

Exactly Solvable Dynamics for Rolling D-brane

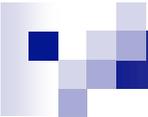
Yu Nakayama (Tokyo univ.)

**Based on :hep-th/0605013 with S.J. Rey, Y.
Sugawara,**

**hep-th/0507040 with S.J. R, Y. S,
hep-th/0412038 with K. Panigrahi, S.J. R,
and**

H. Takayanagi

hep-th/0406173 with Y. S, H.T



Purpose of the talk

- **Large charge** (+BPS) vs. **Small Charge** (+non-BPS)
 - Black hole / String phase transition
 - Hawking temperature vs. Hagedorn temperature
 - Is 2D (pure) Black hole really black?
- **Analyticity** vs. **Non-analyticity**
 - Universality of Tachyon-Radion correspondence
 - Wick rotation in curved space
- **Unitarity** vs. **Open/Closed duality**
 - Optical theorem
 - Lorentzian world-sheet vs. Euclidean world-sheet

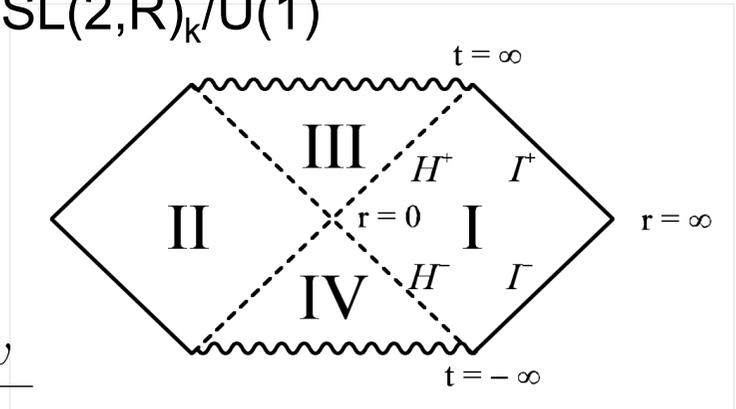
2D (fermionic) Black Hole

- 2D black hole is the simplest black hole geometry as an **exact string background** (Witten; Mandal, Sengupta, Wadia) $SL(2, \mathbb{R})_k/U(1)$

$$ds^2 = 2k(d\rho^2 - \tanh^2 \rho dt^2) \quad e^\Phi = \frac{e^{\Phi_0}}{\cosh \rho}$$

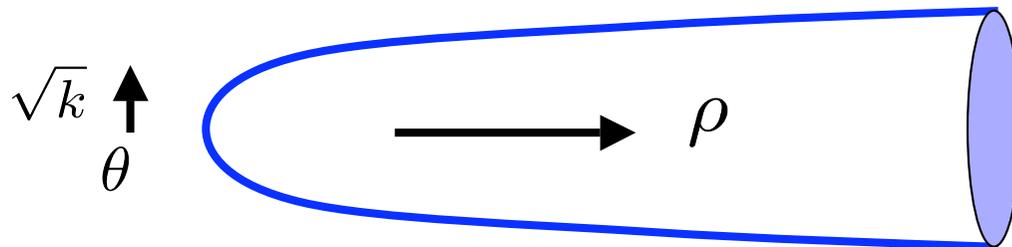
- Global metric looks **Schwarzschild-like**

$$u = \sinh \rho e^t \quad v = -\sinh \rho e^{-t} \quad ds^2 = -2k \frac{dudv}{1 - uv}$$



- In **Euclidean** geometry, 2D black hole is **cigar geometry**:

$$ds^2 = 2k(d\rho^2 + \tanh^2 \rho d\theta^2) \quad e^\Phi = \frac{e^{\Phi_0}}{\cosh \rho}$$



