

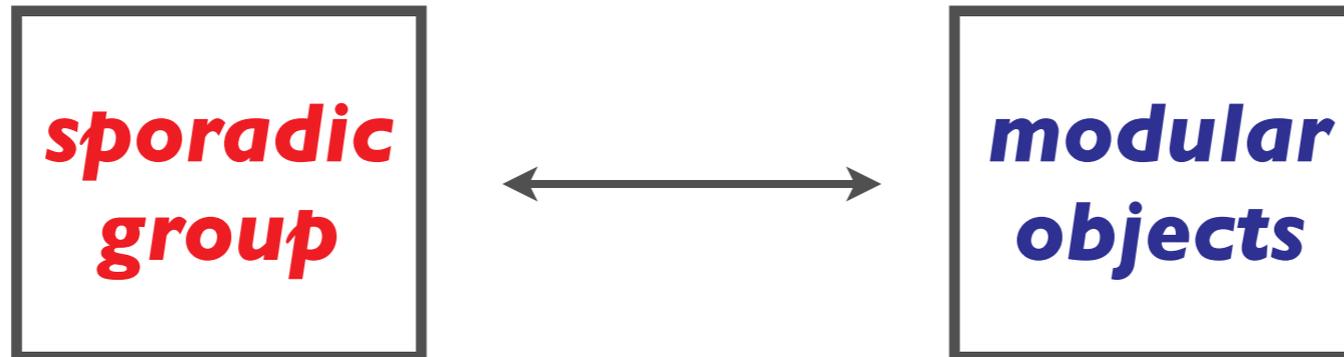
Strings 2011, Uppsala

***K3, M<sub>24</sub>, and  
Holographic Moonshines***

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# Introduction

# The Moonshine Phenomenon



$$J(\tau) = J(\tau + 1) = J(-\frac{1}{\tau}) \quad (q = e^{2\pi i\tau})$$

$$= q^{-1} + 196884q + 21493760q^2 + \dots$$

$$\begin{array}{ccc} \parallel & & \parallel \\ 1 + 196883 & 1 + 196883 + & 21296876 \end{array}$$

dim of irreps of Monster

[McKay late 70's]

