

# Infrared behaviour of Duality cascades and warped throats

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Based on:

S. Franco, A. Hanany, A.M.U, hep-th/0502113

J.F.G. Cascales, F. Saad, A.M.U, hep-th/0503079

S. Franco, A. Hanany, F. Saad, A.M.U, hep-th/0505040

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## Motivation

- AdS/CFT correspondence between quiver gauge theory on D3-branes at CY singularity  $X_6$  and IIB theory on  $\text{AdS}_5 \times Y_5$
- Going non-conformal by addition of fractional branes in the quiver side and 3-form fluxes on 3-cycles of  $Y_5$  on the IIB side
  - Warped throats and cascading gauge theories provide non-trivial examples of holography for non-conformal cascading and IR confining gauge theories. [Klebanov, Strassler]
- Story is well-understood for conifold, where confinement of gauge theory related to complex deformation of the throat
- Cascades and throats with fluxes are known for other singularities
  - [Franco, He, Herzog, Walcher; Ejaz, Herzog, Klebanov]
  - Extend IR understanding to other singularities

## The conifold case

[Klebanov, Strassler]

- The gauge theory  $SU(N) \times SU(N + M)$  suffers a cascade of Seiberg dualities as the RG scale runs to the infrared

$$SU(N) \times SU(N + M) \rightarrow SU(N) \times SU(N - M) \rightarrow \\ \rightarrow SU(N - 2M) \times SU(N - M) \rightarrow \dots$$

- For  $N = KM$ , this proceeds for  $K$  steps, the left over IR theory is  $SU(M)$  with no charged matter, confines and develops a gaugino condensate.

