

The large charge expansion

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Introduction

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Study subsectors with large charge Q .

Large charge Q becomes controlling parameter in a perturbative expansion!

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Consider systems with large quantum number
many degrees of freedom

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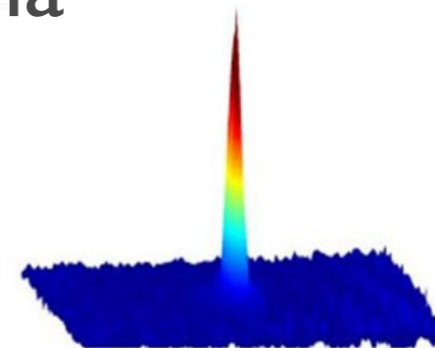
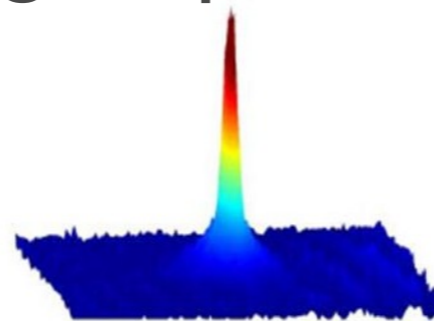


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Consider systems with large quantum number
many degrees of freedom



emergent phenomena

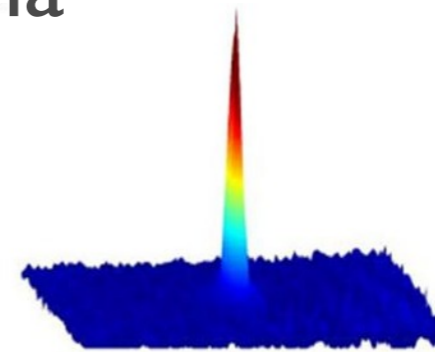
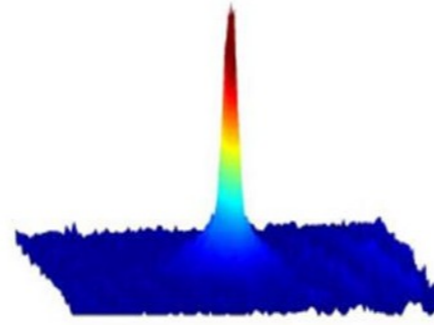


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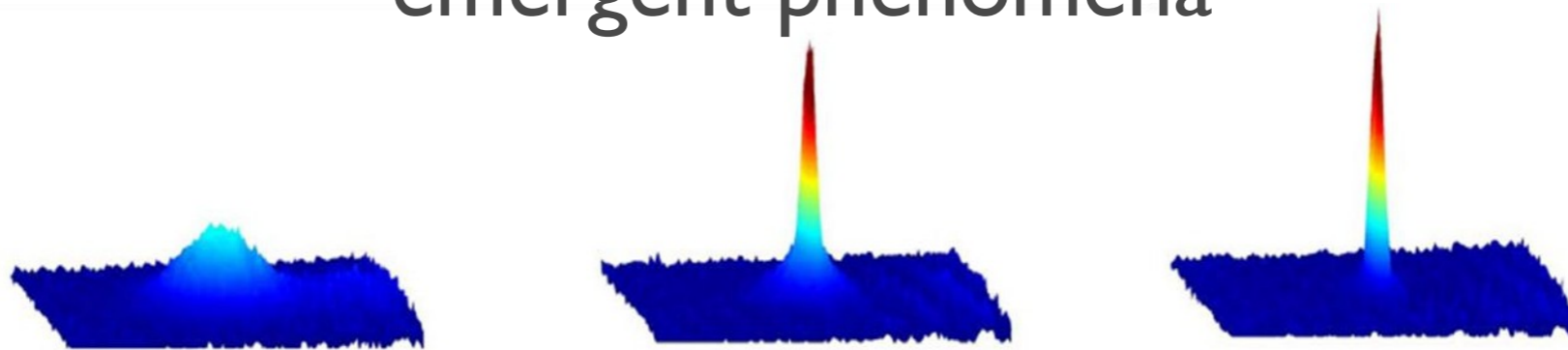


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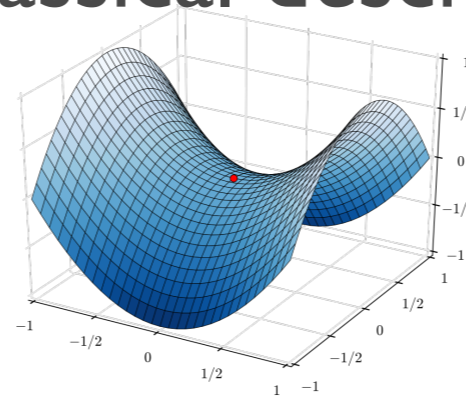
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emergent phenomena



semiclassical description

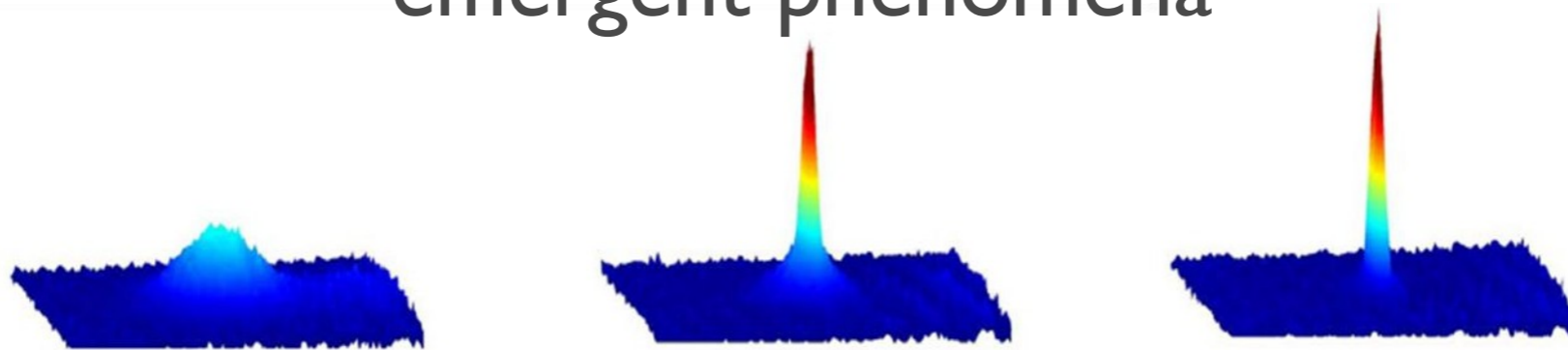


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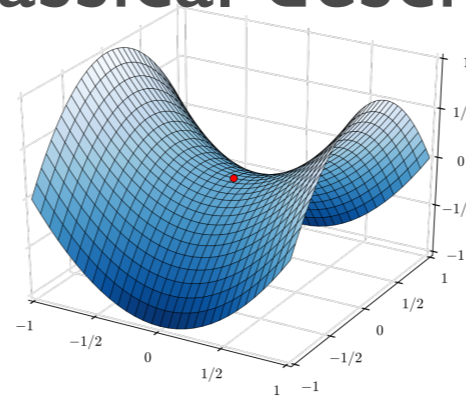
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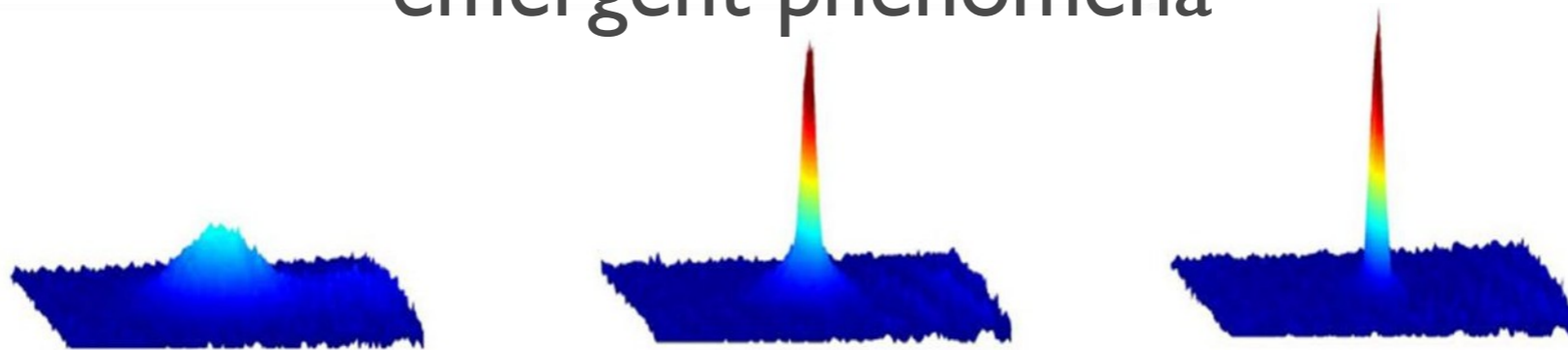
e.g. superfluid

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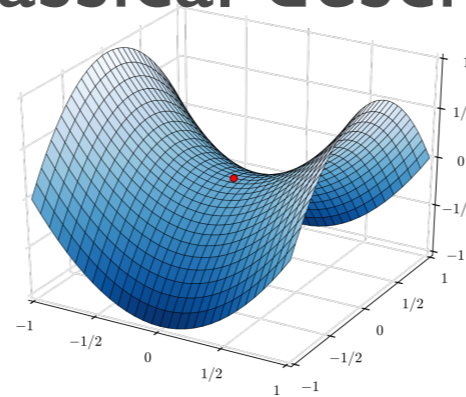
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e.g. superfluid

works also for strongly coupled systems!

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accessible?

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no



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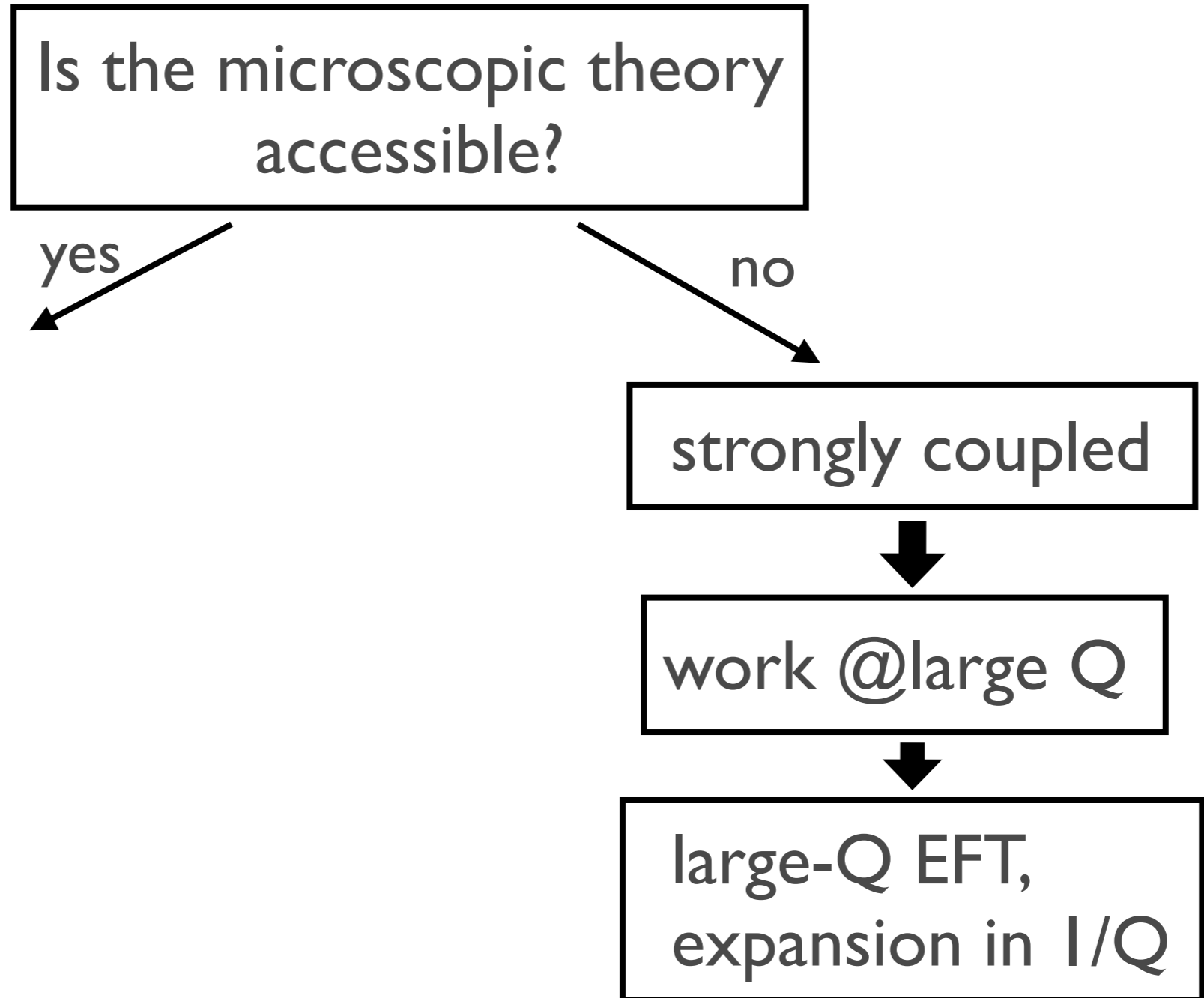
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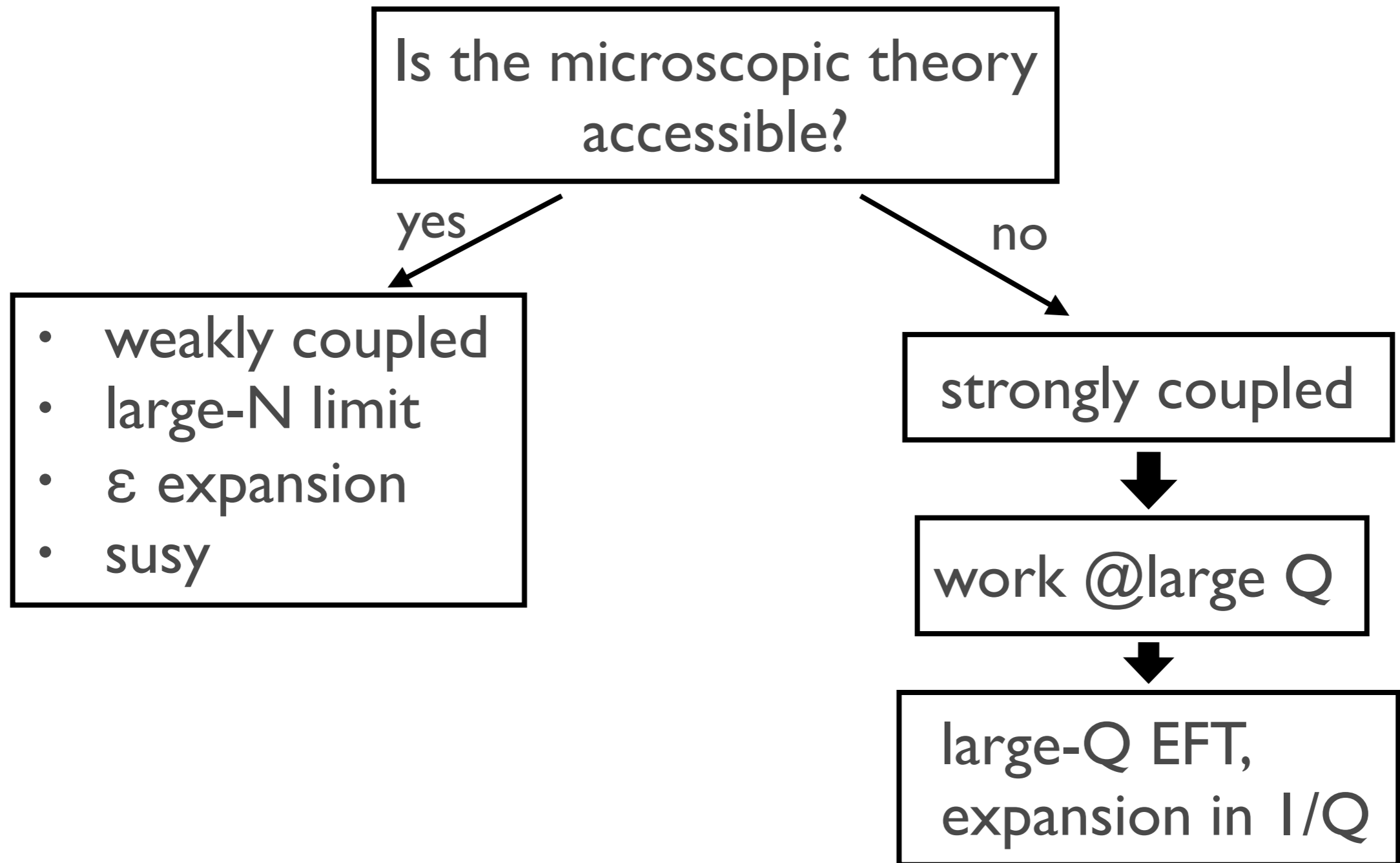
work @large Q

