

超弦理論入門

立川裕二 (IPMU & IAS)

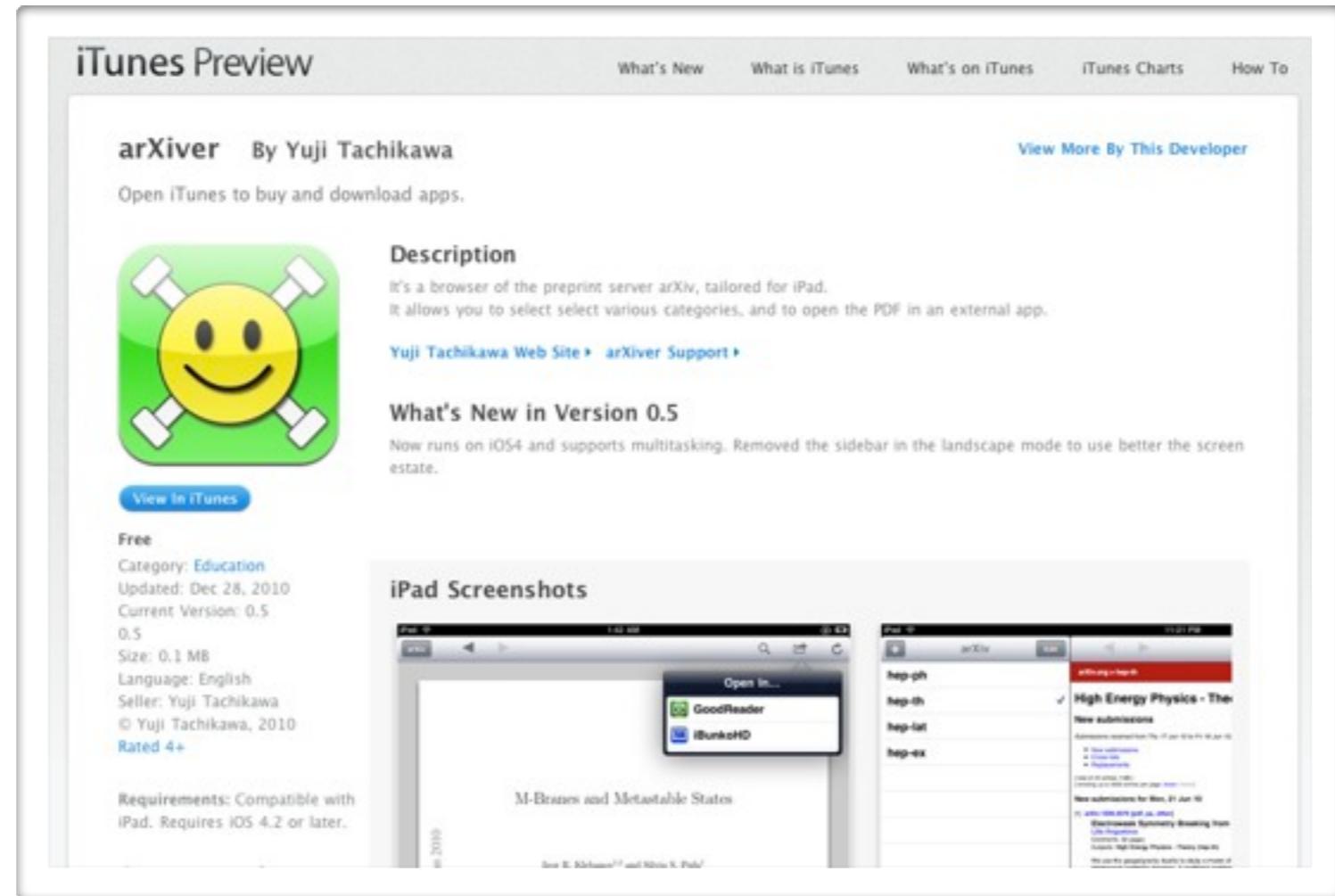
at the 17th ICEPP Winter Symposium



Institute for the Physics and Mathematics of the Universe

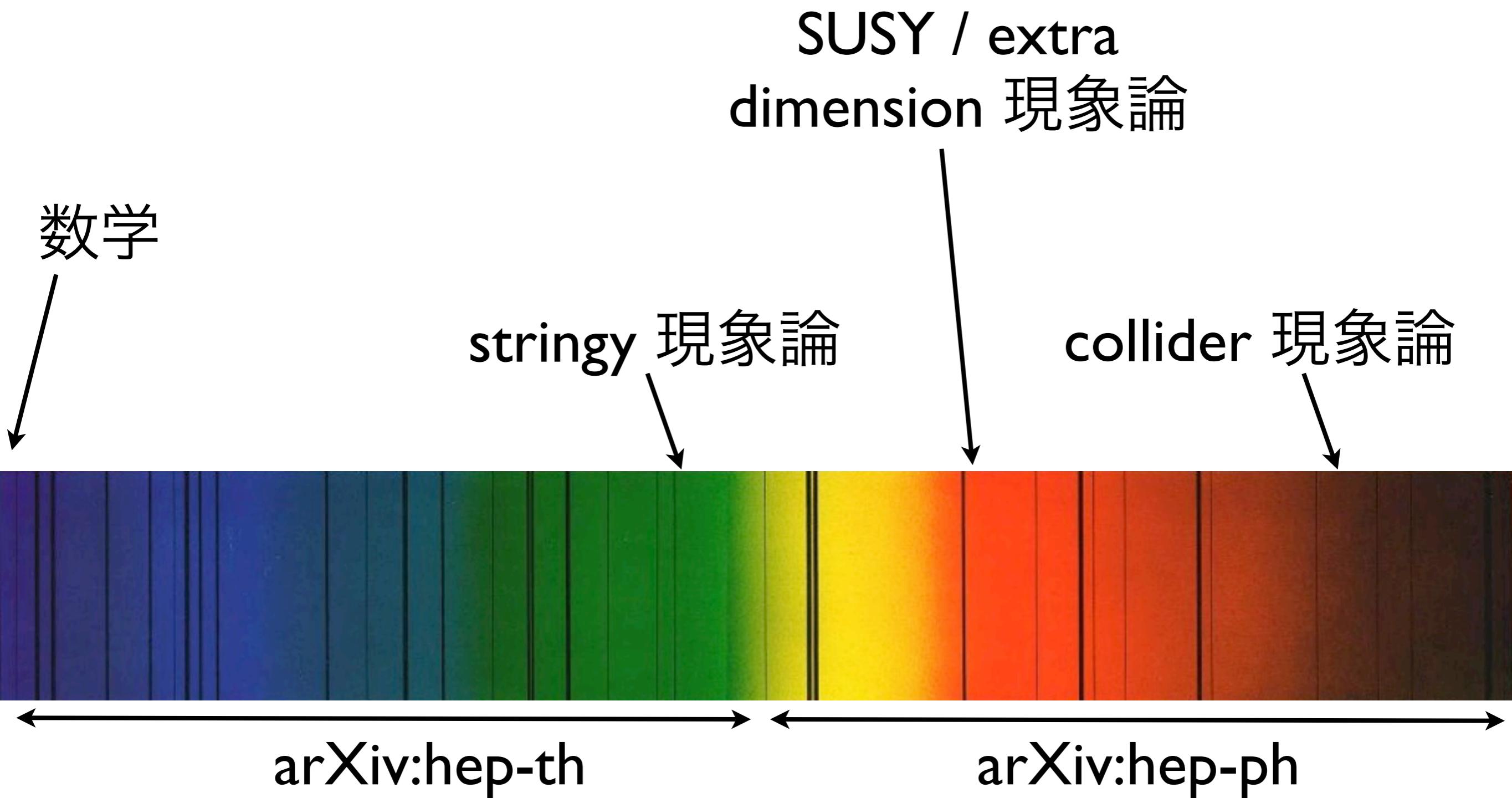
数物連携宇宙研究機構

数年前できました。宜しくお願ひします。

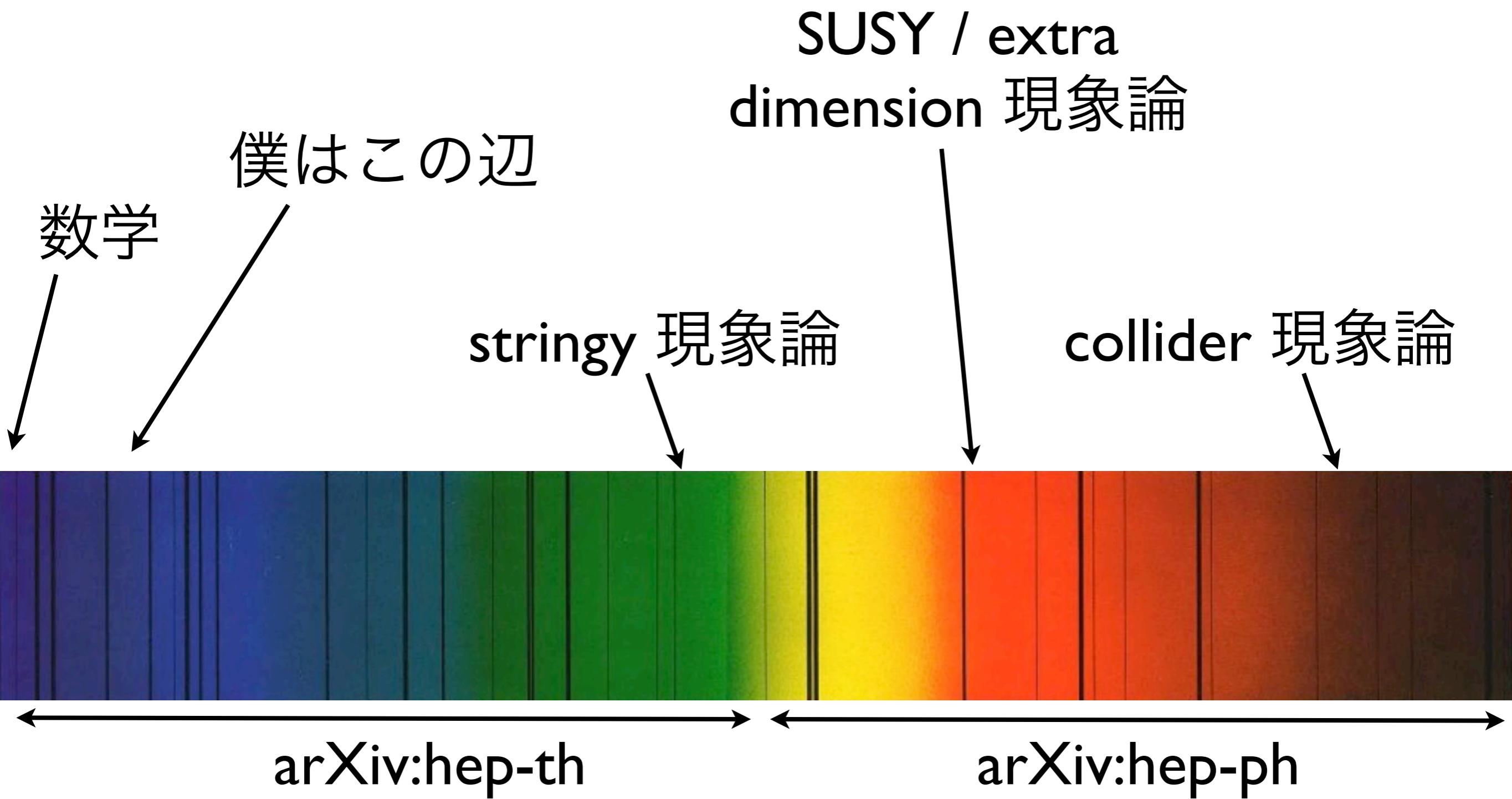


- プログラミングが趣味なので、同好の方はよろしくおねがいします。
- Mac / iPhone / iPad で Objective-C を使います。

Spectrum of High Energy Physics Theory



Spectrum of High Energy Physics Theory



- というわけで、超弦理論だとするとどういうのが実験で見えるかもしれない、という話は出来ません...
- ごめんなさい。

- そのかわり、なぜ理論家が超弦理論が
自然な枠組みだと思っているかを伝え
られたらと思います。

- 場の量子論:
標準模型の書かれている言葉
- 点粒子を相対論的に量子化したもの
- 案外難しかった。
 - 電磁気: 戦後すぐ、
 - 弱い力/強い力: ~1970年代
 - 重力: 出来ていない

- (超)弦理論
重力を量子化できる。
- 弦を相対論的に量子化したもの
- かなり難しかった。
はじまり: 1974 ~
第一次革命: 1984 ~
第二次革命: 1995 ~