**M**



- Probe the dual geometry by using e.g.  $M$  additional D3-brane probes

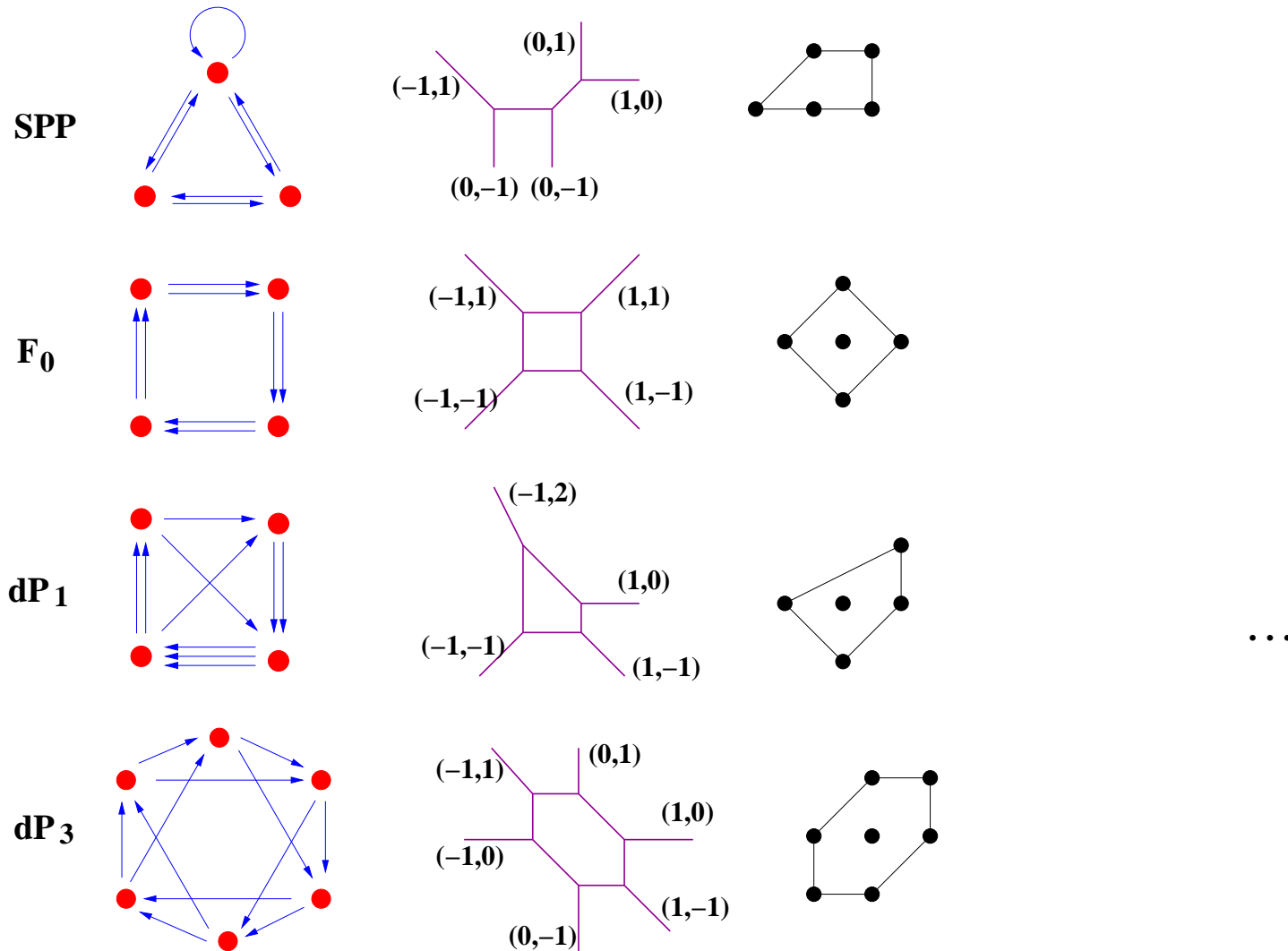


For  $N_f = N_c$  there is a quantum deformed moduli space. D3-brane probe moduli space corresponds to mesonic branch

$$\det \mathcal{M} = M_{11}M_{22} - M_{12}M_{21} = \epsilon \rightarrow \text{Deformed conifold}$$

# Toric singularities

- Diverse techniques to obtain quiver for toric singularities  
Partial resolution, mirror symmetry, exceptional collections, ...
- Some examples of quivers and toric/web diagrams



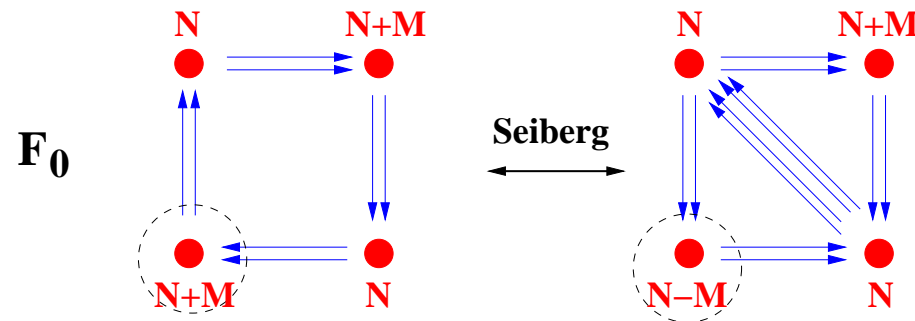
# Seiberg duality



- For a single node in a quiver gauge theory: Reverse arrows, introduce mesons, and take into account superpotential

[Beasley, Plesser; Feng, Hanany, He, A.U.]

- Example of  $F_0$



- Cascades of Seiberg dualities for several quiver gauge theories with fractional branes  $\rightarrow F_0$ , SPP,  $dP_n$ 's,  $Y^{p,q}$ , ...

[Franco, He, Herzog, Walcher; Ejaz, Herzog, Klebanov; Franco, Hanany, A.U;...]

- Supported by singular warped conical throats (of Klebanov-Tseytlin type) for  $dP_n$ 's,  $Y^{p,q}$  [Franco, He, Herzog, Walcher; Ejaz, Herzog, Klebanov]

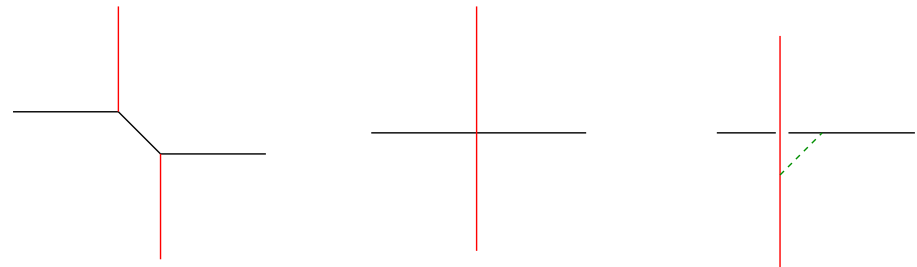
# Complex deformations

- Complex deformations are easily identified in web diagram as recombinations of external legs into subwebs