large- Q EFT,
expansion in $1/Q$

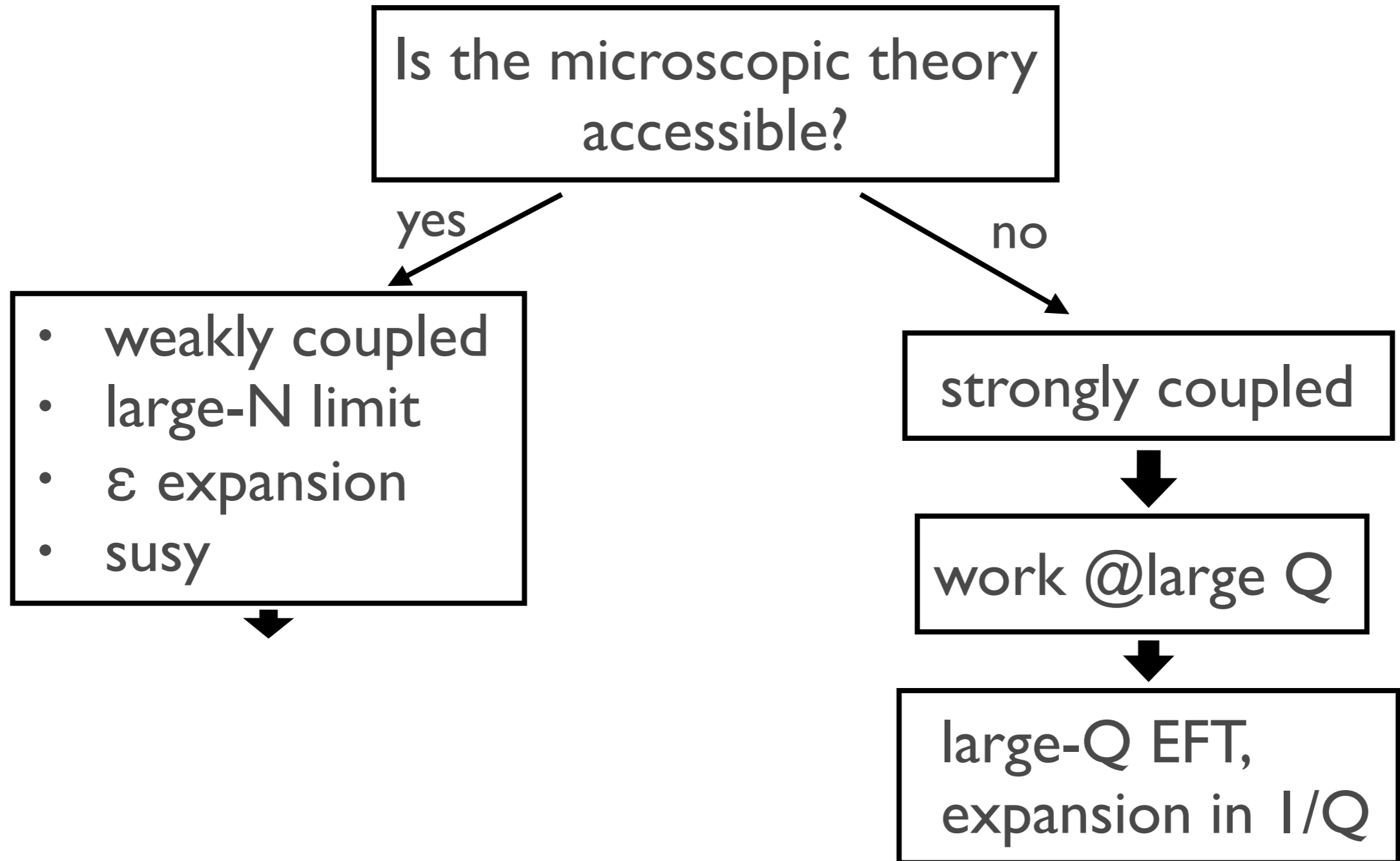
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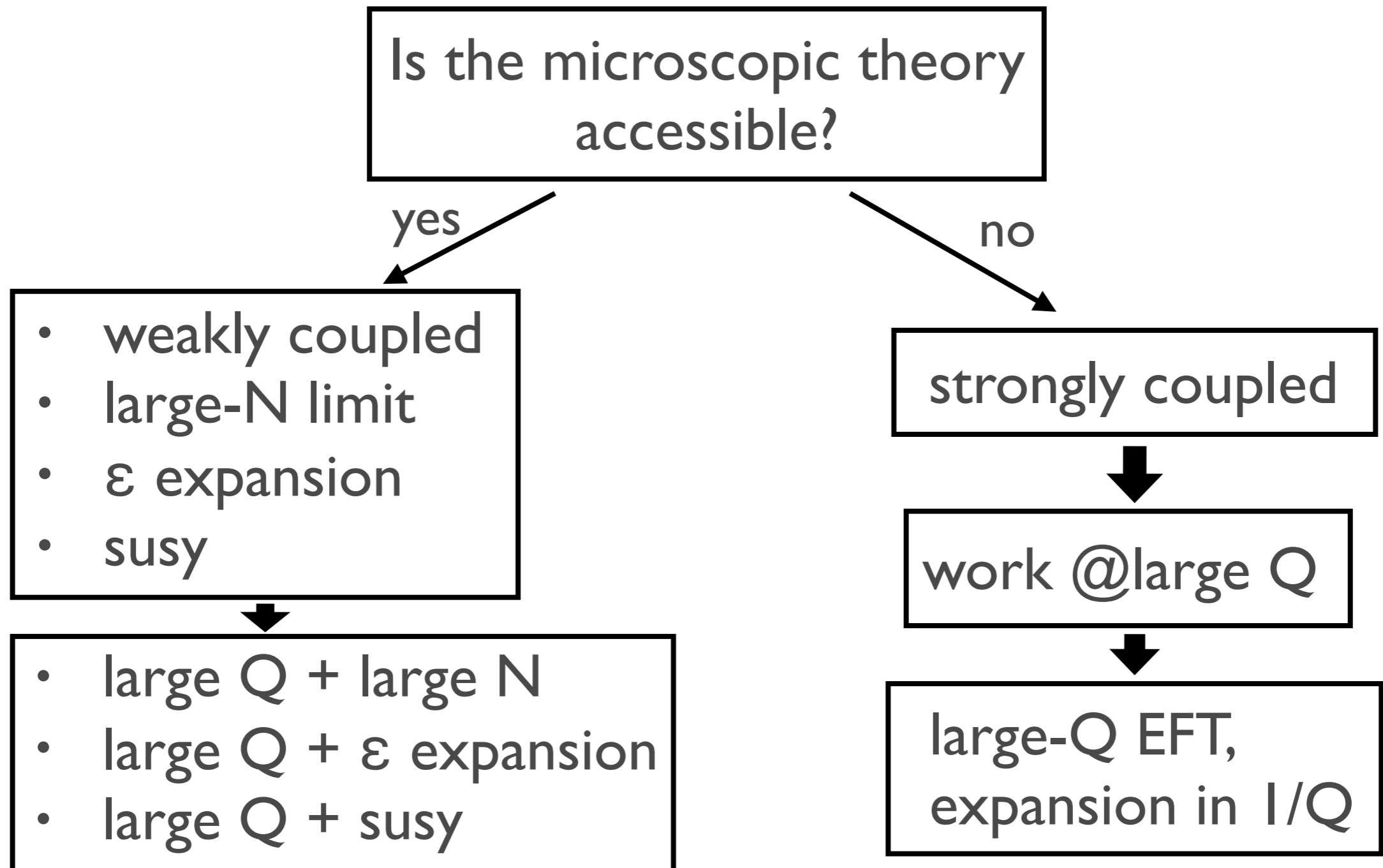
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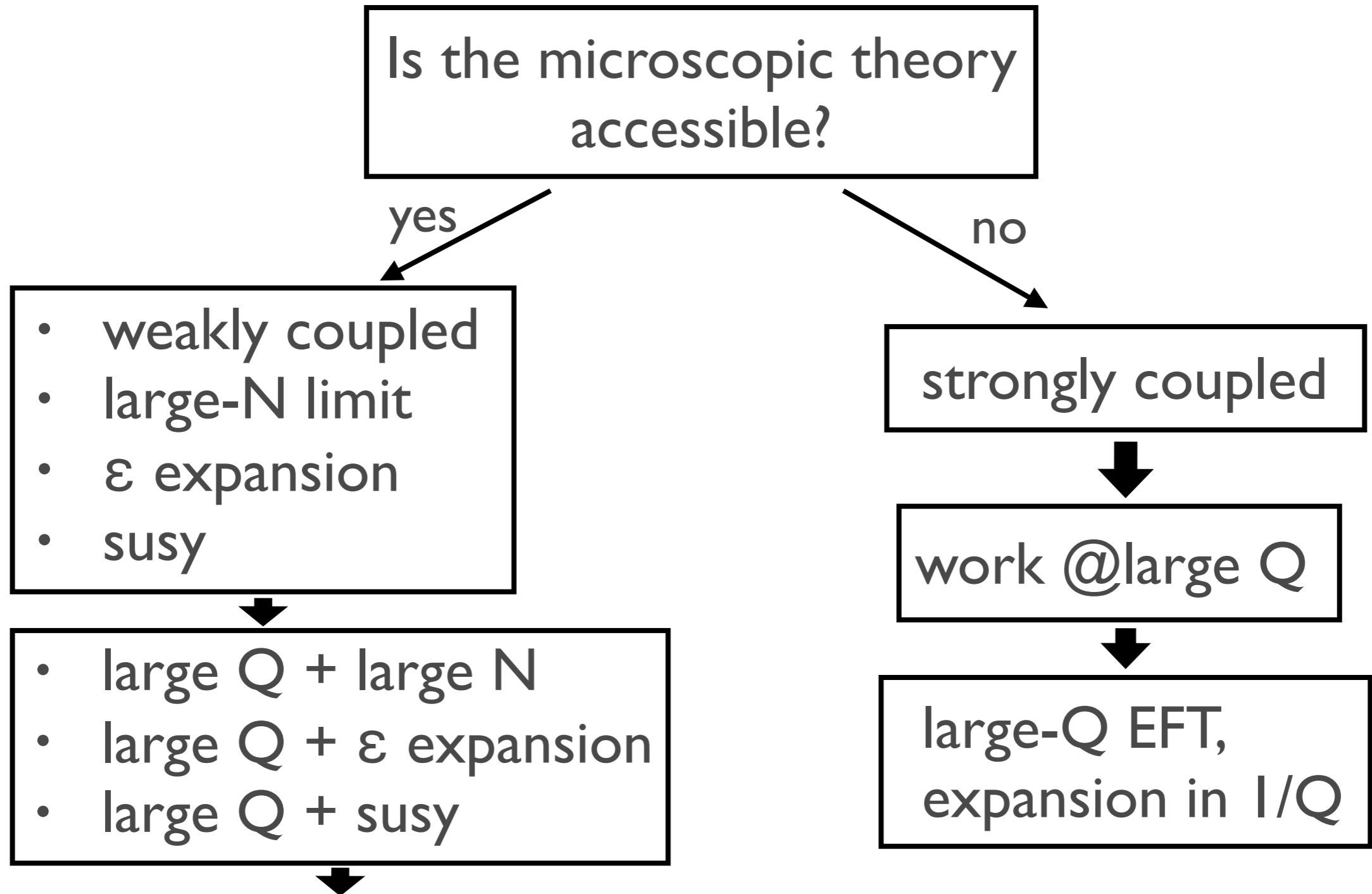
Introduction



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Introduction

Is the microscopic theory accessible?

yes

no

- weakly coupled
- large-N limit
- ϵ expansion
- susy

strongly coupled

- large Q + large N
- large Q + ϵ expansion
- large Q + susy

work @large Q

large- Q EFT,
expansion in $1/Q$

go beyond perturbation theory in $1/Q$, calculate non-perturbative (exponential) corrections!

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Superfluid

isolated vacuum

- Wilson-Fisher CFT
- NRCFT (unitary Fermi gas)
- $N=2$ SCFT in 3d
- asymptotically safe model in 4d
- Gross-Neveu model

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EFT of the moduli space

moduli space of vacua

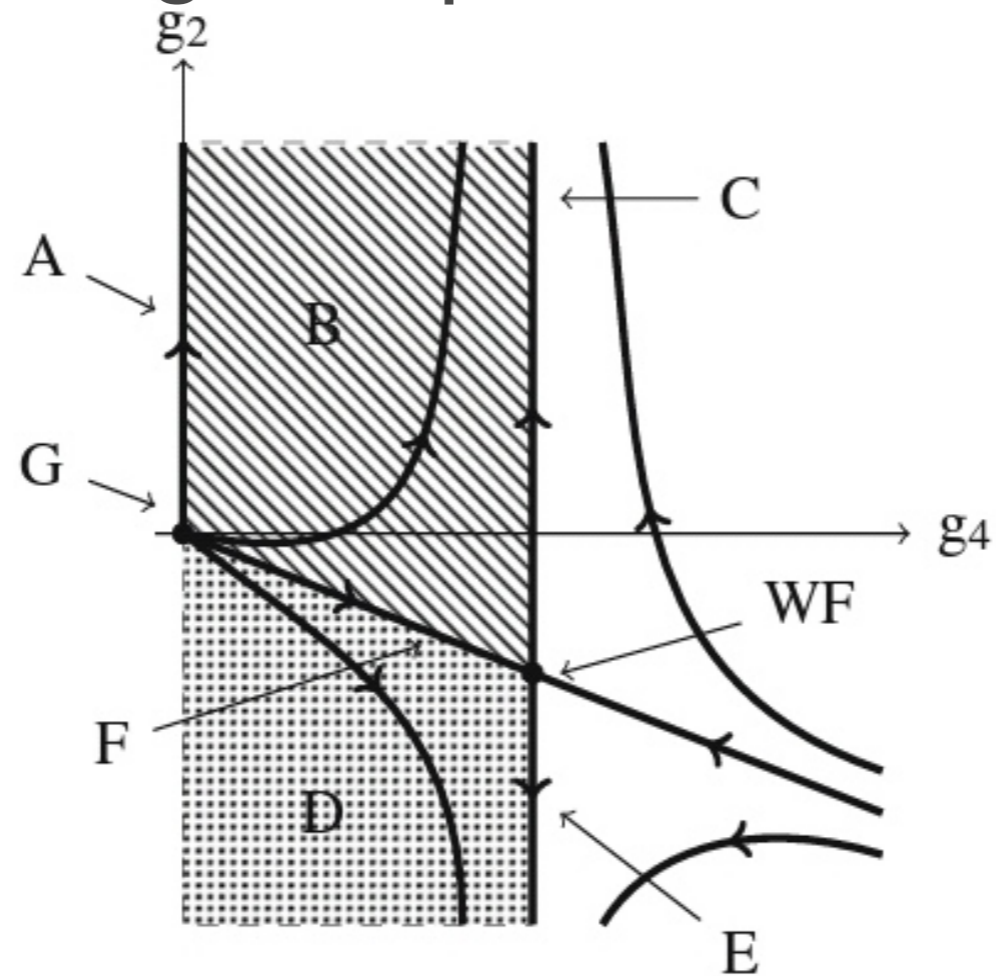
- free boson
- $N=2$ theories in 4d

Introduction

Example: Scalar field theories in $2 < D < 4$ have a strongly-coupled interacting fixed point, the **Wilson-Fisher FP**.

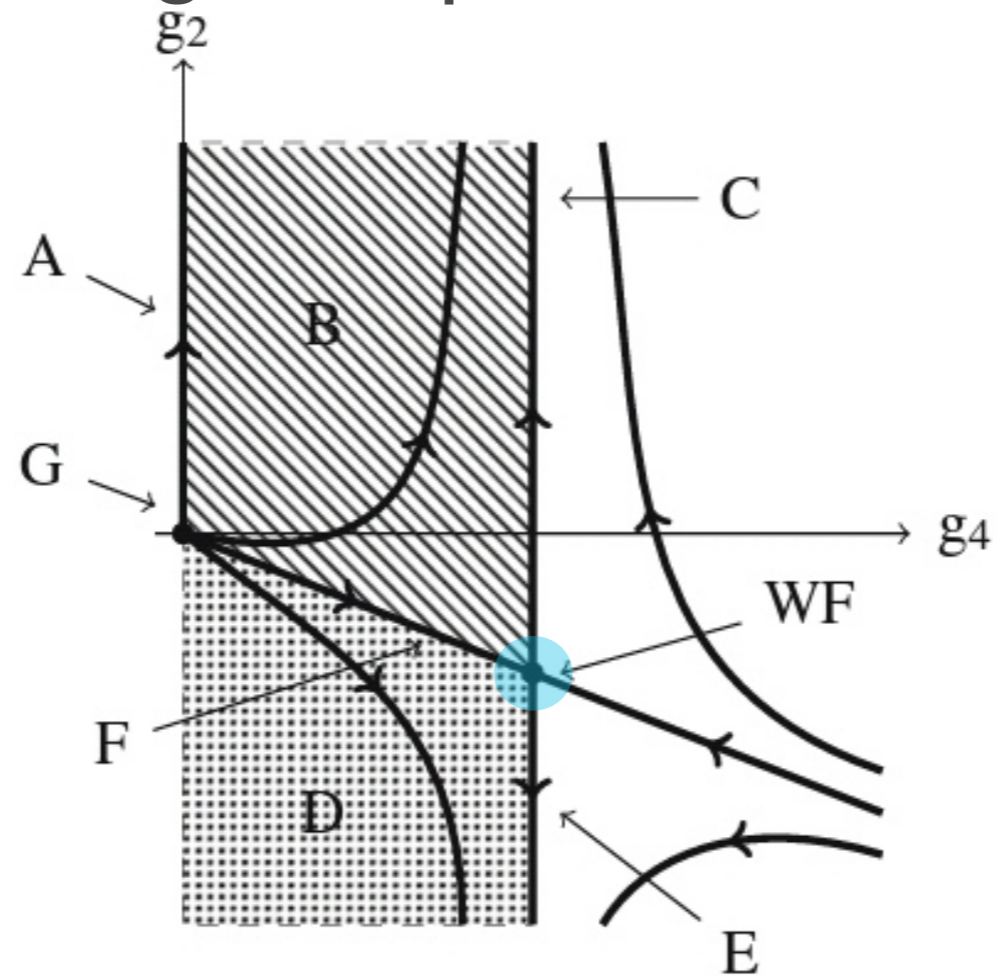
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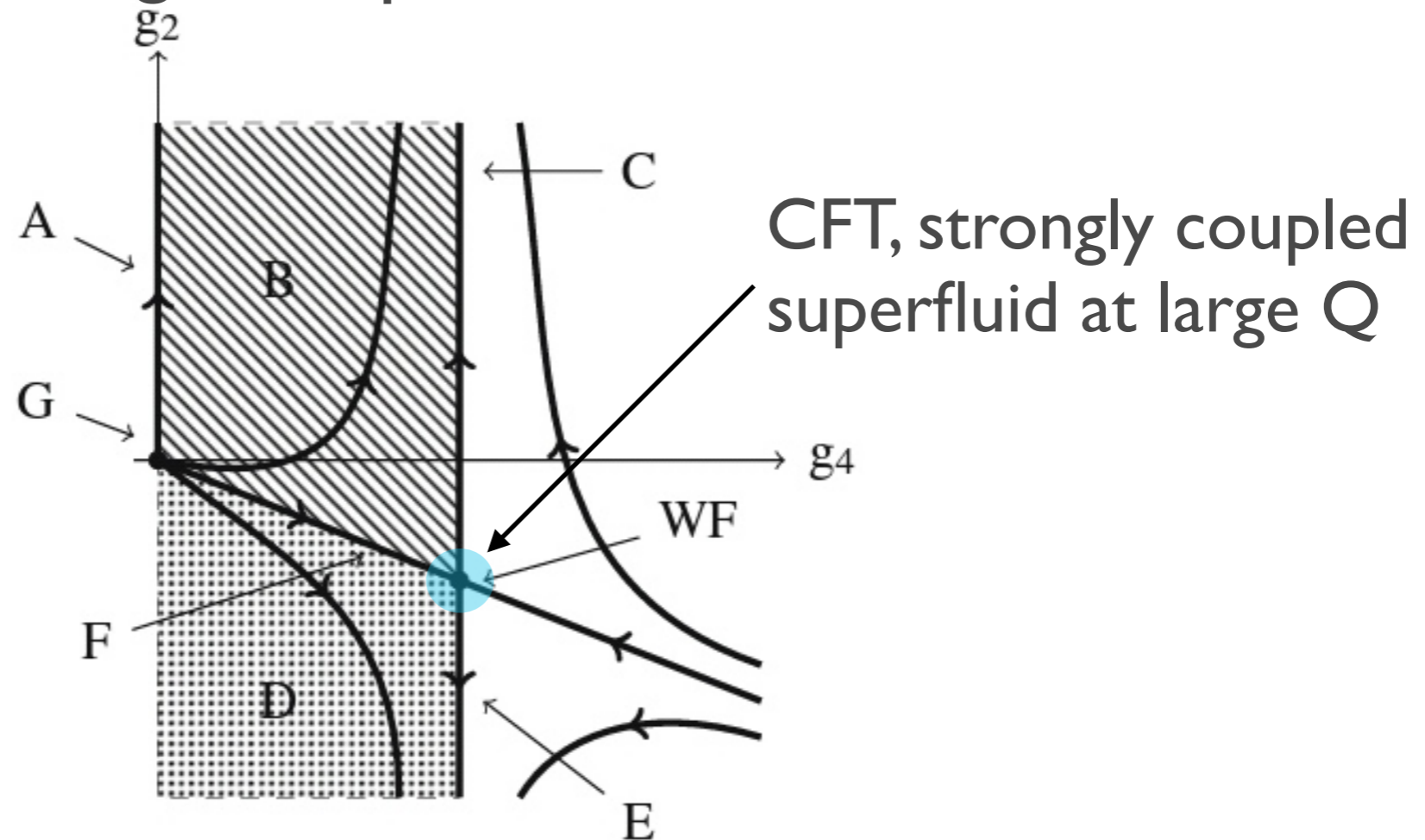
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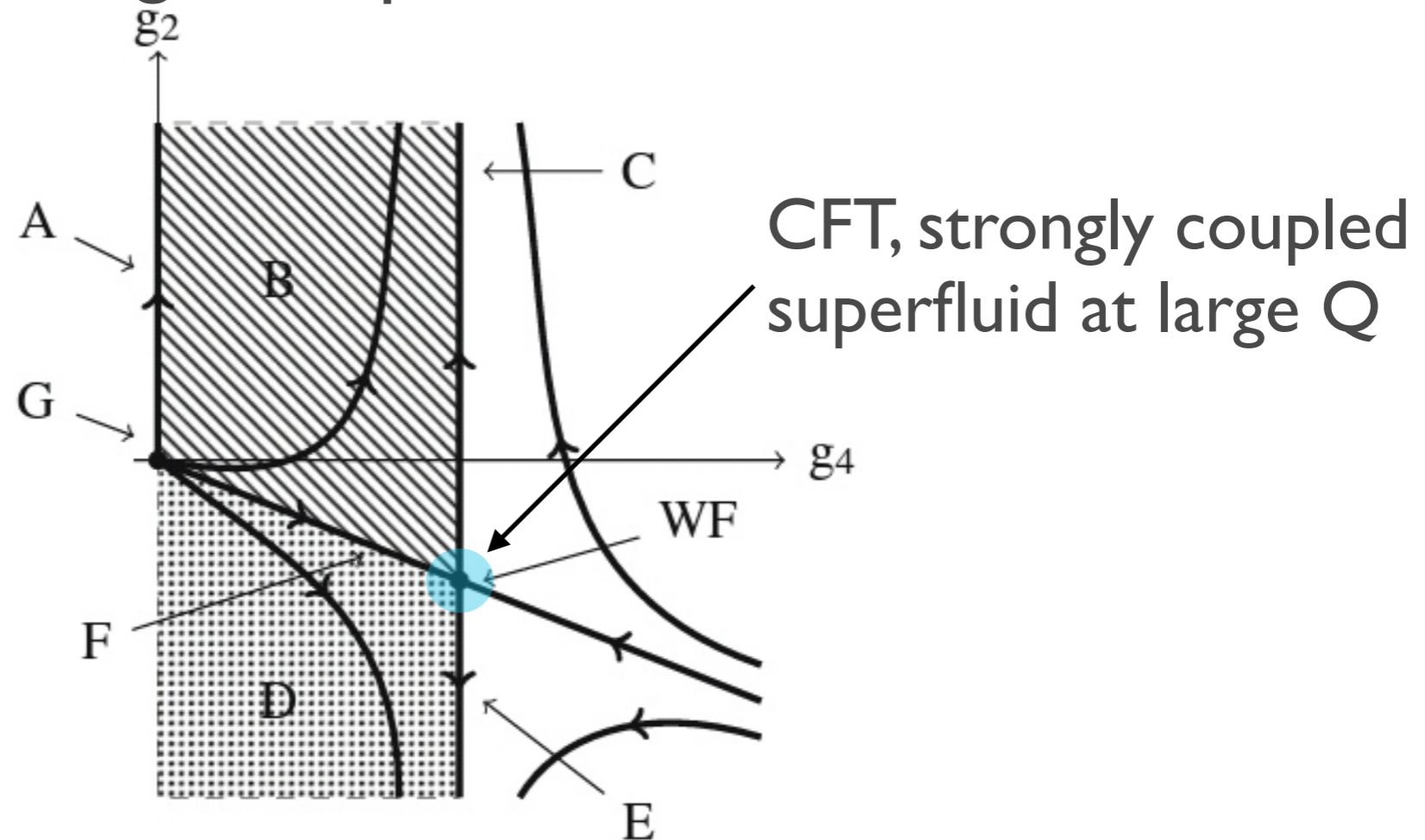
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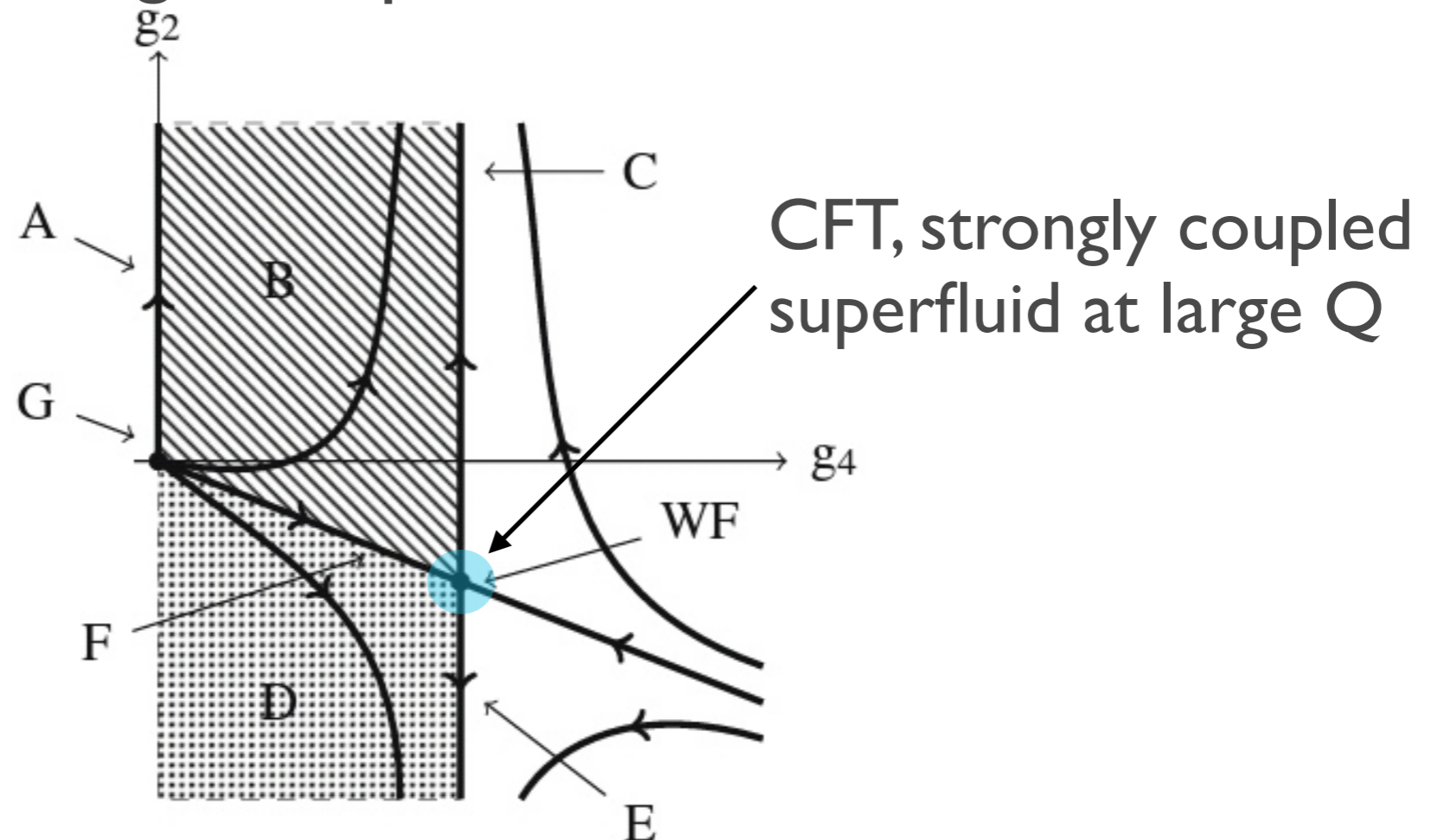
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$O(2N)$ vector model in $D=3$:

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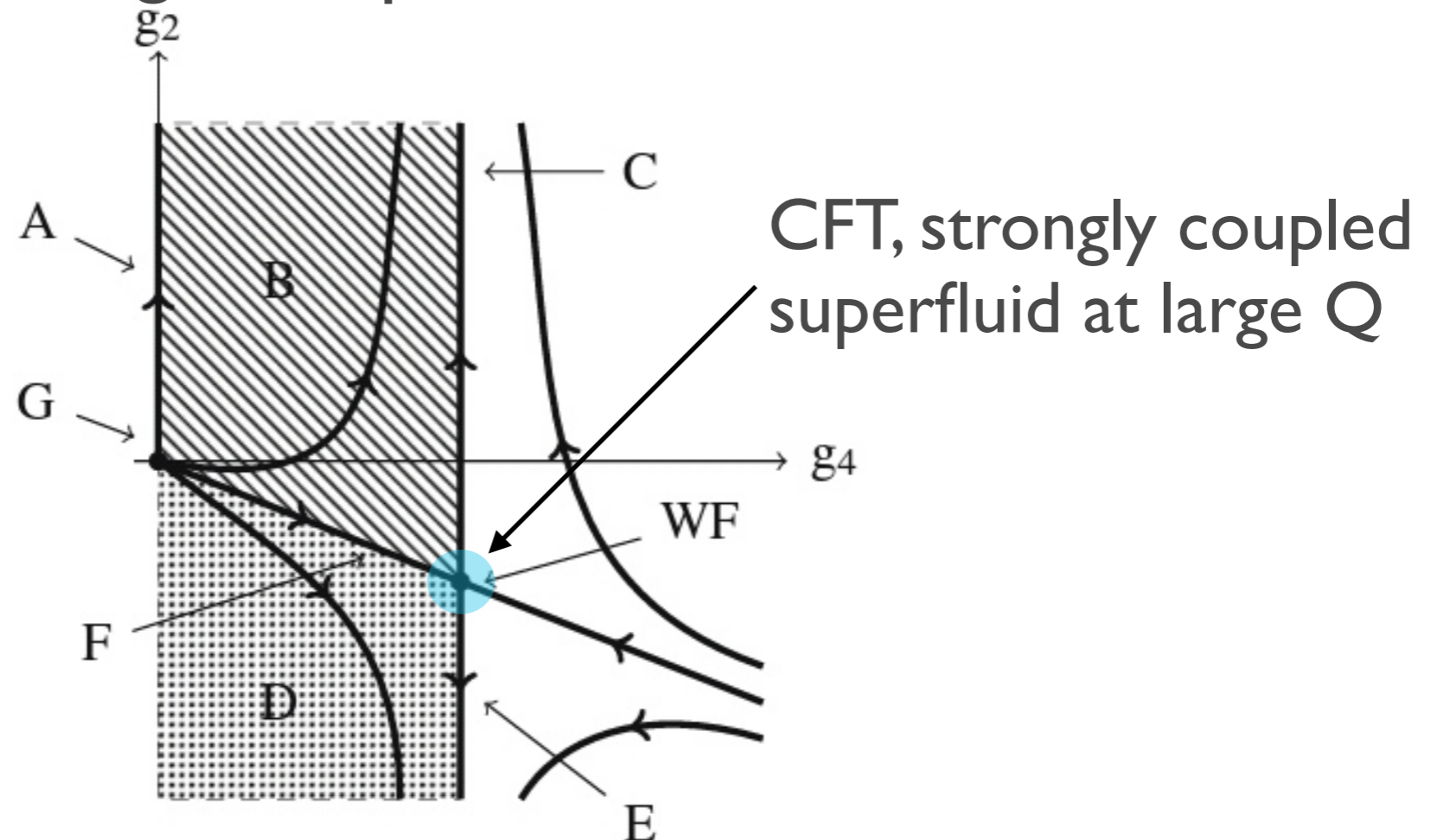


O(2N) vector model in D=3:

$$S[\phi_i] = \sum_{i=1}^N \int dt d\Sigma \left[g^{\mu\nu} (\partial_\mu \phi_i)^\dagger (\partial_\nu \phi_i) + r (\phi_i^\dagger \phi_i) + \frac{u}{2N} (\phi_i^\dagger \phi_i)^2 \right]$$

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For $r=R/8$, this flows to the WF fixed pt in the IR, $u \rightarrow \infty$

The $O(2)$ model

Simplest example: $O(2)$ model in $(2+1)$ dimensions

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$$\mathcal{L}_{UV} = \partial_\mu \phi^* \partial^\mu \phi - g^2 (\phi^* \phi)^2$$

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Assume that also the IR DOF are encoded by **cplx scalar**

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Look at scales: put system in box (2-sphere) of scale R
Second scale given by U(1) charge Q:

$$\rho^{1/2} \sim Q^{1/2} / R$$

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8


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Fixing the charge breaks symmetries:

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Beyond LO: use **dimensional analysis, parity and scale invariance** to determine (tree-level) operators in effective action (Lorentz scalars of scaling dimension 3, including couplings to geometric invariants)

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Use ρ -scaling to determine which terms are not suppressed:

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$$\mathcal{L} = k_{3/2}(\partial_\mu\chi\partial^\mu\chi)^{3/2} + k_{1/2}R(\partial_\mu\chi\partial^\mu\chi)^{1/2} + \mathcal{O}(Q^{-1/2})$$

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↙ LO Lagrangian ↙ curvature coupling

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↖ suppressed by inverse powers of Q

The O(2) model

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 $\partial\chi \sim \rho^{1/2}, \quad \partial \dots \partial\chi \sim \rho^{-1/4}$

Result for NLSM action in D=3:

$$\mathcal{L} = \kappa_{3/2}(\partial_\mu\chi\partial^\mu\chi)^{3/2} + \kappa_{1/2}R(\partial_\mu\chi\partial^\mu\chi)^{1/2} + \mathcal{O}(Q^{-1/2})$$

Annotations:

- LO Lagrangian (points to $\kappa_{3/2}(\partial_\mu\chi\partial^\mu\chi)^{3/2}$)
- curvature coupling (points to $\kappa_{1/2}R(\partial_\mu\chi\partial^\mu\chi)^{1/2}$)
- dimensionless parameters (points to $\kappa_{3/2}$ and $\kappa_{1/2}$)
- suppressed by inverse powers of Q (points to $\mathcal{O}(Q^{-1/2})$)

The O(2) model

Use ρ -scaling to determine which terms are not suppressed:
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Result for NLSM action in D=3:

$$\mathcal{L} = k_{3/2}(\partial_\mu\chi\partial^\mu\chi)^{3/2} + k_{1/2}R(\partial_\mu\chi\partial^\mu\chi)^{1/2} + \mathcal{O}(Q^{-1/2})$$

Annotations for the equation above:

- LO Lagrangian (points to the first two terms)
- curvature coupling (points to the R term)
- dimensionless parameters (points to $k_{3/2}$ and $k_{1/2}$)
- suppressed by inverse powers of Q (points to $\mathcal{O}(Q^{-1/2})$)

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LO Lagrangian

curvature coupling

dimensionless parameters

suppressed by inverse powers of Q

Energy of classical ground state at fixed charge:

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Use ρ -scaling to determine which terms are not suppressed:
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vacuum + Goldstone + 1/Q-suppressed corrections

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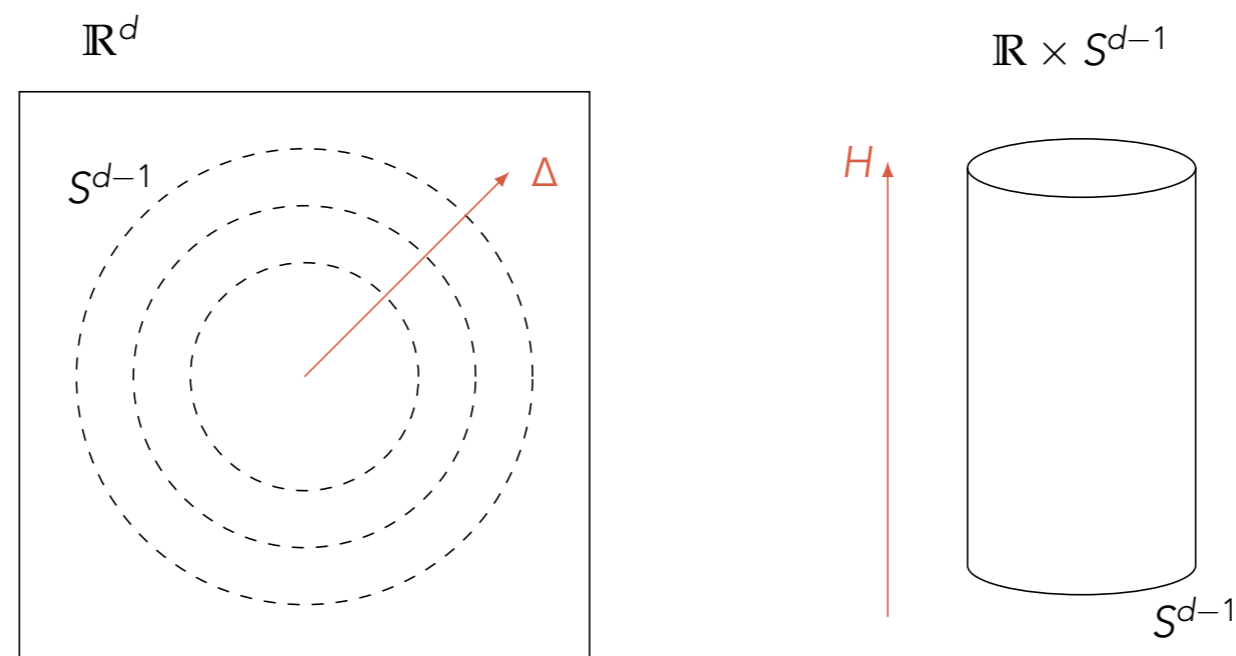
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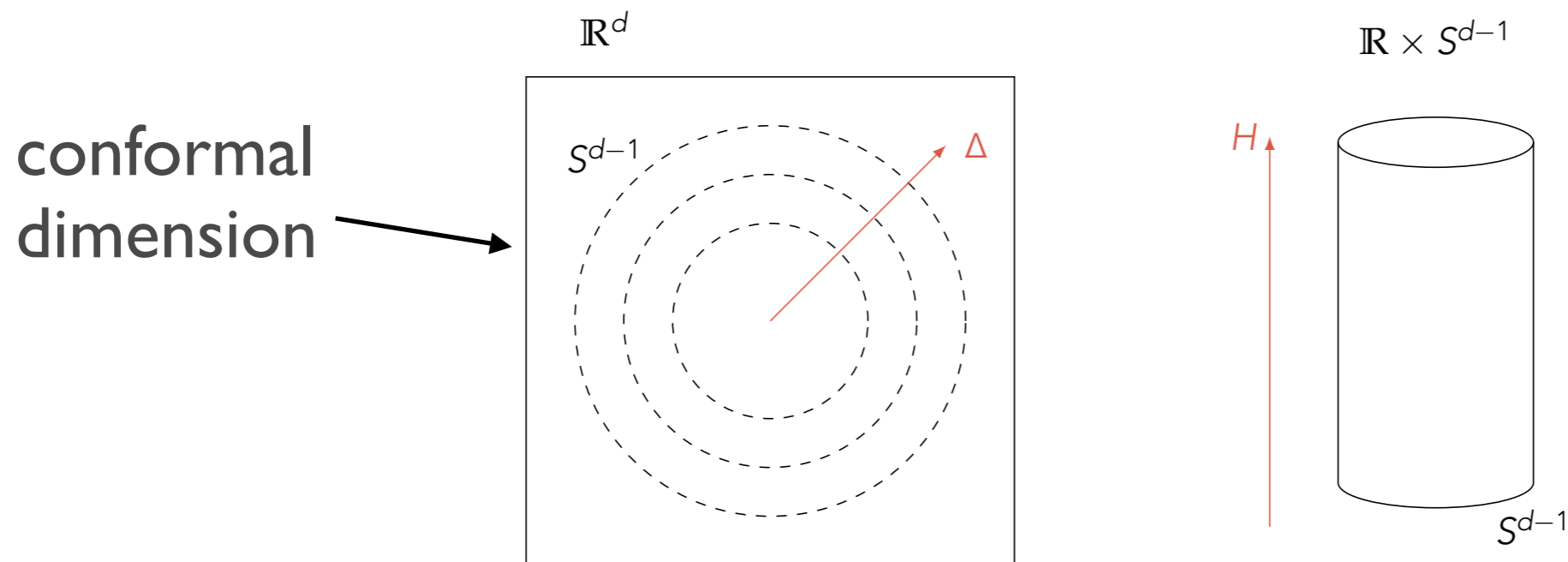
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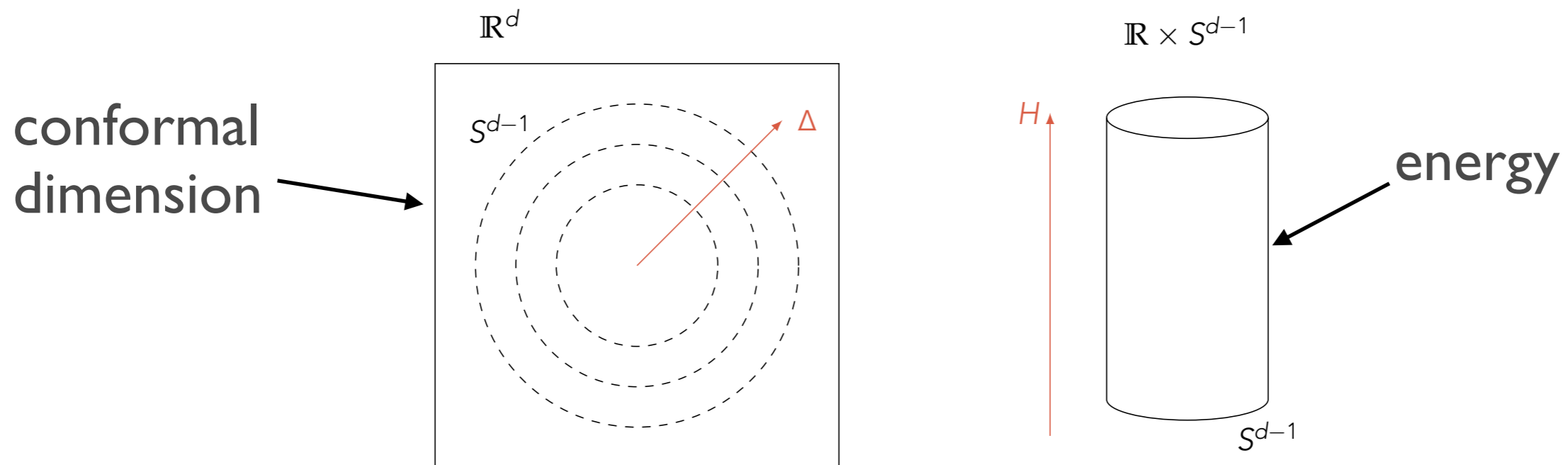
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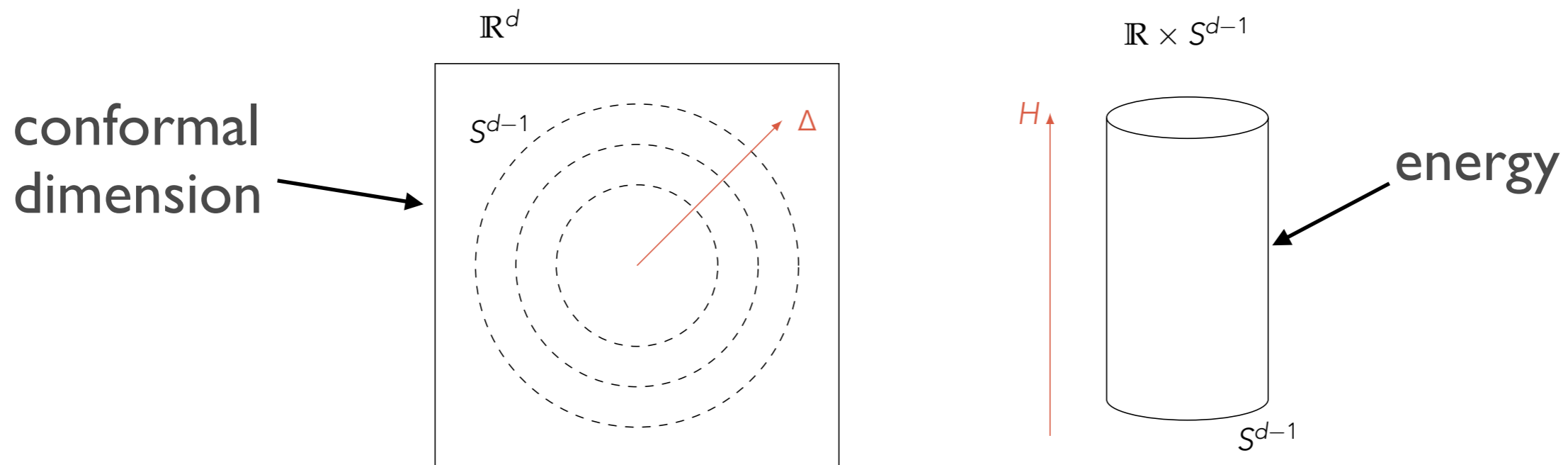
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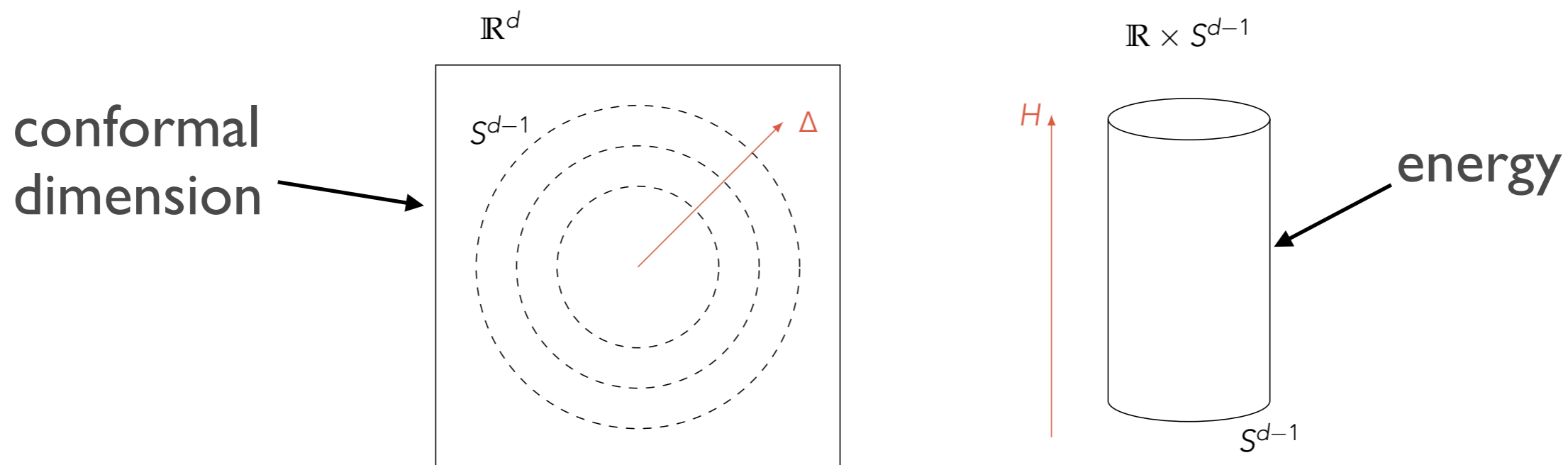