黎明期



Veneziano



Nambu

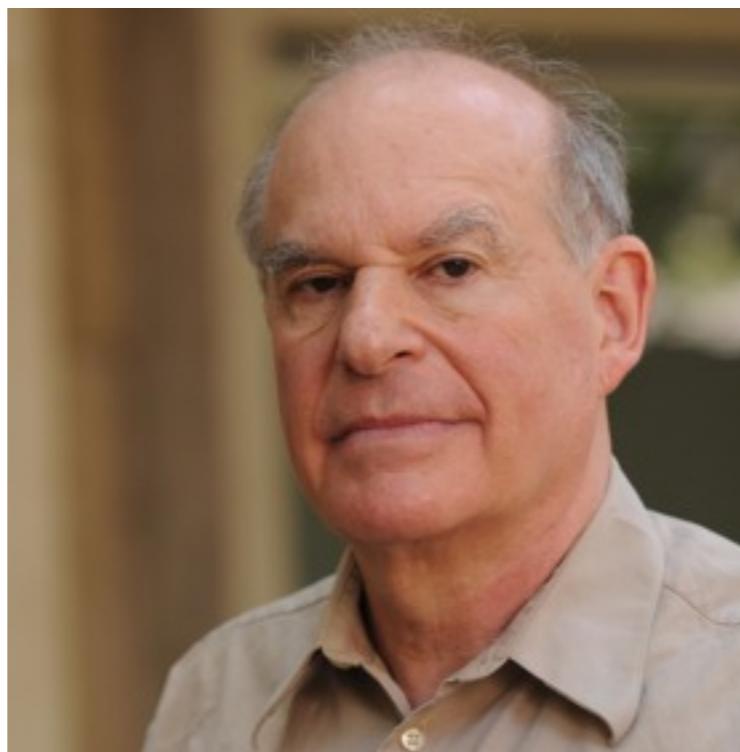


Polyakov

| 1984~



Green



Schwarz



Witten

1995~



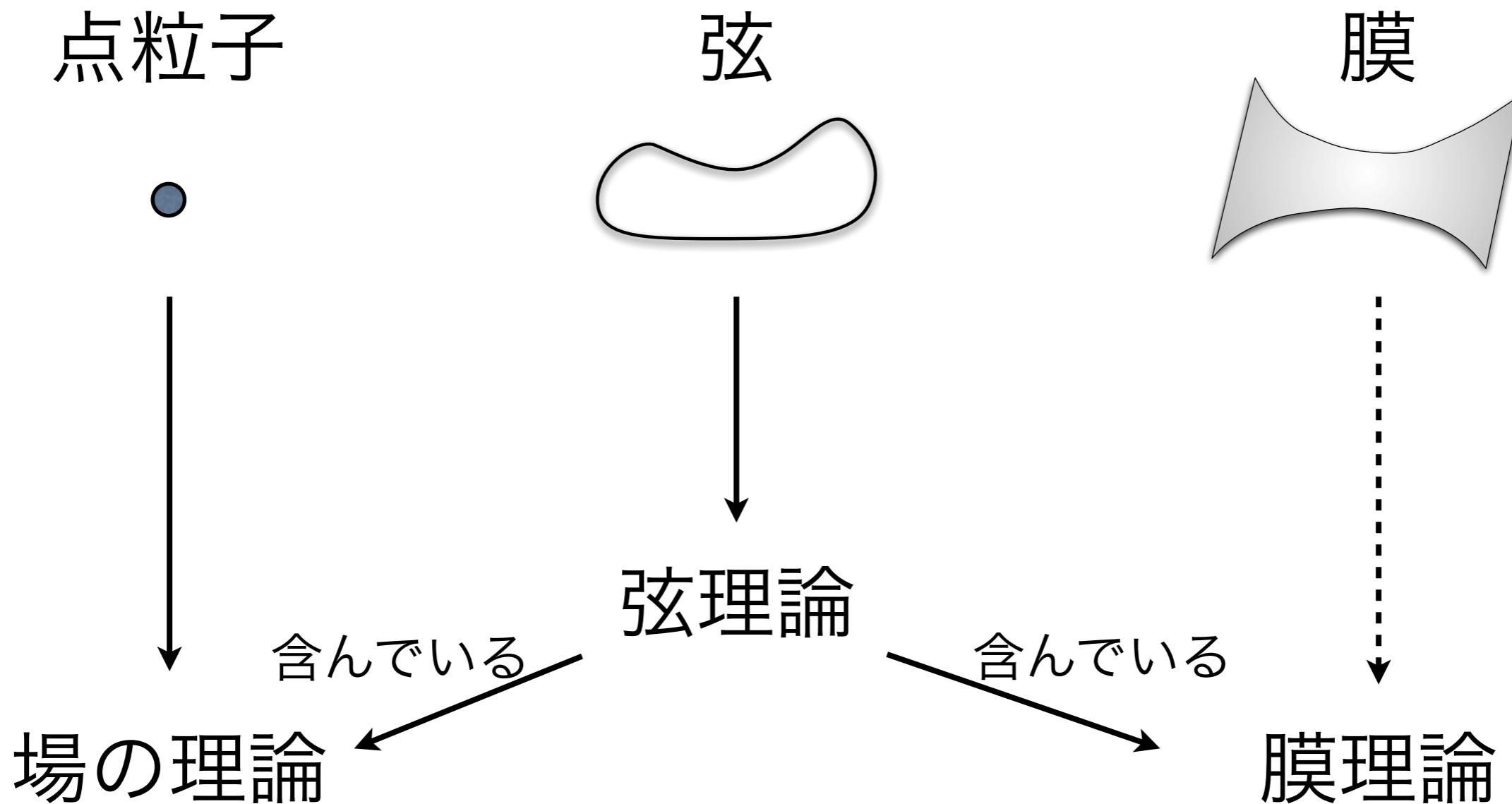
Polchinski



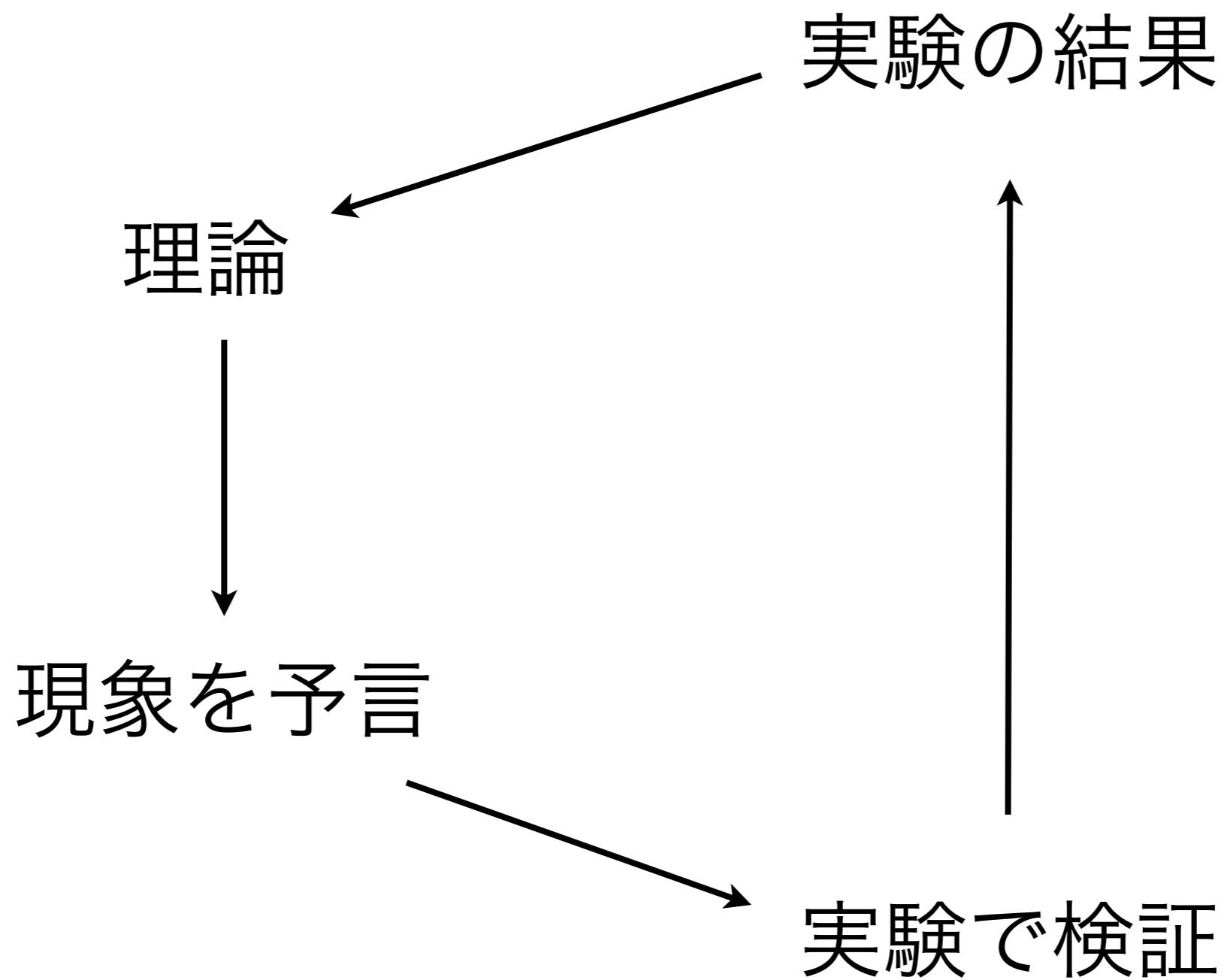
Maldacena

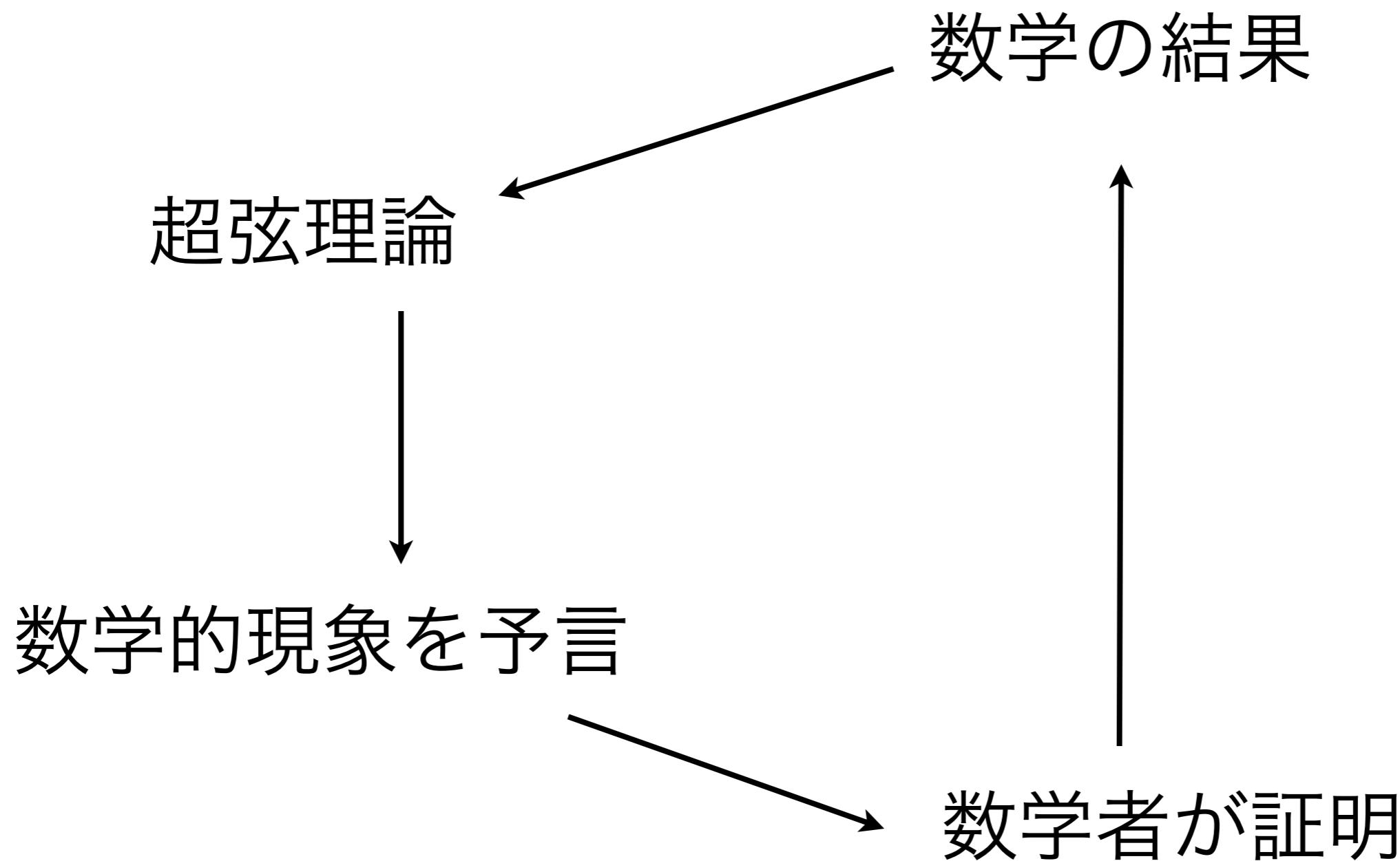