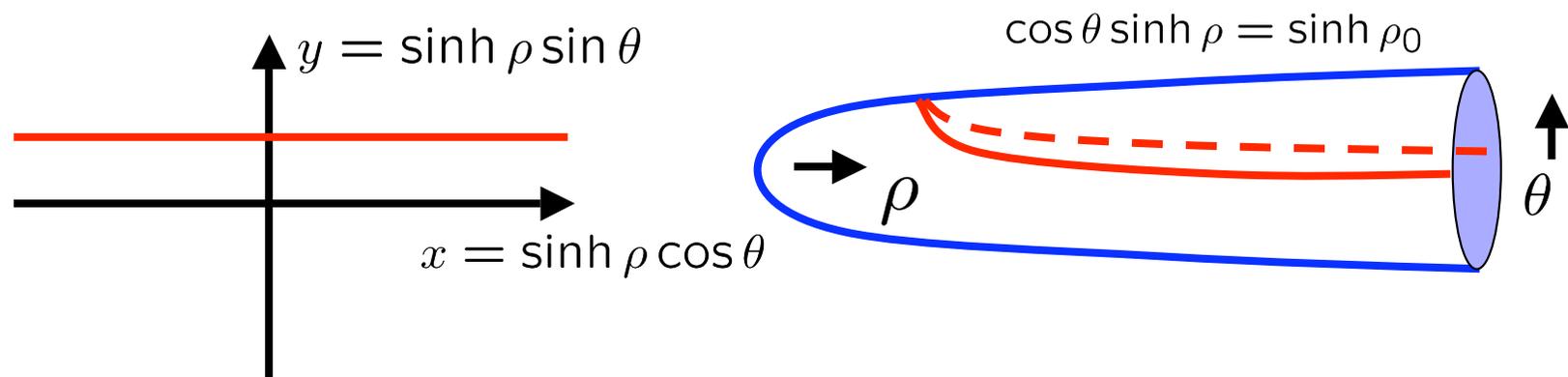
$$T_{Hw} = \frac{1}{\beta_{Hw}} = \frac{1}{2\pi\sqrt{2k}}$$

$$T_{Hg} = \frac{1}{\beta_{Hg}} = \frac{1}{4\pi\sqrt{1 - \frac{1}{2k}}}$$

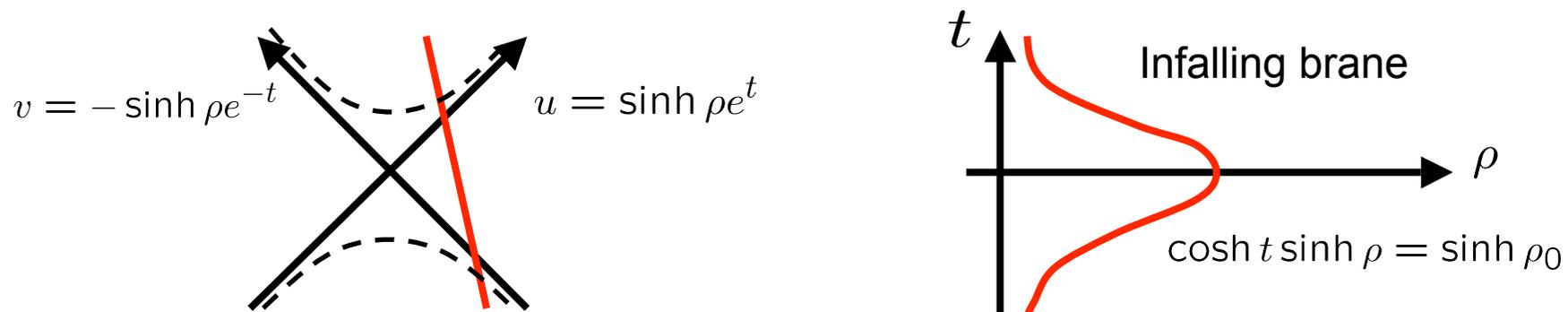
Branes in 2D Black Hole geometries

Classical D-branes are classified by solutions of **DBI action**.

- Euclidean case (Class 2' brane "hairpin brane": Ribault-Schomerus. see also Eguchi-Sugawara, Ahn-Stanishkov-Yamamoto, Lukyanov-Vitcev-Zamolodchikov)



- Lorentzian case ("Rolling D-brane": Kutasov, NST, Yogendran, NRS)