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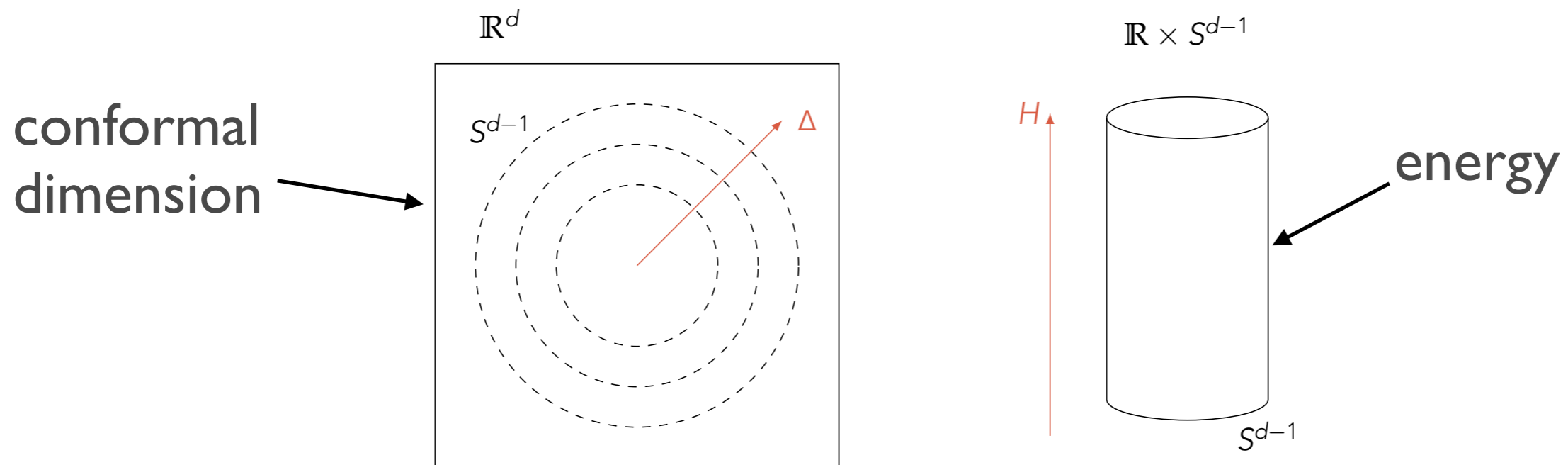
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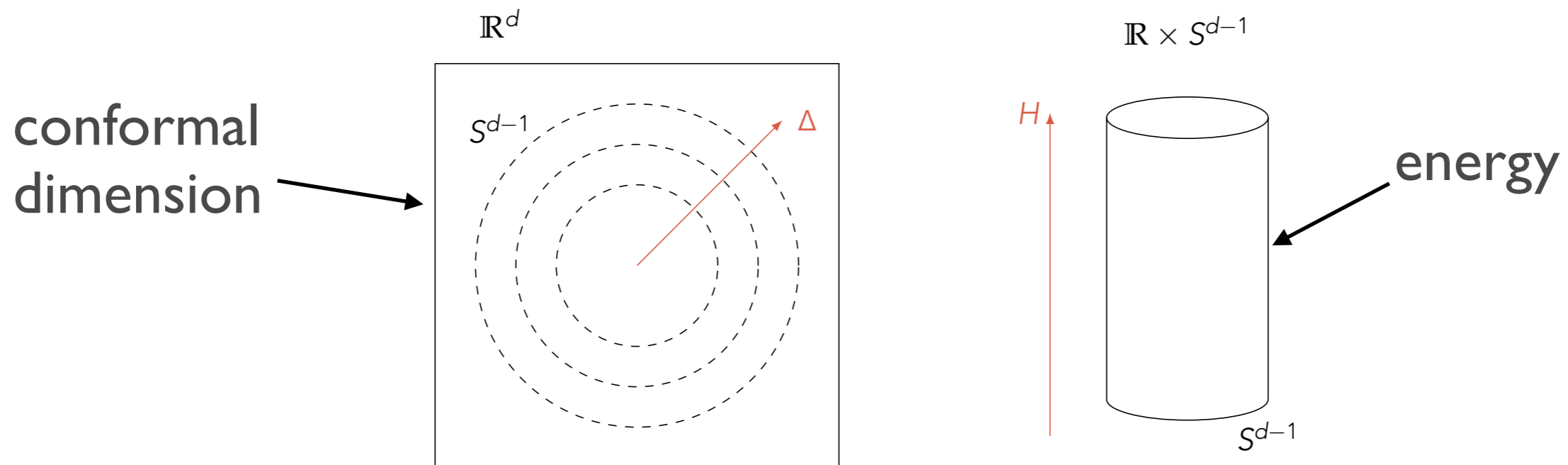
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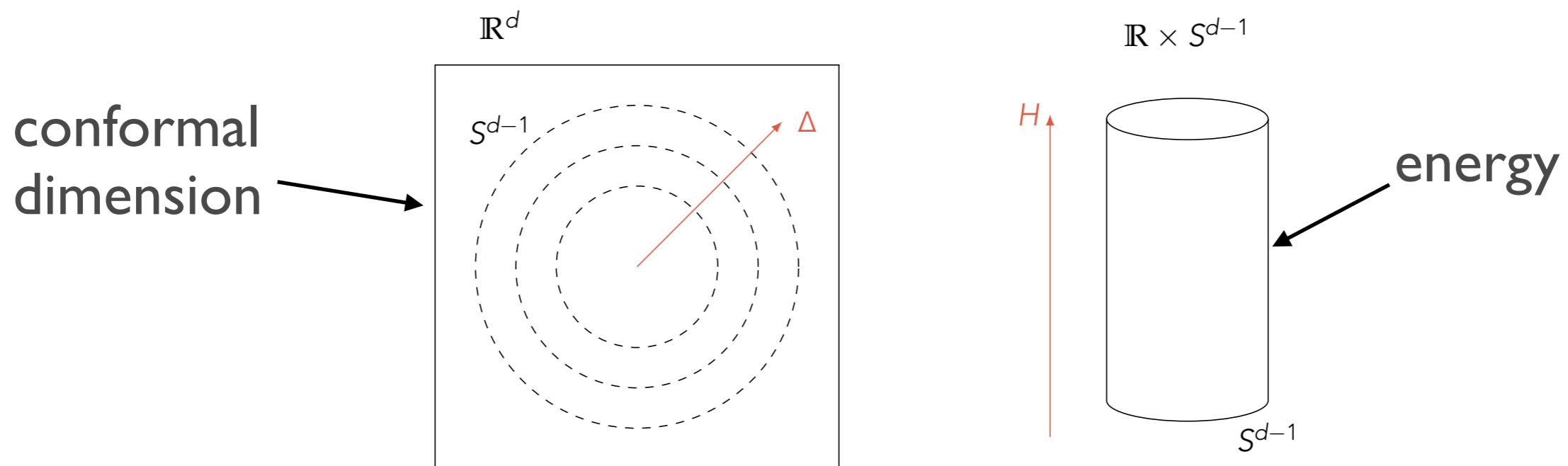
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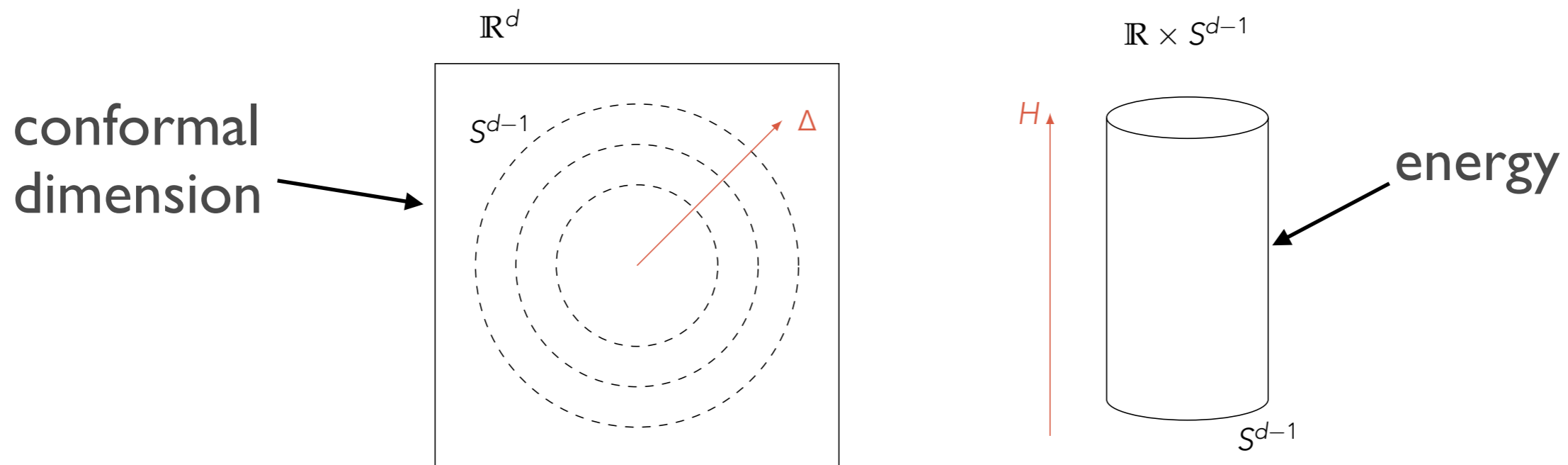
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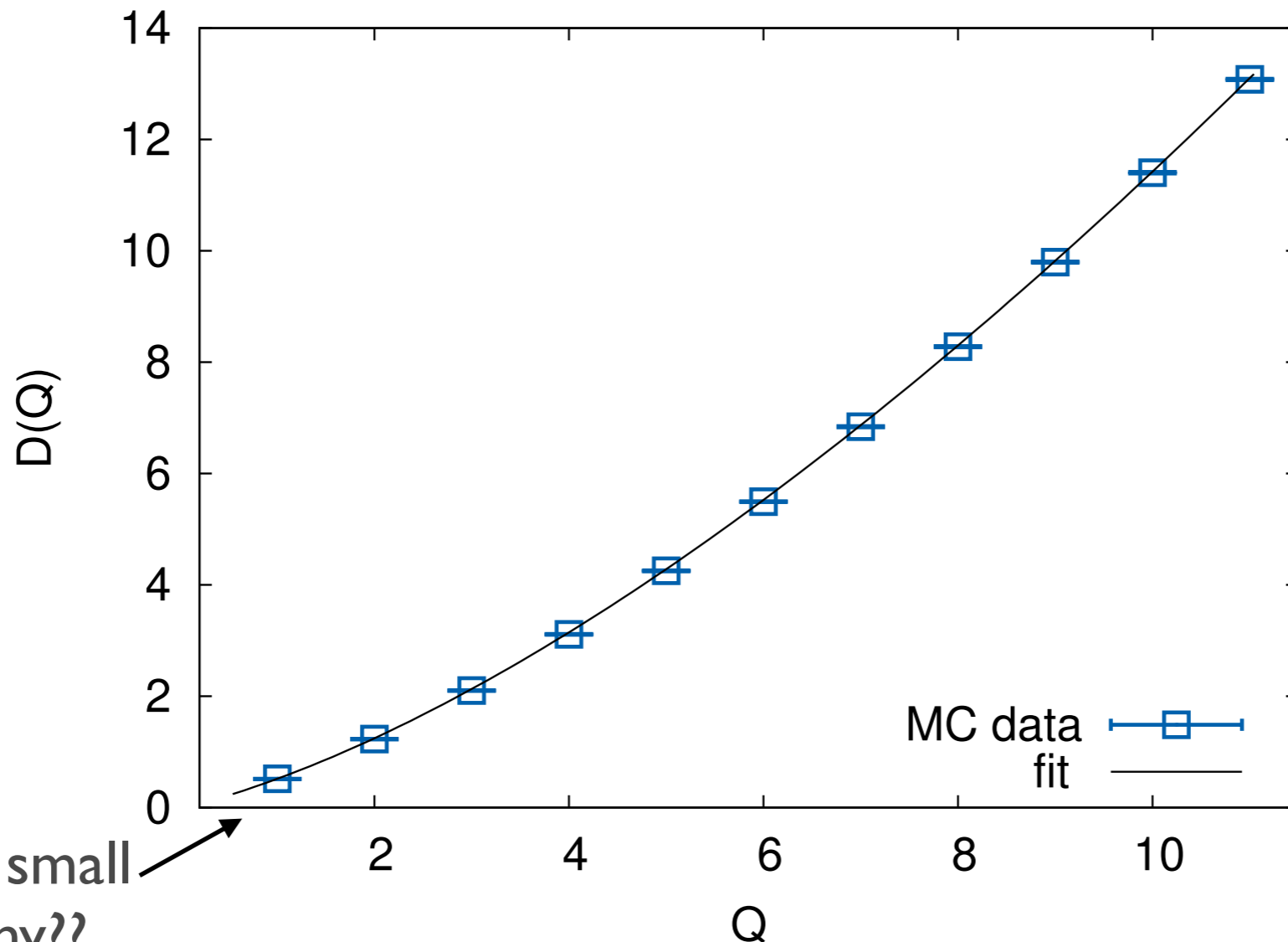
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Independent calculation on the lattice:



Excellent agreement!!

$$c_{3/2} = 1.195(10)$$

$$c_{1/2} = 0.075(10)$$

works for small charge. Why??

D. Banerjee, Sh. Chandrasekharan, D. Orlando [hep-th/1707.00711]

Large-charge expansion works extremely well for O(2).

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On top of the conformal Goldstone of $O(2)$, a new sector with $N-1$ non-relativistic type II Goldstones and $N-1$ massive modes with $m=2\mu$ appears.