- (超)弦理論
重力を量子化できる。
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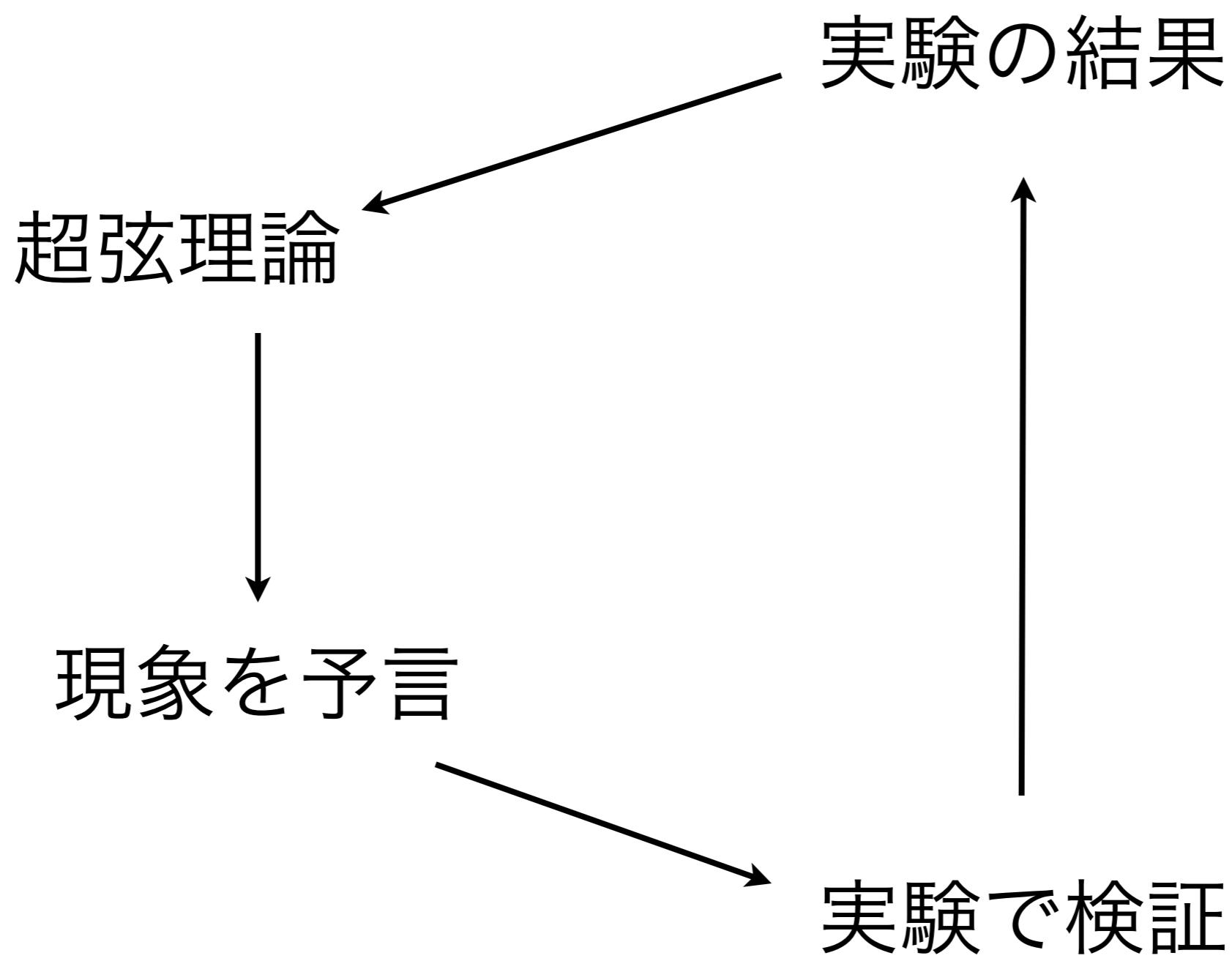
- 2000年ごろには枠組みはほぼ完成。



- 超弦理論は自動的に量子重力を含む !!
- 残念ながら、低エネルギー (TeV)
では場の量子論とほとんど変わらない...
- (でも数学には沢山応用がある)

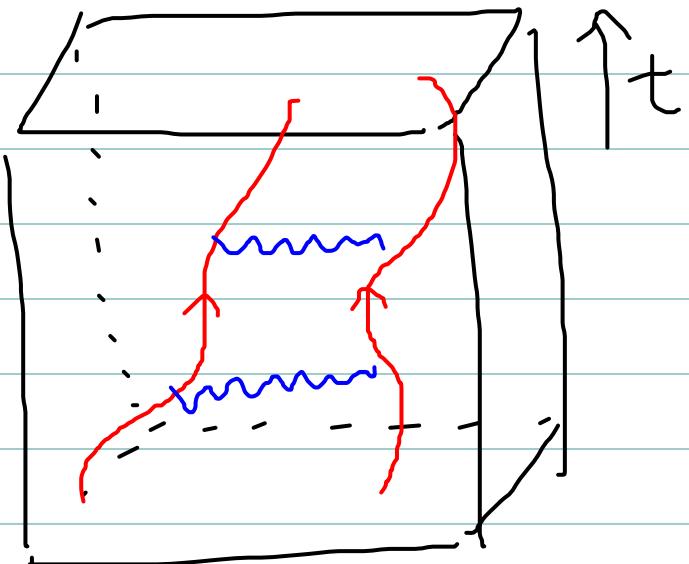






と行きたいものです...