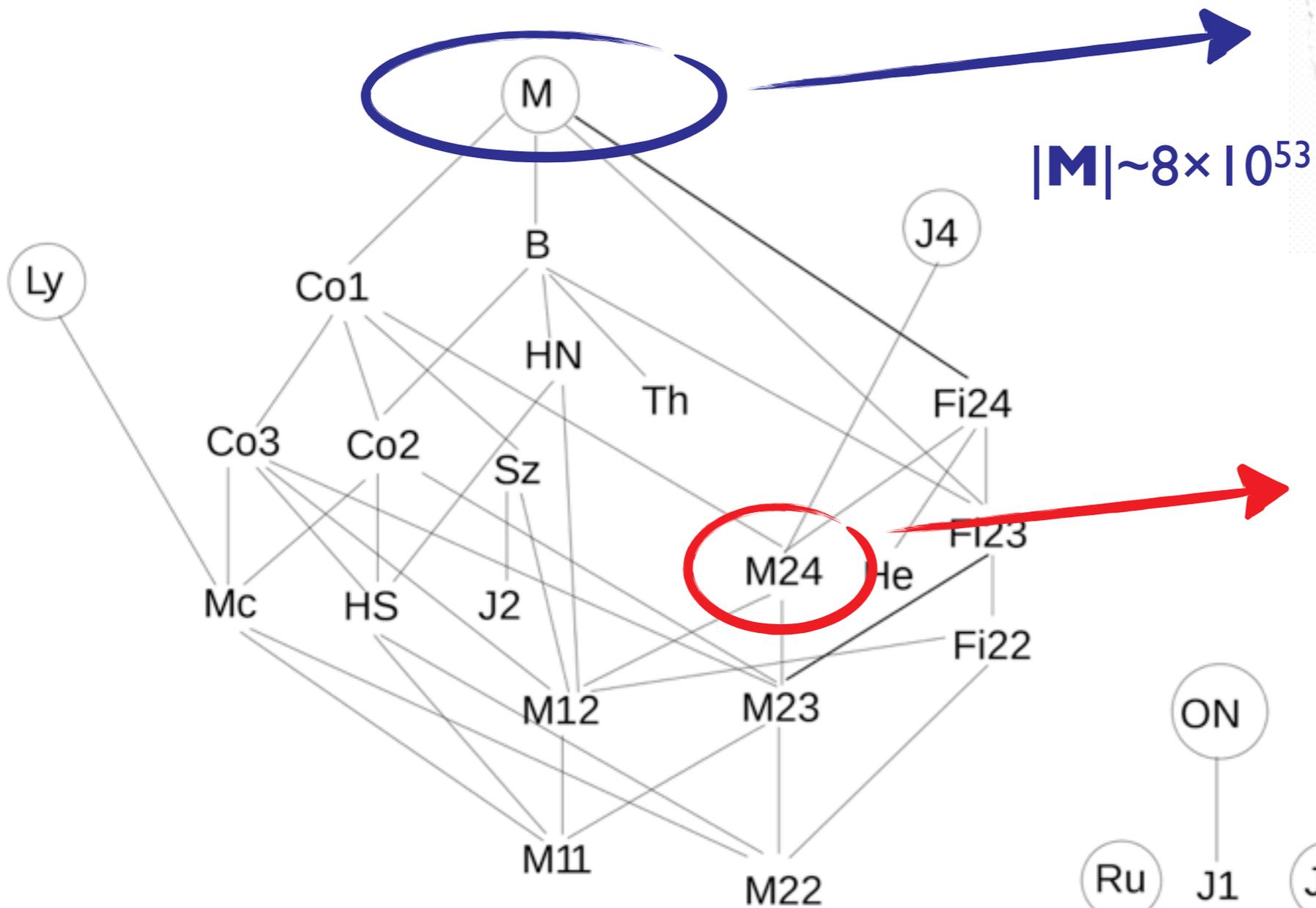
The Largest  
Sporadic Group



The Most Natural  
Modular Function

# Sporadic Groups

**The 26 finite simple groups that don't come in  $\infty$ -families.**



largest Mathieu group  
 $\sim 2 \times 10^9$

# **Mathieu 24**

It has a natural 24-dim representation, on which

$$M_{24} \subset S_{24}.$$

$N$  : a 24-dim Niemeier lattice



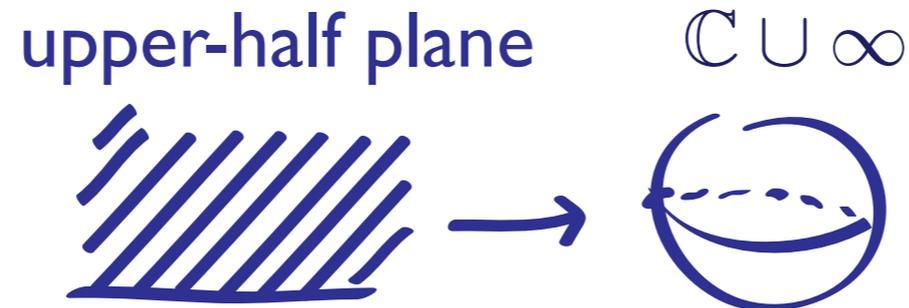
$M_{24}$

All symmetries  $G$  of  $K3$  manifold preserving the hyper-Kähler structure satisfy

$$G \subset M_{23} \subset M_{24}.$$

# Modular $\subset$ Automorphic Forms

A modular form  $f(\tau)$  transforms “covariantly” under a subgroup  $\Gamma$  of  $SL(2, \mathbb{R})$ .



An “automorphic form” transforms under more general groups (a higher dimensional version of a modular form).

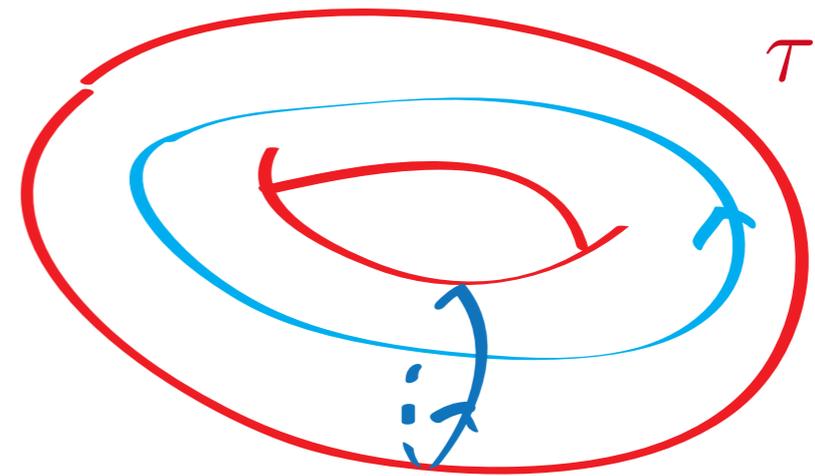
In physics, this “automorphism” reflects an underlying symmetry of the problem.

# ***String theory is good at producing automorphic forms!***

World-sheet Symmetries  
(mapping class group of  $\Sigma$ )

eg.  $Z_{2d} \text{CFT}(\tau)$  transforms under  $SL(2, \mathbb{Z})$ .

Space-Time Symmetries  
(Such as  $T$ -,  $S$ -dualities)



All symmetries have to be reflected in suitable partition functions.

# String Theory and Moonshine

e.g. *Monstrous moonshine* (partially) explained by CFT and proven using bosonic string theory.

$$J(\tau) = q^{-1} + 196884 q + 21493760 q^2 + \dots$$
$$= \text{Tr}_{\mathcal{H}} q^{L_0 - \frac{c}{24}} \left( \begin{array}{l} \text{partition function of a } c=24 \text{ 2d} \\ \text{chiral CFT with a Monster symmetry} \end{array} \right)$$

[’88 Frenkel-Lepowsky-Meurmann  
/Dixon-Ginsparg-Harvey]

The “moonshine conjecture” was proven by introducing generalised Kac-Moody algebras and considering the lift of  $J(\tau)$  into an automorphic form. [R. Borcherds ’92]

Physically, this “lift” corresponds to considering the **full 26-dim bosonic string**.

