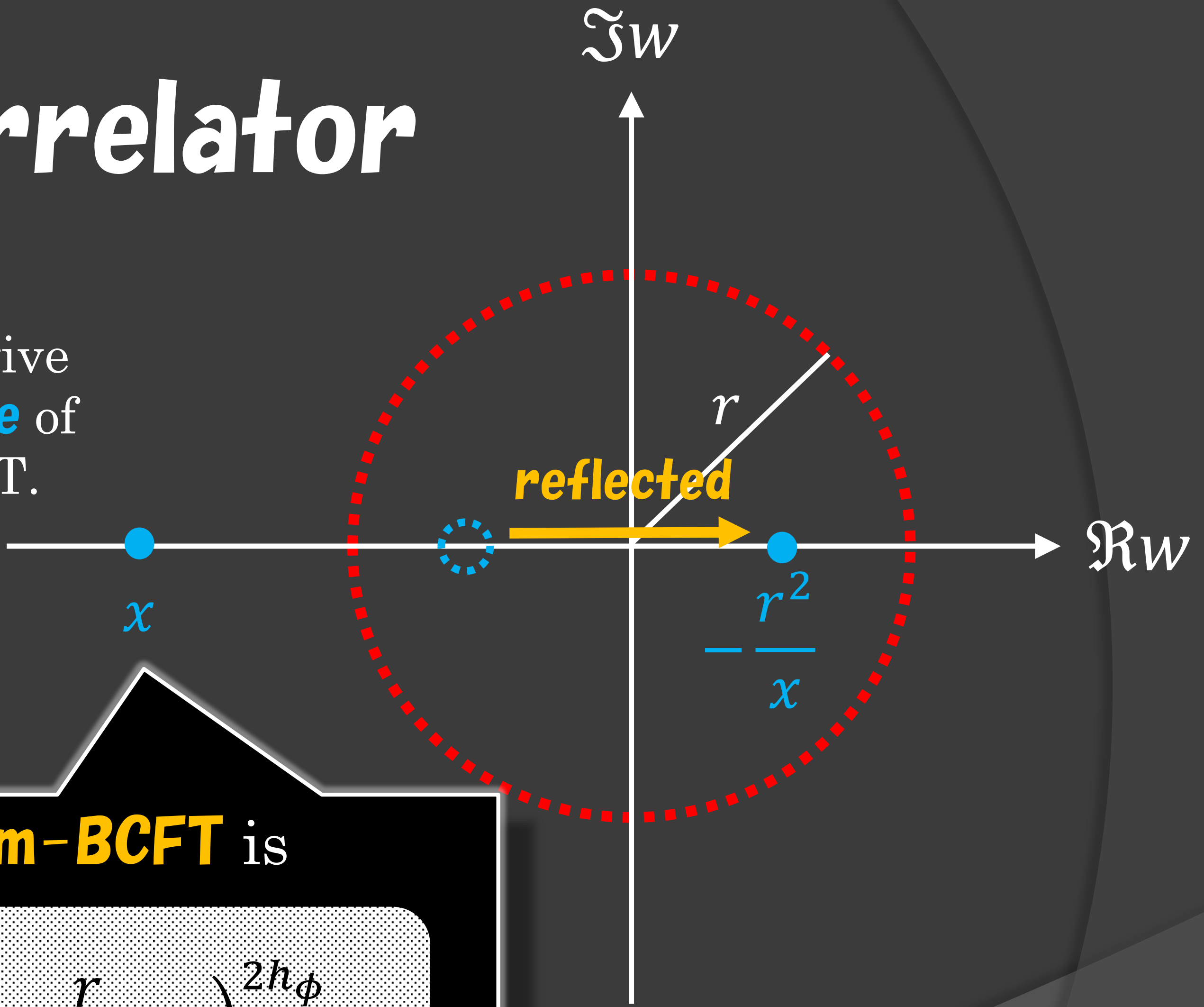
$$S_A = \lim_{n\to 1}\frac{1}{1-n}\log\langle\sigma_n(a,t)\overline{\sigma_n}(b,t)\rangle_{Im-disk}$$

$$= \begin{cases} \frac{c}{6}\log\frac{a^2+r^2-t^2}{r\epsilon} + \frac{c}{6}\log\frac{b^2+r^2-t^2}{r\epsilon} + 2g_B & \text{: disconnected} \\ \frac{c}{3}\log\frac{b-a}{\epsilon} & \text{: connected} \end{cases}$$

CFT calculation based on our prescription completely matches with the gravity calculation!

