
A Universal Pattern at Infinite Field Distance



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In collaboration with Etheredge, Heidenreich, McNamara, Rudelius and Ruiz: [2306.16440](#) and ongoing
Castellano and Ruiz: 23xx.xxxx
Montero: 23xx.xxxx

Strings 2023, Perimeter, July 2023

Given an Effective Field Theory (EFT) coupled to Einstein gravity,

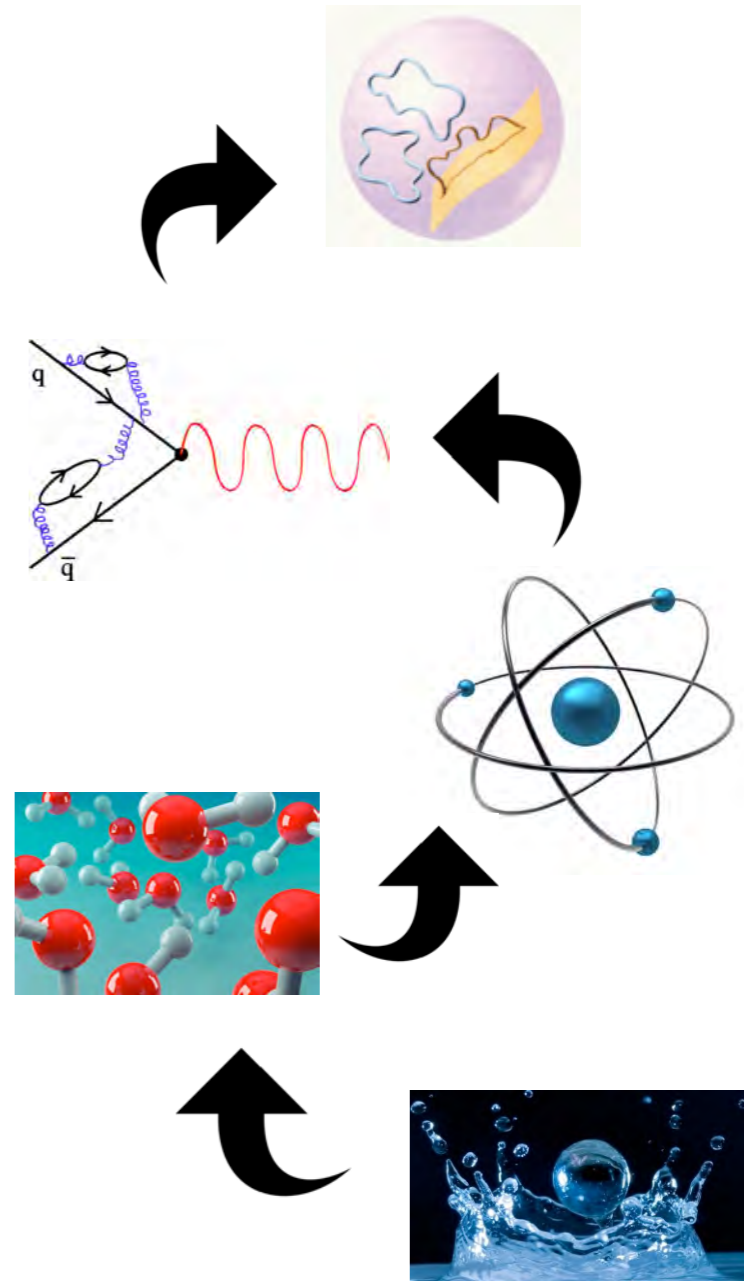
Energy



Quantum Gravity

cut-off scale Λ

EFT coupled to classical gravity



Given an Effective Field Theory (EFT) coupled to Einstein gravity, what is the cut-off at which semiclassical gravity breaks down **and how**?

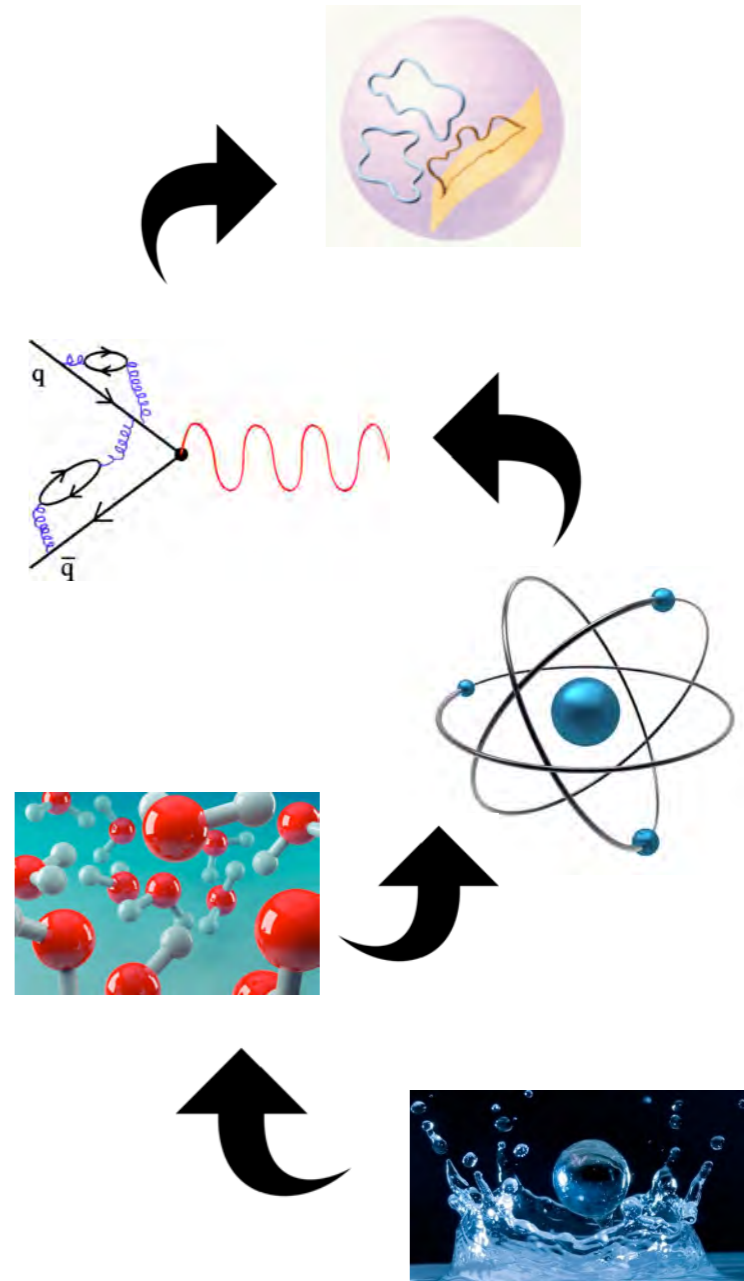
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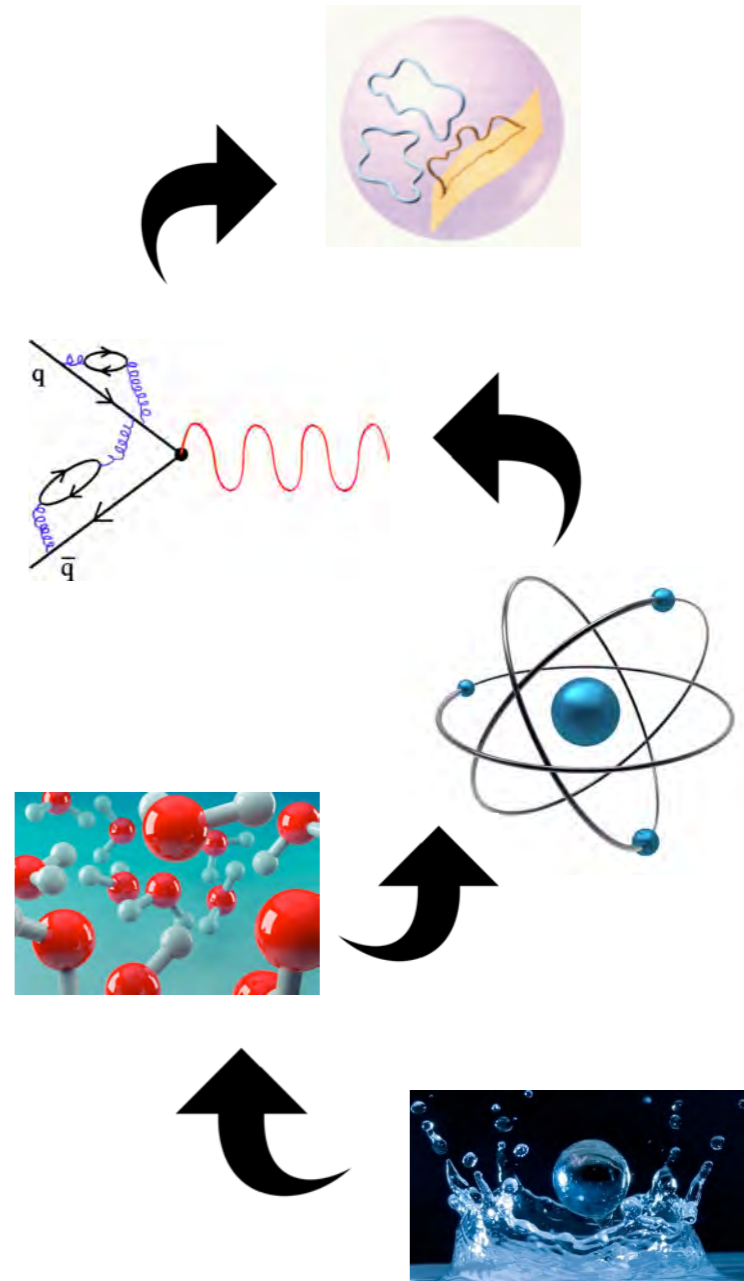


Is it M_p ?