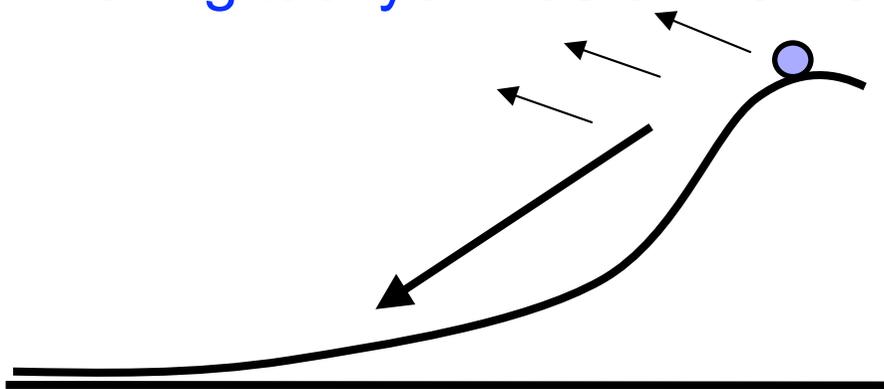
Tachyon-Radion correspondence

- D-brane near NS5-brane shows resemblance to **rolling tachyon** (Kutasov): **Rolling D-brane**

$$L_{D0} = -e^{-\Phi} \sqrt{\left(\frac{ds}{dt}\right)^2} = -V(X) \sqrt{1 - \dot{X}^2}$$

where $V(X) = M_0 e^{\sqrt{\frac{2}{k}}X}$.

- **Rolling tachyon** has similar form (Sen ...).



$$L_T = -V(T) \sqrt{1 - \dot{T}^2}$$

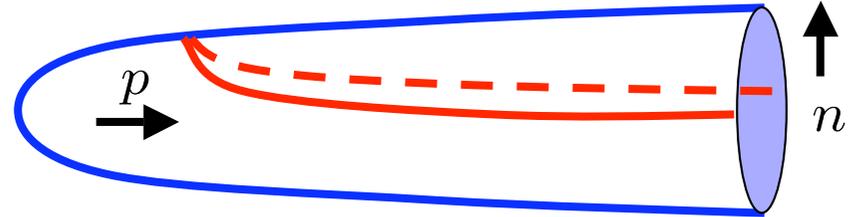
$$V(T) \sim \frac{1}{\cosh T}$$

Is tachyon-radion correspondence **universal**?
Artifact at the level of **effective action**?

Euclidean boundary states

- Class 2' boundary states in Euclidean BH

$$\cos \theta \sinh \rho = \sinh \rho_0$$



$$|B; \rho_0\rangle = \int_0^\infty \frac{dp}{2\pi} \sum_{n \in \mathbb{Z}} \Psi_{D1}(\rho_0; p, n) |p, n\rangle,$$

$$\Psi_{D1}(\rho_0; p, n) = \mathcal{N}(k) \frac{\Gamma(ip)\Gamma\left(1 + \frac{ip}{k}\right)}{\Gamma\left(\frac{1}{2} + \frac{ip+n}{2}\right)\Gamma\left(\frac{1}{2} + \frac{ip-n}{2}\right)} [e^{-ip\rho_0} + (-1)^n e^{ip\rho_0}].$$

- Effect of 1/k correction

- Delta function localized trajectory → smeared wavefunction

Poisson distribution:

$$\delta(\phi) \sim \sqrt{\frac{1}{k-1}}$$

- The steeper the hairpin, the wider the trajectory (NRPT).

Wick rotation: rolling D-brane boundary states

Naïve momentum space Wick rotation $n \rightarrow i\omega$ does **not** work.

- Performing Wick rotation in **coordinate space**, or choosing the **contour integral** properly,

$$\Psi(\rho_0; p, \omega) = \frac{\Gamma(\frac{1}{2} - i\frac{p+\omega}{2})\Gamma(\frac{1}{2} - i\frac{p-\omega}{2})\Gamma(1 + \frac{ip}{k})}{\Gamma(1 - ip)} \left[e^{-ip\rho_0} - \frac{\cosh\left(\pi\frac{p-\omega}{2}\right)}{\cosh\left(\pi\frac{p+\omega}{2}\right)} e^{ip\rho_0} \right]$$

$$|falling\rangle = \int \frac{d\omega}{2\pi} \frac{dp}{2\pi} \Psi(\rho_0; p, \omega) |U; p, \omega\rangle$$