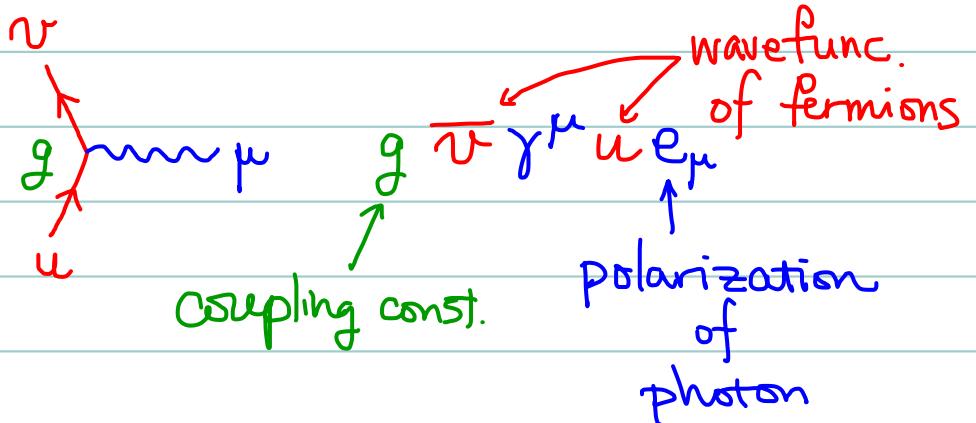
- 今日
場の理論のおさらい
粒子じゃなくて弦にすると？
- 明日午前
弦理論は膜理論を含んでいる
- 明日午後
弦理論から標準模型を出すには？



electron
photon

Need to sum over all possible paths of particles.

$x_{1,2,3}$

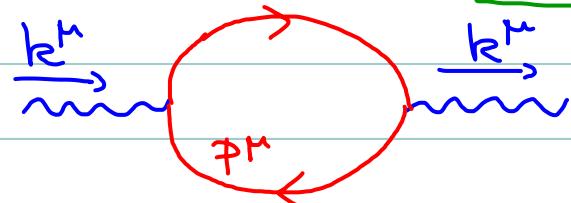


What was the problem back then?

* Compatibility of Lorentz invariance
& Positivity of probability.

~~> Gauge principle

* Treatment of infinities ~~> Renormalization



$$\sim \int d^4 p \frac{k_\mu k_\nu}{p^4} \sim \log(\text{cutoff.}) + f(|k|) \underset{\text{measurable}}{\cancel{}}$$

Counterterm

A horizontal wavy blue line representing a propagator. A red 'X' is placed on the line, indicating a subtraction point. A blue arrow points from this subtraction point down to the text $\sim -\log(\text{cutoff.})$, with the word "cancel" written next to it, indicating the cancellation of the divergent part of the loop diagram.

$$\sim -\log(\text{cutoff.})$$

Positivity of Probability ~ Unitarity

State $|\Psi\rangle = c_1|\uparrow\rangle + c_2|\downarrow\rangle$, $\langle\Psi|\Psi\rangle = \langle\uparrow|\uparrow\rangle = \langle\downarrow|\downarrow\rangle = 1$

$\sim |\uparrow\rangle$ "happens" with prob. $|c_1|^2$, $|\downarrow\rangle$ with prob. $|c_2|^2$.

WE NEED $\langle\Psi|\Psi\rangle \geq 0$ for all $|\Psi\rangle$.

Let $|0\rangle$: vacuum. $\langle 0|0\rangle = 1$.

Create a photon

$$A^\mu|0\rangle$$

$$\langle 0|A^{\mu\dagger} A^\nu|0\rangle \sim -\eta^{\mu\nu} = \begin{pmatrix} -1 & + & + & + \\ + & + & + & + \end{pmatrix}$$

Lorentz inv.

$$|\psi\rangle = c_t |e_t\rangle + c_x |e_x\rangle + c_y |e_y\rangle + c_z |e_z\rangle$$


 $e_t \ e_x \ e_y \ e_z$
 $\uparrow P^M = (E, E, 0, 0)$

$$\langle \psi | \psi \rangle = \underbrace{|c_t|^2 + |c_x|^2}_{\text{red}} + |c_y|^2 + |c_z|^2$$

If always $c_t = c_x$

$$\langle \psi | \psi \rangle = |c_y|^2 + |c_z|^2$$

$\rightsquigarrow |e_y\rangle$ "happens" $\frac{|c_y|^2}{|c_y|^2 + |c_z|^2}$, $|e_z\rangle$ with $\frac{|c_z|^2}{|c_y|^2 + |c_z|^2}$.

$$|\psi\rangle = c_t (\underbrace{|e_t\rangle + |e_x\rangle}_{\text{unphysical mode}}) + c_y |e_y\rangle + c_z |e_z\rangle$$

Mode with polarization $\sim p_\mu$ is unphysical.

$A_\mu(x)$ and $A_\mu(x) + \partial_\mu \lambda(x)$
 are physically THE SAME.

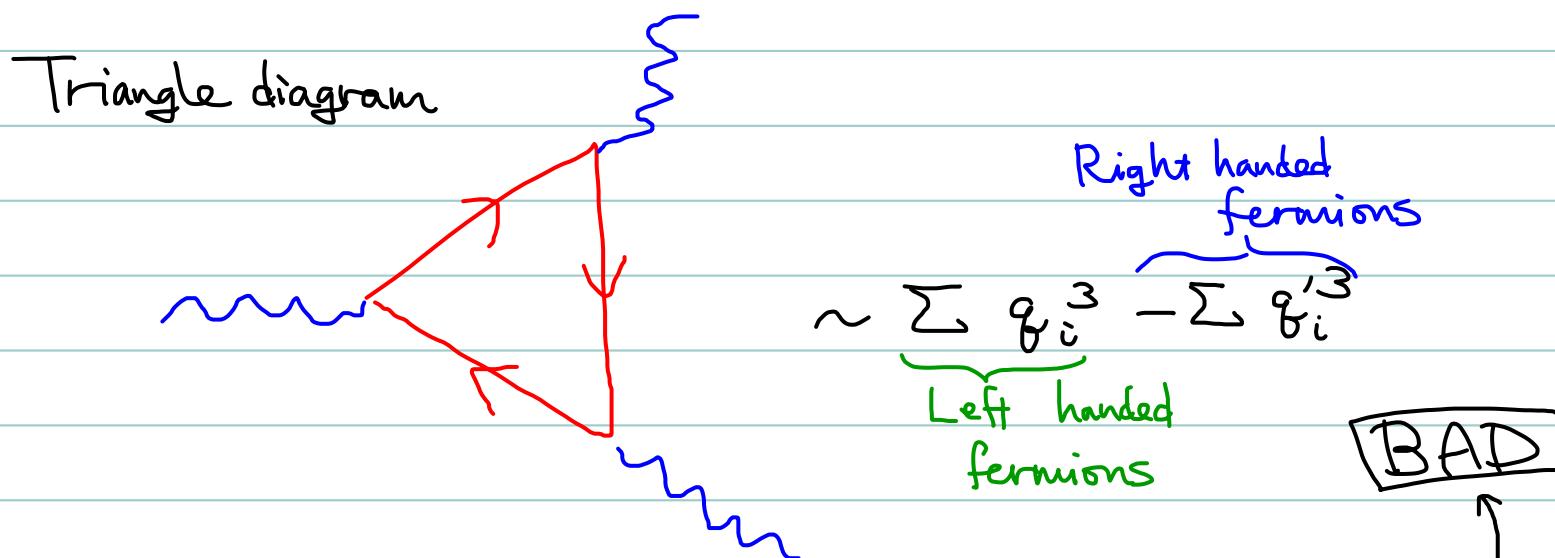
Gauge
Symmetry

This settled the problem for 20 years.