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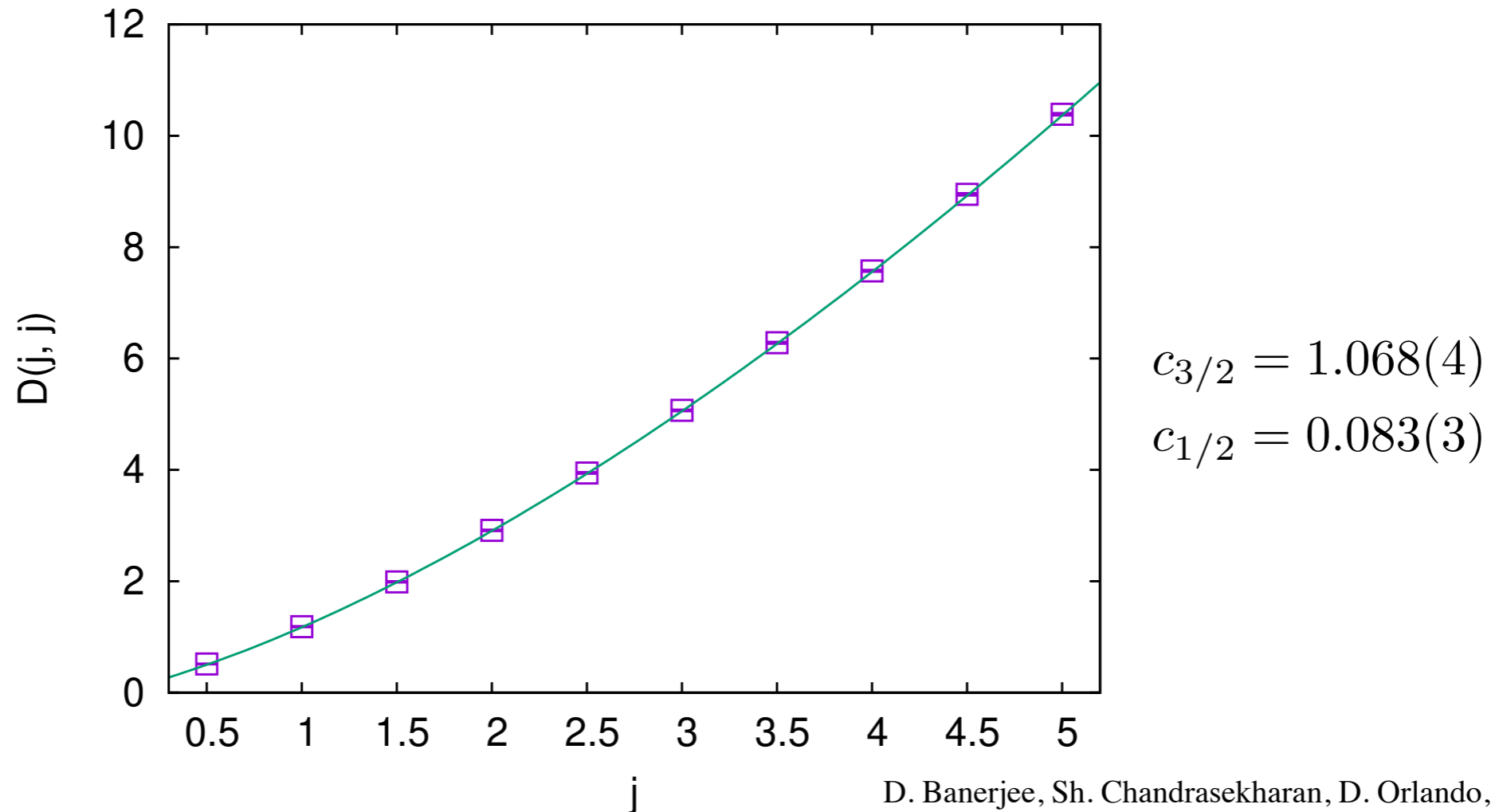
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New **lattice data** for $O(4)$ model:



Again excellent agreement with large- Q prediction!

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Adding the chemical potential gives us more structure to work with!

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Leading order in N :