Quantum Gravity

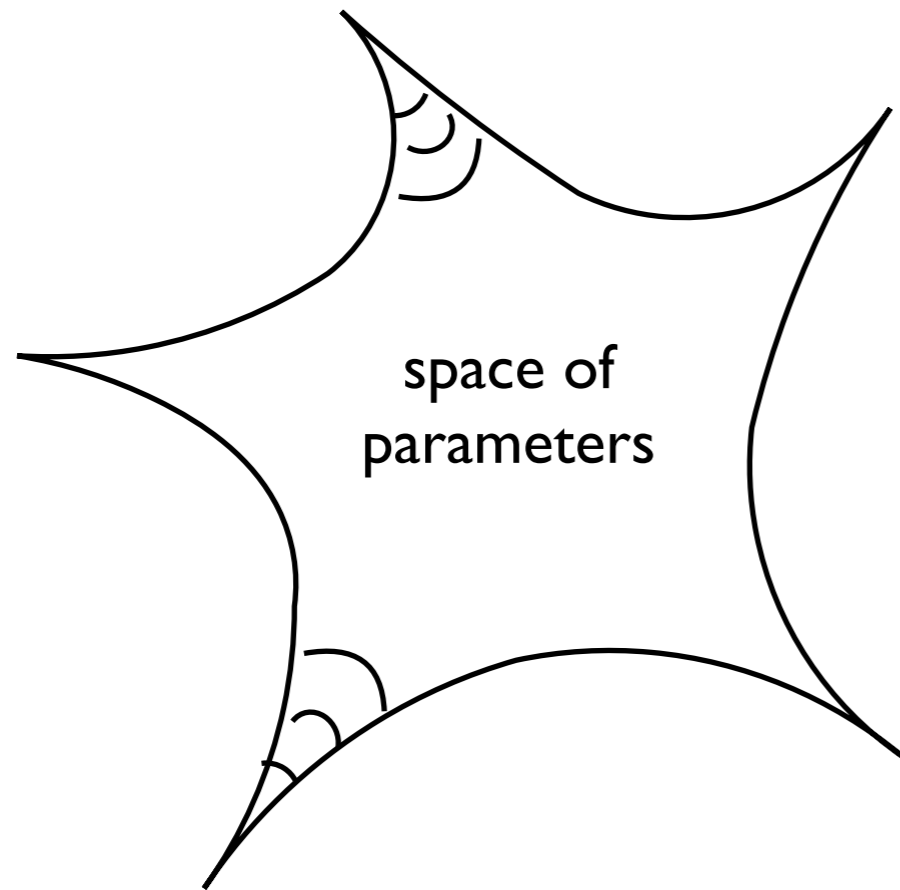
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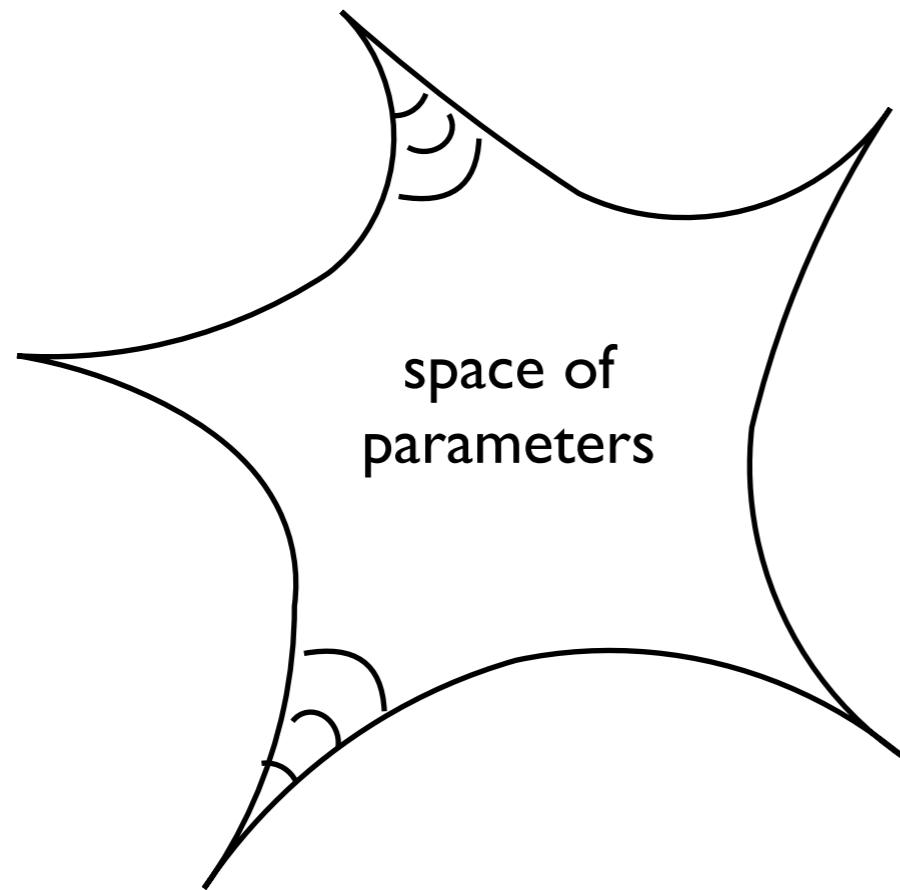
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(for certain regions of the space of EFT parameters)



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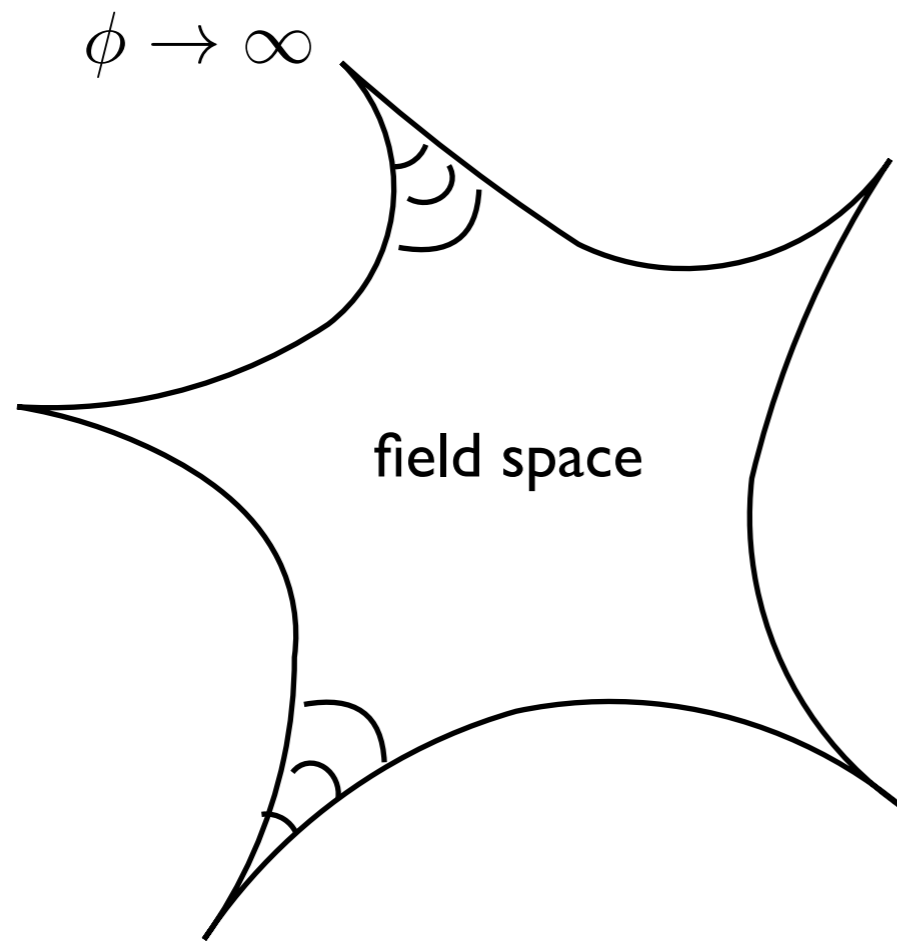


In String Theory:

EFT parameters are given by vacuum expectation values of scalar fields

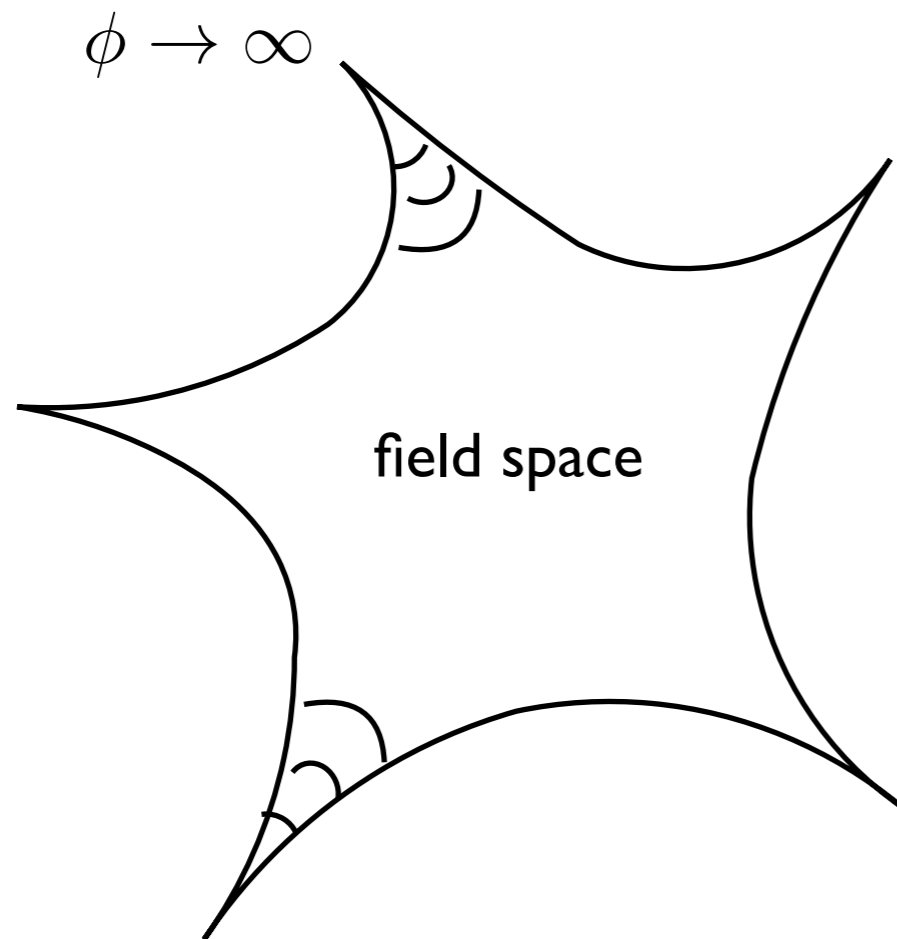
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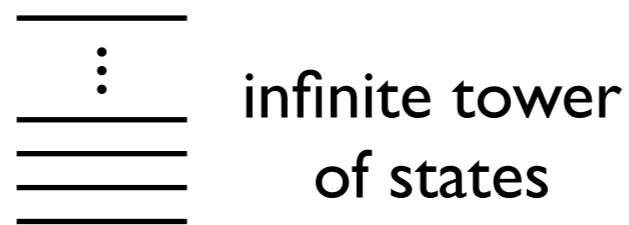