But it was realized that weak interaction is chiral.

$$\begin{array}{c} \mu \\ \text{---} \\ \text{h}\psi_R \end{array} \xrightarrow{g} \begin{array}{c} \mu \\ \text{---} \\ \text{h}\psi_L \end{array}$$
$$\gamma^5 |\psi_L\rangle = + |\psi_L\rangle$$

Triangle diagram



$$\sim \underbrace{\sum g_i^3}_{\text{Left handed fermions}} - \underbrace{\sum g'_i{}^3}_{\text{Right handed fermions}}$$

BAD

produces modes with $C_t + C_x$! \rightsquigarrow Negative Prob.

Summary

* For vector bosons, it's hard to have
Unitarity & Lorentz Invariance.

* Gauge symmetry saves the day.

* If chiral, it's necessary that

$$\sum g_L^3 - \sum g_R^3 = 0 \quad \dots (\#)$$

called **ANOMALY FREE CONDITION.**

Question

the Standard Model is chiral.

Is (#) satisfied?

	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$
(u_L, d_L)	3	2	1/3
\overline{u}_R	3	1	-4/3
\overline{d}_R	3	1	2/3
(e_L, ν_L)	1	2	-1
\overline{e}_R	1	1	+2

$$\sum g_L^3 - \sum g_R^3 = 3 \times 2 \times \frac{1}{27} + 3 \times 1 \times \frac{-64}{27} + 3 \times 1 \times \frac{8}{27} \\ + 1 \times 2 \times (-1)^3 + 1 \times 1 \times 2^3$$

$$= \frac{1}{9}(2 - 64 + 8 - 18 + 72)$$

$$= 0 !!$$

That looks totally accidental.

GUT:

$$SU(3) \times SU(2) \times U(1)$$



$$SU(5)$$



$$SO(10)$$

one generation

$$5 + 10$$

anomaly is
cancelled
between them.

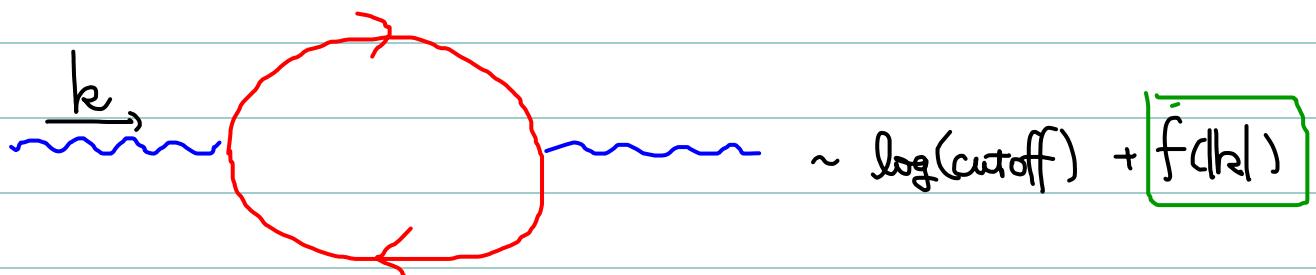
$$16$$

no anomaly!

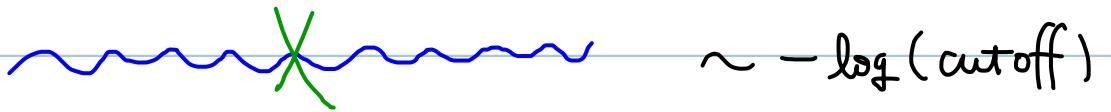
$SO(10)$ doesn't have anomalies.

But why are there 3 generations?

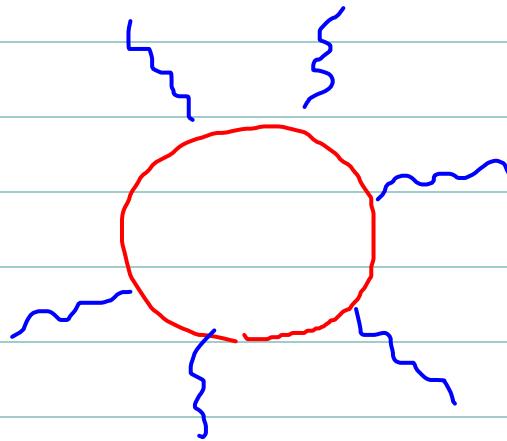
Renormalizability



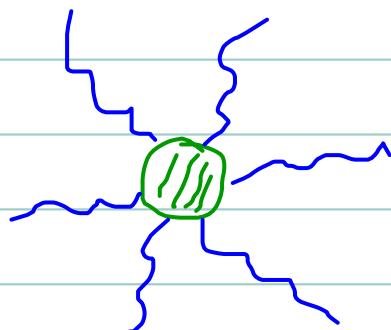
Counter term



How about



do we need counterterm for this?



For gauge theories, NO.

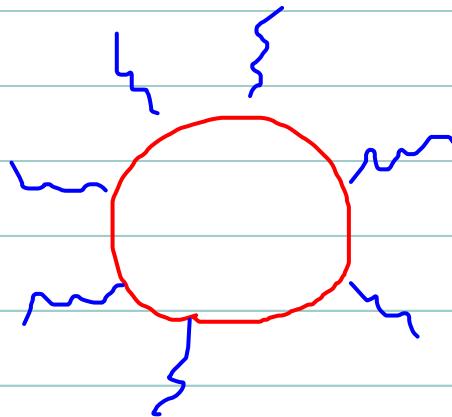
$$\begin{array}{c} \text{---} \\ \text{---} \end{array} \quad \sim \text{const.}$$

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \sim \int d^4 p \frac{1}{(p)^{\# \text{ 足数}}} \sim \text{finite.}$$

Only a few types of counterterms are necessary.

For gravity, it's much worse.

$$\begin{array}{c} \downarrow k_\mu \\ \xrightarrow{\hspace{2cm}} \\ \xrightarrow{\hspace{1cm}} p_\mu \end{array}$$
$$\sim p_\mu k_\nu$$



$$\sim \int d^4 p \left(\frac{p \cdot k}{p'} \right)^{\# \text{ 路数}} \sim \text{infinite!}$$

Need infinite types of counterterms. **NONRENORMALIZABLE.**

ASIDE Do we really need to quantize gravity ?

Is there anything wrong with

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \langle T_{\mu\nu} \rangle$$

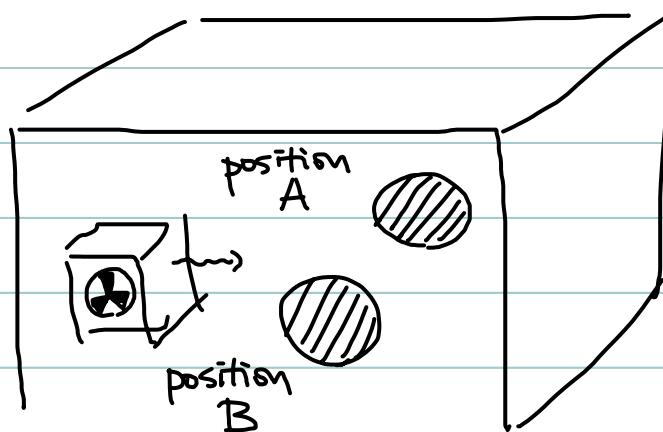
Classical equation Quantum mechanical average

$$\Delta\phi(x) = \langle \rho(x) \rangle$$

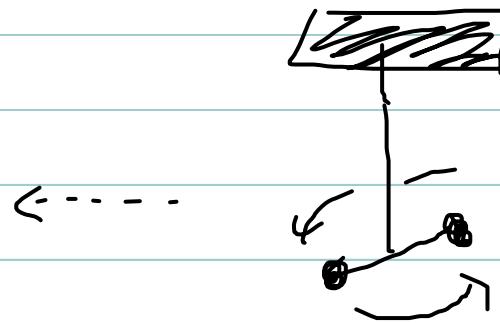
But this leads to various absurdities.