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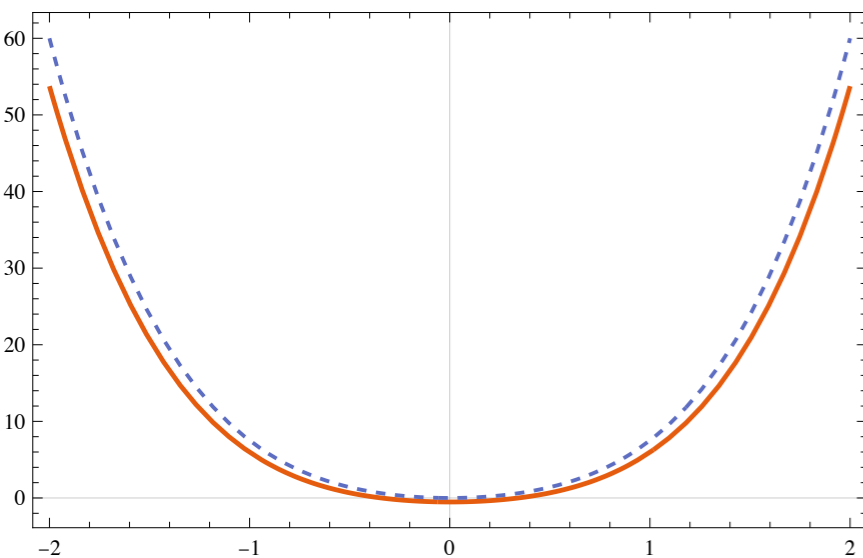
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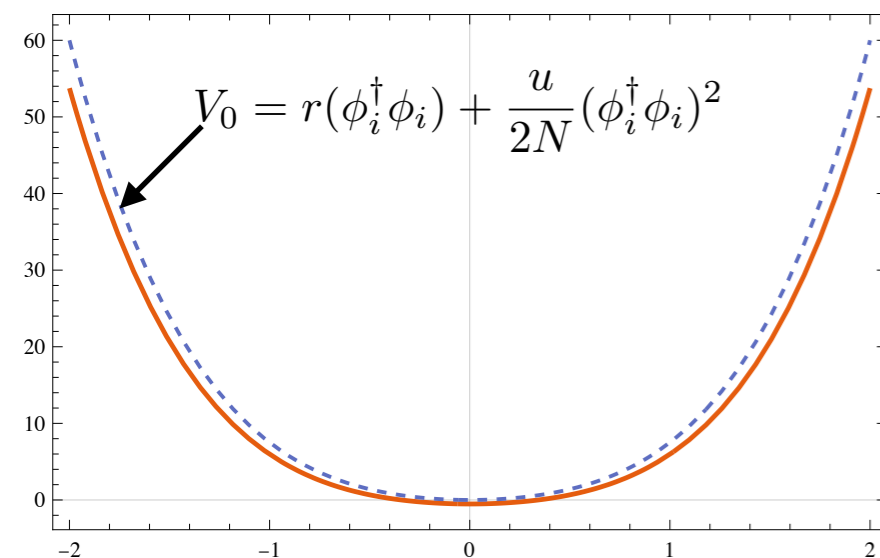
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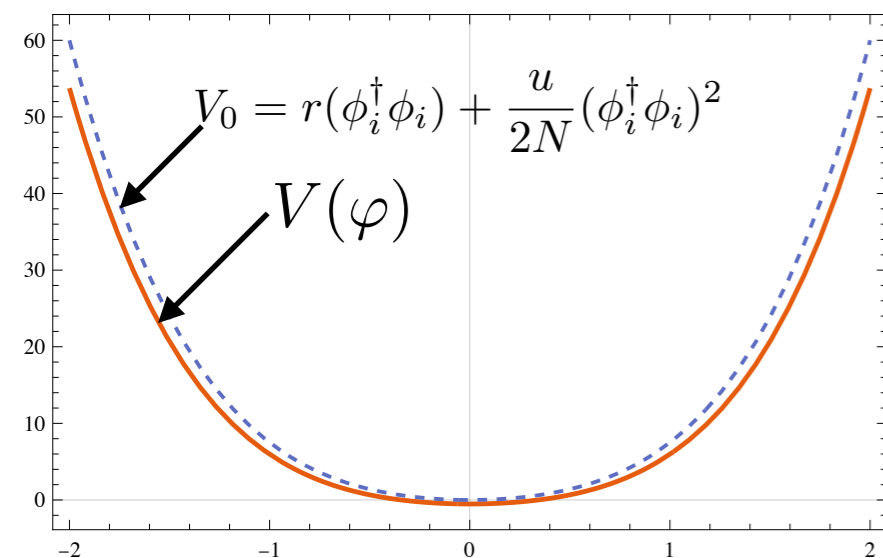
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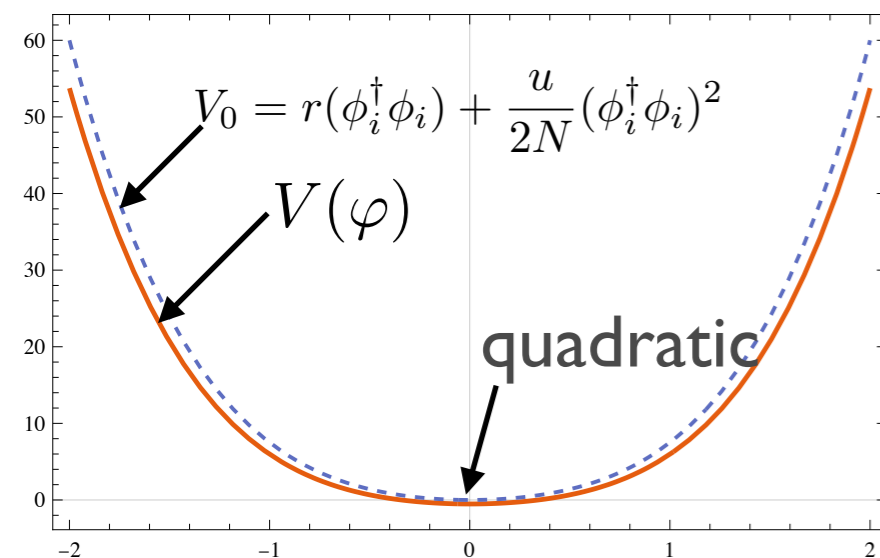
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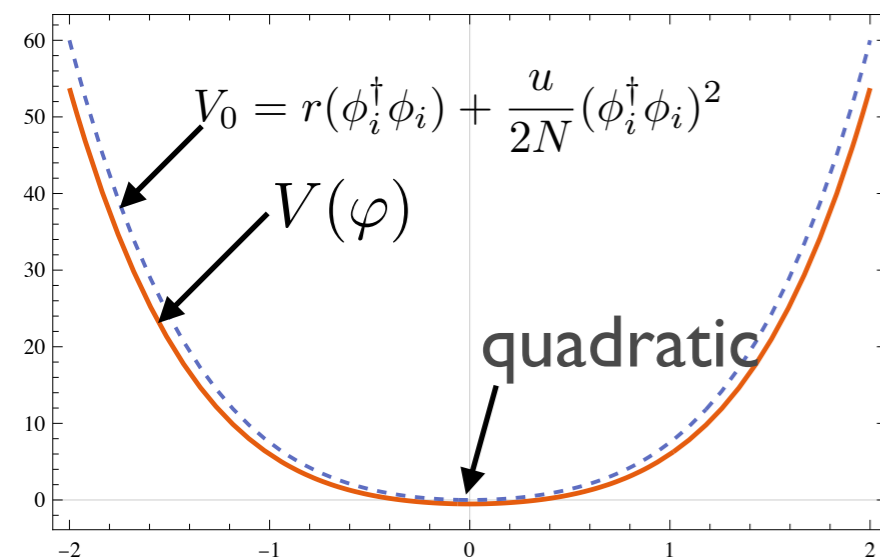
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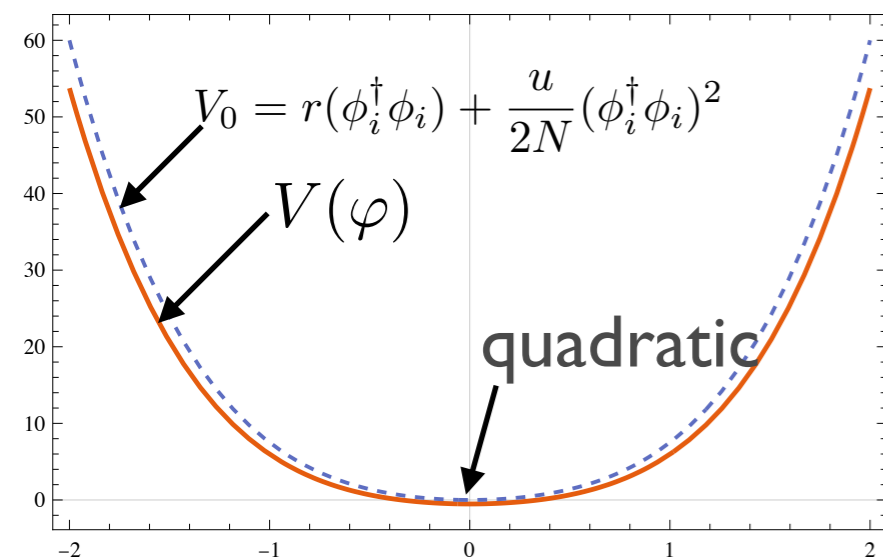
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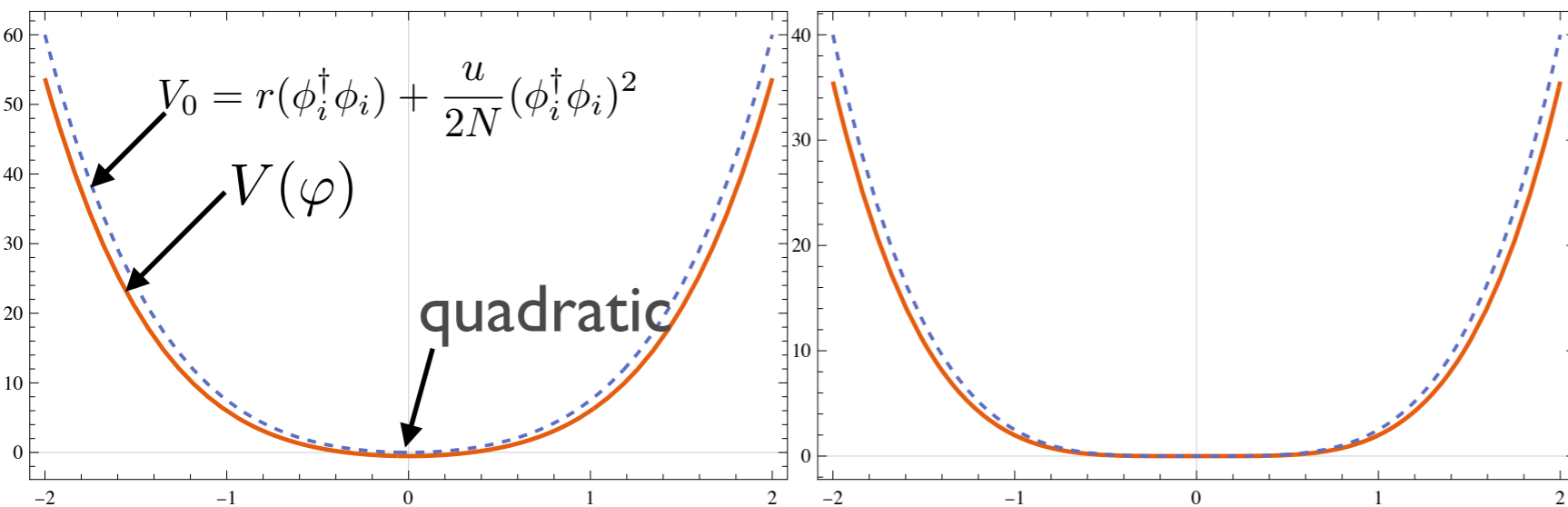
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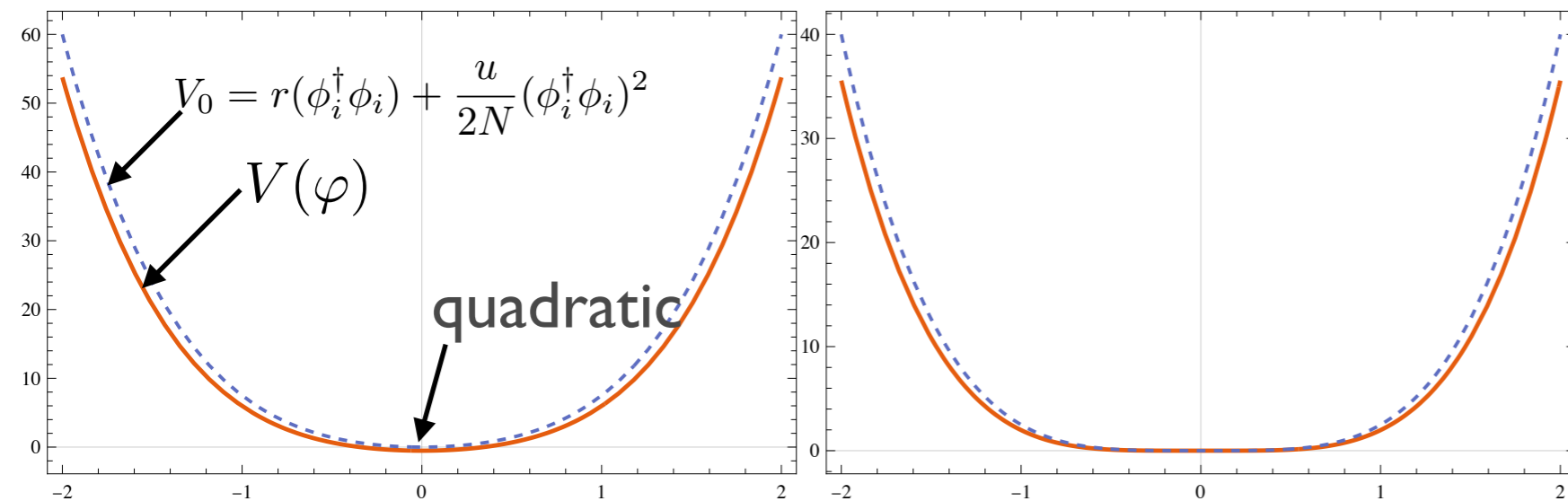
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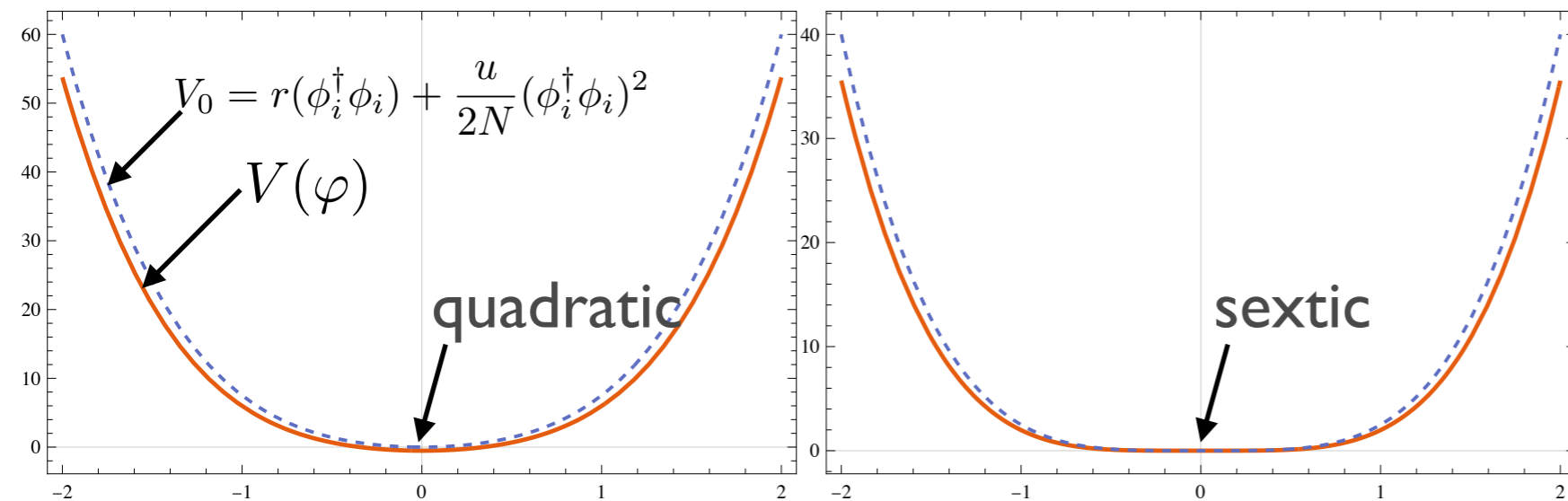
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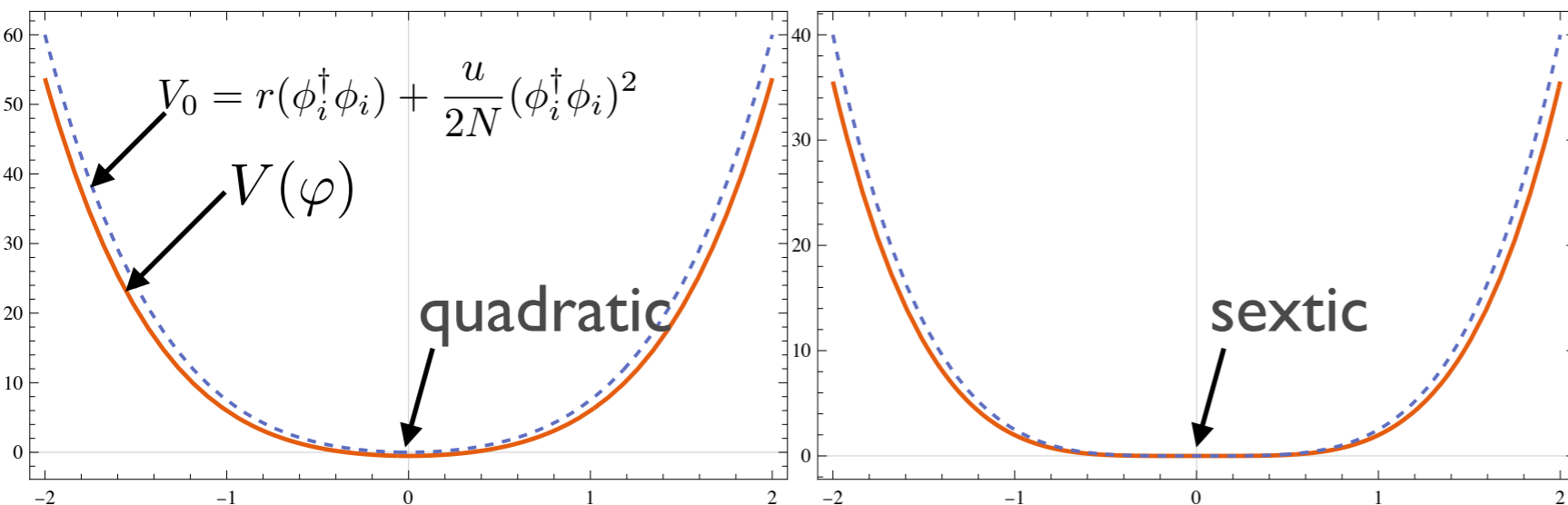
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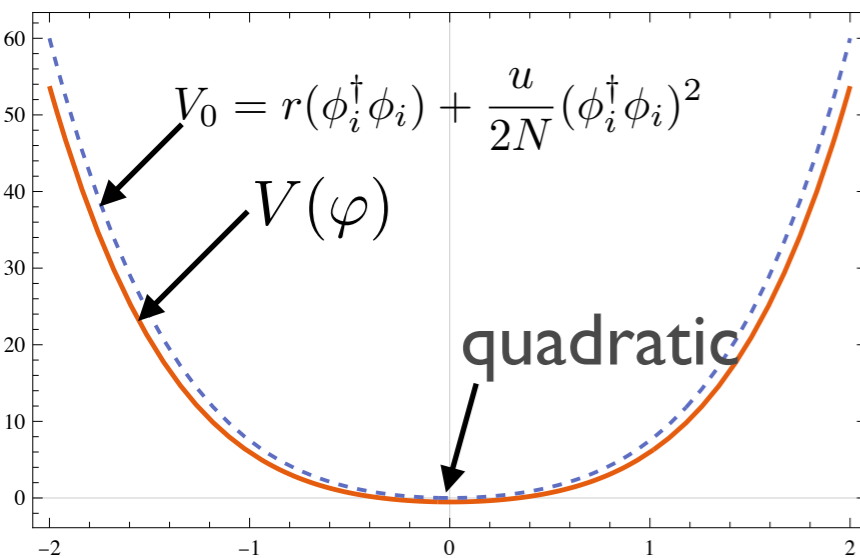
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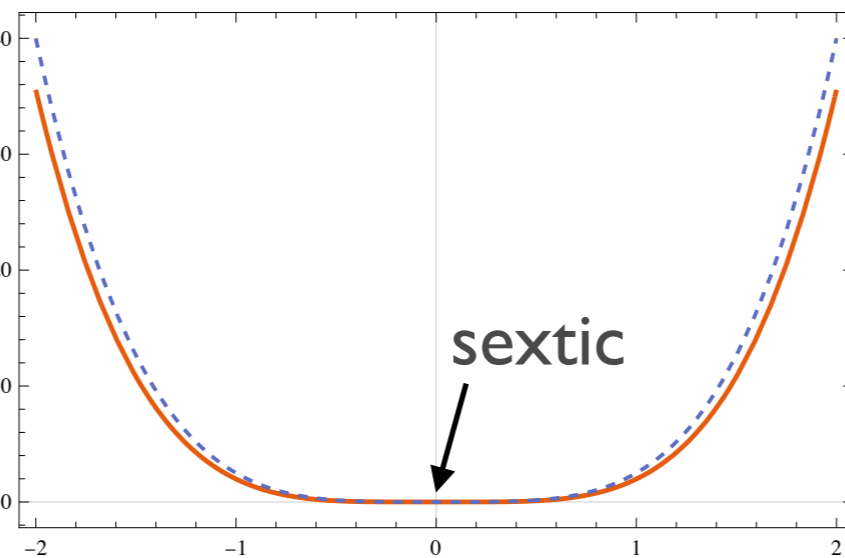
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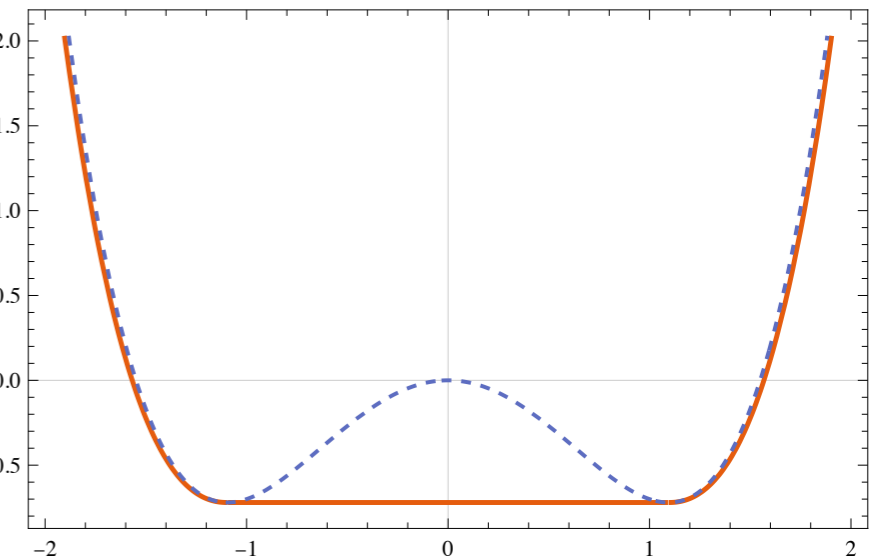
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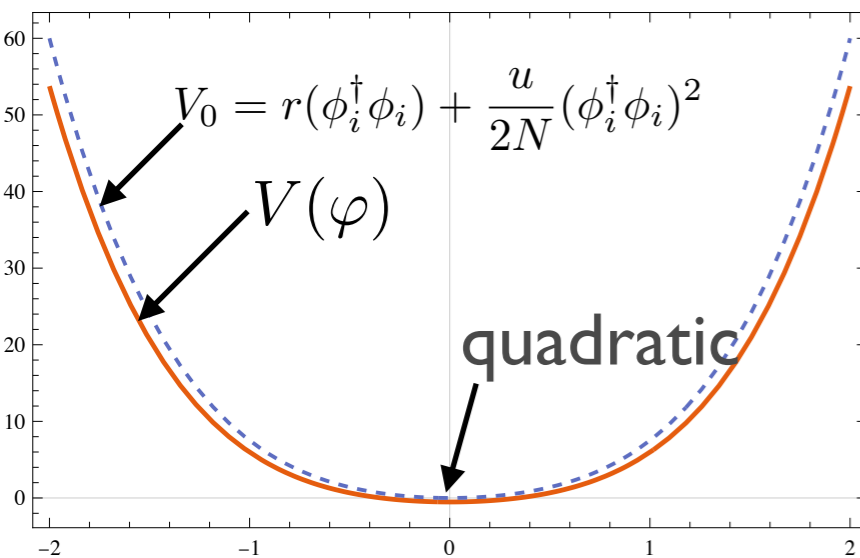
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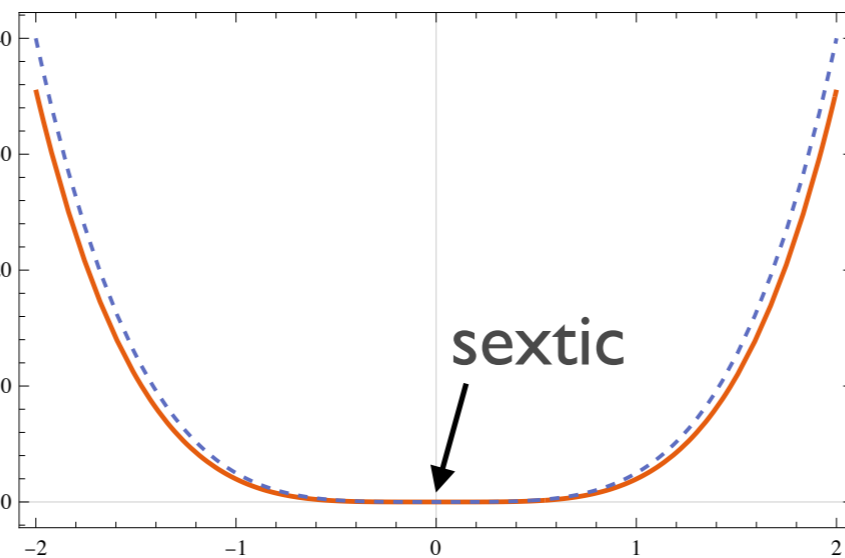
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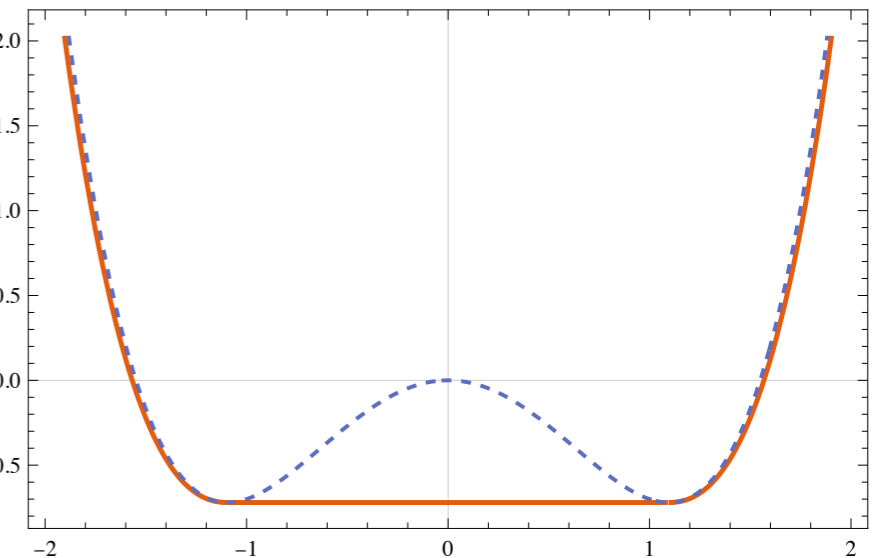
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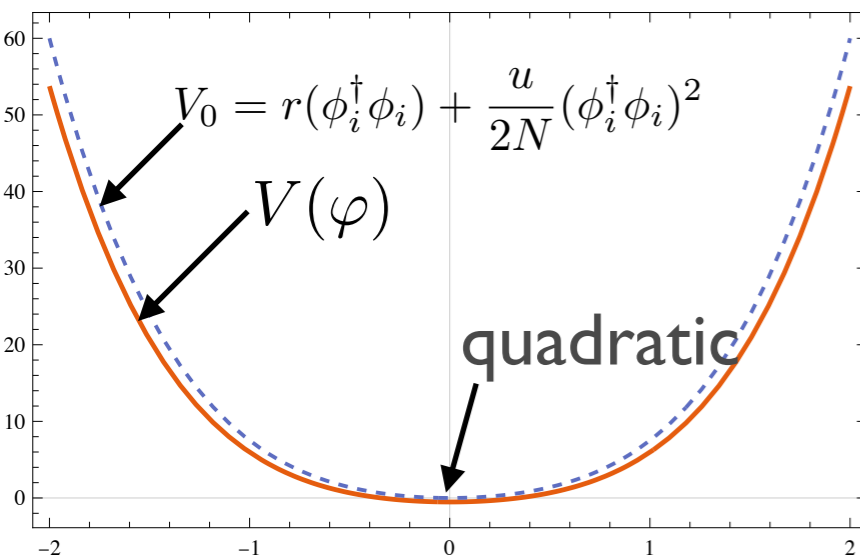
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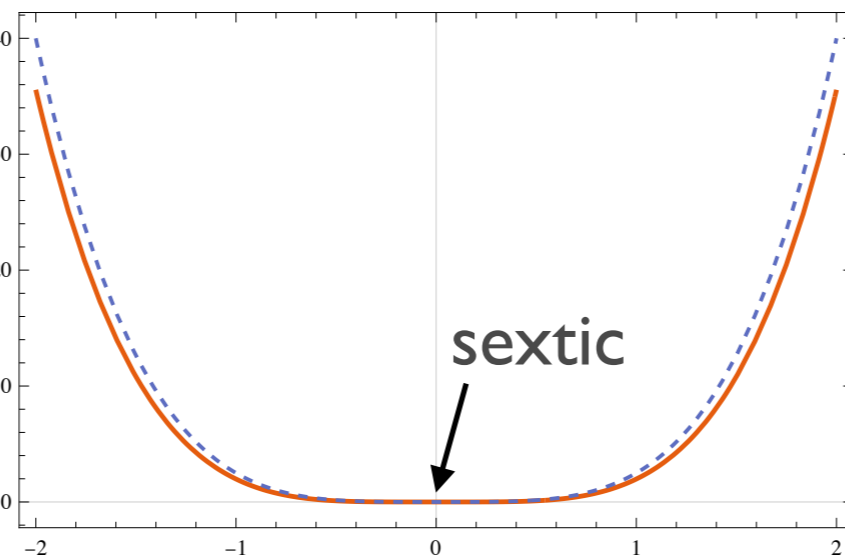
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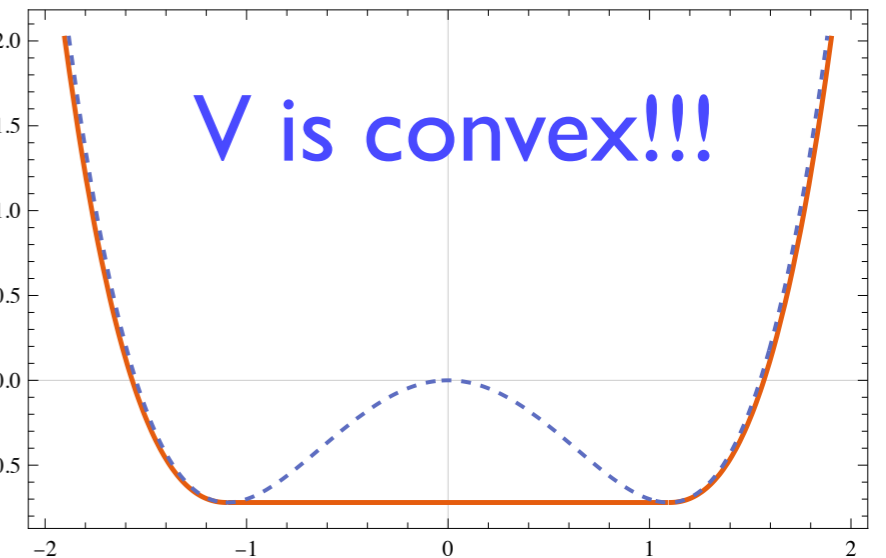
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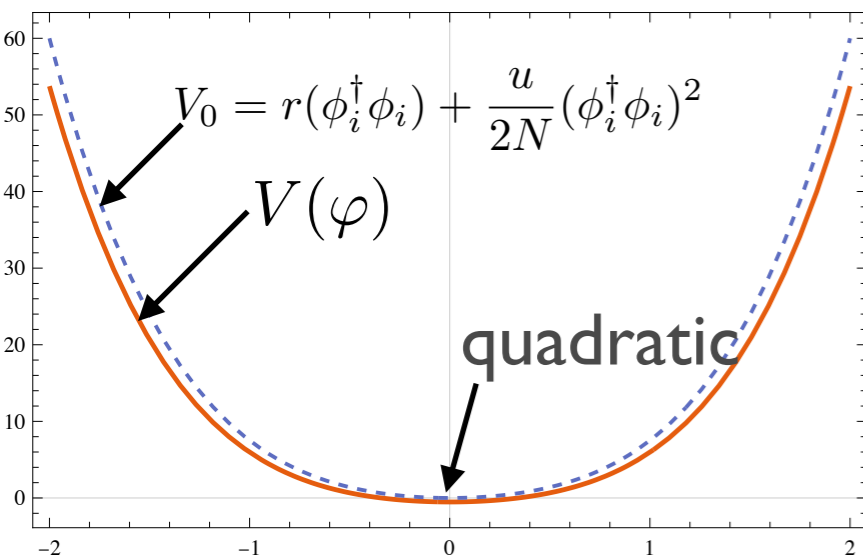
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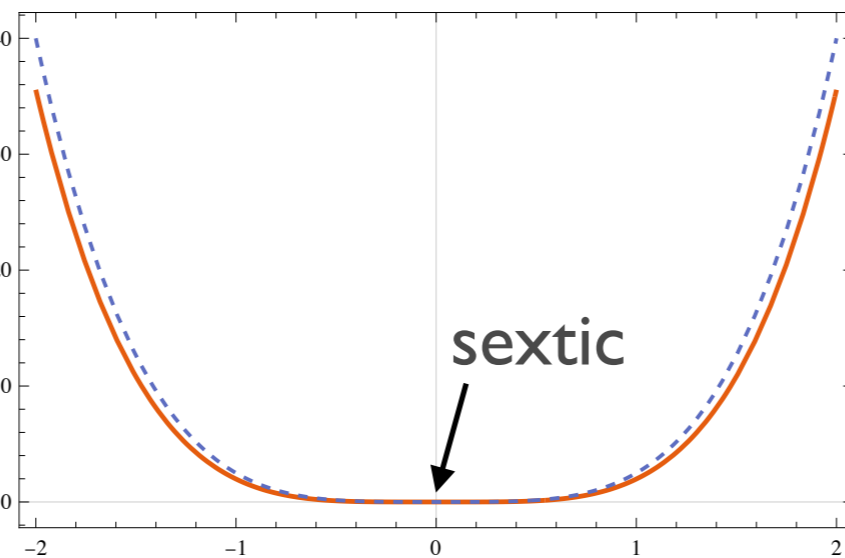
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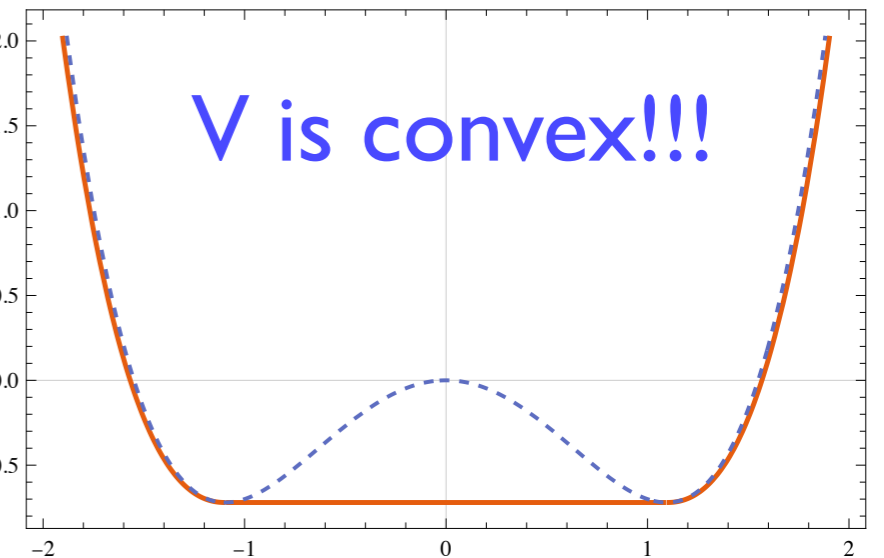
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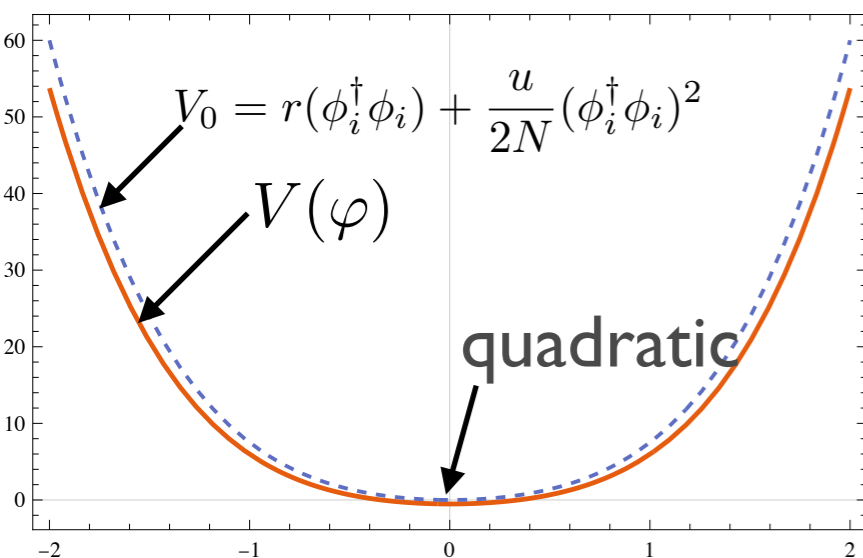
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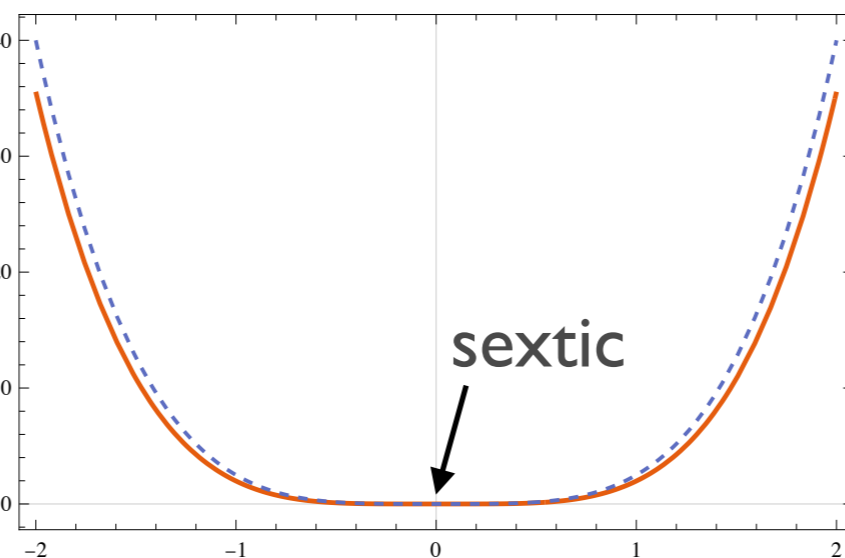
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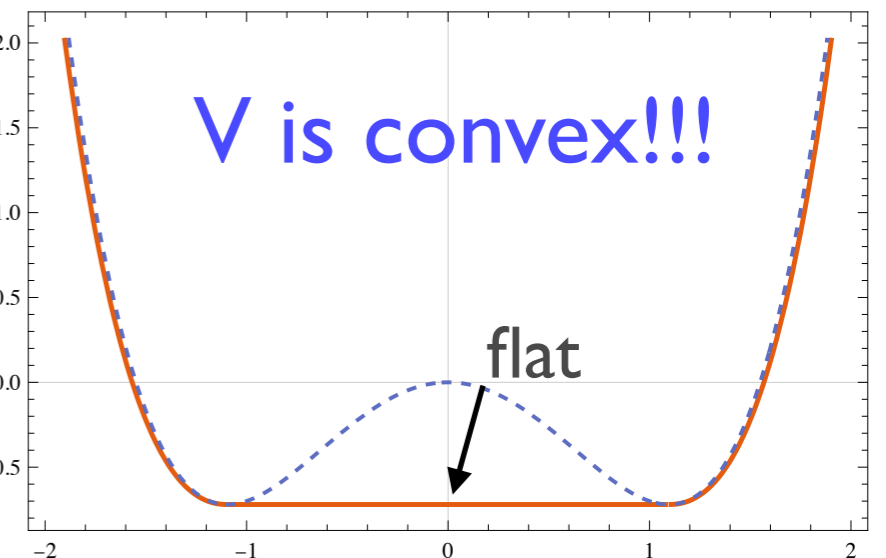
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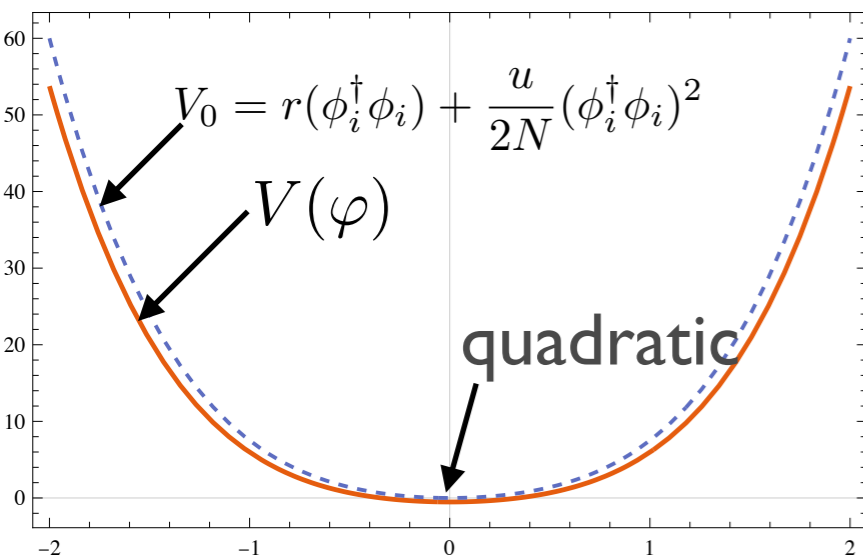
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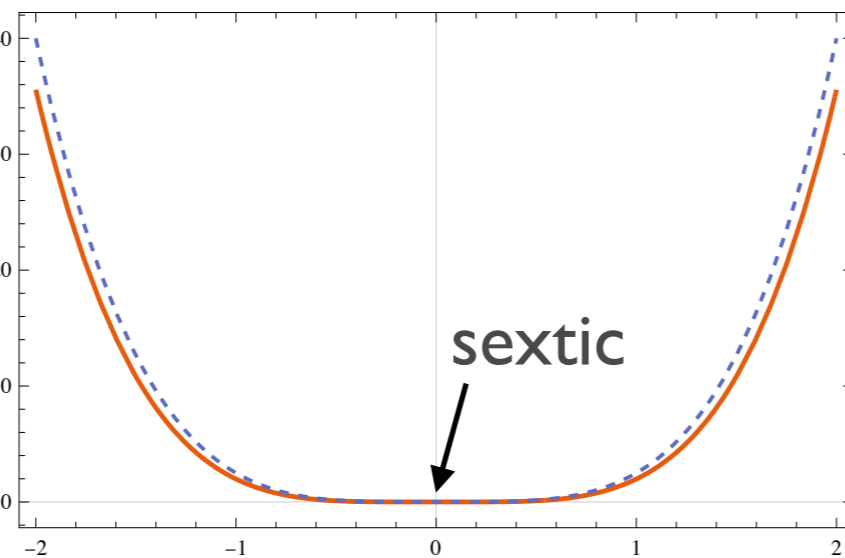
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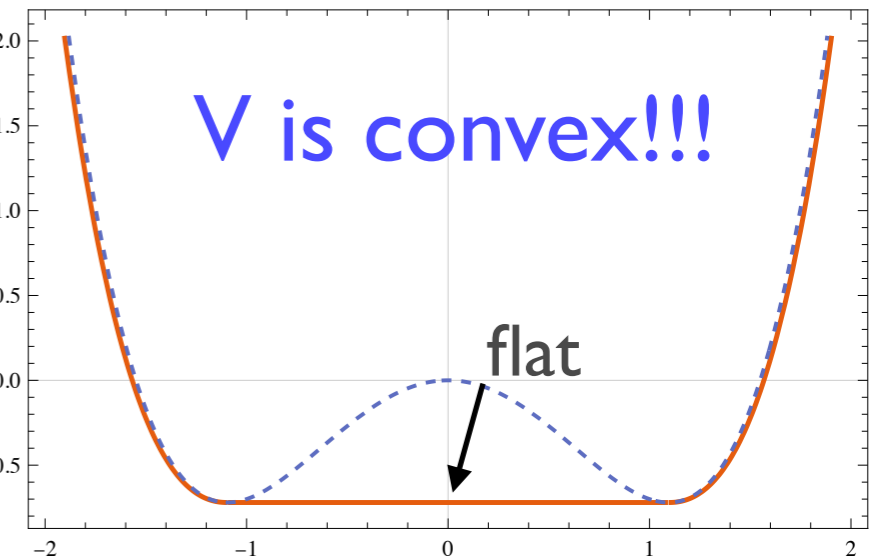
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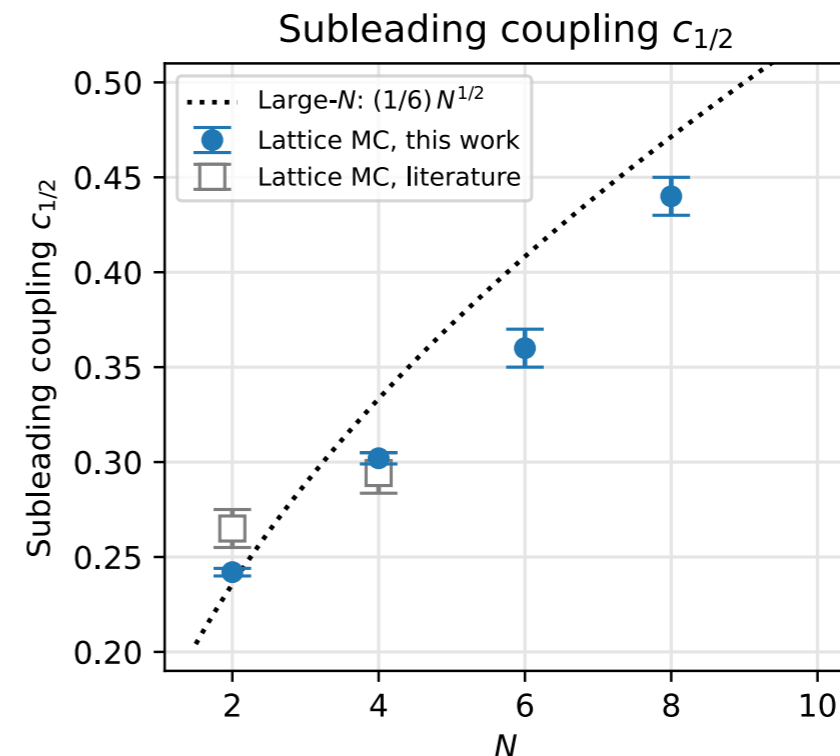
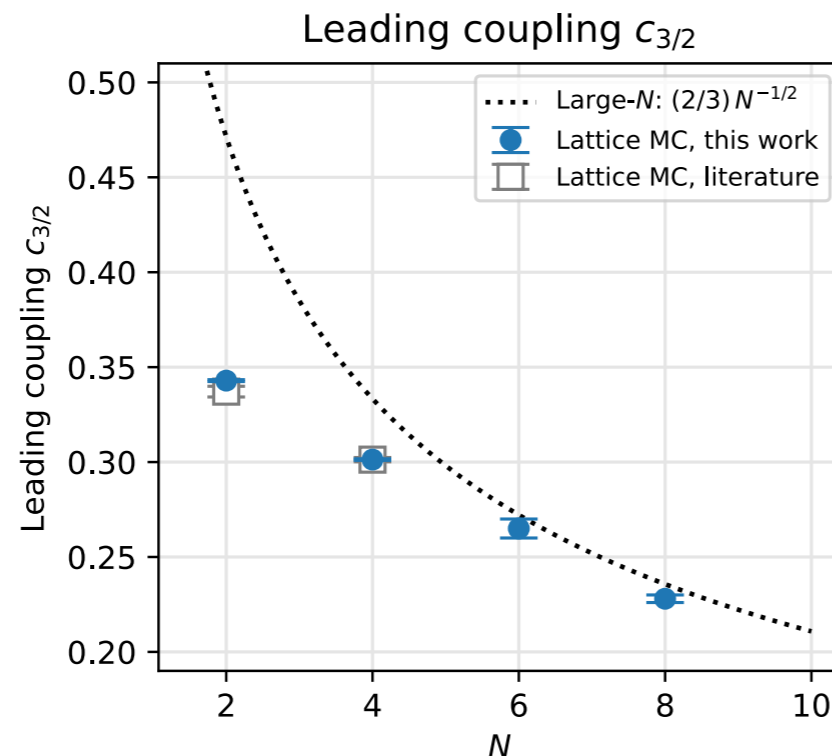
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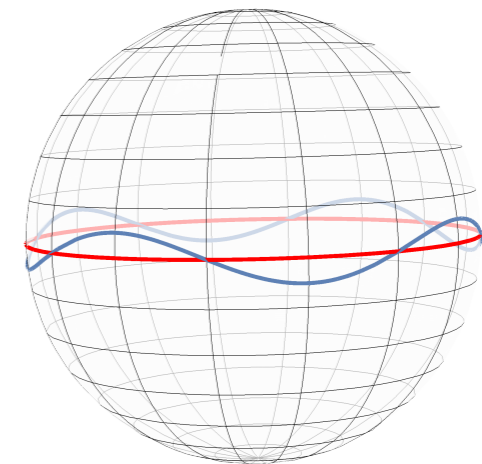
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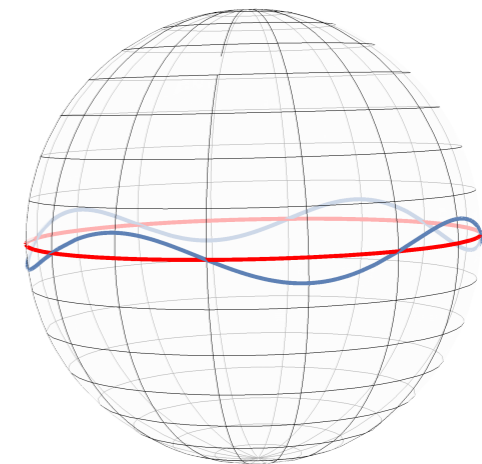
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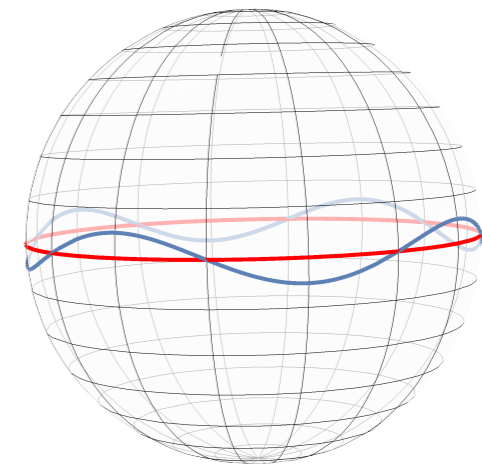
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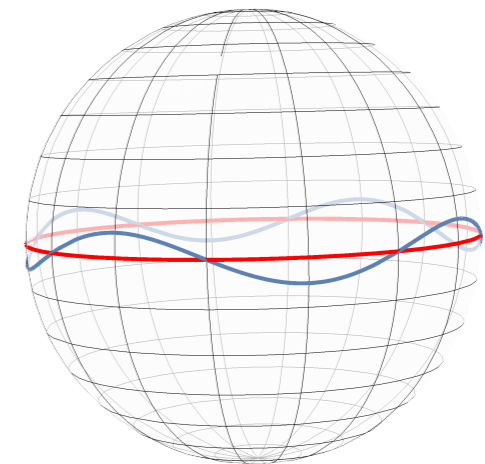
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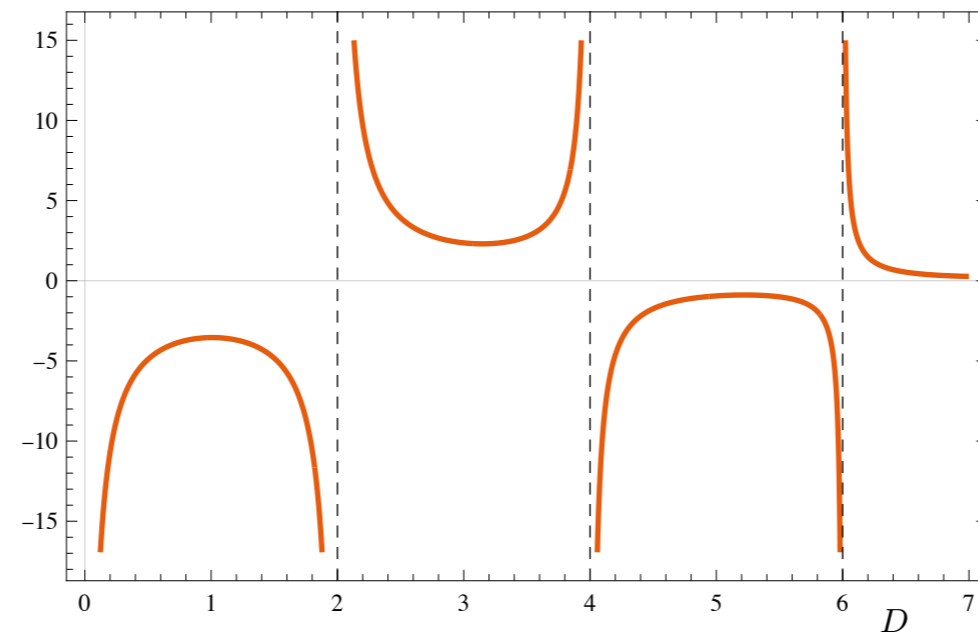
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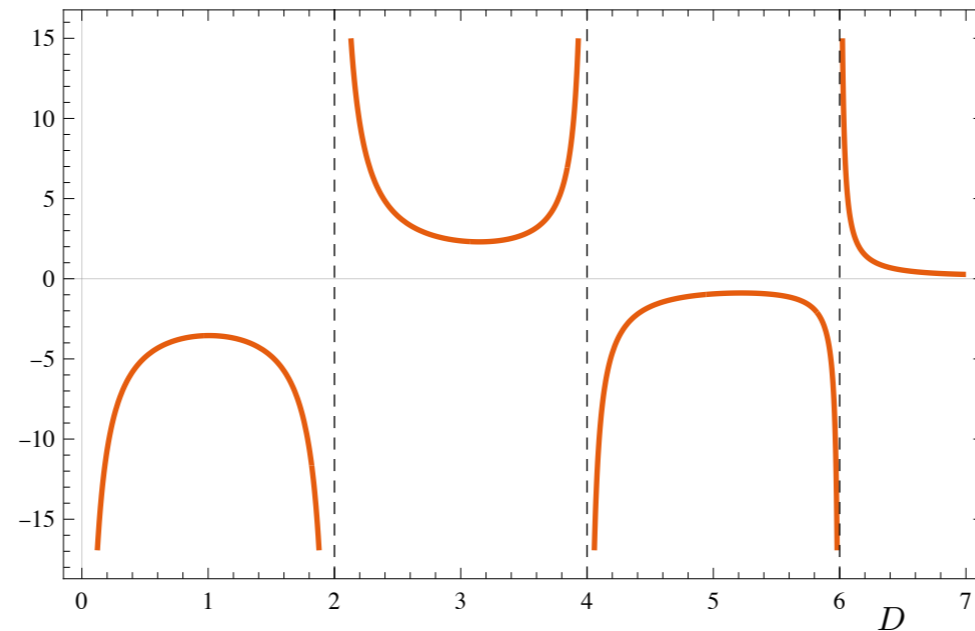
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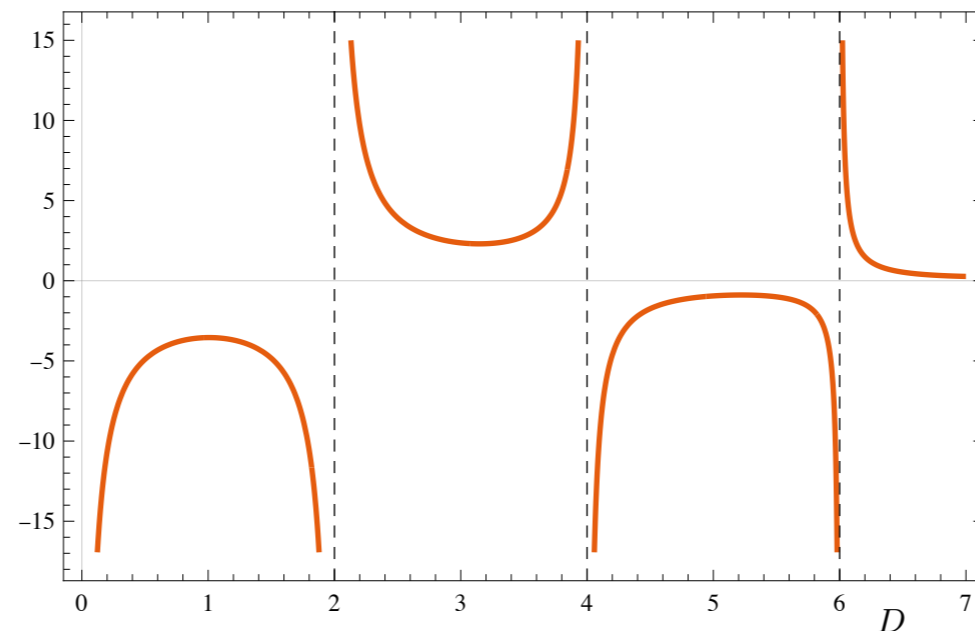