spectrum generating algebra

Borcherds-Kac-Moody Algebra

symmetry

denominator

Sporadic Groups

Automorphic Forms

symmetry

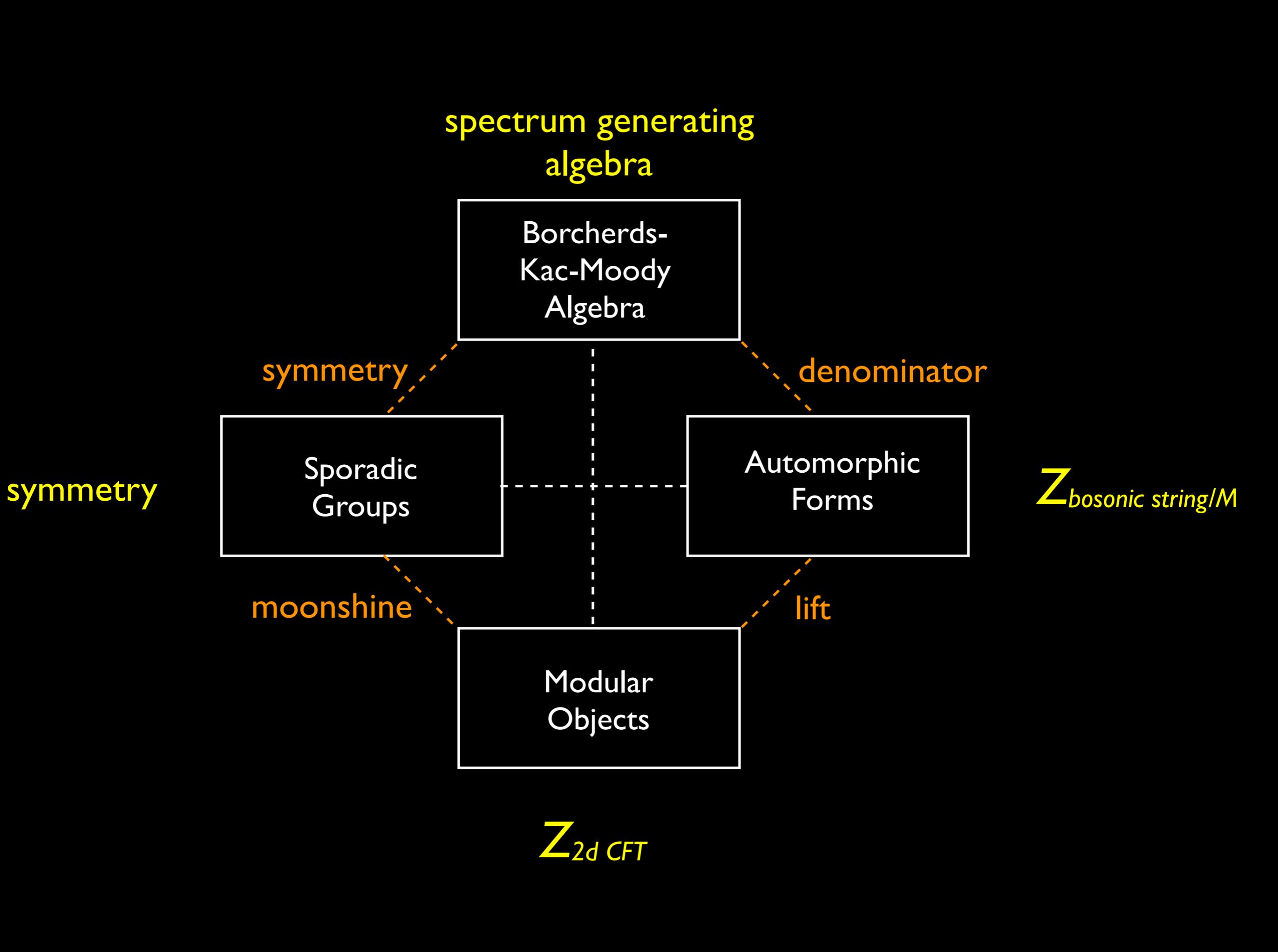
$Z_{\text{bosonic string}/M}$

moonshine

lift

Modular Objects

$Z_{2d \text{ CFT}}$



# ***K3 and M<sub>24</sub>***

Last year, Eguchi-Ooguri-Tachikawa made a striking observation on the BPS-sector of the 2d SCFT with *K3* target:

$\mathcal{Z}(\tau, z; K3) = K3$  elliptic genus

[Eguchi-Ooguri-Taormina-Yang '89]

$$= \frac{\theta_1^2(\tau, z)}{\eta^3(\tau)} \left( 24 \mu(\tau, z) + 2q^{-1/8} \left( -1 + \textcircled{45}q + \textcircled{231}q^2 + \textcircled{770}q^3 + \dots \right) \right)$$

number of massive  $N=4$  multiplets

**also dimensions of irreps of  $M_{24}$ !!!**

# ***A New $M_{24}$ Moonshine?***

## ***WHY?***

## ***WHAT?***

- What precisely should we conjecture? In particular, what should be the moonshine diagram in this case?
- What are the consequences? Are they consistent with what we know?
- What are the evidences?

**(First part)**



# ***A Long-Standing Puzzle***

Along the way we were confronted with an old puzzle of a number theoretic nature (the genus zero property).

Surprisingly, considerations from AdS/CFT give a way out by providing natural explanations for some crucial properties of the modular groups appearing in all moonshines known so far.



***Holographic Modularity of Moonshines***

(Second part)

***K3 and M<sub>24</sub>***

# K3 Sigma-model

In a 2d  $N=(4,4)$  SCFT, BPS states are counted by elliptic genus.

They are modular and are called Jacobi forms.

$$\begin{aligned}\mathcal{Z}(\tau, z; K3) &= \text{Tr}_{\mathcal{H}_{\text{RR}}} \left( (-1)^{J_0 + \bar{J}_0} y^{J_0} q^{L_0 - c/24} \bar{q}^{\bar{L}_0 - c/24} \right) = \sum_{\substack{n \geq 0, \\ \ell \in \mathbb{Z}}} c(4n - \ell^2) q^n y^\ell \\ &= \frac{\theta_1^2(\tau, z)}{\eta^3(\tau)} \left( 24 \mu(\tau, z) + 2q^{-1/8} (-1 + 45q + 231q^2 + 770q^3 + \dots) \right)\end{aligned}$$

Counting the number of massive  $N=4$  multiplets:

$$H(\tau) = 2q^{-1/8} (-1 + 45q + 231q^2 + 770q^3 + \dots)$$

Mock Modular Form

[Zwegers '02/  
Eguchi-Hikami'10]

# ***1/4-BPS States in Type II/K3×T<sup>2</sup>***

Further compactifying down to 4d is achieved by “Borchers-lifting” the elliptic genus.