BCFT Correlator

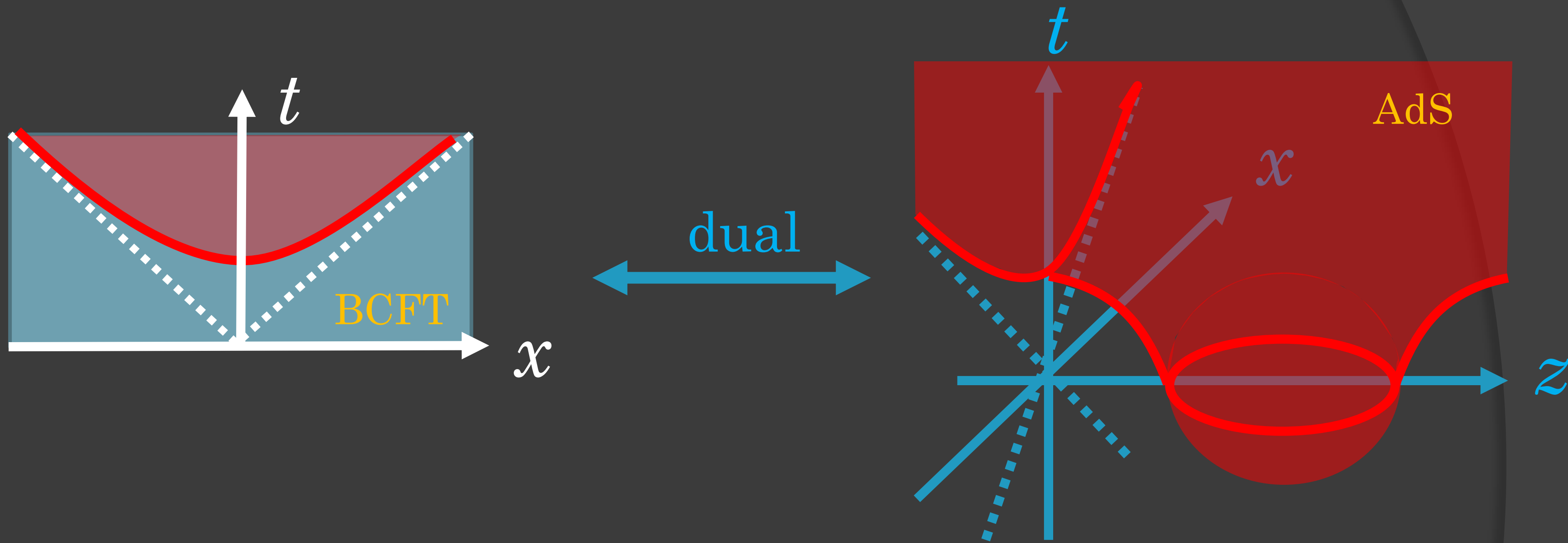
Instead of technical consideration, we can give a ***simple calculation rule*** of a correlator in Im-BCFT.



1-pt. correlator in **Im-BCFT** is

$$\langle \phi(x) \rangle_{disk} = \left(\frac{r}{|w|^2 + r^2} \right)^{2h_\phi}$$

Lorentzian Im-BCFT



Analytic continuation to Lorentzian signature provides time-like boundary in CFT.

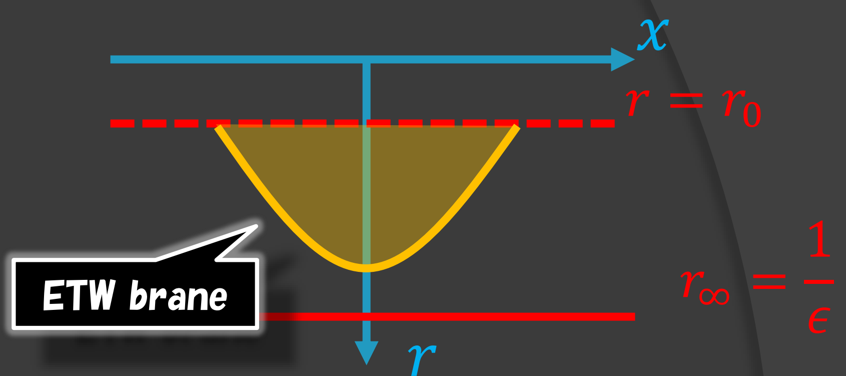
This setup is interesting for two reasons:

- ***universe creation setup***
- ***dS braneworld*** (as mentioned before) by CFT description as boundary state

Contents

Summary②

dS brane/Imaginary BCFT duality



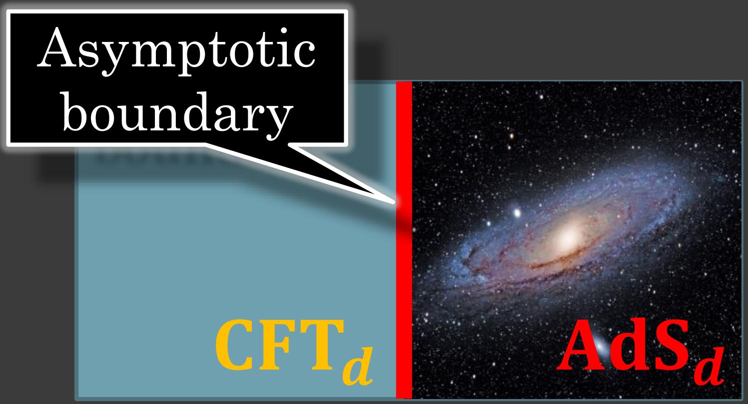
$T > 1/R$
dS brane

Future direction

- Employing this new laboratory
- More understanding of dS braneworld holography
- Universe creation setup
- More justification

Appendix
(extra comments)

Developments from AdS/BCFT



Setup:
AdS$_d$ & CFT$_d$ are glued along the (asymptotic) boundary

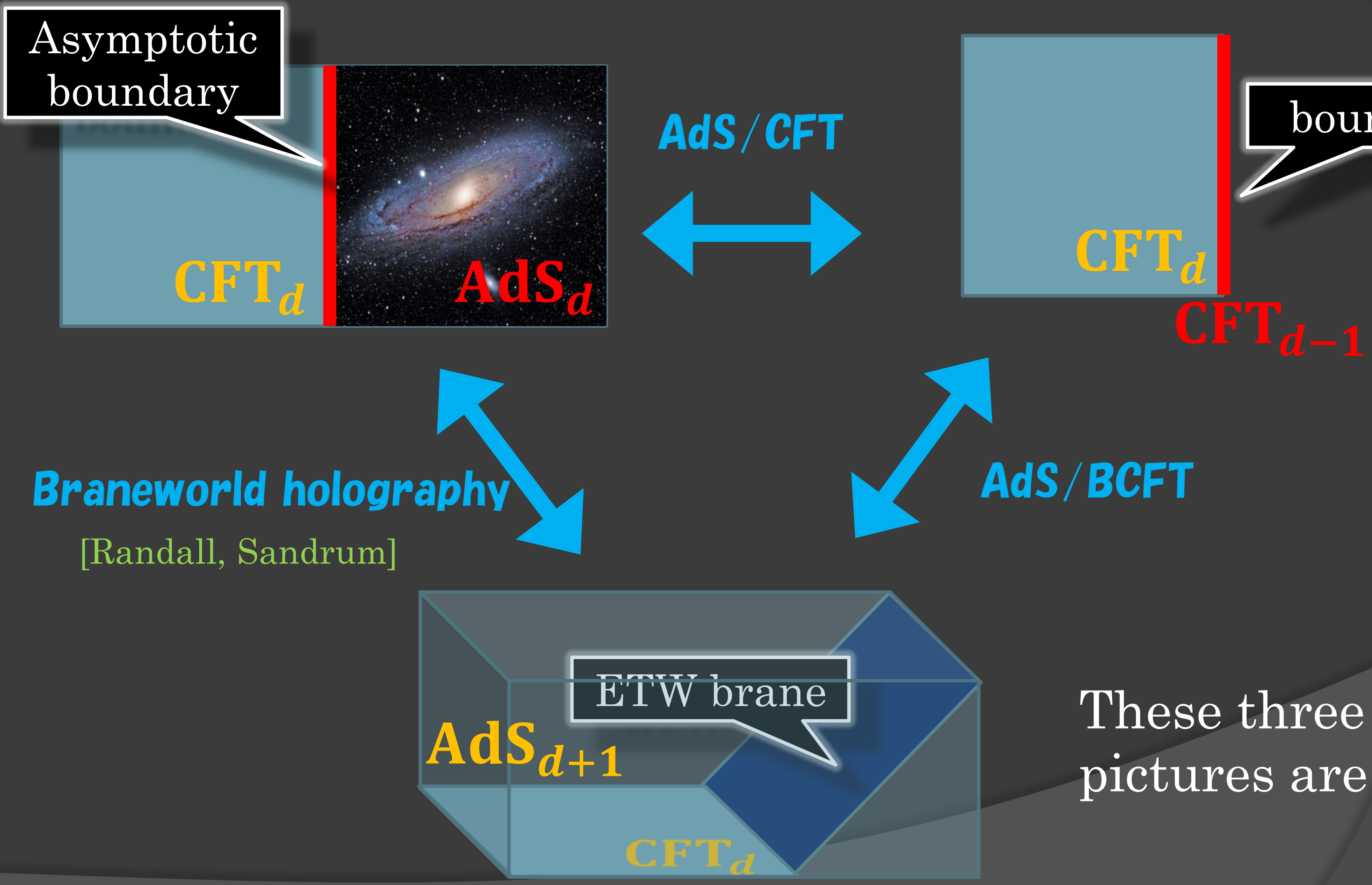
This AdS$_d$ is dynamical.
Light can go through asymptotic boundary.
We can discuss the Page curve in this setup.