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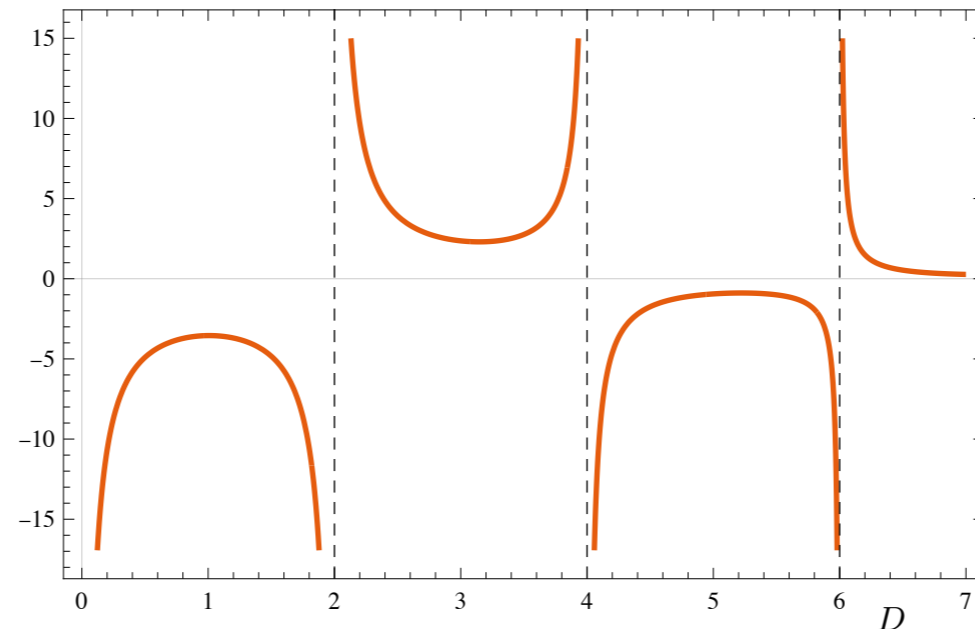
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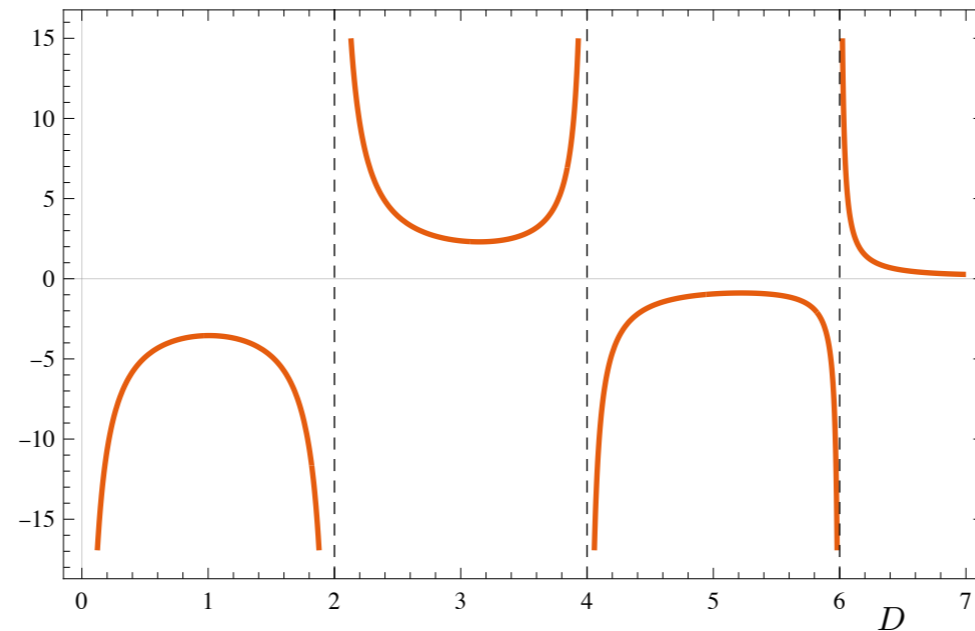
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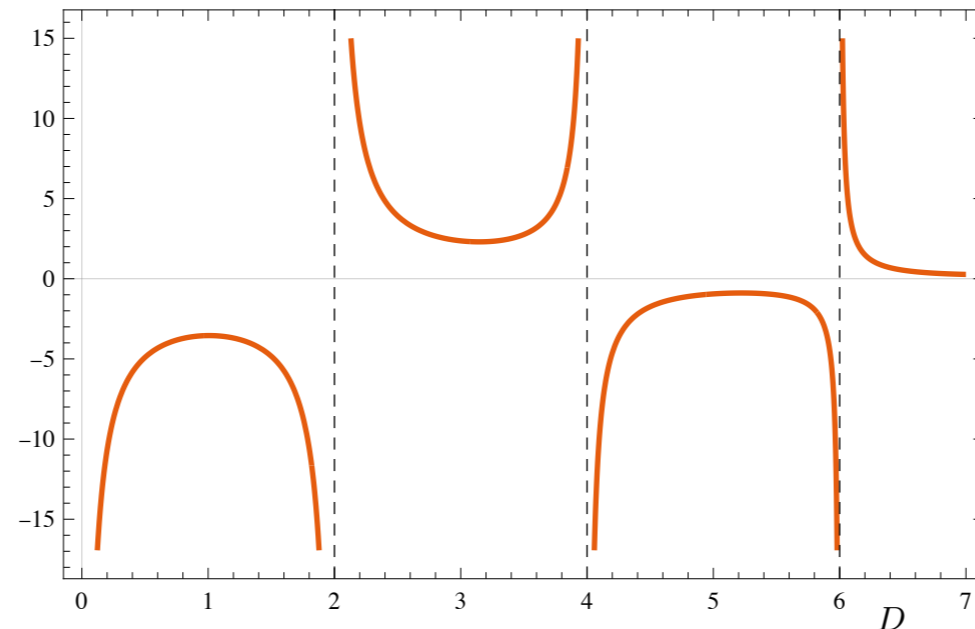
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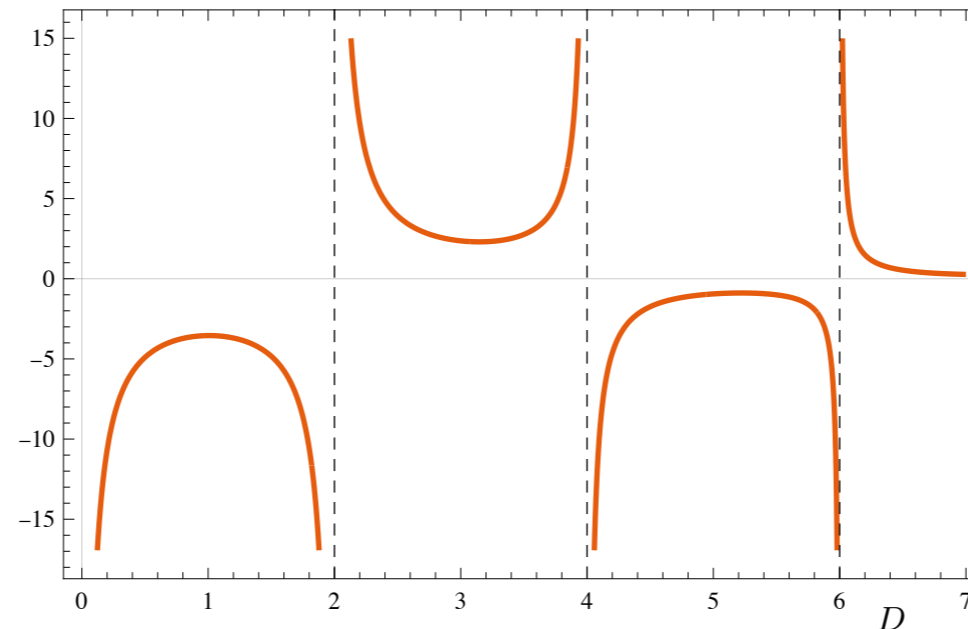
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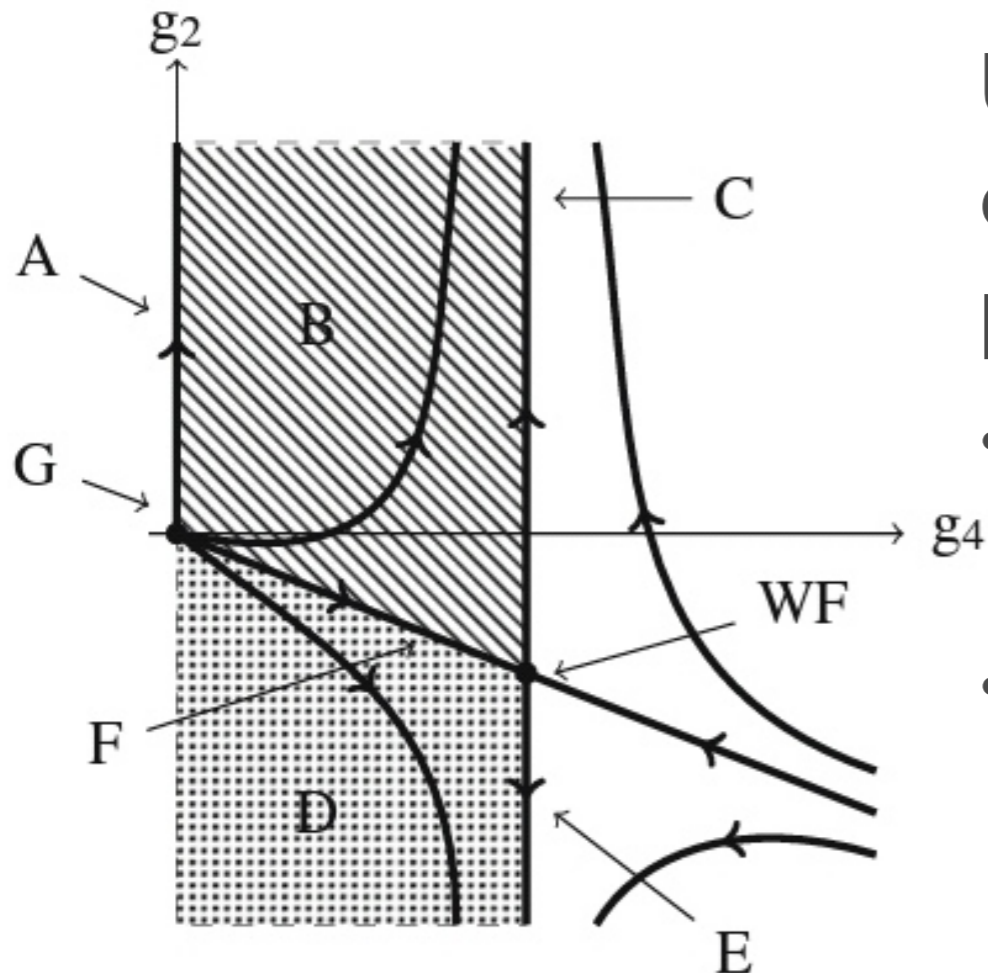
Giombi, Hyman;
Moser, Orlando, Reffert 2110.07617