The resulting automorphic form  $\Phi$  counts the 1/4-BPS states of the  $N=4, d=4$  theory: [Dijkgraaf-Verlinde<sup>2</sup> '97]

$$\Phi(\tau, z, \sigma) = \frac{1}{pqy} \prod_{n,m,\ell} \left( \frac{1}{1 - p^n q^m y^\ell} \right)^{c(4nm - \ell^2)}$$

Fourier coeff. of  $\mathcal{Z}(\tau, z; K3)$

$$= \sum_{P,Q} D(P, Q) p^{Q \cdot Q/2} q^{P \cdot P/2} y^{P \cdot Q}$$

This automorphic form also defines an algebraic structure underlying the 1/4-BPS spectrum.

Dyonic 4d BH  
charge  $(Q, P)$   
 $S_{BH} \sim \log D(P, Q)$

See A. Sen, Strings '07, '08, '10

# ***1/2-BPS States in Type II/K3×T<sup>2</sup>***

$$\frac{1}{\eta^{24}(\tau)} = \sum d(Q) q^{\frac{Q \cdot Q}{2}} \Leftarrow \begin{array}{c} \updownarrow \\ \text{heterotic}/T^6 \end{array}$$

$$= \frac{1}{q \prod_{n \geq 1} (1 - q^n)^{24}}$$

[Dabholkar-Harvey '89]

= partition function of 24 chiral bosons

***1/2-BPS Hilbert space of Type II/K3×T<sup>2</sup> forms a M<sub>24</sub>-representation!***

***M<sub>24</sub> maps 1/2-BPS states to each other.***

# Dyonic Bound States in Type II/K3×T<sup>2</sup>

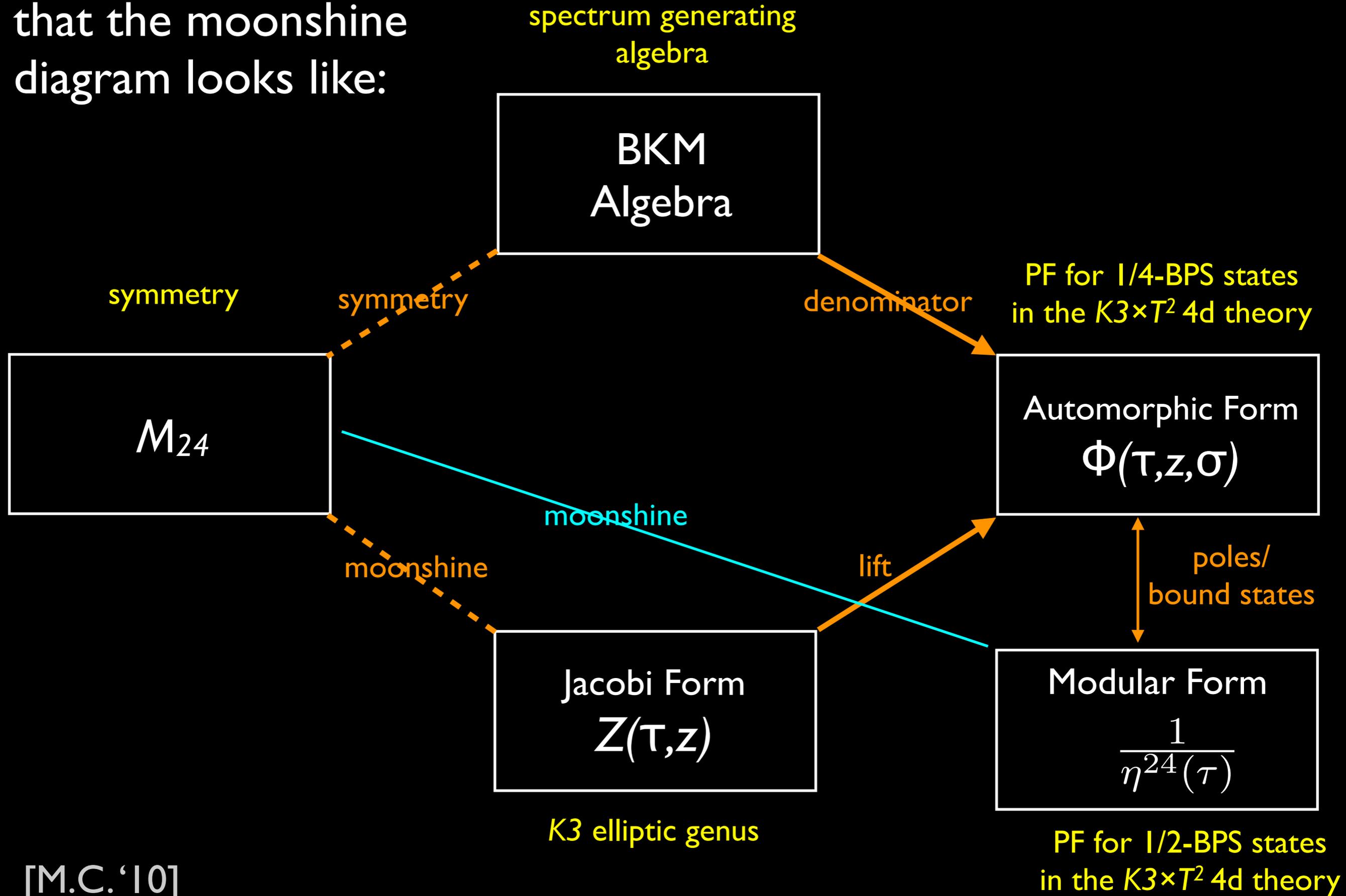
1/4-BPS spectrum has to know about the 1/2-BPS spectrum too!