Developments from AdS/BCFT

Setup:
AdS_d & CFT_d are glued along the (asymptotic) boundary

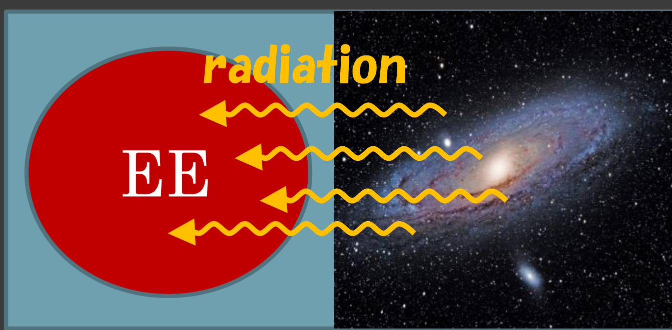
AdS/CFT correspondence:
$\mathrm{AdS}_d = \mathrm{CFT}_{d-1}$
This CFT_{d-1} can be thought of as boundary object of CFT_d

Developments from AdS/BCFT



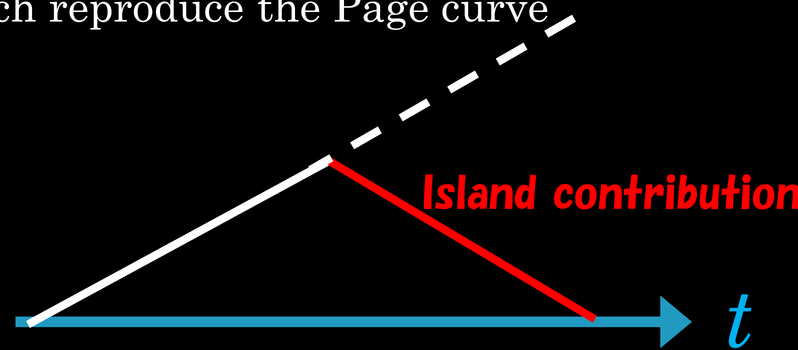
These three pictures are same

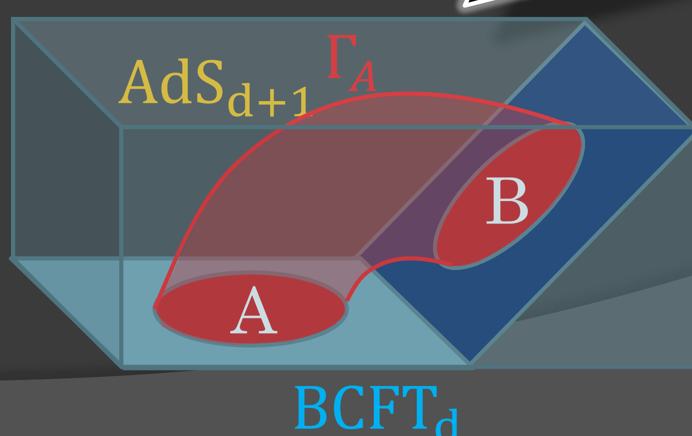
Developments from AdS/BCFT



Braneworld holography

By utilizing **AdS/BCFT**,
We can find the contribution from **Island**
(effective quantum contribution),
which reproduce the Page curve





double holography
[Almheri, Mahajan, Maldacena, Zhao]