→ Higgs branch of 5d theory [Seiberg, Morrison; Aharony, Hanany, Kol]

→ Decomposition of toric polygon as Minkowski sum [Altmann]

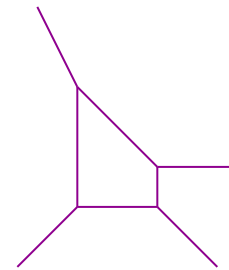
- Familiar example: Conifold



- But other examples too! Ex: SPP



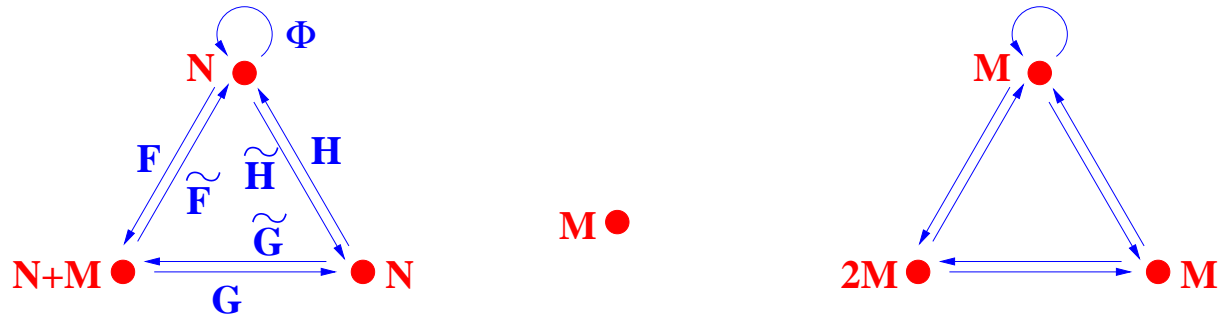
- Also cases with no deformation:  $Y^{p,q}$



- Result: One to one map with gauge theory behaviour!

## The SPP example

- Duality cascade triggered by fractional brane, ending in confining theory.
- Web suggests complex deformation. Probe using additional D3's



- Deformation of mesonic branch

$$\det \mathcal{M} = \epsilon \quad , \quad \mathcal{M} = \begin{pmatrix} FG & F\tilde{F} \\ \tilde{G}G & \tilde{G}\tilde{F} \end{pmatrix}$$

Node  $2M$  confines, vevs for  $\mathcal{M} \propto \mathbf{1}_{2M}$  break to  $N = 4$   $U(N)$

