

Irreversibility, QNEC, and defects

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Irreversibility of the renormalization group is a key property of nonperturbative QFT.

Goal: find renormalization group charges that partially characterize conformal field theories, and that decrease under the RG.

▶ Two directions where irreversibility has been developed:

- 1) Irreversibility theorems in unitary relativistic QFT have been established in $d=2, 3$ and 4 space-time dimensions (C, F, A)
- 2) Irreversibility for defect RG flows, embedded in CFTs. Started with the g-theorem.

▶ Framework for generalizing this:

- D -dimensional CFT coupled to d -dimensional planar static defect
- relevant interactions on defect trigger RG flow between UV and IR fixed points. Bulk remains conformal
- QFTs without defect: special case $D=d$



Overview of irreversibility inequalities

So far, several proofs for different (D, d) . Those based on properties of energy-momentum tensor are:

$d \setminus D$	2	3	4	5	...
1	reflection positivity for stress tensor	reflection positivity for stress tensor	reflection positivity for stress tensor	reflection positivity for stress tensor	reflection positivity for stress tensor
2	reflection positivity for stress tensor	reflection positivity of dilaton	reflection positivity of dilaton	reflection positivity of dilaton	reflection positivity of dilaton
3		No proof	No proof	No proof	No proof
4			unitarity dilaton scattering	unitarity dilaton scattering	unitarity dilaton scattering

$D=d=2$: [Zamolodchikov]'s original C-thm; $D=d=4$: [Komargodski, Schwimmer] (dilaton)

Later on: $d=2$, any D [Jensen, OBannon], $d=4$, any D : [Wang]. Dilaton methods

$d=1$, $D=2$: [Friedan, Konechny].

And recently: $d=1$, any D : [Cuomo, Komargodski, Raviv-Moshe]. Generalized to $d=2$, any D by [Sachar, Sinha, Smolkin].

In parallel, irreversibility theorems have also been obtained using methods from quantum information theory:

$d \setminus D$	2	3	4	5	...
1	positivity of relative entropy	positivity of relative entropy	positivity of relative entropy	positivity of relative entropy	positivity of relative entropy
2	SSA or positivity of relative entropy	SSA + QNEC or positivity of relative entropy	SSA + QNEC or positivity of relative entropy	SSA + QNEC or positivity of relative entropy	SSA + QNEC or positivity of relative entropy
3		SSA	SSA + QNEC	SSA + QNEC	SSA + QNEC
4			SSA	SSA + QNEC	SSA + QNEC

$D=d=2$ entropic C-thm by [Casini, Huerta];

extended by [Casini, Huerta] to $D=d=3$ (not available via correlators/dilaton)

$D=d=4$ by [Casini, Teste, GT]. And unifies $d=2,3,4$ (Markov prop)

$d=1$, any D [Casini, Salazar, GT]; $d=2$, $D=3$ [Casini, Salazar, GT]

... This covers almost 40 years of developments! But a simple and general understanding is still lacking.

