

A holographic dual for Krylov complexity

or

Measuring the wormhole with Krylov complexity

Ruth Shir

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work with:

Eliezer Rabinovici, Adrián Sánchez-Garrido and Julian Sonner



Based on:

- ▶ E. Rabinovici, A. Sánchez-Garrido, RS and J. Sonner, “A bulk manifestation of Krylov complexity,” arXiv:2305.04355[hep-th]
- ▶ E. Rabinovici, A. Sánchez-Garrido, RS and J. Sonner, “Krylov complexity from integrability to chaos,” JHEP 07 (2022), 151
- ▶ E. Rabinovici, A. Sánchez-Garrido, RS and J. Sonner, “Krylov localization and suppression of complexity,” JHEP 03 (2022), 211
- ▶ E. Rabinovici, A. Sánchez-Garrido, RS and J. Sonner, “Operator complexity: a journey to the edge of Krylov space,” JHEP 06 (2021), 062
- ▶ J. L. F. Barbón, E. Rabinovici , RS and R. Sinha, “On The Evolution Of Operator Complexity Beyond Scrambling,” JHEP 10 (2019), 264

Plan

- ▶ Why complexity? Which notion of complexity?
- ▶ Krylov complexity
- ▶ A holographic dual for Krylov complexity
- ▶ Open questions and future directions

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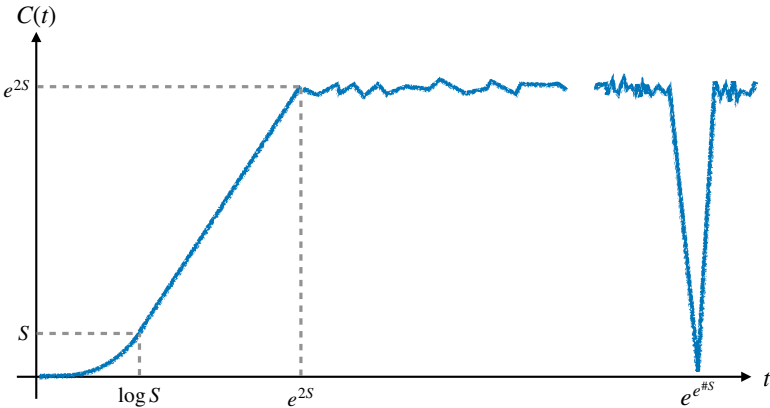
Answer

Quantum complexity???

[Susskind 2014 –]

General time-dependent profile of complexity

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Different notions of quantum complexity

- ▶ **Circuit complexity:** The minimal number of local gates g_i needed to construct an operator U starting from an initial operator U_0 , up to a tolerance parameter ϵ
- ▶ **Geometric complexity:** Choose a penalty metric on space of unitaries; define complexity as shortest path to reach U from U_0
- ▶ **Krylov complexity:** Defined using the system's Hamiltonian and initial state/operator

Can we find a precise bulk-boundary correspondence?

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- ▶ A dense operator which has non-zero projection on every eigenstate of the Liouvillian, $|E_a\rangle\langle E_b|$
- ▶ No degeneracies in the spectrum of the Liouvillian $\omega_{ab} = E_a - E_b$ except for the D -fold degeneracy of zero frequencies $\omega_{aa} = E_a - E_a = 0$

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where K is the Krylov space dimension

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- ▶ Define **position** operator over Krylov basis

$$\hat{n} = \sum_{n=0}^{K-1} n |\mathcal{O}_n\rangle \langle \mathcal{O}_n| \quad (8)$$

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a probe of operator time evolution at all time scales

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- ▶ and as a measure of operator complexity at *all time scales* for *finite systems* in [Barbón Rabinovici RS Sinha 2019]

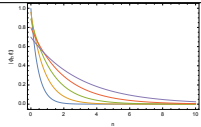
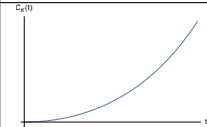
Dynamics of Krylov complexity