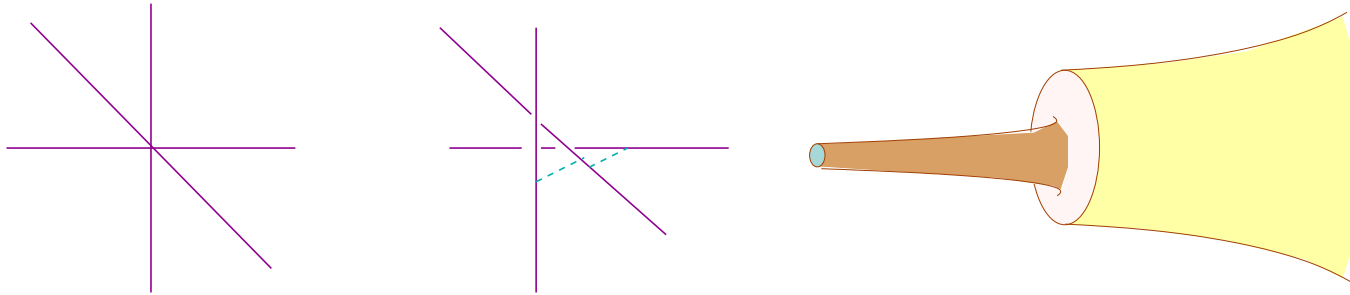
- Moduli space deformed to smooth

$$xy - zw^2 = 0 \quad \rightarrow \quad xy - zw^2 = \epsilon w$$

→ Complex deformation to smooth geometry

## The cone over $dP_3$

- The geometry has a two-dimensional branch of complex deformations

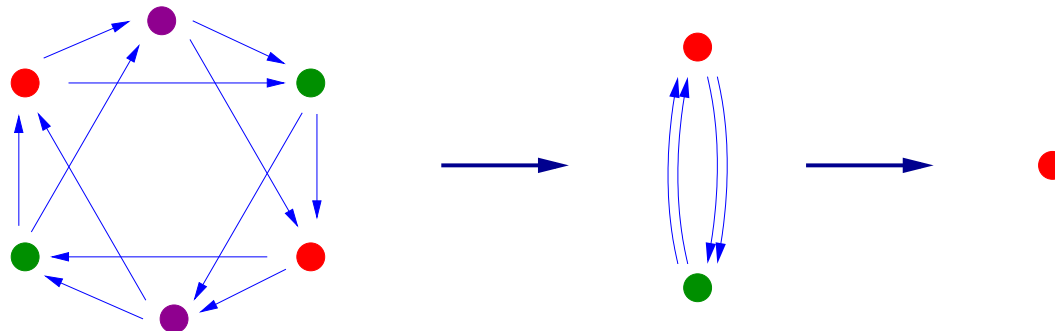


- In the regime of hierarchically different sizes, sequential deformation:  
cone over  $dP_3 \rightarrow$  conifold  $\rightarrow$  smooth

- Suggests two stages of infrared partial confinement

- Full agreement with gauge theory pattern:

Duality cascade  $\rightarrow$  Partial confinement. to conifold th.  $\rightarrow$  Subsequent cascade  $\rightarrow$  Confinement





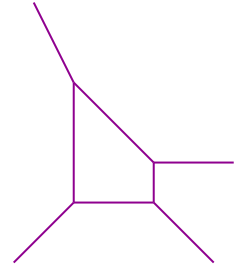
## Cases with no deformation

- Many geometries do not admit complex deformations

For instance,  $Y_{p,q}$  (first order deformation in

[Burrington, Liu, Mahato, Pando-Zayas]

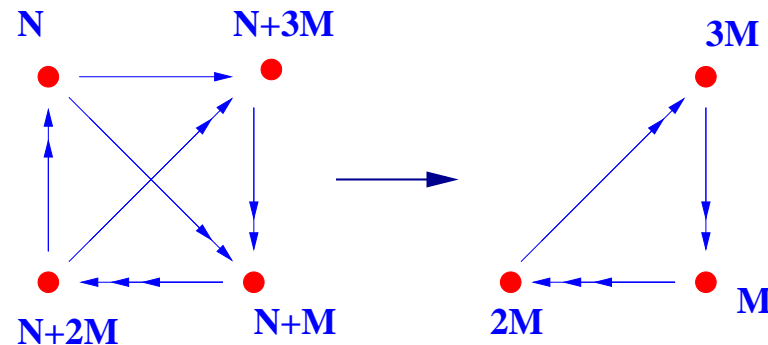
is obstructed at second order)



- But UV cascades, and dual singular throats with fluxes exist.

What is the IR behaviour? Smooth out the supergravity singularity?

- Analyze using gauge theory! [Berenstein, Herzog, Ouyang, Pinansky; Franco, Hanany, Saad, A.U; Bertolini, Bigazzi, Cotrone]



→ No vacuum  $N_f < N_c$  ADS F-term pushes vevs to infinity

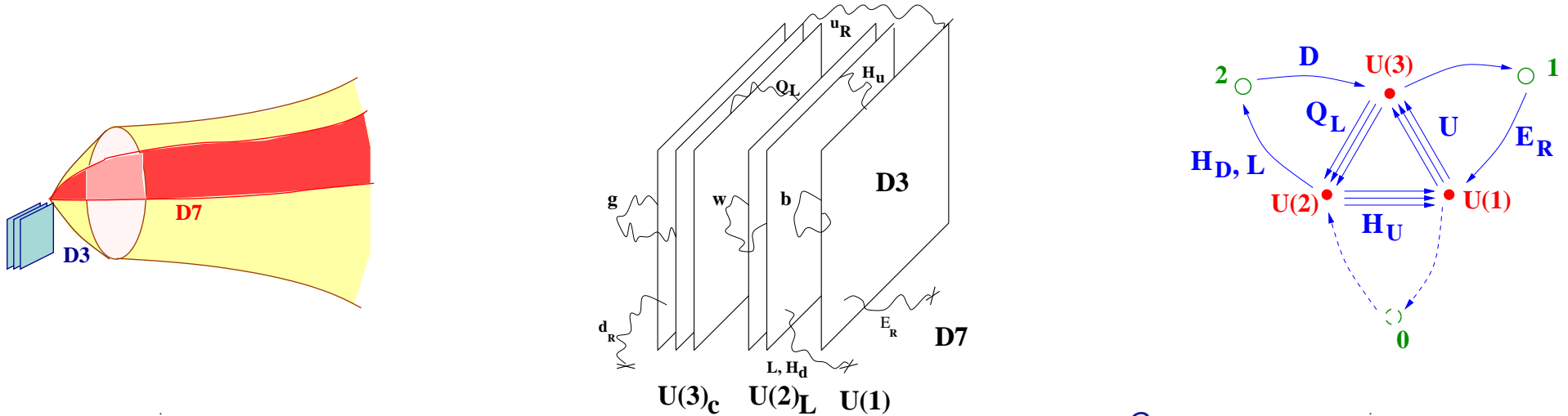
(D-flatness does not forbid it: absence of  $U(1)$ 's or dynamical FI's)