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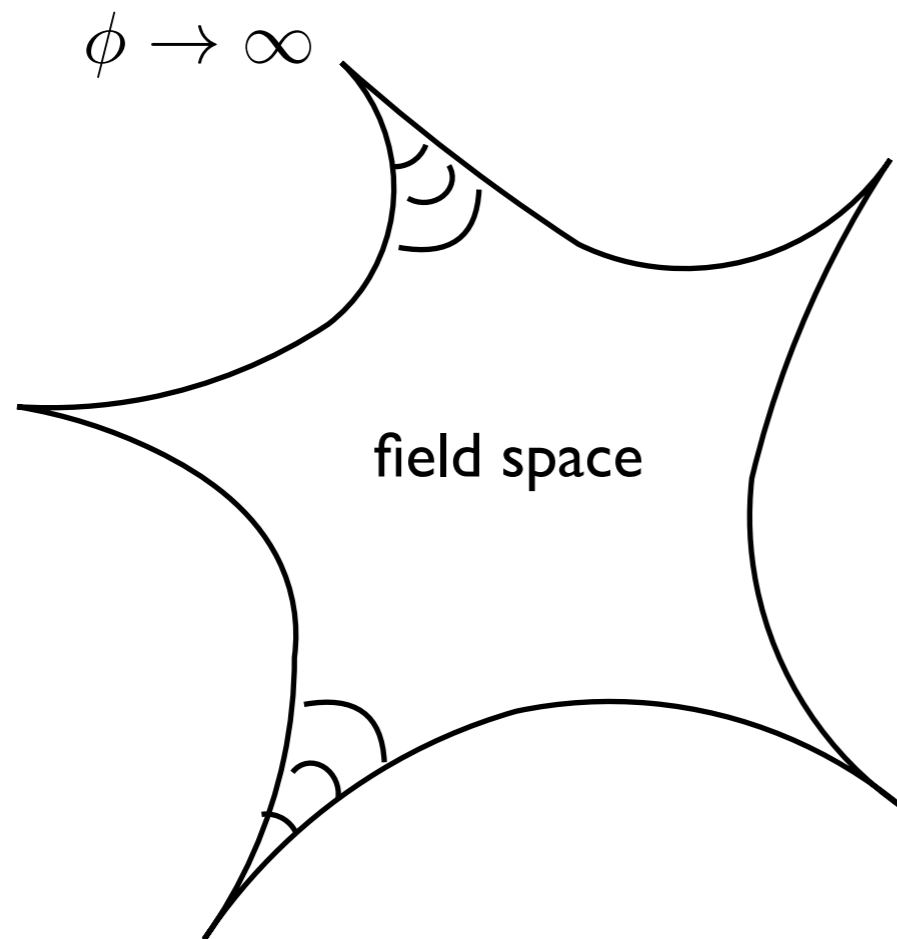


The cut-off goes to zero asymptotically:

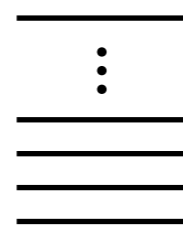


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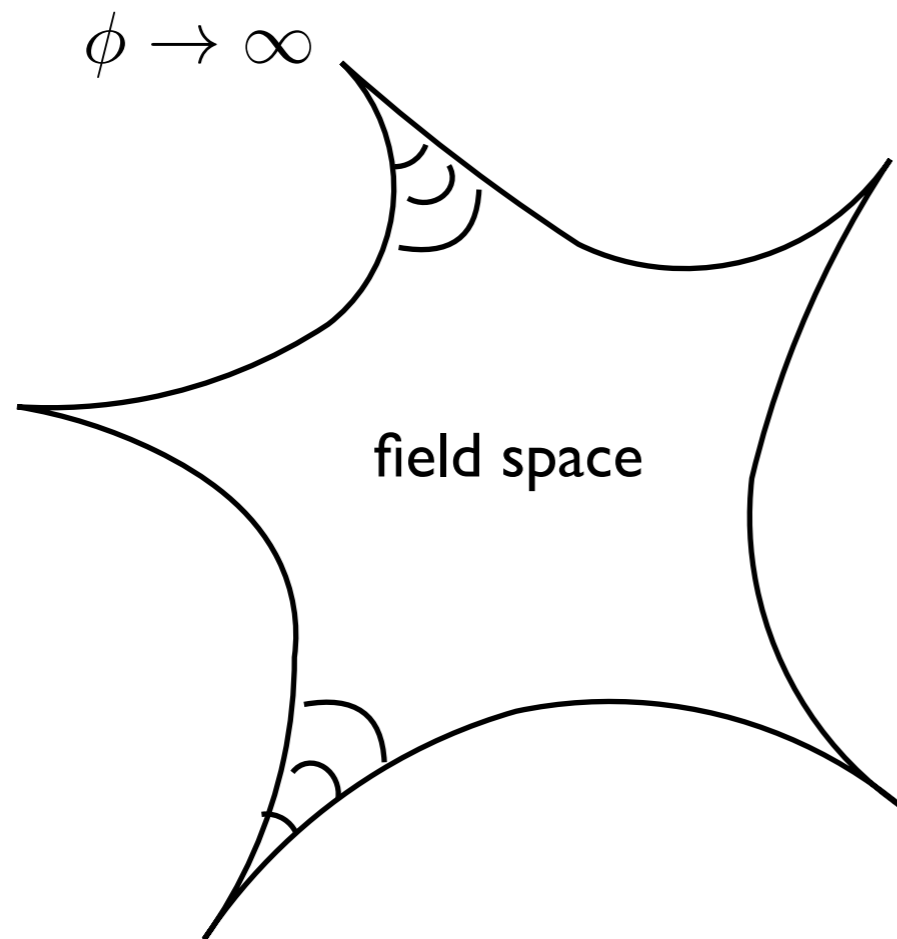
infinite tower
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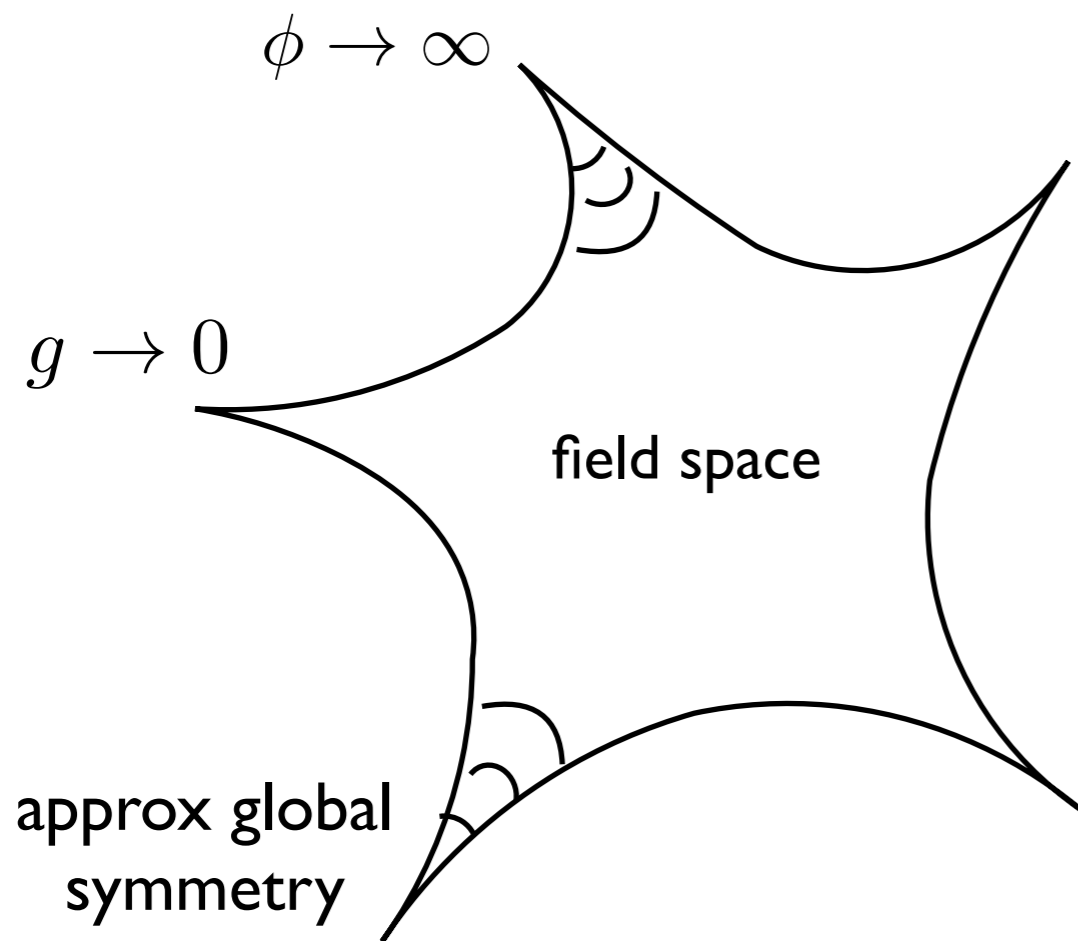
Swampland/quantum gravity constraint:

Distance Conjecture

[Ooguri-Vafa'06]

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From EFT perspective, these are e.g. weak coupling limits for gauge theories, approximate global symmetry limits, ... $\rightarrow \Lambda \ll M_p$

Outline

- 1) Review of Distance Conjecture and drop-off of QG cut-off

- 2) Universal pattern underlying all known string theory examples
 - ➔ To sharpen the conjecture and provide a quantitative bound on the mass of the tower

- 3) Beyond moduli spaces: Implications for accelerated expansion and AdS scale separation
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Distance Conjecture [Ooguri-Vafa'06]

Given an EFT coupled to gravity, with a moduli space parametrized by the vacuum expectation value of some scalar fields:

There is an **infinite tower of states** becoming **exponentially light** at every **infinite field distance** limit of the moduli space

$$m(P) \sim m(Q)e^{-\alpha\Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty$$