Summary

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Concrete examples where a strongly-coupled CFT simplifies at large charge.



$O(2N)$ model in 3d: in the limit of large $U(1)$ charge Q , we computed the conformal dimensions in a controlled perturbative expansion:

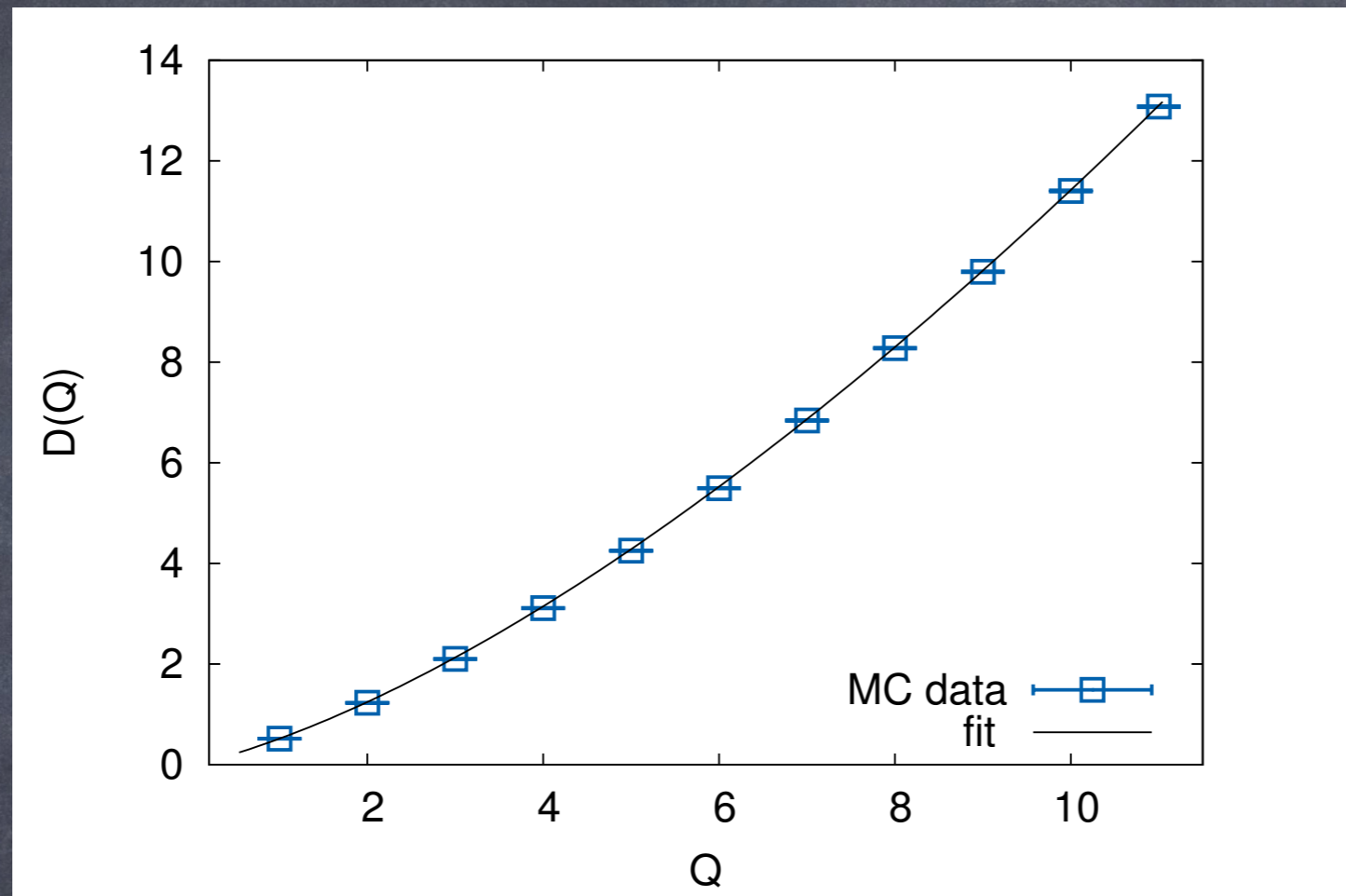
- Excellent agreement with lattice results for $O(2)$, $O(4)$
- large Q and large N : path integral at saddle pt., more control than in EFT, can calculate coefficients
- can follow the flow away from conformal point
- find the full effective potential

Further directions

- Further study of supersymmetric models at large R-charge (higher-dim. moduli spaces) Hellerman, Maeda, Orlando, Reffert, Watanabe; Argyres et al.
- Connection to holography (gravity duals) Loukas, Orlando, Reffert, Sarkar; De la Fuente, Zosso; Giombi, Komatsu, Offertaler.
- Operators with spin; connection to large-spin results Cuomo, de la Fuente, Monin, Pirtskhalava, Rattazzi; Cuomo
- Understanding dualities semi-classically at large charge
- Use/check large-charge results in conformal bootstrap Jafferis and Zhiboedov
- Further lattice simulations: inhomogeneous sector, general $O(N)$ Chandrasekharan et al.
- CFTs in other dimensions (2, 5, 6) Komargodski, Mezei, Pal, Raviv-Moshe; Araujo, Celikbas, Reffert, Orlando; Moser, Orlando, Reffert

Further directions

- Chern-Simons matter theories @large charge Watanabe
- 4- ϵ expansion @large charge Arias-Tamargo, Rodriguez-Gomez, Russo; Badel, Cuomo, Monin, Rattazzi; Watanabe; Antipin et al.
- going away from the conformal point Orlando, Reffert, Sannino; Orlando, Pellizzani, Reffert
- non-relativistic CFTs Favrod, Orlando, Reffert; Kravec, Pal; Orlando, Pellizzani, Reffert; Hellerman, Swanson; Pellizzani
- Boundary CFTs at large Q Cuomo, Mezei, Raviv-Moshe
- Weak gravity conjecture Aharony, Palti; Antipin et al.
- Study fermionic theories. Can large-charge approach be used for QCD (e.g. large baryon number)? Komargodski, Mezei, Pal, Raviv-Moshe



Thank you for your
attention!