[Parker et al 2018] [Barbón Rabinovici RS Sinha 2019] [Rabinovici
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n	b_n	wavefunction	K-complexity	time scales
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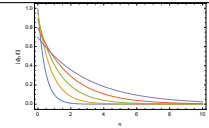
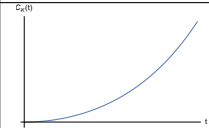
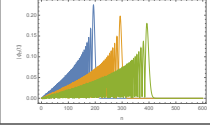
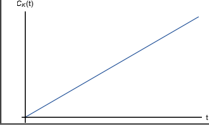
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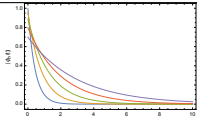
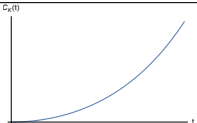
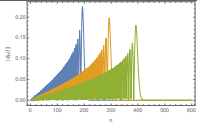
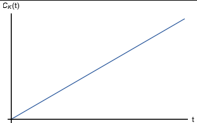
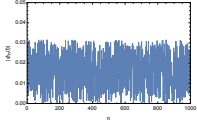
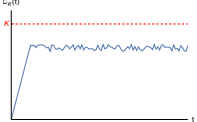
Dynamics of Krylov complexity

[Parker et al 2018] [Barbón Rabinovici RS Sinha 2019] [Rabinovici Sánchez-Garrido RS Sonner 2020]

n	b_n	wavefunction	K-complexity	time scales
$n \ll S$	$b_n \sim n$			$t \lesssim \log S$
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$n \sim e^{2S}$	“descent”			$t \gtrsim e^{2S}$

Numerical results: SYK₄ setup

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- ▶ SYK₄ is a maximally chaotic many-body system. Consider complex SYK₄ with L fermions

$$H_{\text{SYK}} = \sum_{i,j,k,l}^L J_{ij,kl} c_i^\dagger c_j^\dagger c_k c_l \quad (10)$$

where $\{c_i, c_j^\dagger\} = \delta_{ij}$ and $\{c_i, c_j\} = 0 = \{c_i^\dagger, c_j^\dagger\}$

- ▶ The coupling constants are taken from a Gaussian distribution with

$$\overline{J_{ij,kl}} = 0 \quad \overline{|J_{ij,kl}|^2} = \frac{3!J^2}{L^3} \quad (11)$$

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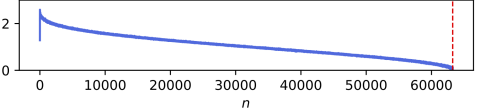
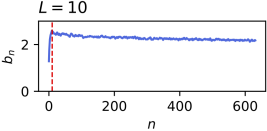
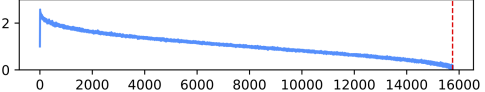
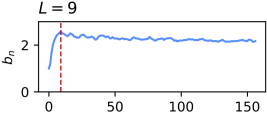
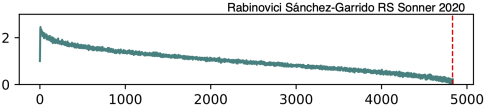
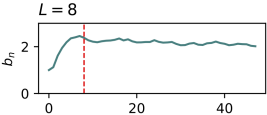
- ▶ Operator taken to be the hopping operator

$$\mathcal{O} = c_{L-1}^\dagger c_L + c_L^\dagger c_{L-1} \quad (12)$$

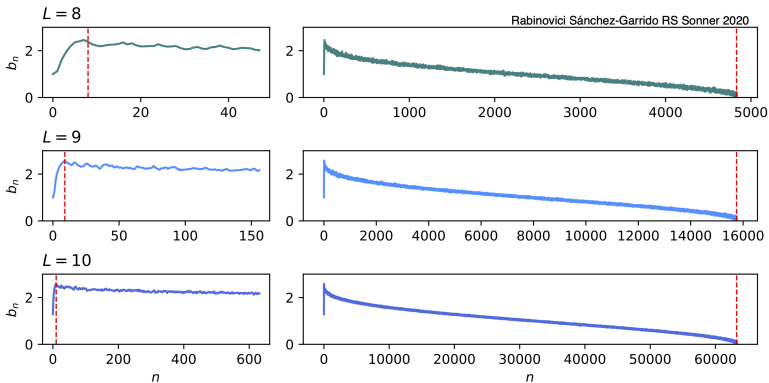
shown in [\[Sonner Vielma 2017\]](#) to satisfy ETH

