

# Solving a 40-year-old Problem: 11D Superfields

Yangrui Hu and S.-N. Hazel Mak

Brown University

Strings 2021 Gong Show

In collaboration with: S. James Gates, Jr.

[arXiv: [1911.00807](https://arxiv.org/abs/1911.00807), [2002.08502](https://arxiv.org/abs/2002.08502), [2006.03609](https://arxiv.org/abs/2006.03609), [2007.05097](https://arxiv.org/abs/2007.05097), [2007.07390](https://arxiv.org/abs/2007.07390)]

# Motivation

Lie algebras & Mathematica<sup>1</sup>



- 1974: first 4D superfield written down! [[Salam, Strathdee 1974](#)]
- 1978: first 11D on-shell supergravity! [[Cremmer, Julia, Scherk 1978](#)]
- ⋮
- 2020: 11D superfield!!! [[Gates, YH, SNHM, arXiv:2002.08502](#)]

---

<sup>1</sup>LieART [[Feger, Kephart, 2012](#)], SUSYno [[Fonseca, 2011](#)]

# The Question

In 11D, we have 32 Grassmann coordinate  $\theta$ 's. Since  $\theta^2 = 0$ ,

$$\mathcal{V}(x, \theta) = v^{(0)}(x) + \sum_{n=1}^{32} v_{\alpha_1 \dots \alpha_n}^{(n)}(x) \theta^{\alpha_1} \dots \theta^{\alpha_n}$$

## Question:

How to write **components** /  **$\theta$ -monomials** in  $\mathfrak{so}(1, 10)$  irreducible representations?

## Difficulties:

- $2^{32} = 4, 294, 967, 296$  total degrees of freedom
- many Fierz identities...
- Lorentz covariant  $\theta$ -monomials at higher levels

# Step 1: Young Tableaux

Totally antisymmetric product as YT:

$$\theta^{\alpha_1} \dots \theta^{\alpha_n} \Leftrightarrow \underbrace{\{32\} \wedge \dots \wedge \{32\}}_{n \text{ times}} \Leftrightarrow \begin{array}{|c|} \hline 32 \\ \hline \vdots \\ \hline 32 \\ -n+1 \\ \hline \end{array}$$

Theorem:

$$\text{YT} \xLeftrightarrow{1-1} \text{Irreducible representation of } \mathfrak{su}(32)$$

## Step 2: Branching Rules

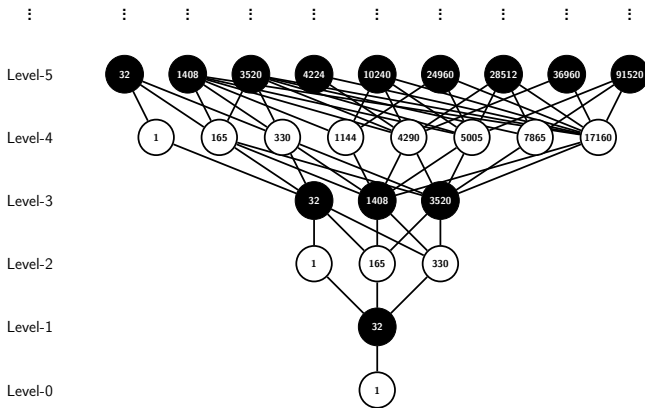
$$\mathfrak{su}(32) \supset \mathfrak{so}(1, 10) \Rightarrow \mathcal{R}_{\mathfrak{su}(32)} \xrightarrow{\text{branching rules}} \bigoplus \mathcal{R}_{\mathfrak{so}(1,10)}$$

Examples:

$$\begin{array}{ccccccc}
 \{496\} & \longrightarrow & \{1\} & \oplus & \{165\} & \oplus & \{330\} \\
 \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} & & \cdot & & \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} & & \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \square \\ \hline \end{array} \\
 & & \Phi & & \Phi_{[abc]} & & \Phi_{[abcd]} \\
 \\ 
 \{4,960\} & \longrightarrow & \{32\} & \oplus & \{1,408\} & \oplus & \{3,520\} \\
 \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \square \\ \hline \end{array} & & \square & & \begin{array}{|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & \text{IR} & \begin{array}{|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} & \text{IR} \\
 & & \Psi_{\alpha} & & \Psi_{[ab]\alpha} & & \Psi_{[abc]\alpha}
 \end{array}$$