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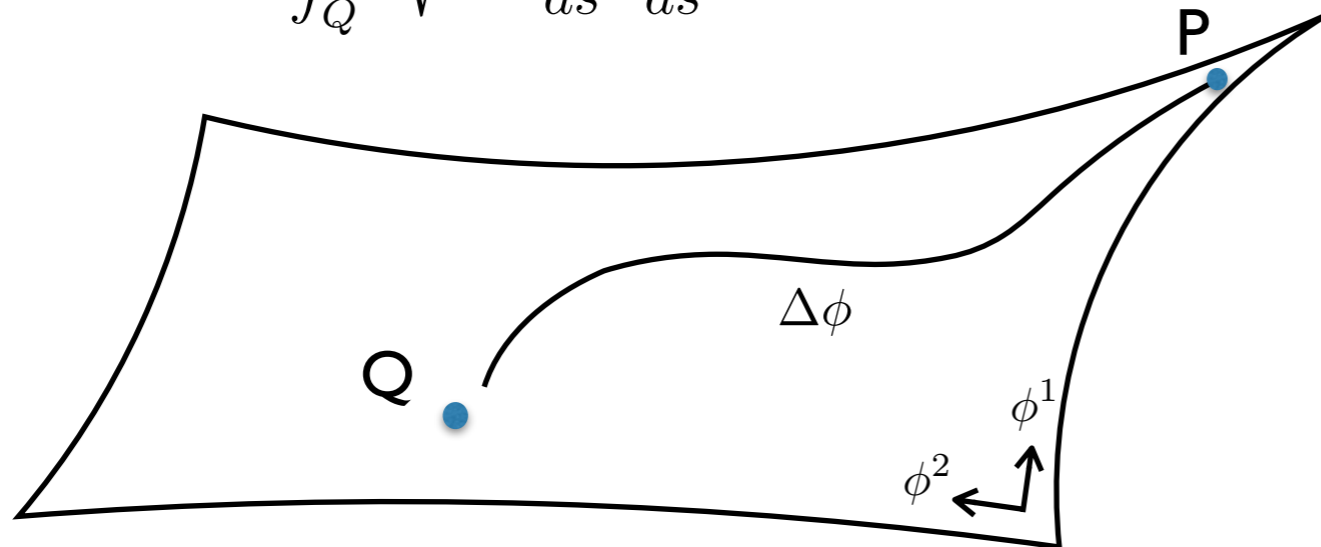
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$\mathcal{L} = g_{ij}(\phi)\partial\phi^i\partial\phi^j \rightarrow$ scalar manifold (moduli space)

$$\Delta\phi = \int_Q^P \sqrt{g_{ij} \frac{d\phi^i}{ds} \frac{d\phi^j}{ds}} ds \equiv \text{geodesic distance}$$



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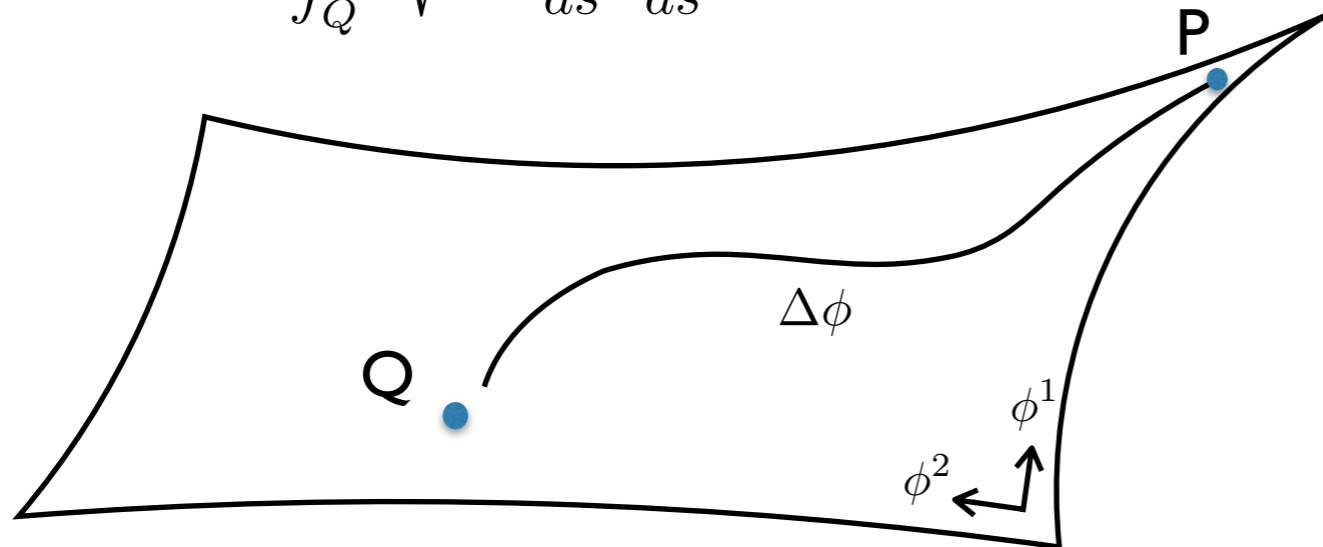
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For example:

- Kaluza-Klein towers as $R \rightarrow \infty$
- winding modes as $R \rightarrow 0$
- string modes as $g_s \rightarrow 0$

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The tower signals a drastic **breakdown** of the effective theory at:

Quantum Gravity (QG) cut-off scale Λ

(above Λ no local field theory description coupled to semiclassical gravity is possible anymore)

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- The QG cut-off is the **string scale** in an weak coupling string limit

$$\Lambda \simeq M_{\text{str}} \ll M_{\text{pl},d}$$

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In general, it is also known as the species scale $\Lambda = \frac{M_p}{N^{1/d-2}}$ [Arkani-Hamed et al'05]
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$$N = \int^{\Lambda} \rho(m) dm \quad (\text{number of single-particle states weakly coupled to gravity})$$

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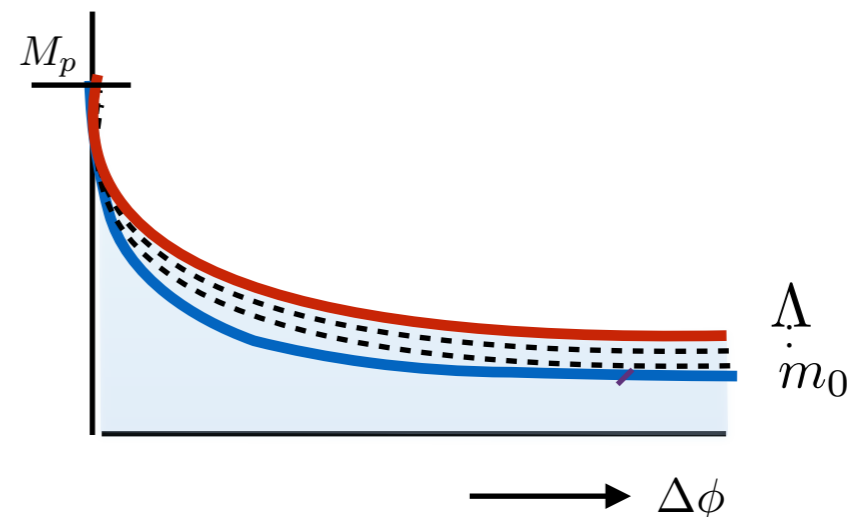
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In the presence of a tower of states $m \sim m_0 \exp(-\alpha \Delta\phi)$:

$$\Lambda \sim M_p \exp(-\lambda \Delta\phi)$$

$$\rightarrow \Delta\phi \lesssim \frac{1}{\lambda} \log \left(\frac{M_p}{\Lambda} \right)$$



Evidence

❖ Plethora of works testing the conjecture in string theory compactifications to flat space:

(M-theory and Heterotic toroidal comp., F-theory and IIB on Calabi-Yau's, M-theory on G2, Type IIA orientifolds, non-SUSY heterotic...)

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Many interesting connections with mathematics (algebraic geometry, etc.)!



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❖ $\text{AdS}_{d+1}/\text{CFT}_d$ All known infinite distance limits in the conformal manifold (for $d > 2$) contain towers of higher spin operators decaying exponentially with the Zamolodchikov distance



CFT Distance Conjecture [Perlmutter, Rastelli, Vafa, IV'21] [Baume, Calderon-Infante'21]

(part of the conjecture recently proven using CFT techniques in [Baume, Calderon-Infante'23])

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What general lessons can we extract from this?

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❖ Provide a bottom-up rationale beyond string theory examples

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- Emergence proposal [Grimm,Palti,IV'18] [Heidenreich et al'18] [Grimm,Palti,IV'18][Gendler,IV'20]
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In terms of the number of species:

“ lightest = less dense ”

$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} N}{N} = -1$$

$$N = \int^{\Lambda} \rho(m) dm \quad \text{with} \quad \Lambda \sim \frac{1}{N^{1/(d-2)}}$$

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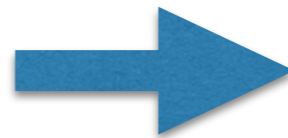
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Fixing the $\mathcal{O}(1)$ factors:

➔ Sharp upper bound on asymptotic field range: $\Delta\phi \leq \frac{1}{\alpha_{\min}} \log \left(\frac{M_p}{\Lambda} \right)$