Page & Geilker, PRL 47 (1981) 979

Consider the following Schrödinger's cat-like setup



Position of a heavy obj.
is chosen QM'ally.



Measure its
position using torsion
balance from the outside.

If $\Delta\phi(x) = \langle \rho(x) \rangle$, the $\langle \rho(x) \rangle$

is the average value of the two positions of the heavy object.

→ $\phi(x)$ measures this.

But of course the torsion balance sees inside the box

$\phi(x)$ is either for position A or B, not the average.

Conclusion It's too naive to say

$$\underbrace{R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R}_{\text{classical}} = \underbrace{\langle T_{\mu\nu} \rangle}_{\text{Quantum}}$$

Summary

- * Fermion content of the SM is fascinating.
- * Quantization of gravity is tough in the standard QFT.

So, what happens if we consider

Strings

instead of

particles ?

What changes if string is used ?

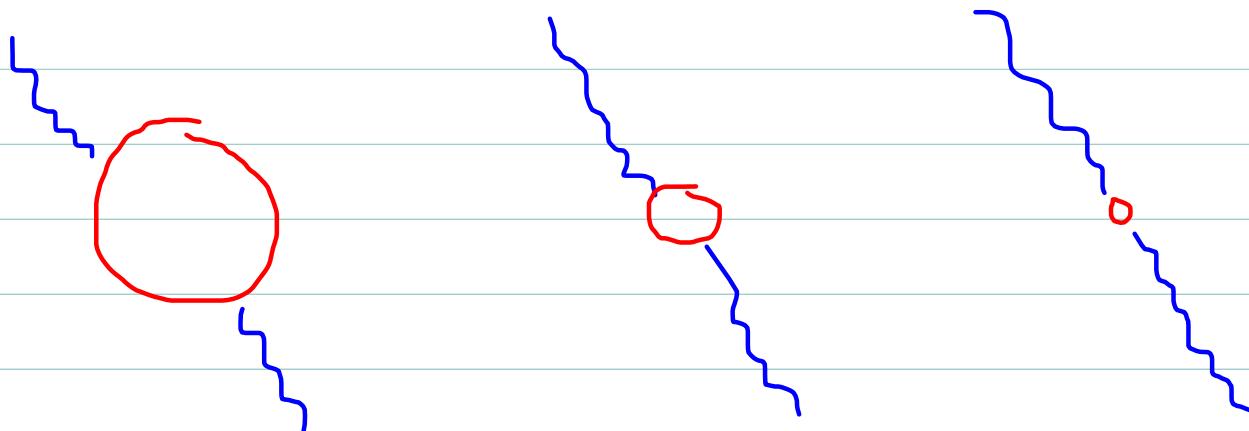
It has a length scale l_s s.t.