- Suggests no smoothing of supergravity throat

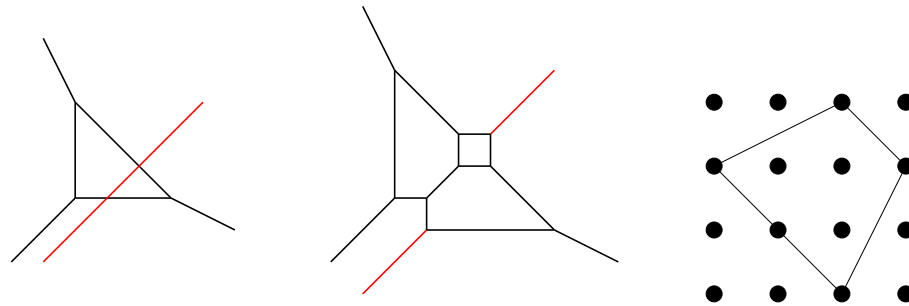
## The SM on the throat

- Model building application:

SM on D-branes at  $C^3/Z_3$  at the end of a throat [Cascales, Saad, A.U.]



- Use a geometry which admits a deformation to  $C^3/Z_3$

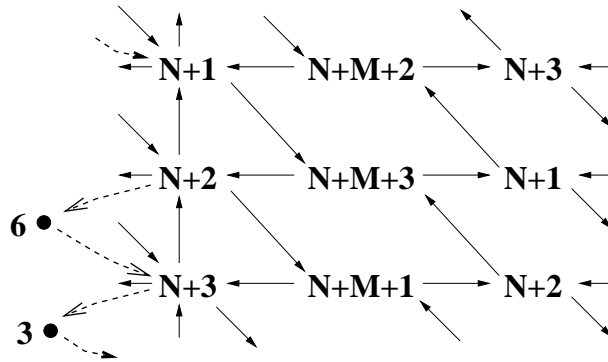


→ Orbifold of the SPP

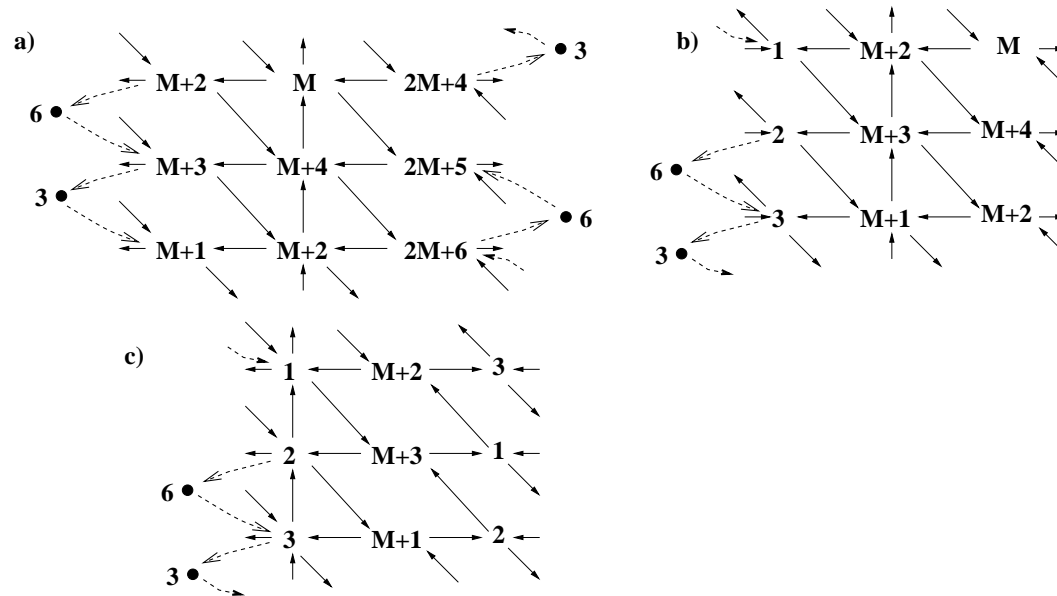
- Quiver is known: **Holographic dual of the SM on the throat**