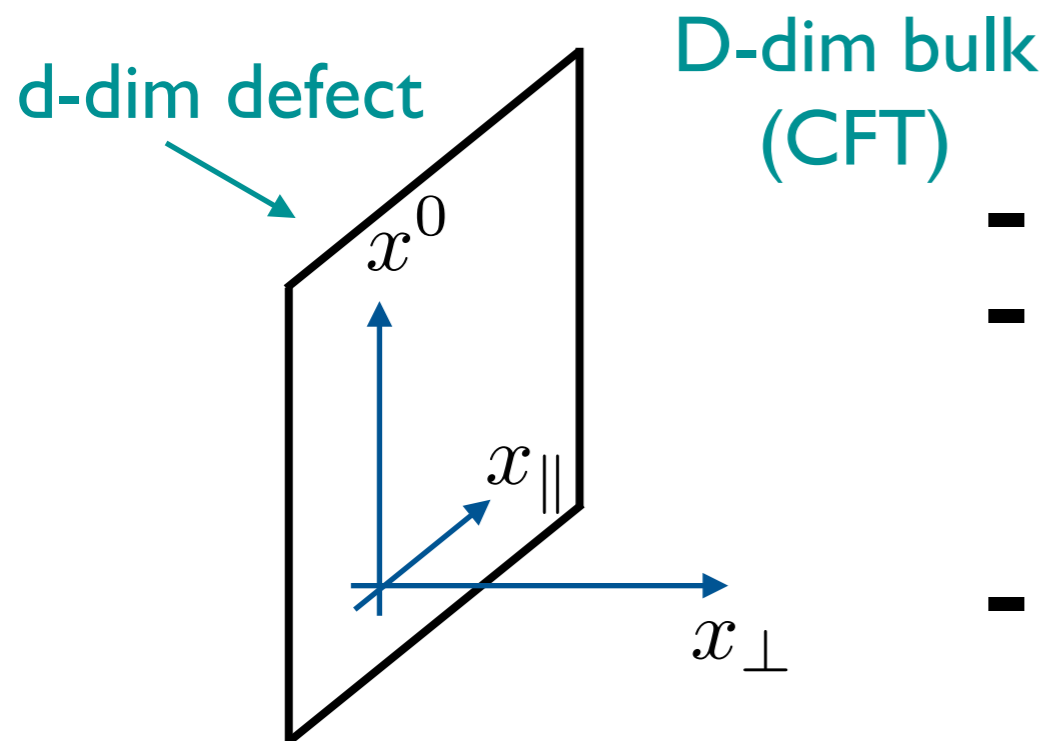
In this talk: we prove an inequality for second derivatives of the Relative Entropy. This establishes the remaining irreversibility thms, and unifies all known theorems for (D, d) . [Casini, Salazar, GT, 2023]

A. Key ingredients



QFT setup

D-dimensional CFT with a d-dimensional planar defect

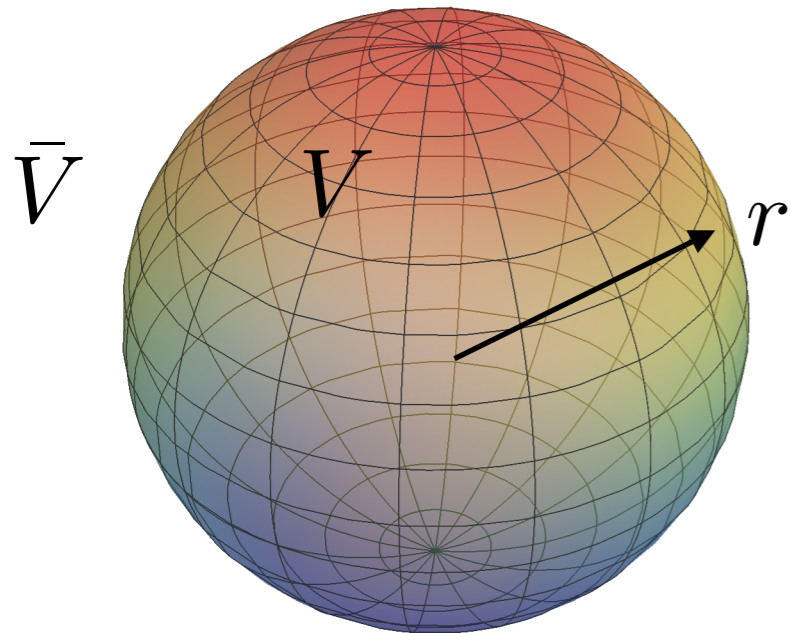


- the defect is conformal in the UV
- turn on relevant deformations on the defect
$$S = S_{UV,CFT} + \int d^d x g \mathcal{O}$$
- triggers an RG flow, which ends on a different IR conformal defect
- the bulk does not flow

Particular case: $d=D$, no bulk. Gives QFT RG flow without defects



Entanglement entropy



density matrix $\rho_r = \text{tr}_{\bar{V}} |0\rangle\langle 0|$

von Neumann entropy $S(r) = -\text{tr}_V (\rho_r \log \rho_r)$

Can probe RG flow by

$$mr \ll 1 \text{ (UV)} \Rightarrow mr \gg 1 \text{ (IR)}$$

- Structure near a fixed point, first without defects ($D=d$):

$$S(r) = S_{local} + S_{non-local}$$

$$S_{local}(r) = \mu_{d-2} r^{d-2} + \mu_{d-4} r^{d-4} + \dots \quad \text{UV divergent}$$

$$S_{non-local}(r) = \begin{cases} (-1)^{\frac{d}{2}-1} 4A \log \frac{r}{\epsilon} & d \text{ even} \\ (-1)^{\frac{d-1}{2}} F & d \text{ odd} \end{cases} \quad \text{universal}$$

RG irreversibility: decrease of A or F; proved for $d = 2, 3, 4$

The universal terms can also be isolated from the partition function on the sphere (free energy in de Sitter), or from the Weyl anomaly if d even.