Dictionary & Graphical Rules in [\[Gates, YH, SNHM, arXiv:2006.03609\]](#)

# 11D Scalar Superfield



- Level-0 to level-32, symmetric about level-16
- Breitenlohner's method:  $\mathcal{V}_{[\mathcal{R}]} = \mathcal{V} \otimes [\mathcal{R}]$  [Breitenlohner 1977]

# 11D Supergravity Surprise!

- $\mathcal{V} \Rightarrow$  linearized Nordström SG [Gates, YH, Jiang, SNHM 2019]
- Semi-prepotential candidate:  $\mathcal{V} = D^\alpha \mathcal{V}_\alpha$

Physical Component	Irrep	Level
graviton $h_{ab}$	$\{65\}, \{1\}$	16
gravitino $\psi_a^\beta$	$\{320\}, \{32\}$	17
3-form $B_{[3]}$	$\{165\}$	16

- Prepotential candidate:  $\mathcal{V}_\alpha$   
contains 2-form  $h_{[ab]} = \{55\}$  at level-17  $\Rightarrow$  Poincaré vielbein

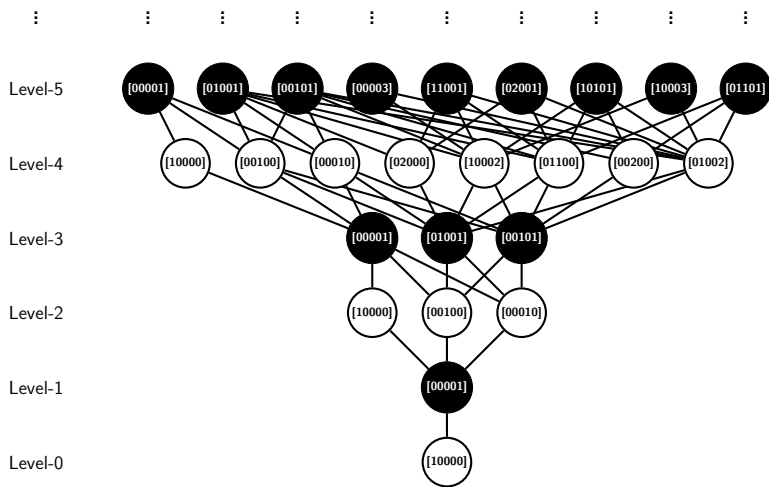
THANK YOU!!

## Backup: More Background

- Our work leads to a formalism demonstrating a *manifest linear* realization of  $\mathfrak{so}(1, 10)$ !
  - GS-formalism [Green and Sethi 1999]: started with a complex chiral spinor in  $\mathfrak{so}(1, 9) \Rightarrow$  coset symmetry  $\mathfrak{so}(1, 10)/\mathfrak{so}(1, 9)$  realized in some *non-linear* manner
  - [Becker and Butter 2003]: coset symmetry  $\mathfrak{so}(1, 10)/\mathfrak{so}(1, 3)$  realized in some *non-linear* manner
- A lot of previous work focused on modified SUGRA covariant derivatives, field strengths, Bianchi identities, etc. They are “orthogonal” to our current efforts.
- [Gates, Vashakidze 1987] showed the relationship between off-shell superfields and higher derivative terms in the *open string* effective action



# Backup: 11D Adynkra



## Backup: Poincaré vielbein & Gravitino

Decompositions of the inverse frame and gravitino fields in 11D yield

$$e_a^m = \{h_{(ab)} + \eta_{ab}h + h_{[ab]}\} \eta^{bm}$$
$$\{121\} \quad \{65\} \quad \{1\} \quad \{55\}$$

where  $h_{(ab)}$  is the conformal graviton,  $h$  is the trace, and  $h_{[ab]}$  is the two-form; and

$$\tilde{\psi}_a^\alpha = \psi_a^\alpha - \frac{1}{11}(\gamma_a)^{\alpha\beta}\psi_\beta$$
$$\{352\} \quad \{320\} \quad \{32\}$$

where  $\psi_a^\alpha$  is the conformal gravitino and  $\psi_\beta \equiv (\gamma^a)_{\alpha\beta}\tilde{\psi}_a^\alpha$  is the  $\gamma$ -trace.