$$\Phi(\tau, z, \sigma) \xrightarrow{z \rightarrow 0} \frac{1}{(2\pi iz)^2} \frac{1}{\eta^{24}(\tau)} \frac{1}{\eta^{24}(\sigma)}$$

[A. Sen/M.C.-Verlinde '07, '08]

Physics suggests  
that the moonshine  
diagram looks like:



# What exactly are we conjecturing?

The mock modular form containing the  $M_{24}$  information in  $\mathcal{Z}(\tau, z; K3)$

$$H(\tau) = q^{-1/8} \left( -2 + 90q + 462q^2 + 1540q^3 + \dots \right)$$

$$\stackrel{?}{=} q^{-1/8} \left( -2 + \sum_{n=1}^{\infty} q^n \dim K_n \right) \quad K_n = M_{24}\text{-representation}$$

True. But meaningless.

Instead we need information about the vector (Hilbert) space!

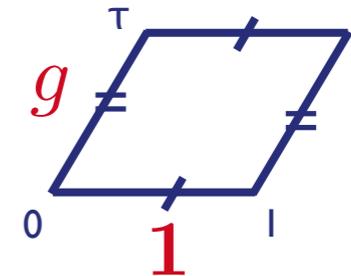
# Twisting the Symmetries

Hilbert space  $\mathcal{H}$



The **twisted** partition function  $Z_g = \text{Tr}_{\mathcal{H}}(g \dots)$  gives finer information about  $\mathcal{H}$  than just the dimension.

e.g. 2d CFT  $Z_g(\tau) = \text{Tr}_{\mathcal{H}}(g q^{\hat{H}} \dots)$



boundary condition:  ~~$SL(2, \mathbb{Z})$~~   $\rightarrow \Gamma_g$

e.g. When there is supersymmetry, we can compute

$$\text{Tr}_{\mathcal{H}}(q^{L_0}) \quad \text{or} \quad \text{Tr}_{\mathcal{H}}((-1)^F q^{L_0})$$

PF

index (twisted PF)

The invariance group has to preserve the spin structure.

# A Concrete Conjecture

In particular, we want to compute the **twisted** K3 elliptic genus  $\mathcal{Z}_g(\tau, z; K3)$  and extract the interesting part  $H_g(\tau)$ .

## Conjecture:

For all  $g \in M_{24}$  It should be a (mock) modular form of a discrete subgroup  $\Gamma_g \subset SL(2, R)$  and satisfy

$$H_g(\tau) = q^{-1/8} \left( -2 + \sum_{n=1}^{\infty} q^n \text{Tr}_{K_n} g \right)$$

**Status:** The candidate (mock) modular forms  $H_g(\tau)$  have been proposed for all  $[g] \in M_{24}$ .

[M.C. /Gaberdiel-Hohenneger-Volpato/Eguchi-Hikami '10]

## Checks:

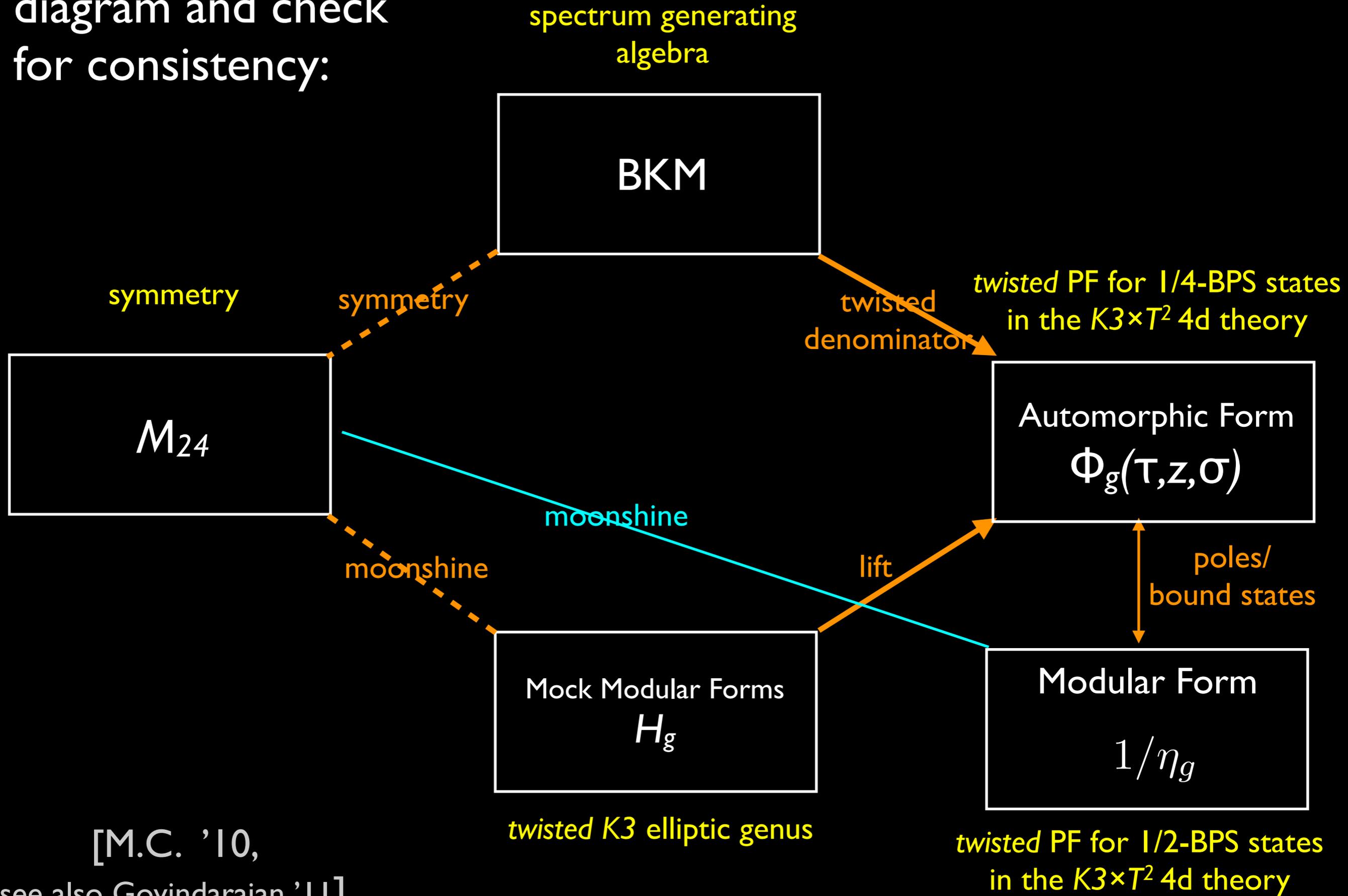
1)  $K_{n < 1000}$  constructed!