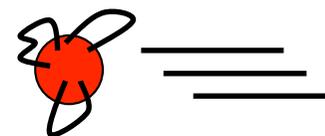
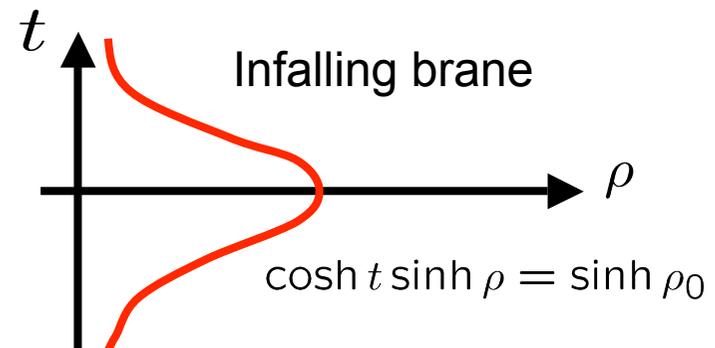
- Finite k correction:

- Trajectory is **smeared** (NPRT)

$$\delta(\phi) \sim \sqrt{\frac{1}{k-1}}$$

- Rolling D-brane gathers moss (Kutasov)

→ analytic continuation of **winding tachyon**?



D-brane radiation

From boundary states, we can **compute closed string emission** from falling D-branes (NST,NRS see also Sahakyan).

- Let us assume $k > 1$. For fixed mass level, $\omega^2(p, M) = p^2 + 2kM^2$

$$\begin{aligned} N(M) &\sim \int \frac{d^{5-p}\mathbf{k}_\perp}{(2\pi)^{5-p}} \int_0^\infty \frac{dp}{2\pi} \frac{1}{2\omega(p, M)} |\Psi(\rho_0; p, \omega(p, M))|^2 \\ &\sim \frac{1}{M} \int \frac{d^{5-p}\mathbf{k}_\perp}{(2\pi)^{5-p}} \int_0^\infty dp e^{\pi(1-\frac{1}{k})p - \pi\sqrt{p^2 + 2k(M^2 + \mathbf{k}_\perp^2)}} \\ &\sim M^{2-\frac{p}{2}} e^{-2\pi M \sqrt{1-\frac{1}{2k}}} = M^{2-\frac{p}{2}} e^{-2\pi M \frac{\beta_{Hg}}{2}} \end{aligned}$$

- **Saddle point approximation** is used as $M \rightarrow \infty$
- **Hagedorn temperature** (with α' correction) appeared in infalling mode!
- We are adding extra directions so that the theory is critical.

Tachyon-radion correspondence

- We can sum over **all the final states**

$$N = \sum_M N(M) \sim \int dM \rho(M) N(M) \sim \int \frac{dM}{M} M^{-p/2}$$

$$\rho(M) \sim \frac{1}{M^3} e^{2\pi M \sqrt{1 - \frac{1}{2k}}} \quad N(M) \sim M^{2 - \frac{p}{2}} e^{-2\pi M \sqrt{1 - \frac{1}{2k}}}$$

- **Density of states exactly cancels with the radiation density** → shows the same behavior in **rolling tachyon**
(Lambert-Liu-Maldacena)

Tachyon-radion correspondence is true at the **stringy level**.

- **Remarkable cancellation** of stringy corrections → **universal property** of rolling (falling) D-brane?

Black hole/ String transition at $k = 1$

Evaluation changes drastically at $k=1$ (BH/String transition)

- There is **no nontrivial saddle point** for $k < 1$

$$\begin{aligned} N(M) &\sim \int_0^\infty dp e^{\pi \left(1 - \frac{1}{k}\right) p - \pi \sqrt{p^2 + 2kM^2}} \\ &\sim e^{-2\pi M \sqrt{\frac{k}{2}}} \end{aligned}$$

- Emission rate is **UV convergent** (BH/String transition)
- Summary:
 - Tachyon-radion correspondence is **universal** ($k > 1$)
 - **BH/String transition** is observed in physical quantities (radiation rate)
 - At $k = 1$, $T_{Hw} = T_{Hg}$ and $k < 1$, BH interpretation is questionable.
 - **$1/k$ corrections are crucial.**

Unitarity and Open/Closed duality (NRS: hep-th/0605013)

- Is **unitarity** consistent with **open/closed duality**?

$$\begin{aligned} \text{Im} Z_{cylinder} &= \langle B | \text{---} \text{---} \text{---} | B \rangle \\ &= \sum_M \int \frac{d^{5-p} \mathbf{k}_\perp}{(2\pi)^{5-p}} \int_0^\infty \frac{dp}{2\pi} \frac{1}{2\omega(p, M)} |\Psi(\rho_0; p, \omega(p, M))|^2 \end{aligned}$$

- Open string channel?

