### Master Space and Hilbert Series

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#### In collaboration with:

A. Zaffaroni, A. Hanany, A. Butti, Y.H. He, D. Vegh

#### Other important related works by:

S.T. Yau, D. Martelli, J. Sparks, S. Benvenuti, B. Feng,...

#### Introduction

I would like give you an overview of the use of the Master Space  $\mathcal{F}^{\flat}$  and the Hilbert Series H(q;...) to analyse SUSY field theories. I will use the example of D3 branes at  $\mathcal{N}=1$  singularities  $\chi$ .

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Natural subclass of d.o.f: gauge invariant operators  $\mathcal{O}$  killed by half of the supercharges  $\bar{\mathcal{D}}\mathcal{O}=0$ : the BPS operators. Non-renormalization properties, useful to study strong coupling dynamics and dualities.

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- ▶ Low energy dynamics: vacuum structure. In SUSY field theory a high degeneracy of vacua (moduli space M) is a typical phenomenon. M is defined by the set of F-flat, D-flat conditions modulo gauge invariance.

▶ It is well known that  $\mathcal{M}$  admits a "quotient" description:

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- ▶ Spectrum  $\rightarrow$  BPS operators  $\leftarrow$  Hilbert Series:  $H(q; \mathcal{M})$
- ▶ There exist a nice interplay between  $\mathcal{F}^{\flat}$  and  $H(q; \mathcal{M})$ .



To every  $\mathcal{N}=1$  4d field theories we can associate two objects:  $\mathcal{F}^{\flat}$ ,  $H(q;\mathcal{M})$ .

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- ▶  $\lim_{t\to 0} H(q=e^{-bt};\mathcal{M}) \sim P(q)/(1-q)^{\dim\mathcal{M}} + ...$
- ► The symmetries G of the variety  $\mathcal{M}$  can be encoded into the Hilbert Series:  $H(q; \mathcal{M}) = \sum_{\vec{l}} \chi_{\vec{l}}^{G}(q)$



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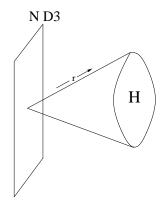


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  - ►  $\lim_{t\to 0} g_{1,\vec{B}}(e^{-bt})/g_{1,0}(e^{-bt}) \sim 1 + t\pi \operatorname{Vol}_{\mathbf{C}_3^{\bar{\mathbf{B}}}}(\mathbf{b})/2\operatorname{Vol}_{\mathbf{H}}(\mathbf{b}) + \dots$

# D3 branes at Singularities



### Generic features: Quiver Gauge Theories

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- ▶ examples:  $\mathbb{C}^3$ ,  $\mathbb{C}^3/\Gamma$ ,  $C(T^{1,1})$ ,  $C(Y^{p,q})$ ,  $C(L^{a,b,c})$ ,...

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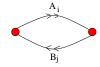
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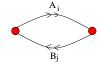
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- ▶ Use the PE function:  $PE_{\nu}[f(t)] := \exp\left(\sum_{k=1}^{\infty} \frac{\nu^k f(t^k)}{k}\right)$ , to pass from N = 1 to N > 1



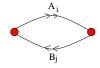


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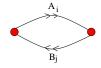
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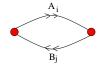
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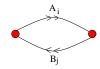
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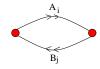




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- ▶ Operators:  $Tr(A_iB_i)$  Mesons; det  $A_i$ , det  $B_i$  Baryons



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# The complex cone over $\mathbb{F}_0$

$$C_{\mathbb{C}}(\mathbb{F}_0) = C_{\mathbb{C}}(\mathbb{P}^1 \times \mathbb{P}^1)$$





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► Explicit symmetries:  $SU(2) \times SU(2) \times U(1)_R \times U(1)_B$ 

