

# A Duality Web in 2 + 1 Dimensions and the Unity of Physics

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#### Based on:

NS and E. Witten, arXiv:1602.04251; NS, T. Senthil, C. Wang, and E. Witten, arXiv:1606.01989; P.-S. Hsin, NS, arXiv:1607.07457

#### Recent related papers:

- A. Karch and D. Tong, arXiv:1606.01893;
- J. Murugan and H. Nastase, arXiv:1606.01912

## Three (almost) independent lines of development – the unity of physics

- The condensed matter, 3d quantum field theory route
  - The supersymmetric route
  - The AdS/CFT, large N route

- Statistical transmutation: a massive particle coupled to a dynamical (statistical, emergent) gauge field with a Chern-Simons term can change its spin and statistics [Wilczek; Polyakov; Jain].
  - Many applications (FQHE, composite fermions, flux attachment, ...)

- Statistical transmutation: a massive particle coupled to a dynamical (statistical, emergent) gauge field with a Chern-Simons term can change its spin and statistics [Wilczek; Polyakov; Jain].
  - Many applications (FQHE, composite fermions, flux attachment, ...)
- This does not mean that a second-quantized theory of massless interacting bosons coupled to a gauge field with a Chern-Simons term is dual to a theory of fermions (or the other way around).

$$|D_B\Phi|^2 - |\Phi|^4 \quad \leftrightarrow \quad |D_b\hat{\Phi}|^2 - |\hat{\Phi}|^4 + \frac{1}{2\pi}Bdb$$

Particle/vortex duality [Peskin; Dasgupta and Halperin]

$$|D_B\Phi|^2 - |\Phi|^4 \quad \leftrightarrow \quad |D_b\hat{\Phi}|^2 - |\hat{\Phi}|^4 + \frac{1}{2\pi}Bdb$$

- LHS is XY, O(2) Wilson-Fisher.
- B is a background field coupled to a global  $U(1)_B$  symmetry.
- RHS is a gauged version of this theory. b is a dynamical field.
- IR duality two different theories flowing to the same IR fixed point.
- $\Phi \leftrightarrow \mathcal{M}_b$  is a monopole operator of b (charged under  $U(1)_B$ ).
- $|\Phi|^2 \leftrightarrow -|\widehat{\Phi}|^2$ . Upon deformation: unbroken  $U(1)_B$  phase is Higgs phase in the RHS; broken  $U(1)_B$  phase massless b.

$$|D_B\Phi|^2 - |\Phi|^4 \quad \leftrightarrow \quad i\bar{\Psi}D_a\Psi + \frac{1}{2\pi}Bda$$

Boson/fermion duality [Chen, Fisher, Wu; Barkeshli, McGreevy]

$$|D_B\Phi|^2 - |\Phi|^4 \quad \leftrightarrow \quad i\bar{\Psi}D_a\Psi + \frac{1}{2\pi}Bda$$

LHS Wilson-Fisher fixed point (B is a background gauge field) RHS QED with gauge field a with a single fermion, a.k.a  $U(1)_{1/2}$ 

- Arguments involve elementary fields with fractional charges and fractional level Chern-Simons terms.
- LHS is T-reversal invariant, while RHS seems like it is not.
- LHS does not need a spin structure, while RHS does. Violating gravitational 't Hooft matching conditions?
- The IR behavior of the RHS is debated.

$$i\bar{\Psi}D\!\!\!/_A\Psi \quad \leftrightarrow \quad i\bar{\chi}D\!\!\!\!/_a\chi + {1\over 4\pi}Ada$$

Fermion/fermion duality [Son; Wang, Senthil; Metlitski, Vishwanath]

$$i\bar{\Psi}D\!\!\!/_A\Psi \quad \leftrightarrow \quad i\bar{\chi}D\!\!\!\!/_a\chi + \frac{1}{4\pi}Ada$$

- Motivated by
  - physics of the lowest Landau level at half-filling [Halperin, Lee, Read]
  - T-Pfaffian state of topological insulators [Chen, Fidkowski, Vishwanath].
- Improperly quantized Chern-Simons term
- LHS is T-reversal invariant (with anomaly) and RHS seems like it is not. Its IR behavior is debated.