$$\text{Im} Z_{cylinder} \stackrel{?}{=} \langle B | \text{---} \text{---} \text{---} | B \rangle$$

- **Euclidean** vs. **Lorentzian** worldsheet

- Gives the same answer in rolling tachyon (Karzcmarek-Liu-Maldacena-Strominger), **but...** (Okuyama-Rozalli, NRS)

Open string computation

- Modular transform is (only) well-defined in **Lorentzian signature world sheet**.

$$\int ds_c \int_0^\infty dp \int_{-\infty}^\infty d\omega \frac{\sinh(\pi p)}{(\cosh(\pi\omega) + \cosh(\pi p)) \sinh(\pi p/k)} q^{\frac{1}{2k}(p^2 - \omega^2)} Z_{osc} \quad \text{closed}$$

$$= \int \frac{ds_o}{s_o} \int_0^\infty dp \int_{-\infty}^\infty d\omega \frac{\sinh(\pi\omega)}{(\cosh(\pi p) + \cosh(\pi\omega)) \sinh(\pi\omega/k)} \tilde{q}^{\frac{1}{2k}(p^2 - \omega^2)} Z_{osc} \quad \text{open}$$

- Imaginary part consists of **two** parts

$$\text{Im} Z_{cyl} = \text{Im} Z_{naive} + \text{Im} Z_{pole}$$

- **Naïve part** corresponds to contribution easily guessed in the Euclidean approach (but not enough)

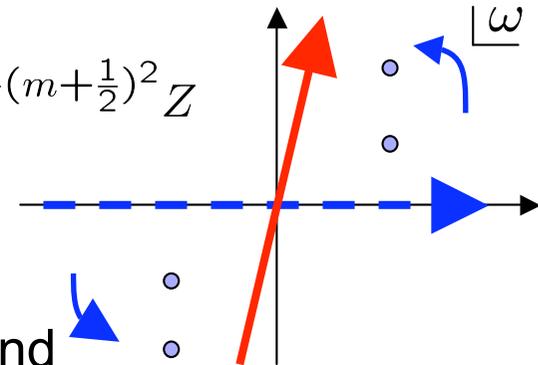
$$Z_{naive} = \int_0^\infty \frac{dt_o}{t_o} \int_0^\infty dp \int_{-\infty}^\infty d\omega \frac{\sin(\pi\omega) \tilde{q}^{\frac{1}{2k}(\omega^2 + p^2)}}{(\cos(\pi\omega) + \cosh(\pi p)) \sin(\pi\omega/k)} Z_{osc}$$

$$\text{Im} Z_{naive} = \sum_{n=1}^{\infty} \int \frac{dt_o}{t_o} \int_{-\infty}^\infty dp \frac{(-1)^{n+1} \sin(\pi nk) e^{-2\pi t_o (\frac{p^2}{2k} + \frac{kn^2}{2})}}{\cos(\pi nk) + \cosh(\pi p)} Z_{osc}$$

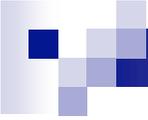
Unitarity meets open/closed duality

- Pole part comes from poles in **Euclidean (Wick) rotation**

$$\text{Im}Z_{pole} = 2 \int_0^\infty \frac{dt_o}{t_o} \sum_{n=1}^\infty \sum_{m=0}^{\lfloor \frac{kn}{2} - \frac{1}{2} \rfloor} (-1)^{n+1} e^{-\pi t_o n(2m+1) + \frac{2\pi t_o}{k} (m + \frac{1}{2})^2} Z$$



- Both contributions are imperative to understand
 - Unitarity
 - Tachyon-Radion correspondence
- **Summary**
 - In Euclidean approach, **no apparent reason** to include/exclude **pole contributions**.
 - **Unitarity** demands its existence, and Lorentzian theory **automatically** knows it.
 - Fortuitously **no pole contribution** in rolling tachyon (in 2D ZZ-brane decay as well).



Summary and Outlook

- **Exact boundary states** for rolling D-brane is constructed.
- **Tachyon-Radion correspondence** is **proved** in α' exact way.
 - Full proof in string field theory?
- **BH/String transition** is observed at $k=1$.
 - Is 2D pure BH really black? Matrix model?
- Consistency between **unitarity** and **open/closed duality** requires careful analytic continuation (Wick rotation).

The shortest path between two truths in the real domain passes through the complex domain. ---- J. Hadamard