Non-perturbative “Descent”

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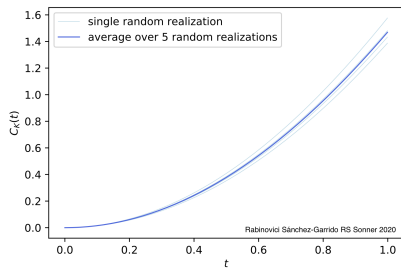
Features a non-perturbative e^{-2S} slope
[Rabinovici Sánchez-Garrido RS Sonner 2020]

K-complexity for SYK₄ with 10 complex fermions

[Rabinovici Sánchez-Garrido RS Sonner 2020]

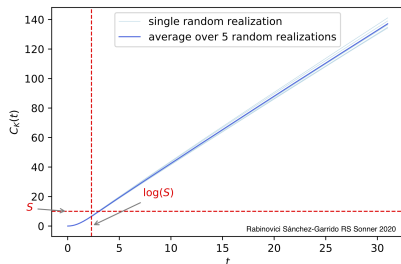
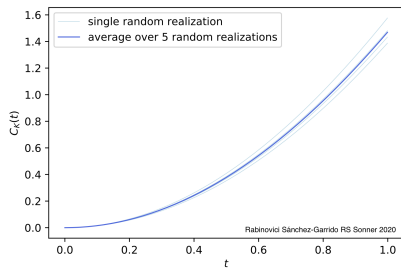
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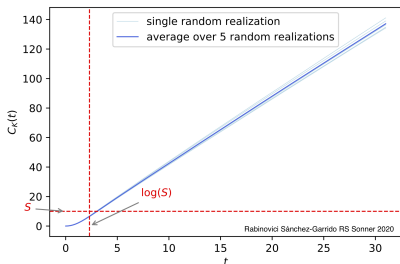
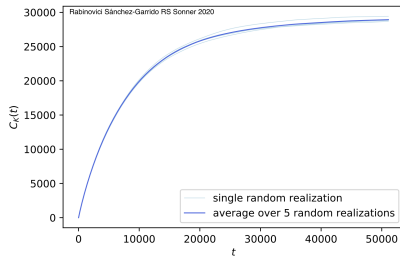
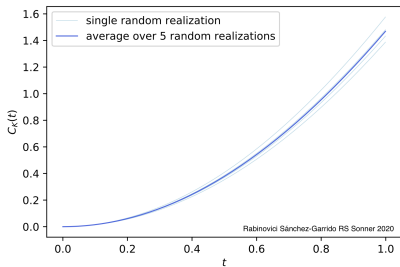
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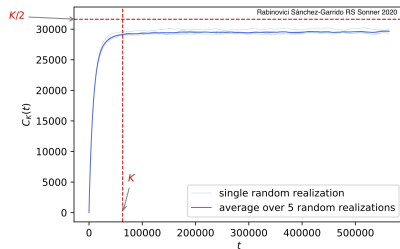
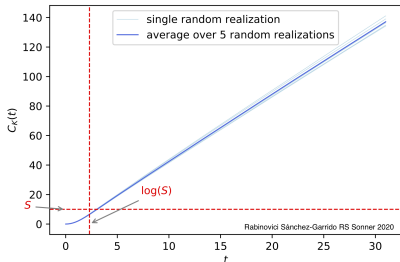
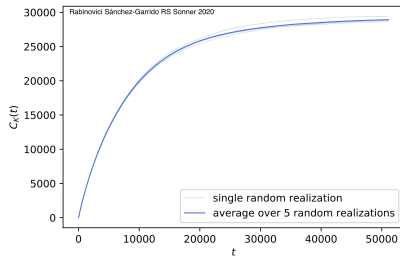
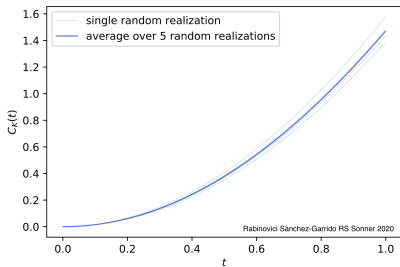
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K-entropy