► For d -dim conformal defect in a D -dim CFT, there are additional terms:

$$S(r) = \mu_{D-2} r^{D-2} + \mu_{D-4} r^{D-4} + \dots + \tilde{\mu}_{d-2} r^{d-2} + \tilde{\mu}_{d-4} r^{d-4} + \dots$$

$$+ \begin{cases} (-1)^{\frac{D-2}{2}} 4A \log \frac{r}{\epsilon} & D \text{ even} \\ (-1)^{\frac{D-1}{2}} F & D \text{ odd} \end{cases} + \begin{cases} (-1)^{\frac{d-2}{2}} 4\tilde{A} \log \frac{r}{\epsilon} & d \text{ even} \\ (-1)^{\frac{d-1}{2}} \tilde{F} & d \text{ odd} \end{cases}$$

For defect RG flows, quantities with tildes flow. One can try to prove irreversibility theorems for the universal \tilde{F} , \tilde{A} , but in general this does not work (exception: $D=d+1$).

[Jensen, O'Bannon]
[Kobayashi, Nishioka, Sato, Watanabe]

One reason: the defect contributes nonzero energy. The universal term in the EE no longer coincides with the one in the free energy. We should look at the quantum-information analog of free energy.



Relative entropy

For two density matrices σ and ρ , the relative entropy is

$$S_{rel}(\rho|\sigma) = \text{tr}(\rho \log \rho - \rho \log \sigma)$$

Introducing the modular Hamiltonian $H = -\log \sigma$

$$S_{rel}(\rho|\sigma) = \text{tr}(\rho \log \rho - \rho \log \sigma + \sigma \log \sigma - \sigma \log \sigma)$$

$$= \langle H \rangle_\rho - \langle H \rangle_\sigma - (S_\rho - S_\sigma)$$

$$= \Delta \langle H \rangle - \Delta S \quad \leftarrow \text{difference of “free energies”}$$

For irreversibility in QFT, we will compare two density matrices:

$\sigma =$ vacuum density matrix for UV fixed point

$\rho =$ vacuum density matrix for QFT w/relevant deformations

The modular Hamiltonian for a CFT on a sphere is

$$H_\sigma = \int_\Sigma d^{D-1}x \eta^\mu \xi^\nu T_{\mu\nu}$$

-conf. transf. of Rindler Hamiltonian
- valid also with conformal defects

η^μ : unit normal to Cauchy surface Σ

$\xi^\nu = \frac{\pi}{R} (R^2 - (x^0)^2 - \vec{x}^2, -2x^0 x^i)$ Killing vector, R: radius of sphere

Note: $\langle H \rangle_\rho$ depends on choice of Cauchy surface. Reason: states evolve with different action.

- for space-like Σ , $\langle H \rangle_\rho \sim R^D$ dominates relative entropy
- for null Σ , $\langle H \rangle_\rho \sim R^{D-2}$ comparable to EE. We choose this limit. It contributes a universal term, proportional to energy of defect. Then

$$-\lim_{R \rightarrow \infty} S_{\text{rel}}(R) = \Delta\mu'_{d-2} R^{d-2} + \Delta\mu'_{d-4} R^{d-4} + \dots + \begin{cases} (-)^{\frac{d-2}{2}} 4 \Delta A' \log(R/\epsilon) & d \text{ even} \\ (-)^{\frac{d-1}{2}} \Delta F' & d \text{ odd} \end{cases}$$



Strong subadditivity and Markov property

Key property of unitary quantum mechanics: strong subadditivity of the EE

[Lieb, Ruskai, 1973]

$$S(A) + S(B) \geq S(A \cup B) + S(A \cap B)$$

For a CFT and for regions with boundary in null cone, **Markov property:**

[Casini, Testé, GT, 2017]