[Y. Tachikawa]

2) Coincides with geometric calculation when the latter is available.

[David-Jatkar-Sen '06]

Twist the moonshine diagram and check for consistency:



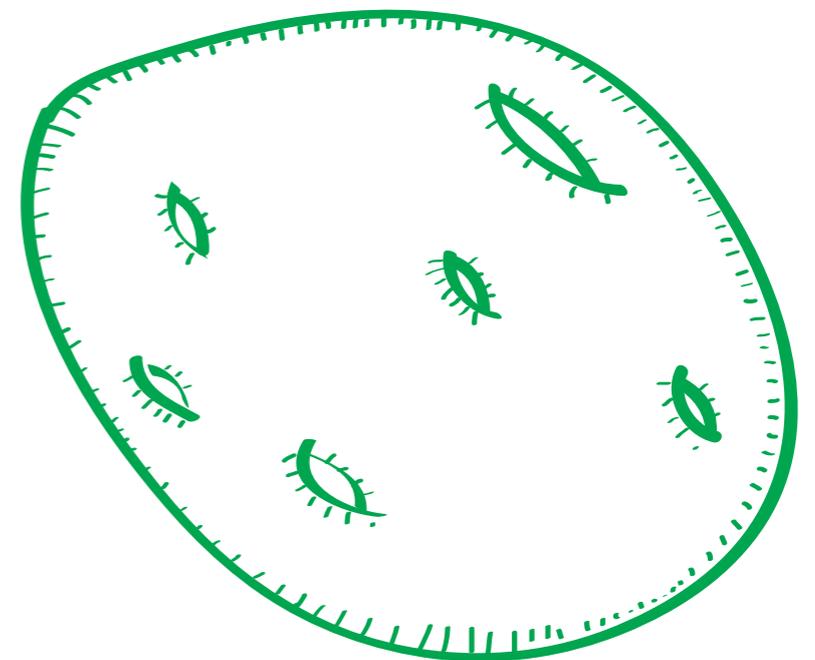
[M.C. '10,  
see also Govindarajan '11]

# **Summary: $M_{24}$ in K3 String Theories**

We don't understand, but we are pretty sure!

Challenge:  
Stringy K3 Symmetries

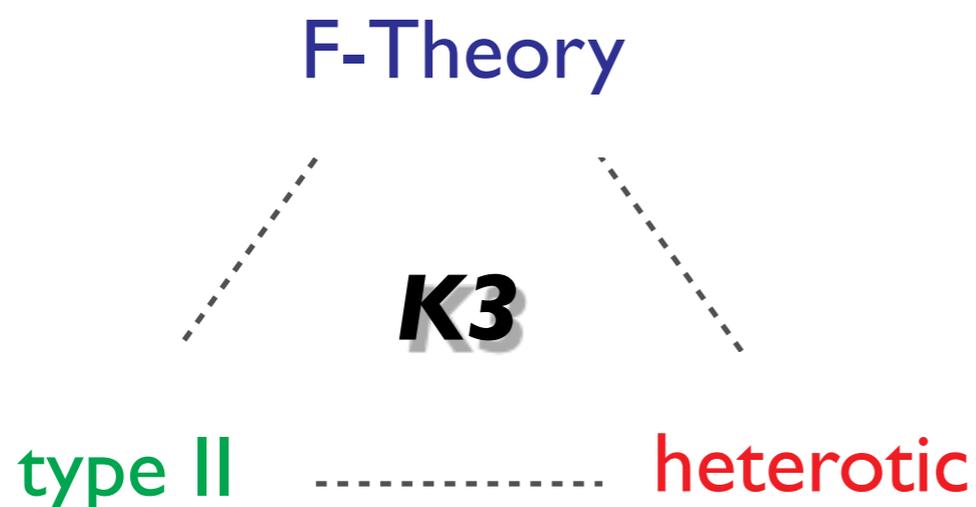
These results are clearly evidence for uncharted symmetries of the stringy geometries of K3. **Note: Classical symmetries are not enough! Stringy symmetries are needed.**



# Summary: $M_{24}$ in K3 String Theories

***This moonshine is physically interesting.***

- We have black holes.
- Since  $K3$  is ubiquitous, it is important to understand these new symmetries.
- We should also explore where else  $M_{24}$  shows its presence.