[Barbón Rabinovici RS Sinha 2019]

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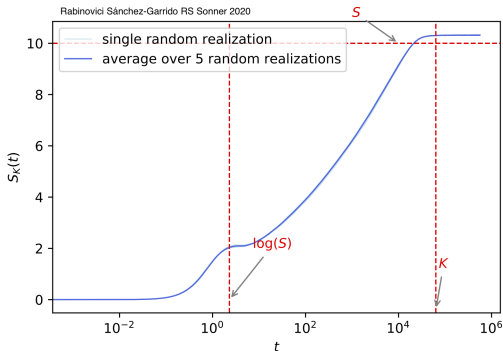
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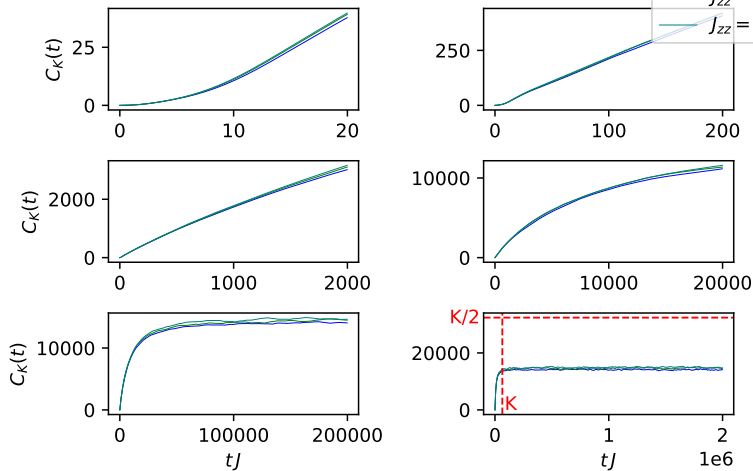
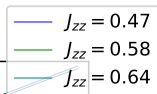
Suppression of K-complexity in XXZ

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XXZ with $N = 12$ $M = 4$ $\mathcal{O} = \sigma_6^z + \sigma_7^z$



Krylov space dimension $K = 64771$

[Rabinovici Sánchez-Garrido RS Sonner 2021]

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- ▶ Given a Hamiltonian H and an initial state $|\Omega\rangle$, the state evolves in time unitarily

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$$H = \begin{pmatrix} a_1 & b_1 & & & & & & & \\ b_1 & a_2 & b_2 & & & & & & \\ & b_2 & a_3 & b_3 & & & & & \\ & & b_3 & a_4 & \ddots & & & & \\ & & & \ddots & \ddots & & & & \\ & & & & & \ddots & b_{K-1} & & \\ & & & & & b_{K-1} & a_K & & \end{pmatrix} \quad (15)$$

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- ▶ Studied in [\[Balasubramanian Caputa Magan Wu 2022\]](#)

A holographic dual for Krylov complexity

The boundary: DSSYK

[Berkooz Narayan Simon 2018][Berkooz Isachenkov Narovlansky Torrents 2018]

- ▶ Start with SYK with N fermions $\{\psi_i, \psi_j\} = 2\delta_{ij}$

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- ▶ Start with SYK with N fermions and p -body interactions

$$H_{\text{SYK}} = i^{p/2} \sum_{1 \leq i_1 \leq i_2 \leq \dots \leq i_p \leq N} J_{i_1 i_2 \dots i_p} \psi_{i_1} \psi_{i_2} \dots \psi_{i_p} \quad (16)$$

where $\overline{J_{i_1 i_2 \dots i_p}} = 0$ and $\overline{J_{i_1 i_2 \dots i_p}^2} = \frac{J^2}{\lambda} \binom{N}{p}^{-1}$