#### The supersymmetric route

- Many dualities of  $4d \mathcal{N} = 1$  theories (IR dualities) [NS; ...]
- They motivated many dualities in 3d
  - $-\mathcal{N}=2$  [Aharony, Hanany, Intriligator, NS, Strassler; Aharony; Giveon, Kutasov; ... Benini, Closset, Cremonesi; Intriligator, NS; Aharony, Razamat, NS, Willett; Park, Park; ...]
  - $-\mathcal{N}=4$  3d mirror symmetry [Intriligator, NS; ...]
- These use particle/vortex duality
- Later derived by compactification of 4d N = 1 dualities on a circle and then flow with relevant operators [Aharony, Razamat, NS, Willett].
  - More checks
  - Leads to many new dualities

## The supersymmetric route

#### The supersymmetric route

- Many checks using supersymmetry and localization
- Related to string duality
- Connected to level/rank duality of 3d topological quantum field theory and 2d RCFT (rigorous [...; Hsin, NS]).
- Can flow from them to non-supersymmetric theories [Jain, Minwalla, Yokoyama; Gur-Ari, Yacoby; ...].
  - This motivates non-supersymmetric dualities.
  - But the flow might not be smooth.

#### The AdS/CFT, large N route

- Same 4d Vasiliev theory is dual to two different 3d field theories [Vasiliev; Sezgin, Sundell; Klebanov, Polyakov; Giombi, Yin; Aharony, Gur-Ari, Yacoby; Giombi, Minwalla, Prakash, Trivedi, Wadia, Yin].
  - Scalars coupled to a Chern-Simons gauge theory
  - Fermions coupled to a Chern-Simons gauge theory
  - Hence, a purely field theoretic duality between them
  - Many explicit checks of this duality at large N [Maldacena, Zhiboedov; Aharony, Giombi, Gur-Ari, Maldacena, Yacoby; Jain, Minwalla, Sharma, Takimi, Wadia, Yokoyama; Minwalla, Yokoyama; Yokoyama; Jain, Mandlik, Minwalla, Takimi, Wadia, Yokoyama; Inbasekar, Jain, Mazumdar, Minwalla, Umesh, Yokoyama; ...]

### Synthesis [Aharony]

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3 [Aharony] + 1 [Hsin, NS] conjectures

 $N_f$  scalars at  $|\Phi|^4$  point coupled to

 $N_f$  fermions coupled to

• 
$$SU(N)_k$$

$$\leftrightarrow$$

$$U(k)_{-N+\frac{N_f}{2},-N+\frac{N_f}{2}}$$

• 
$$U(N)_{k,k}$$

$$\leftrightarrow$$

$$SU(k)_{-N+\frac{N_f}{2},-N+\frac{N_f}{2}}$$

• 
$$U(N)_{k,k+N}$$

$$\leftrightarrow$$

$$U(k)_{-N+\frac{N_f}{2},-N-k+\frac{N_f}{2}}$$

• 
$$U(N)_{k,k-N}$$

$$\leftrightarrow$$

$$U(k)_{-N+\frac{N_f}{2},-N+k+\frac{N_f}{2}}$$

Fits the large N picture  $(N, k \to \infty)$  with finite N/k

Fits the supersymmetric picture

Related to level/rank duality

Baryon and monopole operators match [Radicevic]

#### **Puzzles**

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- Is it true for all N, k, N<sub>f</sub>?
- How can a theory of bosons, which does not need a spin structure be dual to a theory of fermions, which needs it?
- What is the relation to the dualities in the condensed matter literature (with puzzles about quantization of coefficients, Treversal invariance, etc.)?
- What is the precise statement of the dualities (including the coupling to background gauge fields and their Chern-Simons counterterms)?
- Are the assumptions independent? Can we assume some of these dualities and derive others?
- Are there other such dualities?

## Examine $N = k = N_f = 1$

$$i\bar{\Psi}D/A\Psi \quad \leftrightarrow \quad |D_b\Phi|^2 - |\Phi|^4 + \frac{1}{4\pi}bdb + \frac{1}{2\pi}Adb$$

## Examine $N = k = N_f = 1$

Assume: a free fermion coupled to a background A is dual to a gauged Wilson-Fisher fixed point with Chern-Simons interaction for the dynamical field b

$$i\bar{\Psi}D/A\Psi \quad \leftrightarrow \quad |D_b\Phi|^2 - |\Phi|^4 + \frac{1}{4\pi}bdb + \frac{1}{2\pi}Adb$$