► F terms:

$$A_1B_iC_2 = A_2B_iC_1$$
 ,  $B_1C_iD_2 = B_2C_iD_1$   
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▶ for N=1 we can factorize:

$$\mathbf{B}_{i} (\mathbf{A}_{1} \mathbf{C}_{2} - \mathbf{A}_{2} \mathbf{C}_{1}) = 0$$
 ,  $\mathbf{C}_{i} (\mathbf{B}_{1} \mathbf{D}_{2} - \mathbf{B}_{2} \mathbf{D}_{1}) = 0$   
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► For generic N BPS operators and the moduli space  $\mathcal{M}_N$  have  $SU(2)^4 \times U(1)^2$  symmetry

$$\begin{split} &\sum_{\nu=0}^{\infty} \nu^{N} H(t_{1},t_{2},x,y,a_{1},a_{2};\mathcal{M}_{N}) = \\ &\sum_{\beta,\beta'=0}^{\infty} [0,0,\beta,\beta'] \left( PE_{\nu} \left[ \sum_{n=0}^{\infty} [2n+\beta,2n+\beta',0,0] t_{1}^{2n+\beta} t_{2}^{2n+\beta'} \right] \right. \\ &\left. - PE_{\nu} \left[ \sum_{n=1}^{\infty} [2n+\beta,2n+\beta',0,0] t_{1}^{2n+\beta} t_{2}^{2n+\beta'} \right] \right) \end{split}$$

We were almost classical

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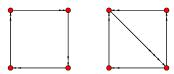
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- ightharpoonup Example  $\mathbb{F}_0$



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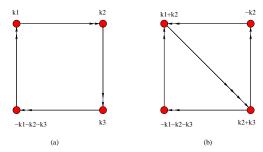
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#### The two three dimensional quiver CS theories



#### with superpotentials

$$W_{(a)} = \epsilon_{ij} \epsilon_{pq} \mathbf{A}_i \mathbf{B}_p \mathbf{C}_j \mathbf{D}_q$$
  

$$W_{(b)} = \epsilon_{ij} \epsilon_{kl} X_{13}^{ik} X_{32}^l X_{21}^j - \epsilon_{ij} \epsilon_{kl} X_{13}^{ik} X_{34}^l X_{41}^j$$

probe the same  $CY_4$  singularities for specific values of  $k_i$ .



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- ► Future applications:
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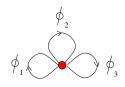
# Thanks for your Attention

• 
$$\mathcal{N}=$$
 4,  $SU(N)\leftrightarrow AdS_5\times S^5$ 

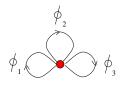


$$\mathcal{N} = 4, \ SU(N) \leftrightarrow AdS_5 \times S^5$$

$$\mathcal{W} = \text{Tr} \left( \phi_1 \phi_2 \phi_3 - \phi_1 \phi_3 \phi_2 \right)$$

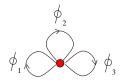


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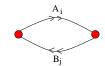


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$$H(q_1,t_2,t_3;\mathbb{C}^3) = \sum_{i,j,k} t_1^i t_2^j t_3^k = \frac{1}{(1-t_1)(1-t_2)(1-t_3)}$$

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Lessons:

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$$\sum_{\nu=0}^{\infty} \nu^{N} H(q,b,a;\mathcal{M}_{N}) = \sum_{\vec{\beta}} Z_{\mathsf{aux}}^{\vec{\beta}}(a;\mathcal{X}) P E_{\nu}[g_{1,\vec{\beta}}(q,b;\mathcal{X})]$$

$$g_{1,\beta,\beta'}(t_1,t_2;\mathbb{F}_0) = \sum_{n=0}^{\infty} (2n+\beta+1)(2n+\beta'+1)t_1^{2n+\beta}t_2^{2n+\beta'}$$

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▶ the  $SU(2)^4$  representation  $[m] \times [n] \times [p] \times [q] = [m, n, p, q]$ .  $g_{1,\beta,\beta'}(t_1,t_2;\mathbb{F}_0)$  can be written explicitely in term of representations of  $SU(2)^4 \times U(1)^2$ :

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 $ightharpoonup Z_{\text{aux}}$  also in rep of the HIDDEN SYMMETRY:

$$\begin{split} Z_{\text{aux}}(t_1, t_2, a_1, a_2; \mathbb{F}_0) &= \\ \sum_{\beta, \beta' = 0}^{\infty} [0, 0, \beta, \beta'] t_1^{\beta} t_2^{\beta'} &- \sum_{\beta, \beta' = 2}^{\infty} [0, 0, \beta - 2, \beta' - 2] t_1^{\beta} t_2^{\beta'} \end{split}$$