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- ▶ Define the ratio parameter $\lambda \equiv \frac{2p^2}{N}$
- ▶ In the limit $N \rightarrow \infty$, $p \rightarrow \infty$ and λ fixed, the **ensemble averaged** effective Hamiltonian is

$$H = \frac{J}{\sqrt{\lambda(1-q)}} \begin{pmatrix} 0 & \sqrt{1-q} & & & \\ \sqrt{1-q} & 0 & \sqrt{1-q^2} & & \\ & \sqrt{1-q^2} & 0 & \sqrt{1-q^3} & \\ & & \sqrt{1-q^3} & 0 & \ddots \\ & & & \ddots & \ddots \end{pmatrix}$$

where $q \equiv e^{-\lambda}$

Krylov complexity for DSSYK

- ▶ In [Rabinovici Sánchez-Garrido RS Sonner 2023] we showed that this Hamiltonian is written in the Krylov basis for H with initial state

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$$C_K(t) = \langle \hat{n}(t) \rangle \quad (19)$$

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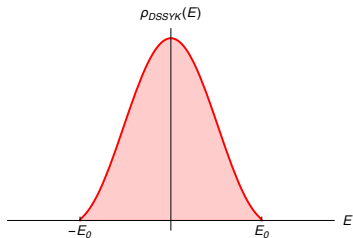
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$$H = -\frac{2J}{\lambda} + 2\lambda J \left(\frac{k^2}{2} + 2e^{-\tilde{l}} \right) \quad (21)$$

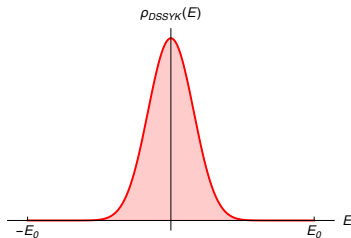
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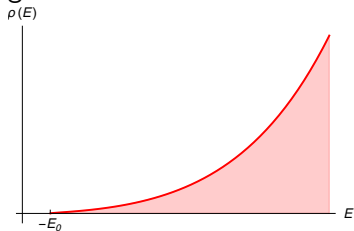
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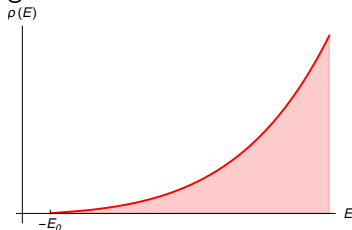
The triple-scaled limit of SYK [Lin 2022]

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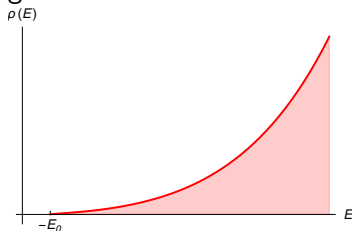
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- ▶ The d.o.s. in the triple-scaled limit is $\rho(E) \propto \sinh(2\pi\sqrt{E})$
- ▶ The **Liouville** form of the triple-scaled Hamiltonian connects DSSYK with the Hamiltonian of JT gravity

K-complexity in the triple-scaled limit of SYK

[Rabinovici Sánchez-Garrido RS Sonner 2023]

- ▶ Recall that $l = \lambda n$ and hence $C_K(t) = \frac{\langle l(t) \rangle}{\lambda}$
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$$\lambda C_K(t) = 2 \log [\cosh (2\lambda J e^{-x_0} t)] + x_0 \quad (22)$$

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The bulk: JT gravity

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- ▶ The action of JT gravity

$$S_{JT} = \int_{\mathcal{M}} d^2x \sqrt{-g} [\Phi_0 R + \Phi(R + 2)] \\ + 2 \int_{\partial\mathcal{M}} dx \sqrt{\gamma} [\Phi_0 K + \Phi(K - 1)] \quad (23)$$

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$$ds^2 \Big|_{\partial\mathcal{M}} = -\frac{dt_b^2}{\epsilon^2}, \quad \Phi \Big|_{\partial\mathcal{M}} = \frac{\phi_b}{\epsilon} \quad (24)$$

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- ▶ Equations of motion

$$0 = R + 2 \implies \text{AdS}_2 \quad (25)$$

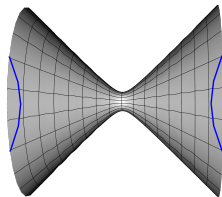
$$0 = (\nabla_\mu \nabla_\nu - g_{\mu\nu})\Phi \quad (26)$$