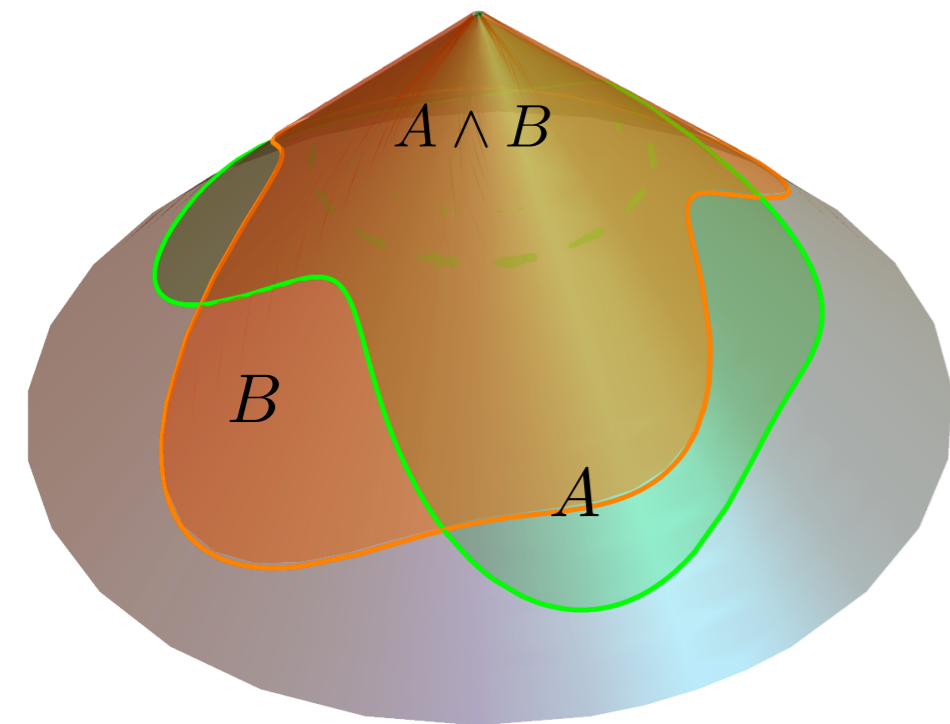
$$S(A \cup B) + S(A \cap B) = S(A) + S(B)$$

$$\Leftrightarrow \log \rho_{A \cup B} = \log \rho_A + \log \rho_B - \log \rho_{A \cap B}$$

- ✓ This is called a quantum Markov state
- ✓ mod Hamiltonian local on null surfaces
- ✓ Tracing out a subsystem becomes a reversible process

Combining SSA w/Markov, we obtain strong superadditivity

$$S_{\text{rel}}(A) + S_{\text{rel}}(B) \leq S_{\text{rel}}(A \cup B) + S_{\text{rel}}(A \cap B)$$



B. Proof of irreversibility inequality

We will explain the proof for a $d=2$ defect, in D -dim CFT.

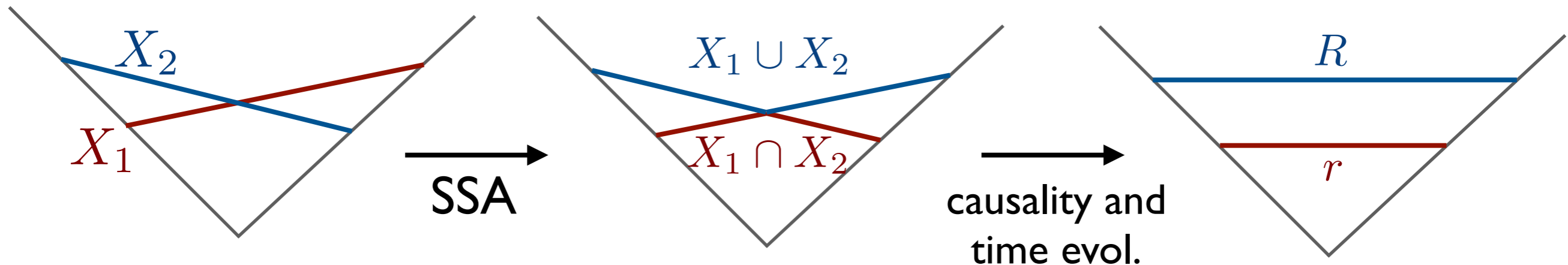
General case: [Casini, Salazar, GT, 2023] and talk at IFQ next week