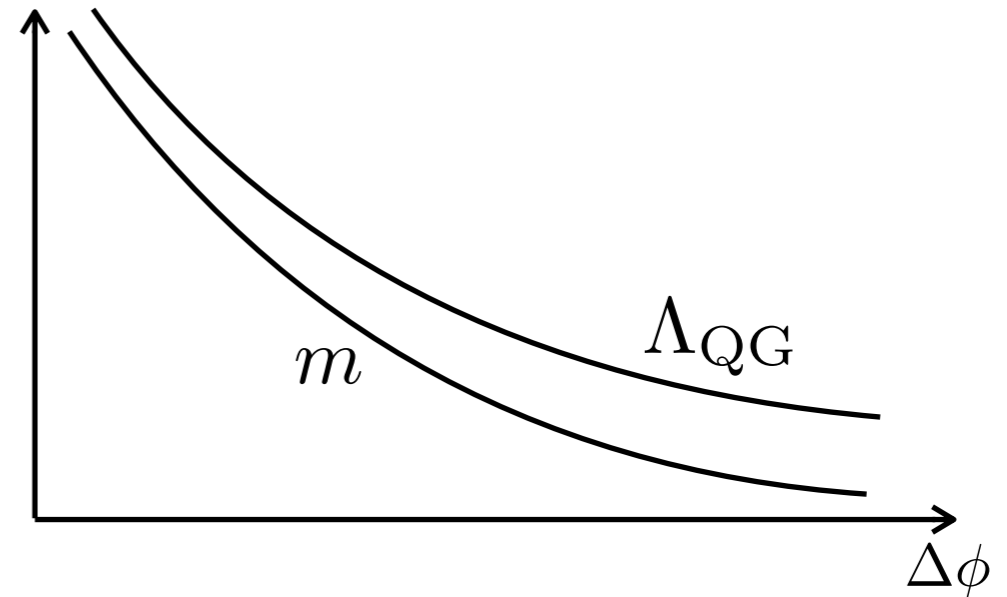
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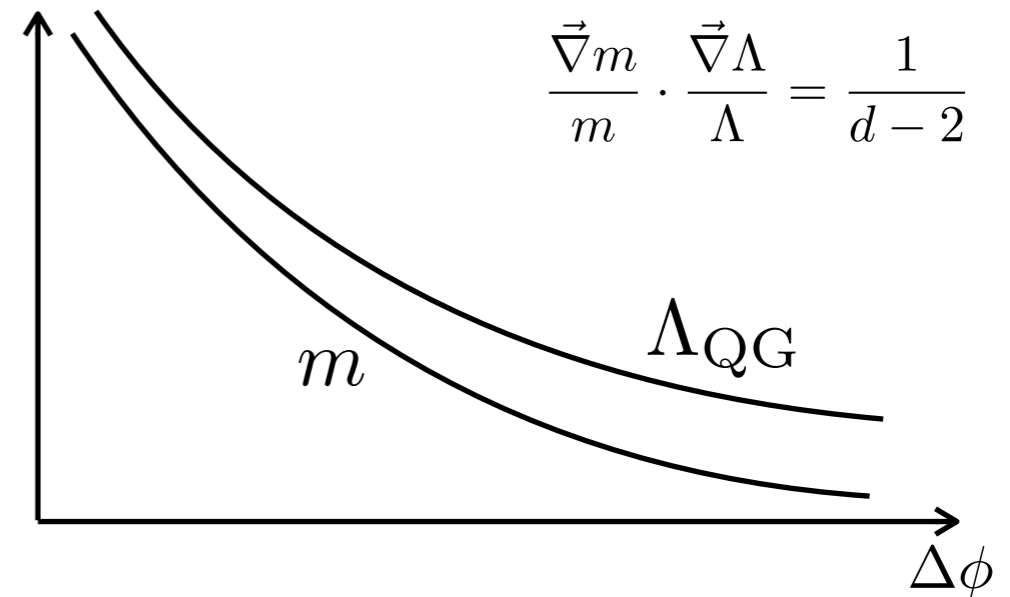
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We find that γ is related to the exponential decay rate α

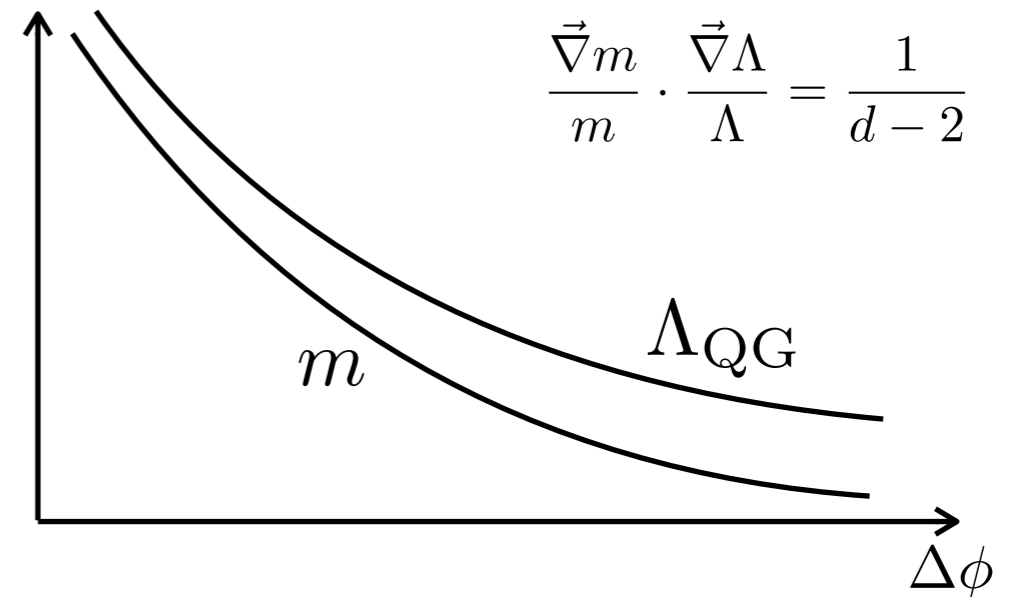
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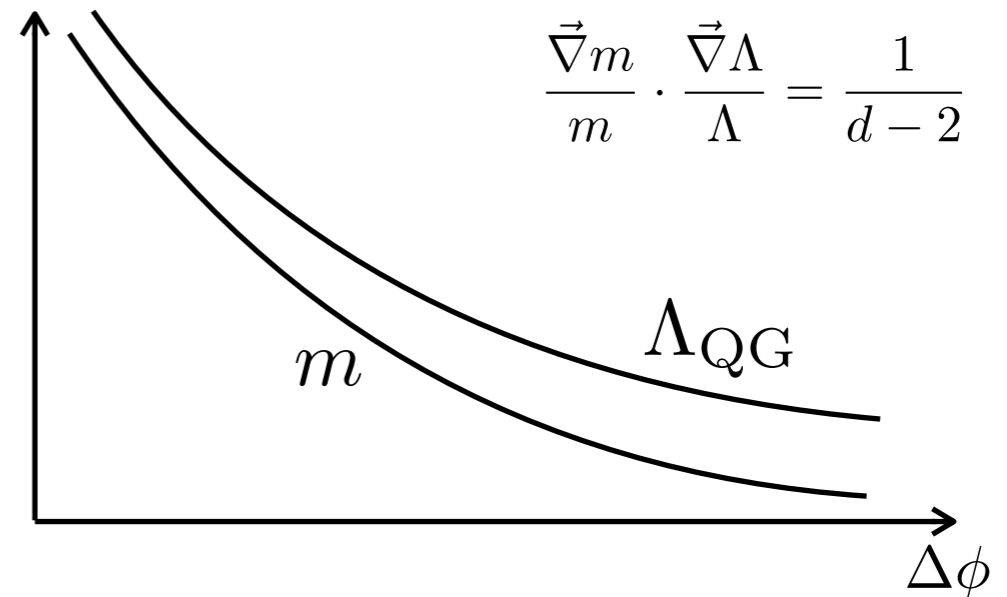
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$$\begin{array}{l}
 \text{---} \\
 \vdots \\
 \text{---} \\
 \text{---} \\
 \text{---}
 \end{array}
 \quad
 \begin{array}{l}
 m_k = k m_{\text{KK}} \\
 \Lambda = M_{p,d+n} \simeq m_{\text{KK}}^{1/(d+n-2)} \\
 \alpha_{\text{KK}} = \sqrt{\frac{d+n-2}{n(d-2)}}
 \end{array}$$

String tower



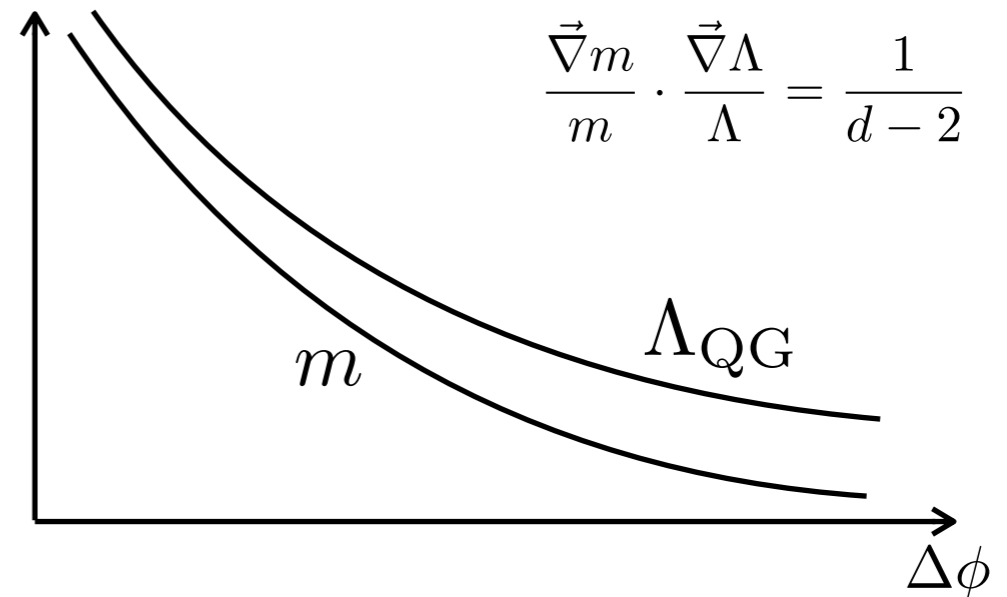
This should surprise you. Why?