Wormhole (ERB) length

[Harlow Jafferis 2018]

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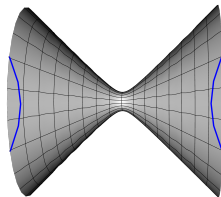
(a) Global coordinates

$$ds^2 = -(1+x^2)d\tau^2 + \frac{dx^2}{1+x^2}$$

$$\Phi(x, \tau) = \Phi_h \sqrt{1+x^2} \cos \tau$$

Wormhole (ERB) length

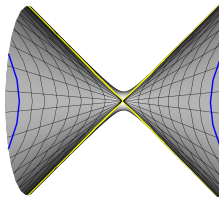
[Harlow Jafferis 2018]



(a) Global coordinates

$$ds^2 = -(1+x^2)d\tau^2 + \frac{dx^2}{1+x^2}$$

$$\Phi(x, \tau) = \phi_h \sqrt{1+x^2} \cos \tau$$



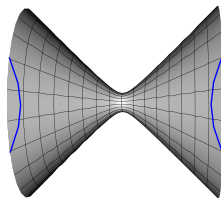
(b) Schwarzschild coordinates

$$ds^2 = -(r^2 - r_s^2)dt^2 + \frac{dr^2}{r^2 - r_s^2}$$

$$\Phi(r, t) = \phi_b r$$

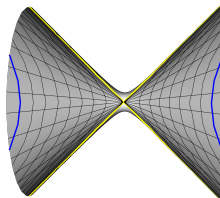
Wormhole (ERB) length

[Harlow Jafferis 2018]



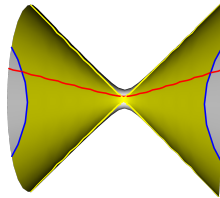
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$$ds^2 = -(1+x^2)d\tau^2 + \frac{dx^2}{1+x^2}$$
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(b) Schwarzschild coordinates

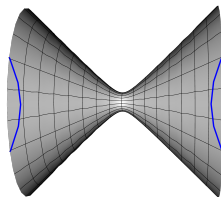
$$ds^2 = -(r^2 - r_s^2)dt^2 + \frac{dr^2}{r^2 - r_s^2}$$
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(c) The **wormhole** length defined as the geodesic distance between two points on the boundaries

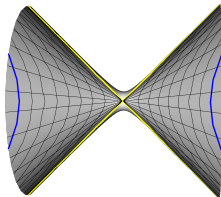
Wormhole (ERB) length

[Harlow Jafferis 2018]



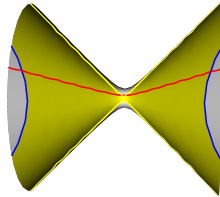
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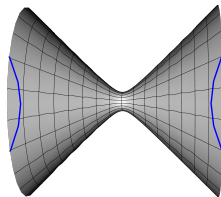
(c) The **wormhole** length defined as the geodesic distance between two points on the boundaries

► Wormhole length

$$\tilde{l} = 2 \log \left(\frac{2\phi_b}{\epsilon} \right) + 2 \log \left[\cosh \left(\frac{\Phi_h}{\phi_b} t_b \right) \right] - 2 \log \Phi_h \quad (27)$$

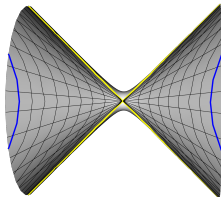
Wormhole (ERB) length

[Harlow Jafferis 2018]



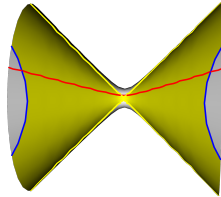
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$$ds^2 = -(1+x^2)d\tau^2 + \frac{dx^2}{1+x^2}$$
$$\Phi(x, \tau) = \phi_h \sqrt{1+x^2} \cos \tau$$



(b) Schwarzschild coordinates

$$ds^2 = -(r^2 - r_s^2)dt^2 + \frac{dr^2}{r^2 - r_s^2}$$
$$\Phi(r, t) = \phi_b r$$