- Same symmetries
  - $-U(1)_{A}$
  - T-reversal invariance (with anomaly) of the RHS follows from particle/vortex duality. T-reversal is a quantum symmetry there. (More below.)
- Conversely, assuming this duality we derive the known particle/vortex duality.

$$i\bar{\Psi}DA\Psi \quad \leftrightarrow \quad |D_b\Phi|^2 - |\Phi|^4 + \frac{1}{4\pi}bdb + \frac{1}{2\pi}Adb$$

#### Mapping of operators

- $\Psi \leftrightarrow \Phi^+ \mathcal{M}_b$  with  $\mathcal{M}_b$  a monopole operator (charged under  $U(1)_A$ )
- Mass term  $\overline{\Psi}\Psi \leftrightarrow |\Phi|^2$
- Mass deformation leads to two phases depending on the sign...

$$i\bar{\Psi}D/A\Psi \quad \leftrightarrow \quad |D_b\Phi|^2 - |\Phi|^4 + \frac{1}{4\pi}bdb + \frac{1}{2\pi}Adb$$

- Mass deformation depends on the sign:
  - $-\Phi$  is massive with spin  $\frac{1}{2}$ . It is charged under  $U(1)_b$  and  $U(1)_A$ . It is mapped to the massive Ψ. (Statistical transmutation of massive particles.)
  - $-U(1)_b$  is Higgsed. Vortex of spin  $-\frac{1}{2}$  is charged under  $U(1)_A$ . It is mapped to the massive Ψ.
- T changes the sign of the fermion mass (= boson mass square) and maps Φ particles to vortices.

$$i\bar{\Psi}D/A\Psi \quad \leftrightarrow \quad |D_b\Phi|^2 - |\Phi|^4 + \frac{1}{4\pi}bdb + \frac{1}{2\pi}Adb$$

- Here both sides of the duality need a spin structure (spinors in the LHS and odd Chern-Simons level in the RHS)
  - But if A is a spin<sub>c</sub> connection,

$$\int \frac{dA}{2\pi} = \frac{1}{2} \int w_2 \operatorname{mod} \mathbf{Z} ,$$

there is no need for a spin structure on either side.

#### Derive many other dualities

Starting with

$$i\bar{\Psi}D/A\Psi \quad \leftrightarrow \quad |D_b\Phi|^2 - |\Phi|^4 + \frac{1}{4\pi}bdb + \frac{1}{2\pi}Adb$$

we can derive other dualities by changing the two sides:

- Add a Chern-Simons counterterm for the classical field A
- Gauge it by turning A into a dynamical field a and adding a new classical field.
- Use other dualities.
- Repeat.

#### Another boson/fermion duality

For example, derive:

$$|D_B\Phi|^2 - |\Phi|^4 \quad \leftrightarrow \quad i\bar{\Psi}D_a\Psi + \frac{1}{2\pi}Bda - \frac{1}{4\pi}BdB$$

LHS Wilson-Fisher fixed point

RHS QED with a single fermion, a.k.a  $U(1)_{1/2}$ 

- Derived from the other duality
- Neither side needs a spin structure when a is a spin<sub>c</sub> connection
- Need a Chern-Simons counterterm for B
- Can map the operators and check the phases
- RHS is T-reversal invariant (quantum symmetry)

#### A fermion/fermion duality

•Derive  $i \bar{\Psi} D\!\!\!\!/_A \Psi \quad \leftrightarrow$ 

$$i\bar{\chi}D\!\!\!/_a\chi - \frac{2}{4\pi}bdb + \frac{1}{2\pi}adb + \frac{1}{2\pi}Adb - \frac{1}{4\pi}AdA$$

LHS free fermion

RHS QED with a single fermion, coupled to  $U(1)_{-2}$  of b.

- If we incorrectly integrate out b, we find the previously mentioned version with improperly quantized Chern-Simons terms.
- No need for a spin structure when a and A are spin<sub>c</sub> connections.
- Can map the operators and check the phases
- T-reversal invariance (with anomaly) is manifest in LHS. It acts non-trivially in the RHS (quantum symmetry).

#### More

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- Many more dualities and relations between them
  - ✓ Add gravitational Chern-Simons counterterms (more checks)
  - ✓ Relation to 4d S-duality in half-space with these 3d theories on the boundary (Witten's S and T operations on 3d field theories)
  - ✓ Generalization to arbitrary N and k
    - Using a precise version of level/rank duality
    - Problem with large  $N_f$
    - Leads to many more dualities
  - Much more can be done