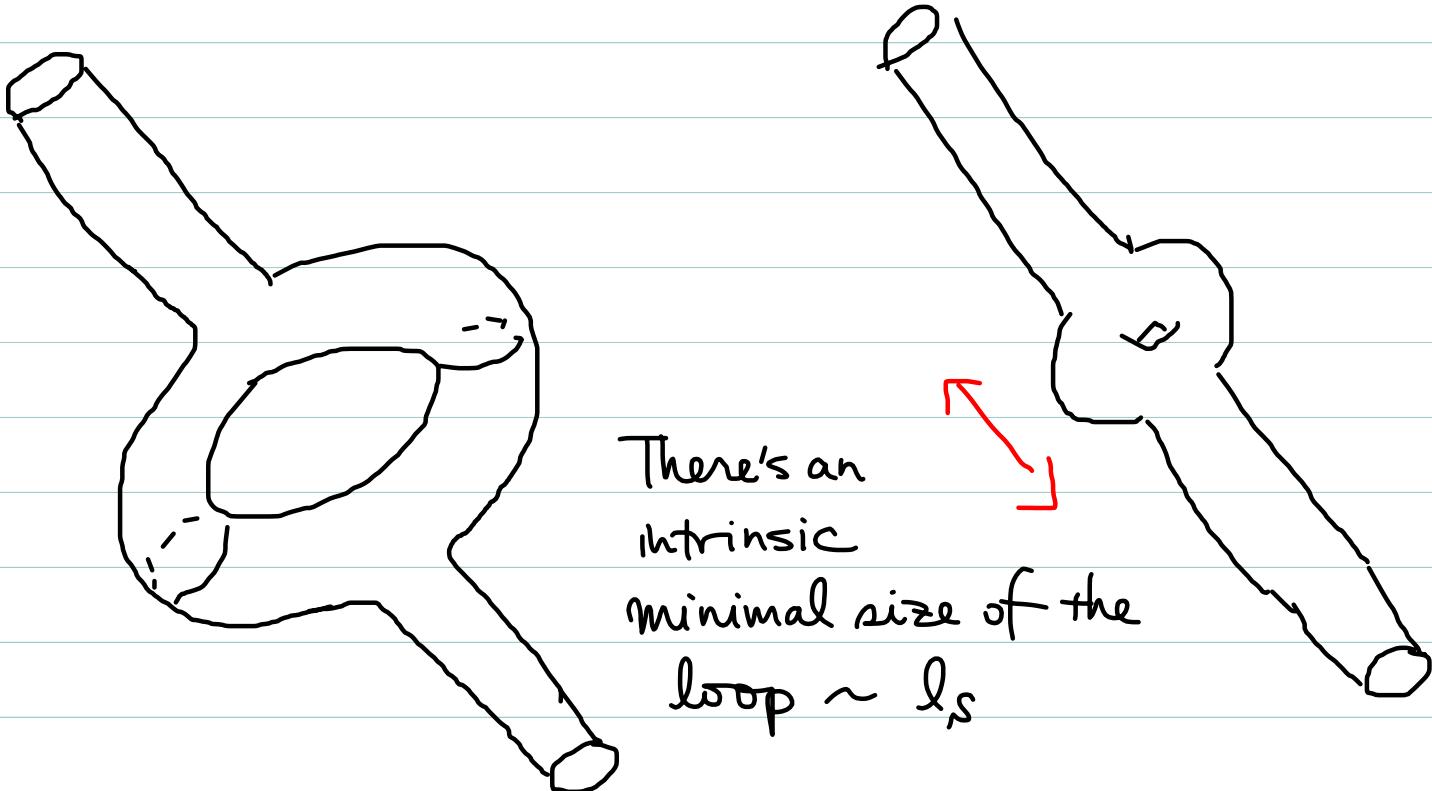
has mass L/l_s^2

The field theory divergence from UV

$\int d^4 p \frac{1}{p^2}$, say, comes from



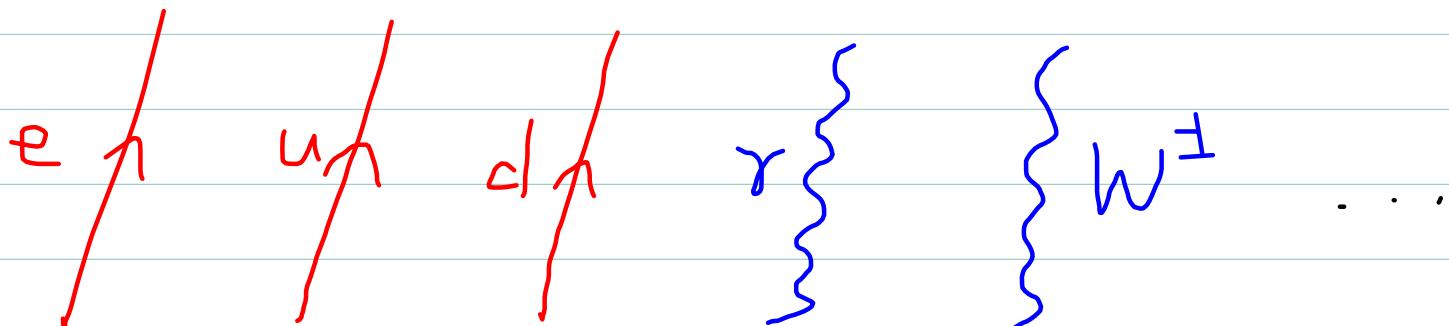
In string theory,



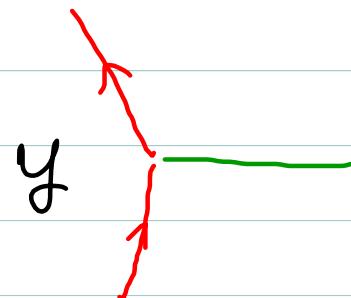
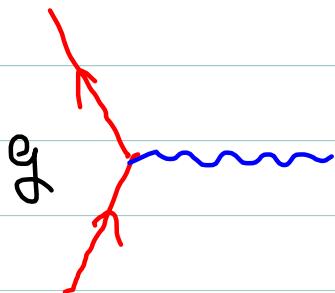
→ Every loop diagram is cut off around
 $|P| \sim 1/l_s$

Another merit: in QFT, we start by specifying the
Lagrangian \sim Gauge fields
Fermions
Scalars, Yukawa couplings ...

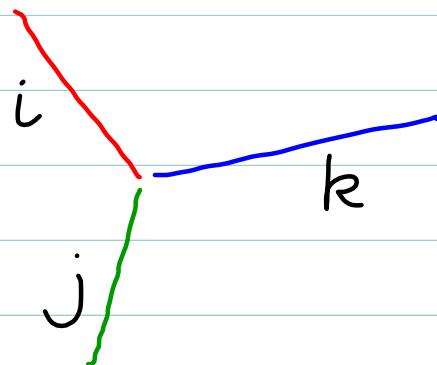
In terms of Feynman diagrams, we have different types of
lines



and various interaction vertices

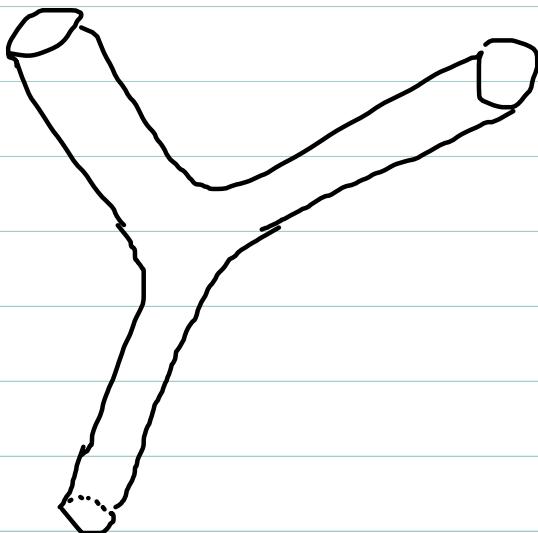


Note g & γ needs to be chosen for each vertex.

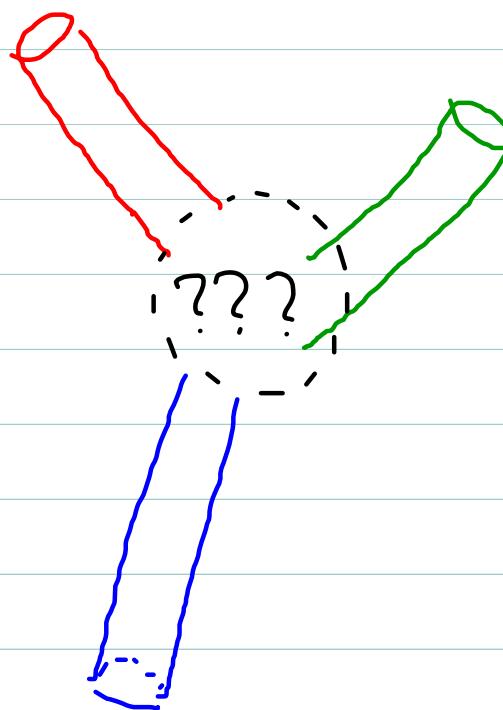


$$\sim \lambda_{ijk}$$

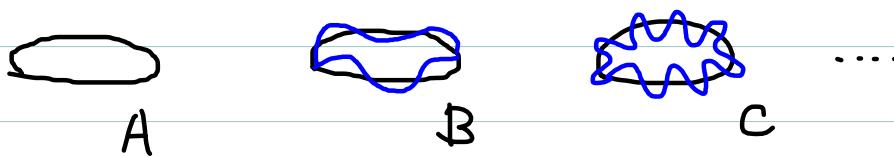
In string theory, you can't join two different types of strings consistently :



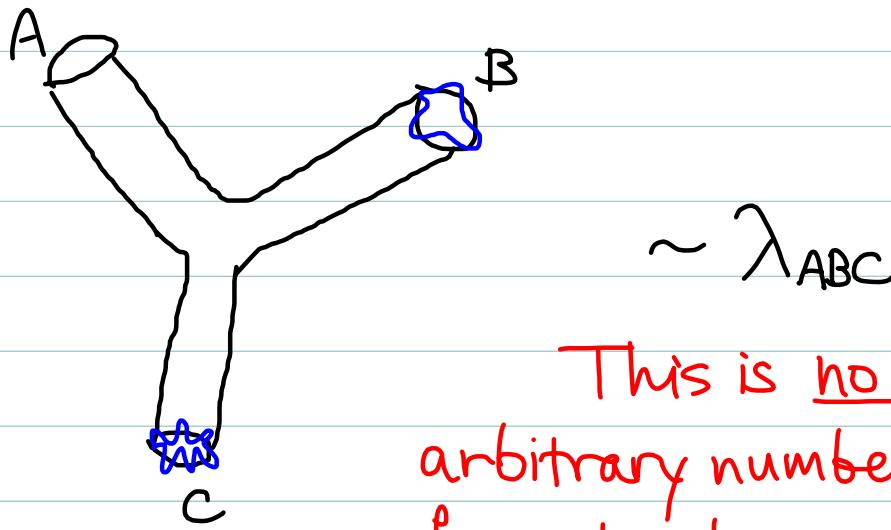
but



Instead, a string can vibrate in different ways:

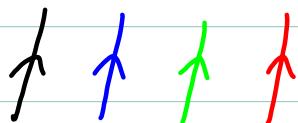


which correspond to various different particles

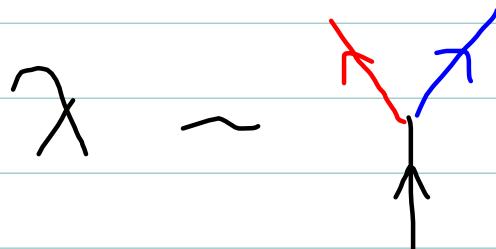


This is no longer an arbitrary numbers. Calculable from the dynamics of strings.

QFT



choose types of particles in a theory.



STRING

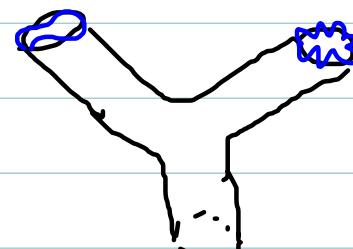


choose a type of string.

study how they vibrate.



interaction is then determined.



There are **two** known consistent types of strings.



Type II string



Heterotic string

Both require $1+9$ dimensional spacetime.
have supersymmetry.
contains quantized gravity.

All these follow from the consistency of
Lorentz inv. & Positivity of probability.