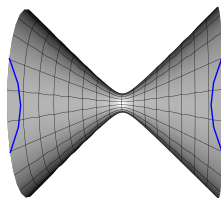


(c) The **wormhole** length defined as the geodesic distance between two points on the boundaries

► **Renormalized** wormhole length

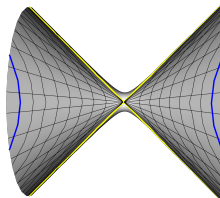
Wormhole (ERB) length

[Harlow Jafferis 2018]



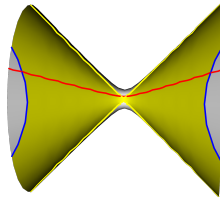
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$$\Phi(r, t) = \phi_b r$$



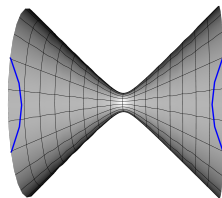
(c) The **wormhole** length defined as the geodesic distance between two points on the boundaries

► Renormalized wormhole length

$$\tilde{l} = 2 \log \left[\cosh \left(\frac{\Phi_h}{\phi_b} t_b \right) \right] - 2 \log \Phi_h \quad (27)$$

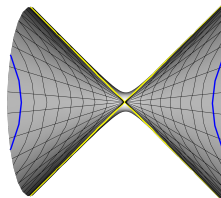
Wormhole (ERB) length

[Harlow Jafferis 2018]



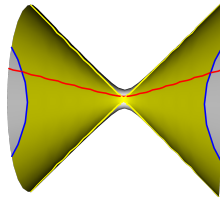
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(c) The **wormhole** length defined as the geodesic distance between two points on the boundaries

► Renormalized wormhole length

$$\tilde{l} = 2 \log \left[\cosh \left(\sqrt{\frac{E}{2\phi_b}} t_b \right) \right] - \log \left(\frac{\phi_b E}{2} \right) \quad (27)$$

Krylov complexity = wormhole length

[Rabinovici Sánchez-Garrido RS Sonner 2023]

K-complexity in triple-scaled SYK

$$\lambda C_K(t) = 2 \log \left[\cosh \left(\sqrt{\lambda J E} t \right) \right] - \log \left(\frac{E}{4\lambda J} \right) \quad (28)$$

Krylov complexity = wormhole length

[Rabinovici Sánchez-Garrido RS Sonner 2023]

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Wormhole length in JT gravity

$$\tilde{l}(t) = 2 \log \left[\cosh \left(\sqrt{\frac{E}{2\phi_b}} t \right) \right] - \log \left(\frac{\phi_b E}{2} \right) \quad (29)$$

Summary, open questions and future directions

Result

Krylov basis in triple-scaled SYK = Wormhole length basis in JT gravity

Summary, open questions and future directions

Result

Krylov basis in triple-scaled SYK = Wormhole length basis in JT gravity

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Krylov basis in triple-scaled SYK = Wormhole length basis in JT gravity

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- ▶ Quantum corrections?

Summary, open questions and future directions

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Krylov basis in triple-scaled SYK = Wormhole length basis in JT gravity

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- ▶ Quantum corrections?
- ▶ Going away from the small λ limit

Summary, open questions and future directions

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- ▶ Saturation of K-complexity?

Summary, open questions and future directions

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Krylov basis in triple-scaled SYK = Wormhole length basis in JT gravity

Krylov complexity in triple-scaled SYK = Wormhole length in JT gravity

- ▶ Quantum corrections?
- ▶ Going away from the small λ limit
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- ▶ The Krylov basis with operators is richer, what is the precise geometric interpretation of operator insertions in DSSYK

Summary, open questions and future directions

Result

Krylov basis in triple-scaled SYK = Wormhole length basis in JT gravity

Krylov complexity in triple-scaled SYK = Wormhole length in JT gravity

- ▶ Quantum corrections?
- ▶ Going away from the small λ limit
- ▶ Saturation of K-complexity?
- ▶ The Krylov basis with operators is richer, what is the precise geometric interpretation of operator insertions in DSSYK
- ▶ Higher dimensional gravity?