Mass of the tower:

$$m \sim m_0 \exp(-\alpha \Delta\phi)$$

Quantum Gravity cut-off:

$$\Lambda_{\text{QG}} \sim m^\gamma \sim M_p \exp(-\lambda \Delta\phi)$$



I) The structure/density of the towers of states fixes γ

We find that γ is related to the exponential decay rate α

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String tower

$$\begin{array}{l} m_k = \sqrt{k} M_{\text{str}} , \rho(k) \sim e^{\sqrt{k}} \\ \Lambda \simeq M_{\text{str}} \\ \alpha_{\text{str}} = \frac{1}{\sqrt{d-2}} \end{array}$$

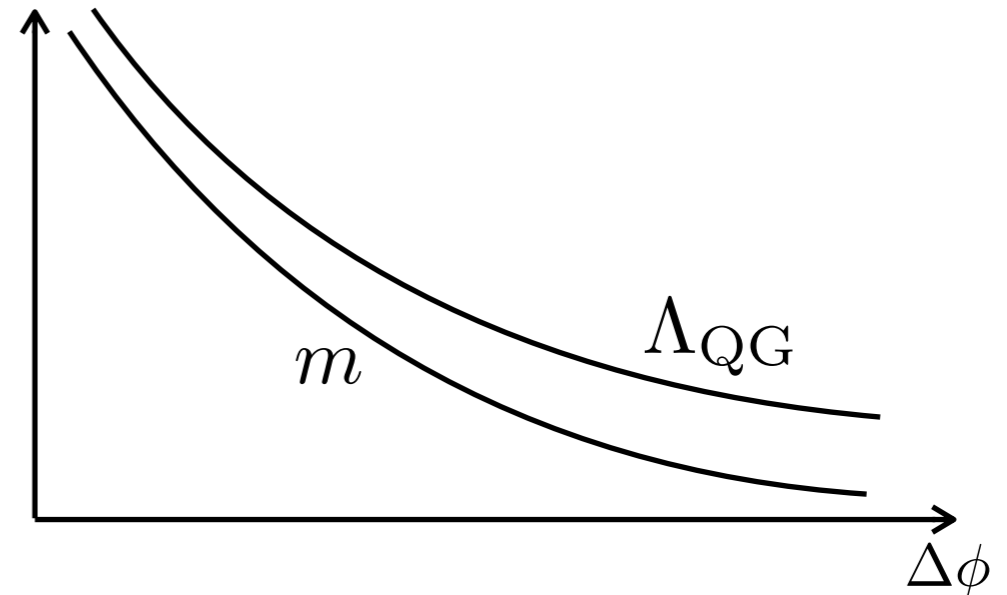
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2) Λ is sensitive to all towers becoming light, and not only the leading one

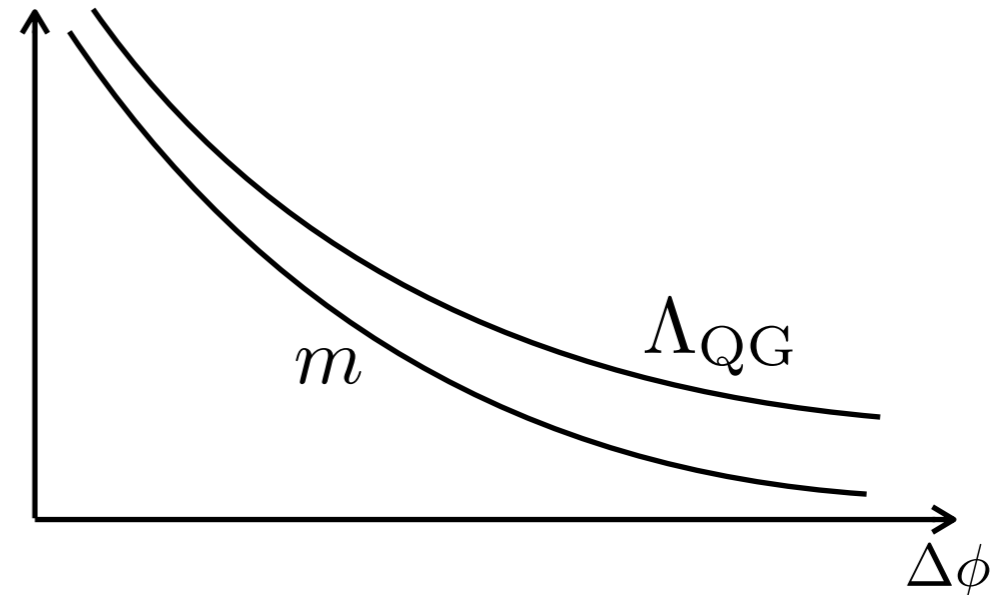
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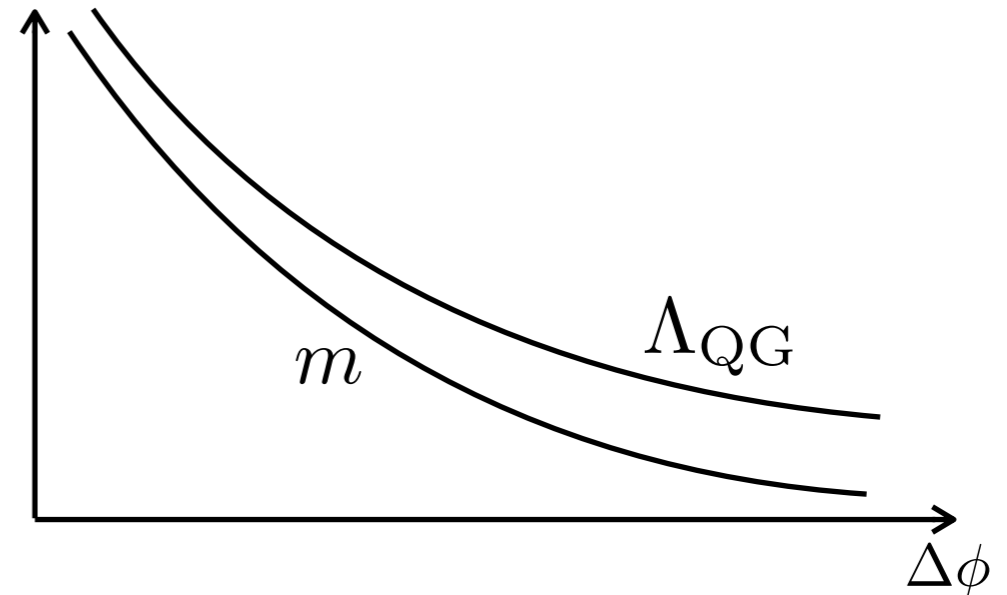
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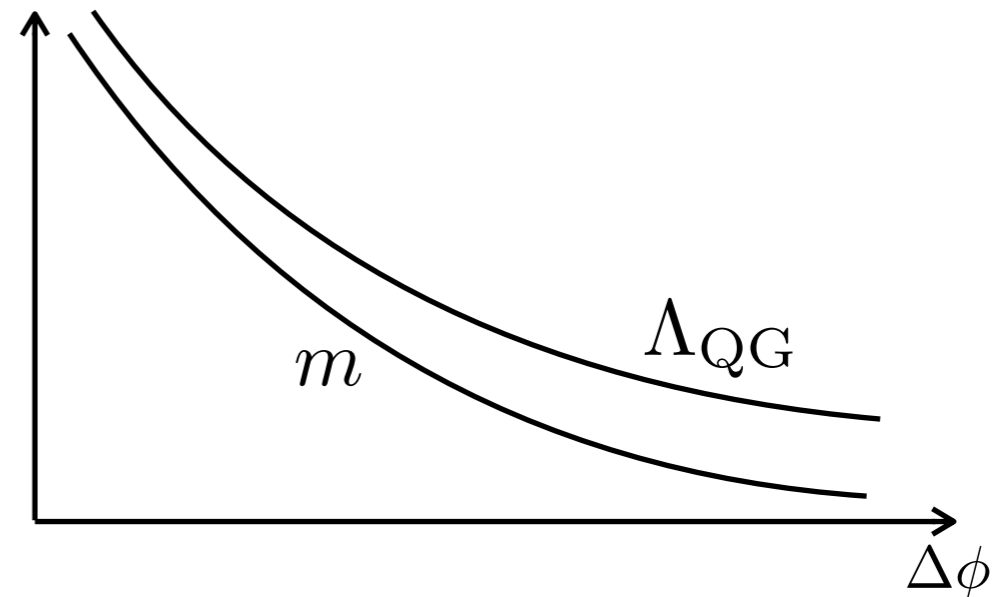
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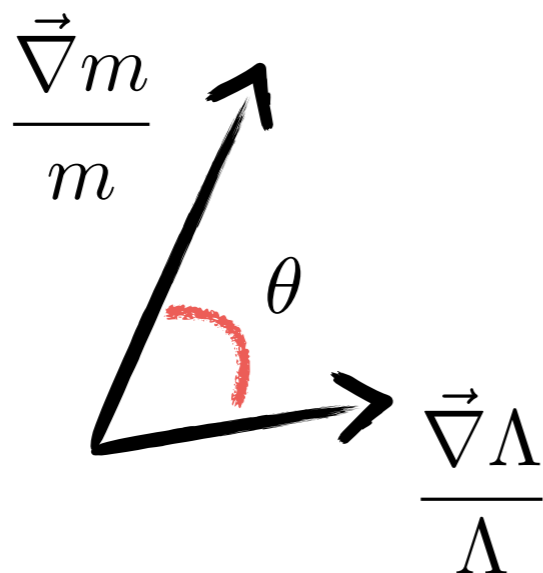
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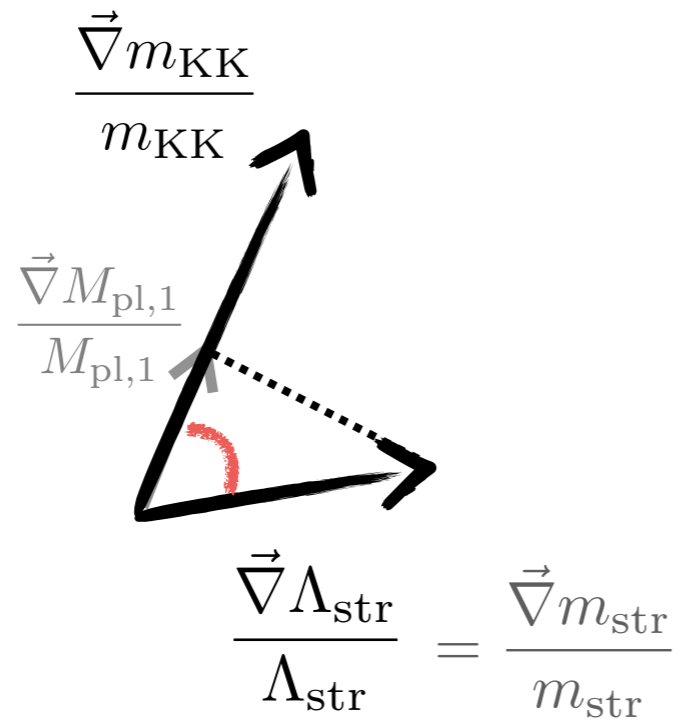
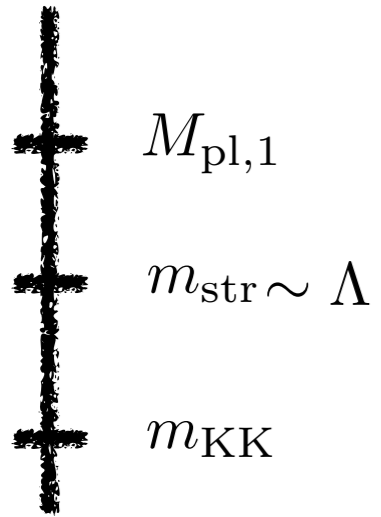
How??