The entropic C-theorem

[Casini, Huerta, '04, '12]

Warm-up without defect, ie $D=d=2$



$$\Delta S(X_1) + \Delta S(X_2) \geq \Delta S(X_1 \cup X_2) + \Delta S(X_1 \cap X_2) = \Delta S(R) + \Delta S(r)$$

$$2\Delta S(\sqrt{rR}) \geq \Delta S(R) + \Delta S(r). \text{ When } r \rightarrow R \Rightarrow (R \Delta S'(R))' \leq 0$$

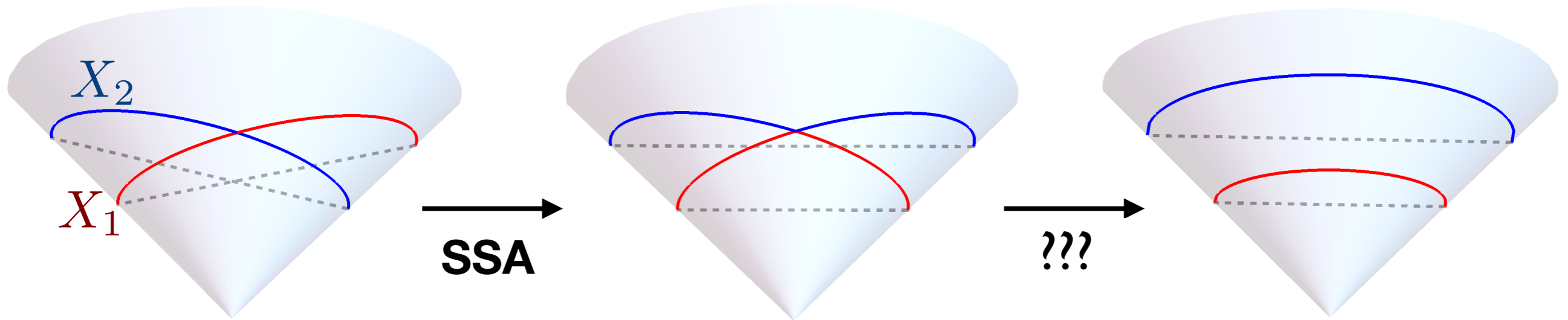
implies the decrease of the central charge $C_{IR} < C_{UV}$



d=2 defect

Let's try the same for a d=2 defect, in a D-dim CFT bulk.

The problem is that now the spheres extend into the bulk, and their causal union and intersection no longer give causal diamonds of spheres.



➔ New ingredient: use QNEC in the bulk (conformal), to bound the entropy of union and intersection by entropies of diamonds (spheres)

As we will see, this requires the relative entropy:

$$S_{\text{rel}}(X_1) + S_{\text{rel}}(X_2) \leq S_{\text{rel}}(X_1 \cup X_2) + S_{\text{rel}}(X_1 \cap X_2) \leq S_{\text{rel}}(R) + S_{\text{rel}}(r)$$