→ **Peculiar technicolor:** Cascade of Seiberg dualities, ending in partial confinement leaving the SM degrees of freedom

- Quiver of the UV theory at a particular point in the Seiberg duality cascade



- Quivers of the theory at the last stages of cascading



# Conclusions

- Branes at singularities offer a tractable way to study the interplay between gauge and string theory dynamics
- We have reviewed the correspondence between duality cascades and warped throats, with an emphasis on
  - Complex deformations and IR confinement
  - Absence of complex deformations and no vacuum
- Progress in understanding the structure of throats / cascades with richer geometry / dynamics at its bottom / infrared
  - Terminal singularities, multiwarp throats, ...
- Model building application: Holographic dual of the Standard Model on the throat
- Hopefully, many other insights and answers to come

# Strings 07



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## Appendix: Sample computation: $dP_3 \rightarrow \text{conifold}$

The superpotential reads

$$\begin{aligned} W = & X_{12}X_{23}X_{34}X_{45}X_{56}X_{61} + X_{13}X_{35}X_{51} + X_{24}X_{46}X_{62} - \\ & - X_{23}X_{35}X_{56}X_{62} - X_{13}X_{34}X_{46}X_{61} - X_{12}X_{24}X_{45}X_{51} \end{aligned}$$

The  $SU(2M)$  nodes 2 and 4 condense, we have mesons

$$\begin{aligned} \mathcal{M} &= \begin{bmatrix} M_{63} & M_{62} \\ M_{53} & M_{52} \end{bmatrix} = \begin{bmatrix} X_{61}X_{13} & X_{61}X_{12} \\ X_{51}X_{13} & X_{51}X_{12} \end{bmatrix} \\ \mathcal{N} &= \begin{bmatrix} N_{36} & N_{35} \\ N_{26} & N_{25} \end{bmatrix} = \begin{bmatrix} X_{34}X_{46} & X_{34}X_{45} \\ X_{24}X_{46} & X_{24}X_{45} \end{bmatrix} \end{aligned}$$

and baryons  $\mathcal{B}$ ,  $\tilde{\mathcal{B}}$ ,  $\mathcal{A}$ ,  $\tilde{\mathcal{A}}$ , with

$$\det \mathcal{M} - \mathcal{B}\tilde{\mathcal{B}} = \Lambda^{4M} \quad ; \quad \det \mathcal{N} - \mathcal{A}\tilde{\mathcal{A}} = \Lambda^{4M}$$

The superpotential reads

$$\begin{aligned} W = & M_{62}X_{23}N_{35}X_{56} + M_{53}X_{35} + N_{26}X_{62} - \\ & - X_{23}X_{35}X_{56}X_{62} - M_{63}N_{36} - M_{52}N_{25} + \\ & + X_1 (\det \mathcal{M} - \mathcal{B}\tilde{\mathcal{B}} - \Lambda^{4M}) + X_2 (\det \mathcal{N} - \mathcal{A}\tilde{\mathcal{A}} - \Lambda^{4M}) \end{aligned}$$

Along the mesonic branch, in the abelian case

$$\begin{aligned} W = & M_{62}X_{23}N_{35}X_{56} - X_{23}X_{35}X_{56}X_{62} - M_{63}N_{36} - M_{52}N_{25} + \\ & + M_{53}X_{35} + N_{26}X_{62} + M_{63}M_{52} - M_{53}M_{62} + N_{36}N_{25} - N_{26}N_{35} \end{aligned}$$

Using the equations of motion

$$W = X_{23}N_{35}X_{56}M_{62} - X_{23}M_{62}X_{56}N_{35}$$

Going non-Abelian, the gauge group is  $SU(M)_{25} \times SU(M)_{36}$ .

The conifold theory