The angle between them is always such that the pattern gets satisfied

Example 1:

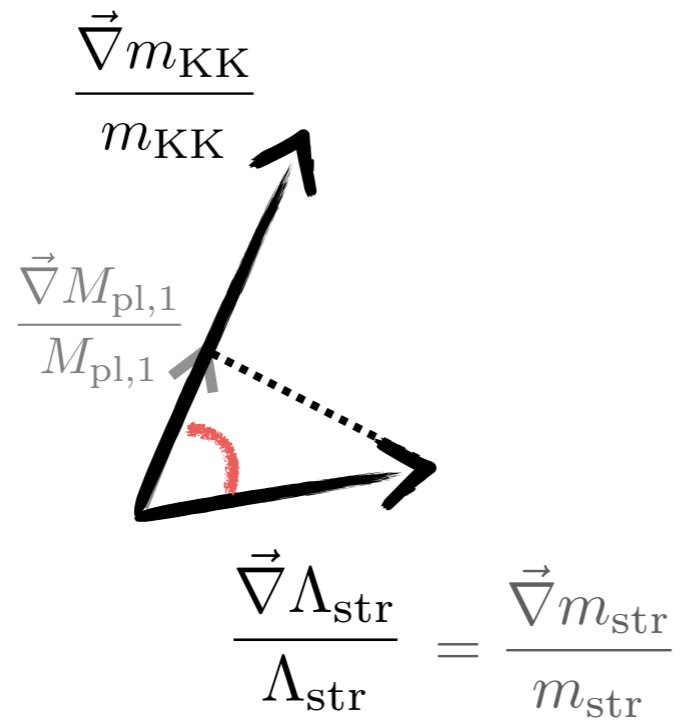
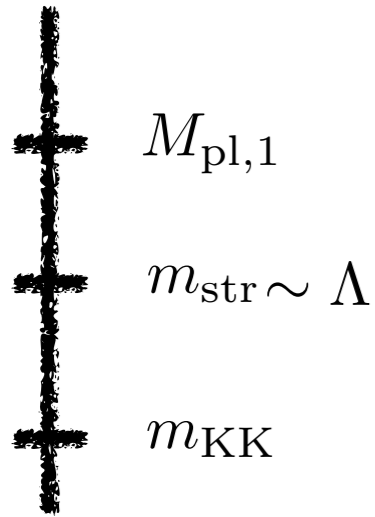
Example 2:

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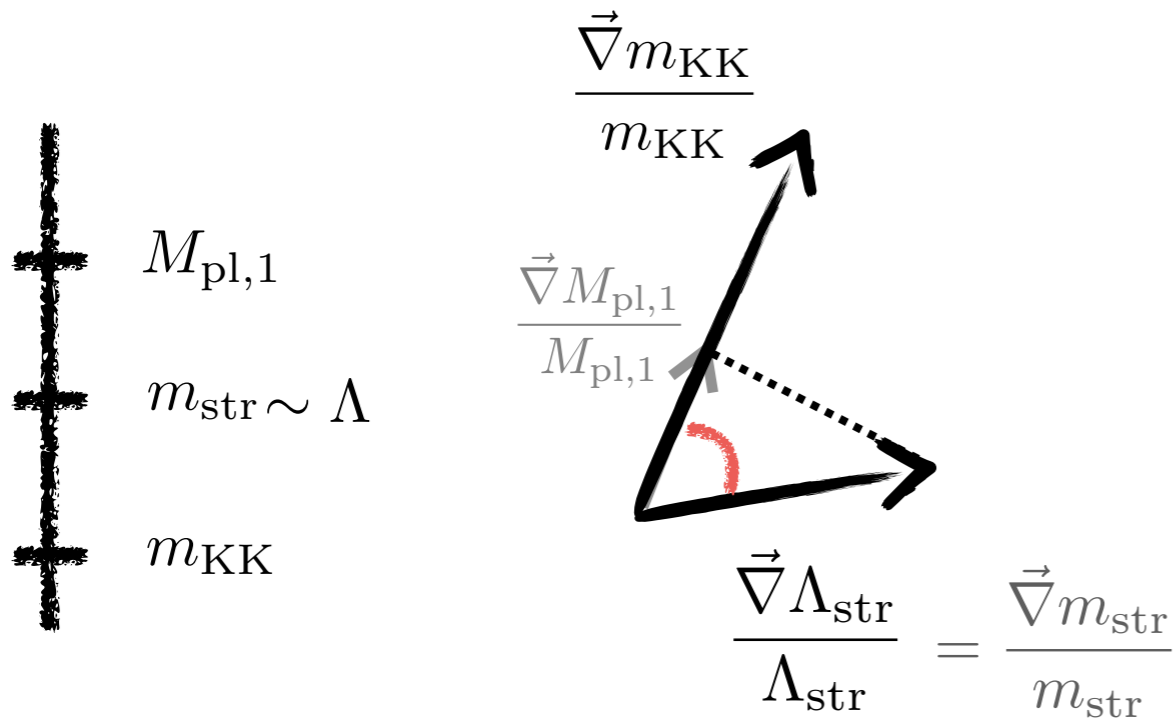