Then we would get defect irreversibility, $(RS'_{\text{rel}}(R))' \geq 0$



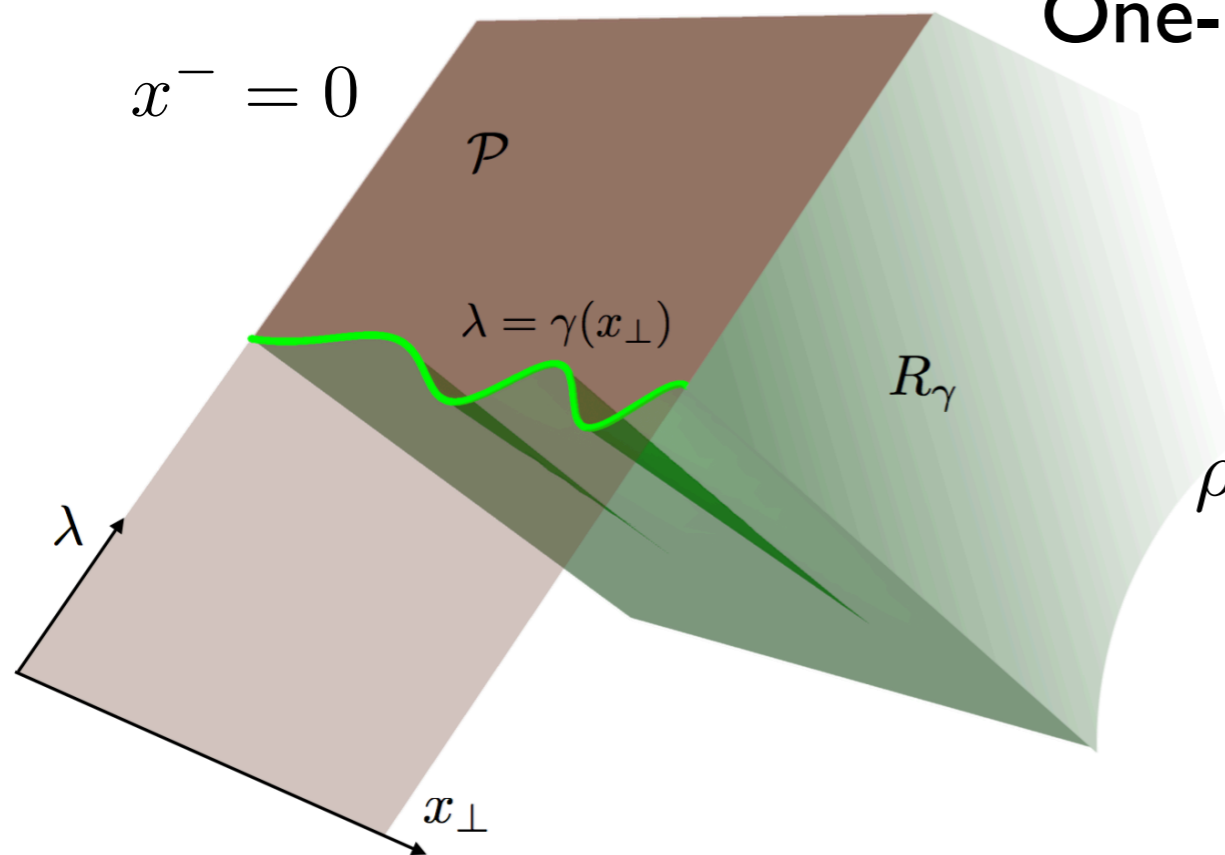
QNEC

[Bousso, Fisher, Leichenauer, Wall]

[Balakrishnan, Faulkner, Khandker, Wang]

[Ceyhan, Faulkner]

Originated from quantum focusing conjecture, but is valid in general QFT. Simplest setup: deformations on null plane



One-parameter deformation of null boundary:

$$x^+ = \gamma_s(x_\perp) = \gamma(x_\perp) + a(x_\perp)s$$

$$(a(x_\perp) \geq 0)$$

Relative entropy between σ (vacuum) and ρ (excited state) is convex under null defs

$$\frac{d^2 S_{\text{rel}}(\gamma_s)}{ds^2} \geq 0$$

QNEC is also valid on light-cone if we have a CFT. It is valid in the CFT bulk of the theory with defect.

$$\Rightarrow [S_{\text{rel}}(R) - S_{\text{rel}}(X_1 \cup X_2)] - [S_{\text{rel}}(X_1 \cap X_2) - S_{\text{rel}}(r)] \geq 0$$

and this establishes the irrev. inequality for $d=2$ and all D .

C. Conclusions

The result for general (d,D) is $RS''_{\text{rel}}(R) - (d - 3)S'_{\text{rel}}(R) \geq 0$

[Casini, Salazar, GT, 2023]

- The result is independent of D ; depends only on defect dim. d
- The theory on defect is nonlocal (it interacts w/bulk). It's remarkable that we get the same inequality as with no bulk ($d=D$)
- Relative entropy appears, required by QNEC.
- Implies **irreversibility of defect RG flows for $d=2,3,4$ and all D .**

Future directions:

- 1) Irreversibility for $d>4$? Requires new tools.
- 2) Theories with less symmetries? Interesting for condensed matter
- 3) Relation to non-entropic results?