Sporadic Quantum  
Black Holes

# **Holographic Modularity of the Moonshines**

with **John Duncan** 1107.xxxx

# A Long-Standing Mystery

For positive every integer  $n$  we can define a certain subgroup  $\Gamma_0(n)^+ \subset SL(2, \mathbb{R})$ .

Q: What are the prime numbers  $p$  such that  $\Gamma_0(p)^+$  is **genus zero**?

A:  $\Gamma_0(p)^+$  is genus zero iff  $p$  divides  **$|M|$** !



[Ogg '73]

**WHY??**



**WHY??**

# An “Explanation” from Monstrous Moonshine

$$J(\tau) = q^{-1} + 196884 q + 21493760 q^2 + \dots$$
$$= \text{Tr}_{\mathcal{H}} q^{L_0 - \frac{c}{24}}$$

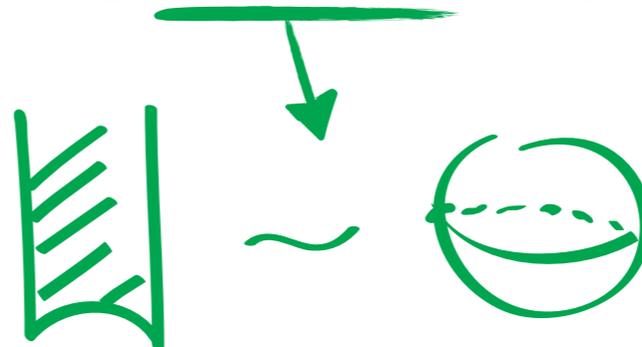
$\mathcal{H}$   
 $\mathcal{M}$

As before, we consider the twisted partition function

$$J_g(\tau) = \text{Tr}_{\mathcal{H}}(gq^{L_0 - \frac{c}{24}}) = \sum_{n \geq -1} c_g(n) q^n \quad \text{for every } g \in \mathbf{M}$$

**Moonshine Conjecture** (Conway-Norton '79):

$J_g(\tau)$  is invariant under some genus zero  $\Gamma_g \subset SL(2, \mathbb{R})$ .



# Genus Zero Property

Genus zero groups  $\Gamma \subset SL(2, R)$  are rare.  
But only genus zero groups seem to appear  
in moonshines.



Not just in Monstrous moonshine,  
this has also been extended to

- The “generalised moonshine”
- Groups other than the Monster

see for instance  
Norton '84/Carnahan '08  
Höhn '03/Duncan '05, '06

***BUT WHY GENUS ZERO??***

# NO Genus Zero for the New $M_{24}$ Moonshine

**Heresy!**

But true, by inspecting

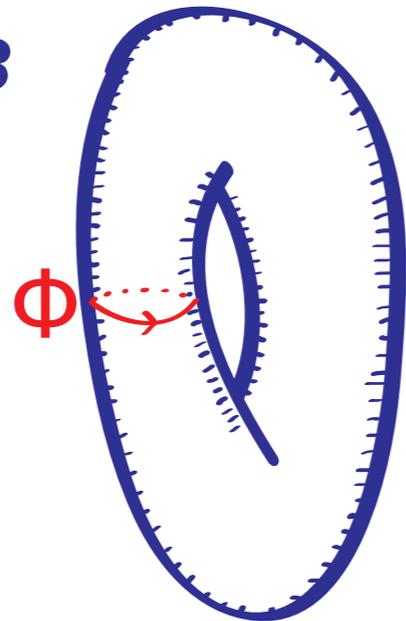
$$H_g(\tau) = q^{-1/8} \left( -2 + \sum_{n=1}^{\infty} \underbrace{c_{g,n}}_{\downarrow} q^n \right)$$

$\ell \backslash [g]$	1A	2A	2B	3A	3B	4A	4B	4C	5A	6A	6B	7A	$\overline{7A}$	8A	10A	11A
1	90	-6	10	0	6	-6	2	2	0	0	-2	-1	-1	-2	0	2
2	462	14	-18	-6	0	-2	-2	6	2	2	0	0	0	-2	2	0
3	1540	-28	20	10	-14	4	-4	-4	0	2	2	0	0	0	0	0
4	4554	42	-38	0	12	-6	2	-6	-6	0	4	4	4	-2	2	0
5	11592	-56	72	-18	0	-8	8	0	2	-2	0	0	0	0	2	-2
6	27830	86	-90	20	-16	6	-2	6	0	-4	0	-2	-2	2	0	0
7	61686	-138	118	0	30	6	-10	-2	6	0	-2	2	2	-2	-2	-2
8	131100	188	-180	-30	0	-4	4	-12	0	2	0	-3	-3	0	0	2
9	265650	-238	258	42	-42	-14	10	10	-10	2	6	0	0	-2	-2	0

# AdS<sub>3</sub>/CFT<sub>2</sub>

Recall: In a Euclidean 3d gravity, the only smooth solution with aAdS boundary conditions and torus boundary is the **solid torus** :

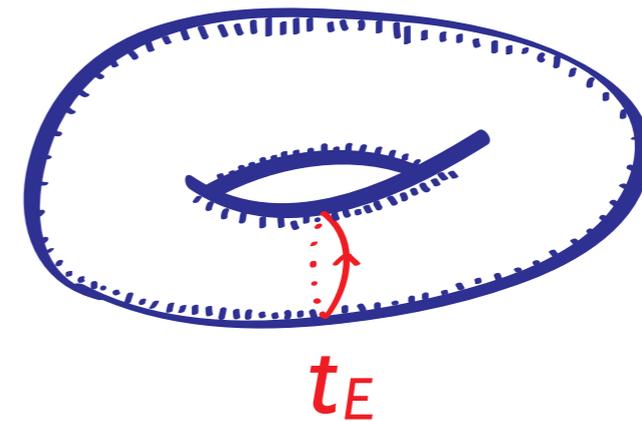
AdS<sub>3</sub>



$$\tau \rightarrow -\frac{1}{\tau}$$

(large diff.)

BTZ



⇒ Saddle points of the gravity path integral are labeled by the **contractible cycle**.

$$\gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_\infty \setminus SL(2, \mathbb{Z}) \quad (\tau \sim \tau + 1)$$

[Maldacena-Strominger '98]

Twisting the boundary condition: ~~SL(2,Z)~~ →  $\Gamma_g$ .

# **$AdS_3/CFT_2$**

**Assuming** a CFT has a dual description given by **semi-classical-like AdS gravity**

⇒ The (twisted) partition function  $Z_g(\tau)$  can also be computed from the gravity side by summing over saddle point contributions

$$\sum_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_\infty \setminus \Gamma_g} f\left(\frac{a\tau + b}{c\tau + d}\right)$$

⇒  $Z_g(\tau)$  has to be **Rademacher-summable!**

See also the “Farey Tail” papers:

Dijkgraaf-Maldacena-Moore-Verlinde '00,

Kraus-Larsen/Dijkgraaf-de Boer-M.C.-Manschot-Verlinde/Denef-Moore/Manschot-Moore '06

# Rademacher-Summability

$$Z_g(\tau) = \text{Reg} \left[ \sum_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma_\infty \setminus \Gamma_g} f\left(\frac{a\tau + b}{c\tau + d}\right) \right]$$

convergent,  
anomaly-free

$$e.g. \quad J(\tau) = e(-\tau) + \lim_{K \rightarrow \infty} \sum_{\substack{0 < c < K \\ -K^2 < d < K^2}} e\left(-\frac{a\tau + b}{c\tau + d}\right) - e\left(-\frac{a}{c}\right) - 12 \quad , \quad e(x) = e^{2\pi i x}$$

$$= \lim_{K \rightarrow \infty} \sum_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \in (\Gamma_\infty \setminus SL(2, \mathbb{Z}))_{<K}} \left( q^{-1} \Big|_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \cdot \tau} - q^{-1} \Big|_{\begin{pmatrix} a & b \\ c & d \end{pmatrix} \cdot i\infty} \right) - 12$$

$$= \text{Reg} \left[ \sum_{\gamma \in (\Gamma_\infty \setminus SL(2, \mathbb{Z}))} q^{-1} \Big|_{\gamma \cdot \tau} \right]$$

# Rademacher-Summability

puts stringent constraints on  $\Gamma_g$ .

e.g. If we want to have a **modular function** (weight 0, weakly holomorphic),

Rademacher-summability

$\Leftrightarrow$

$\Gamma_g$  is genus zero

[Duncan-Frenkel '09]

In particular,  $J_g(\tau) = \text{Reg} \left[ \sum_{\gamma \in (\Gamma_\infty \setminus \Gamma_g)} q^{-1} \Big|_{\gamma \cdot \tau} \right]$  for all  $g \in \mathbf{M}$ !

Note: this holds for higher central charge cases ( $q^n$ ) as well.

# Rademacher-Summability of the $M_{24}$ Moonshine

Here we have  $H_g(\tau)$  = weight 1/2 Mock modular form.

Rademacher-summability  ~~$\iff$~~   $\Gamma_g$  is genus zero

Instead, with the help of previous results of Bringmann-Ono ('06) and Eguchi-Hikami ('09), we show

$$H_g(\tau) = -2 \operatorname{Reg} \left[ \sum_{\gamma \in (\Gamma_\infty \setminus \Gamma_g)} q^{-1/8} \Big|_{\gamma \cdot \tau} \right] \quad \text{for all } [g] \in M_{24}.$$

# ***To Summarise***

We propose a new pattern  
**Rademacher Summability**

$Z_g(\tau)$  = modular functions  
e.g. Monster moonshine

**$g=0$**

$Z_g(\tau)$  = *mock* modular forms  
for the new  $M_{24}$  moonshine

***verified***

Motivation for the new pattern:

All known theories of moonshine have a CFT interpretation.

Assuming the existence of a good dual description  
(for a higher  $c$  version of the sporadic CFT)  
 $\Rightarrow$  All twisted PF have to be Rademacher-summable.



A glass of whiskey for physicists?

AdS/CFT  $\rightarrow \rightarrow$  AdS/CMT, AdS/QCD,  
AdS/Hydrodynamics..... +AdS/NT??

**Thank You!**

# To Summarise

Our result raises many physical questions:

- For  $M_{24}$  moonshine the dual situation is (more or less) clear:  $AdS_3 \times S^3 \times K3$ . How about other sporadic CFT's? Our results can be interpreted as an evidence that there exists 2d chiral CFT's with sporadic symmetries with semi-classical-like (saddle-dominated)  $AdS_3$  duals.
- Relatedly, this brings back the question: What are the criteria for a CFT to have a (weakly coupled)  $AdS_3$  duals?

[see for instance, Heemskerk-Penedones-Polchinski-Sully '09, El-Showk-Papadodimas '11]

- How to understand the twisting from a bulk viewpoint in general? [some progress reported in A. Sen Strings 2011]