e.g. Type IIB on S^1 , $R \rightarrow \infty$, $g_s \rightarrow 0$

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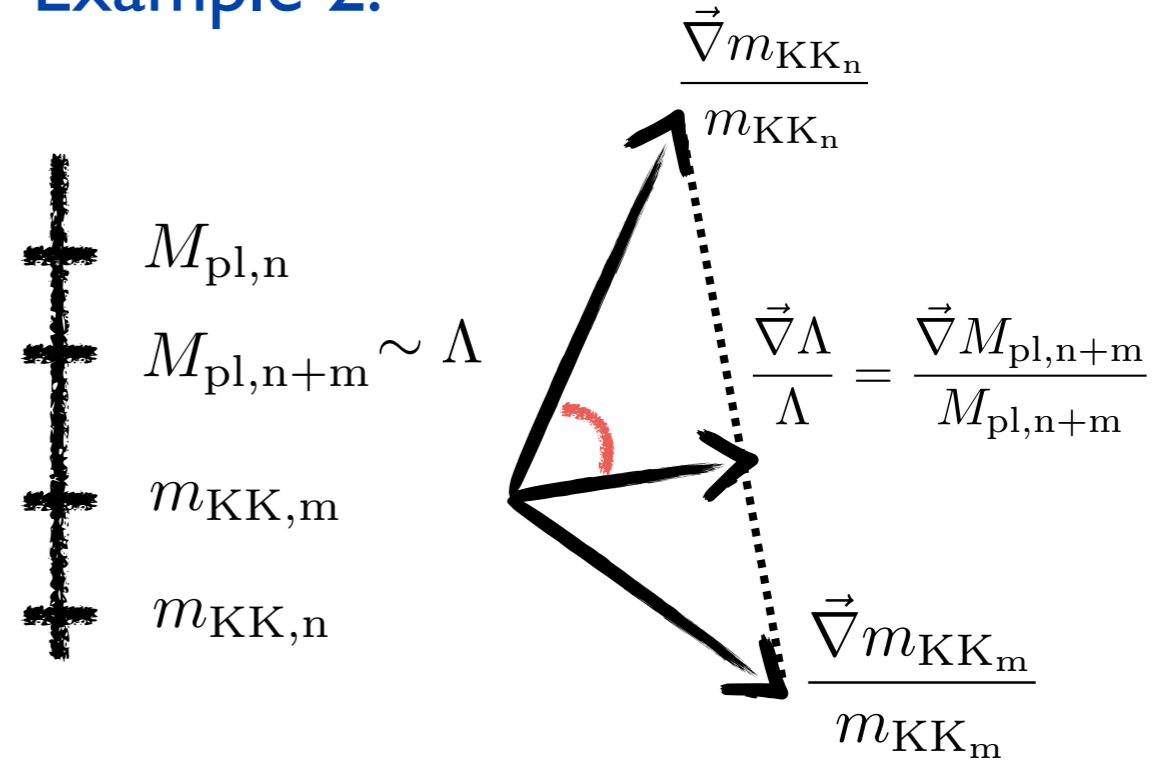
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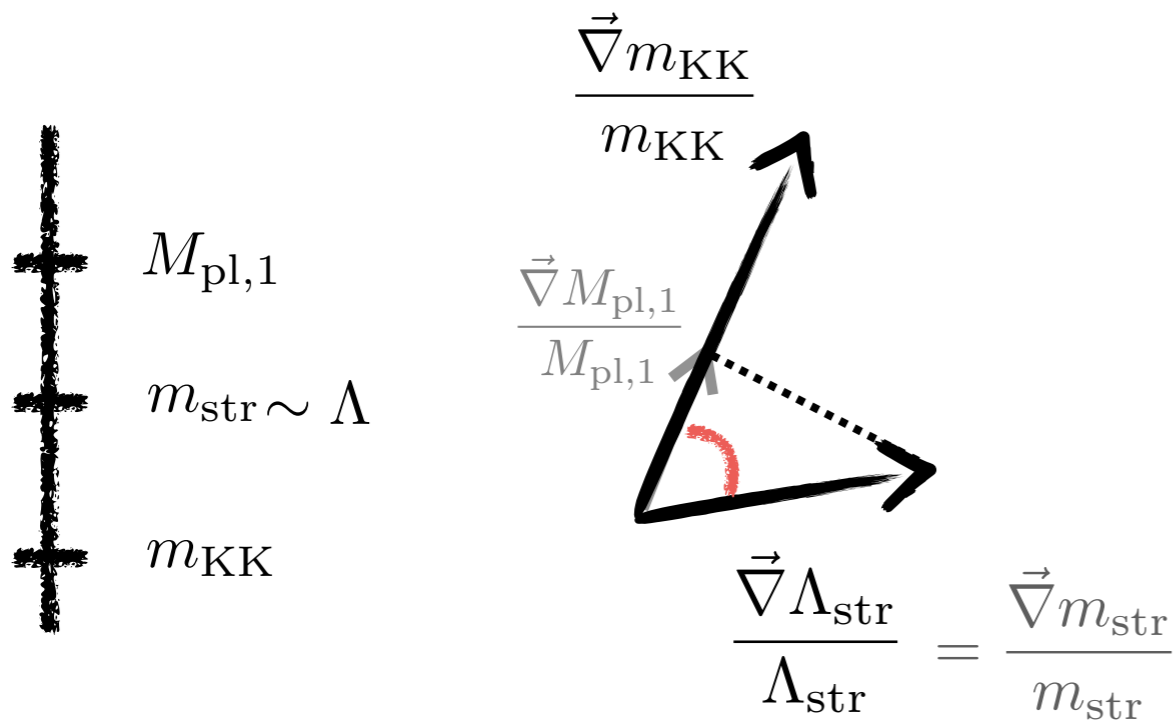
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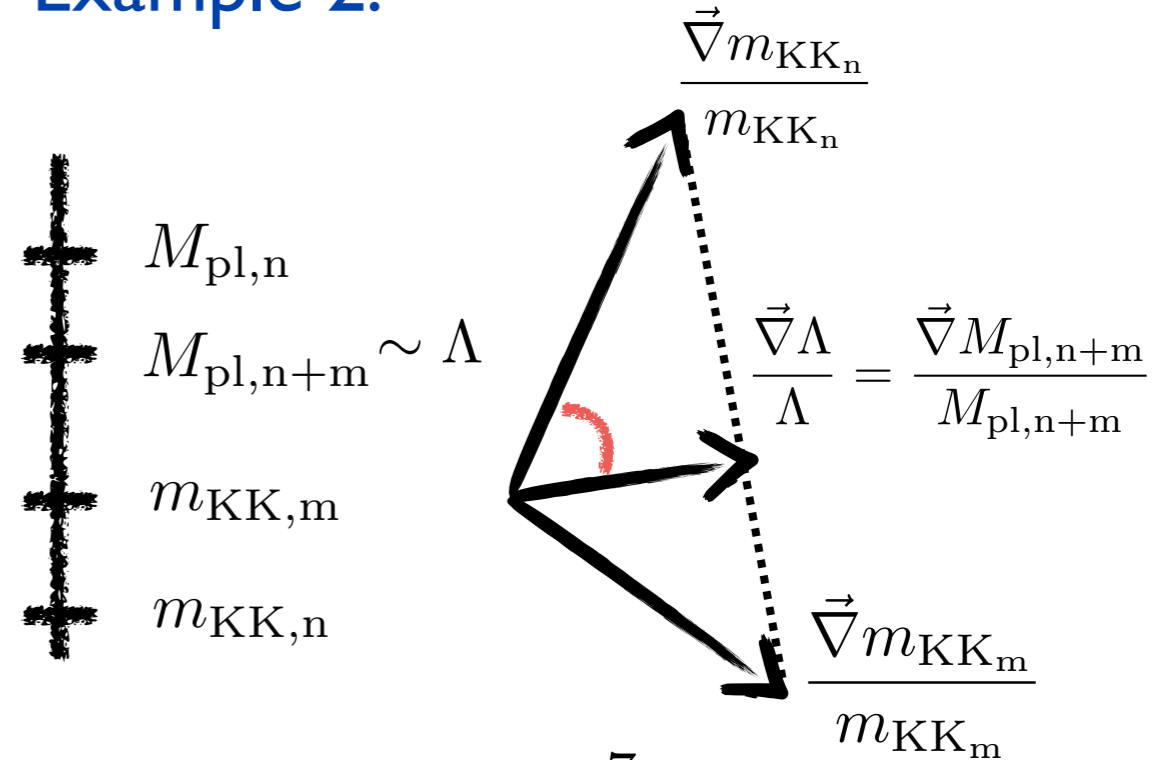
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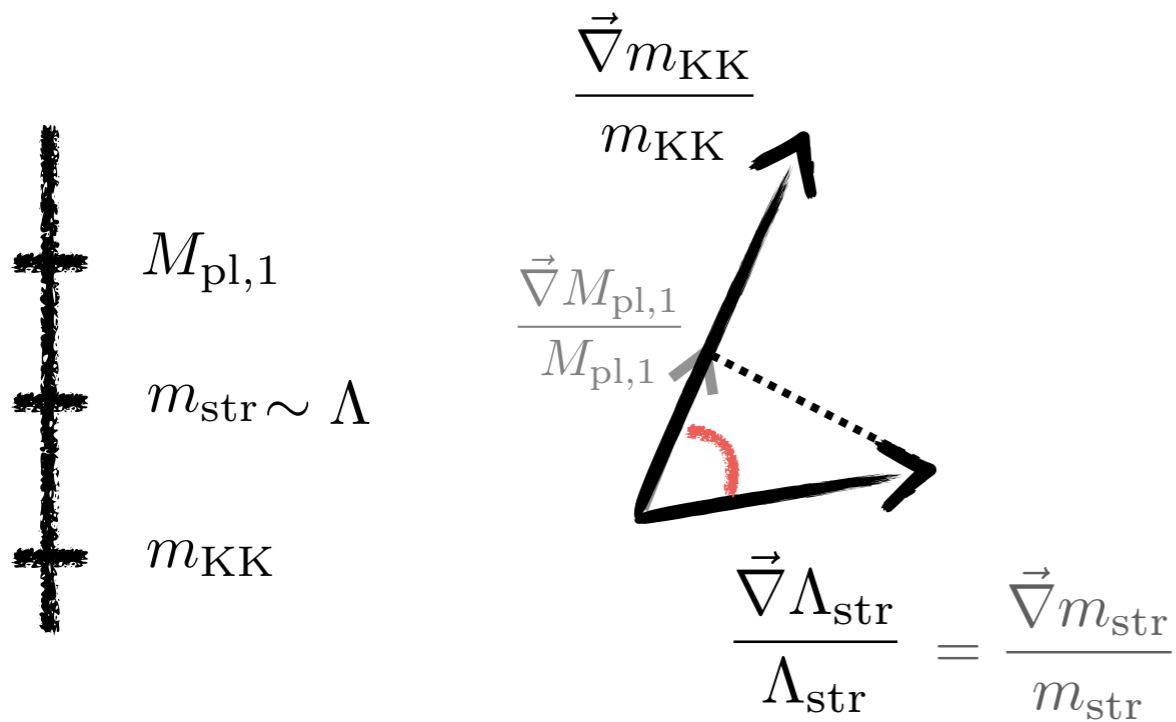
e.g. M-theory on T^7

Leading tower: m_{KK_1}

Subleading towers: m_{KK_i} , $i = 2, \dots, 5$

$$\frac{\vec{\nabla} m_{\text{KK}_1}}{m_{\text{KK}_1}} \cdot \frac{\vec{\nabla} M_{\text{pl},9}}{M_{\text{pl},9}} = \frac{1}{d-2}$$

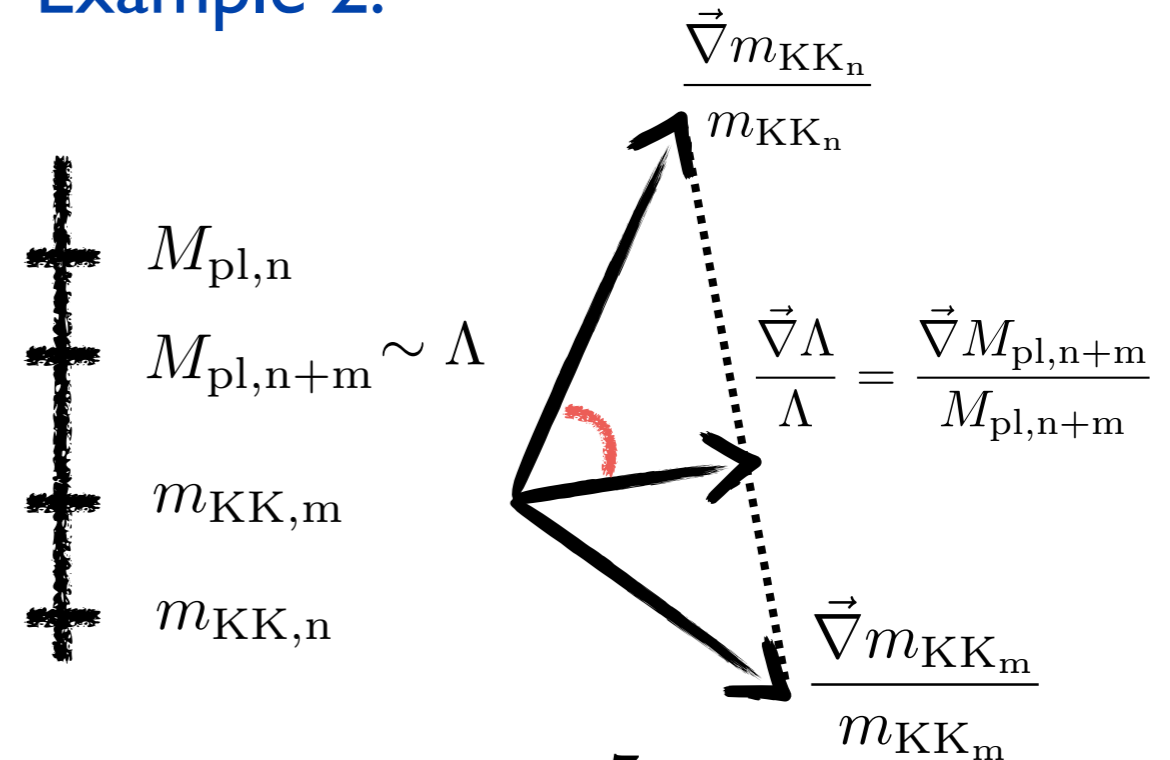
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Even if the leading tower is the same (e.g. a KK tower of one extra dimension), the quantum gravity cut-off can be different depending on the subleading towers

Subleading towers matter!

Evidence for the Pattern

We check it in string theory examples:

See Ignacio Ruiz's poster

- 32 (and 16) supercharges: M-theory on (orientifold) toroidal compactifications
- 8 supercharges: Type II and M-theory Calabi-Yau compactifications
- 4d $N=1$ theories from heterotic, F-theory, Type II...

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Open question: can we find a bottom-up argument?

$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} N}{N} = -1 \quad \text{constrains the variation (with the moduli) of the mass and the density of states}$$

Maybe black hole/entropy bounds?

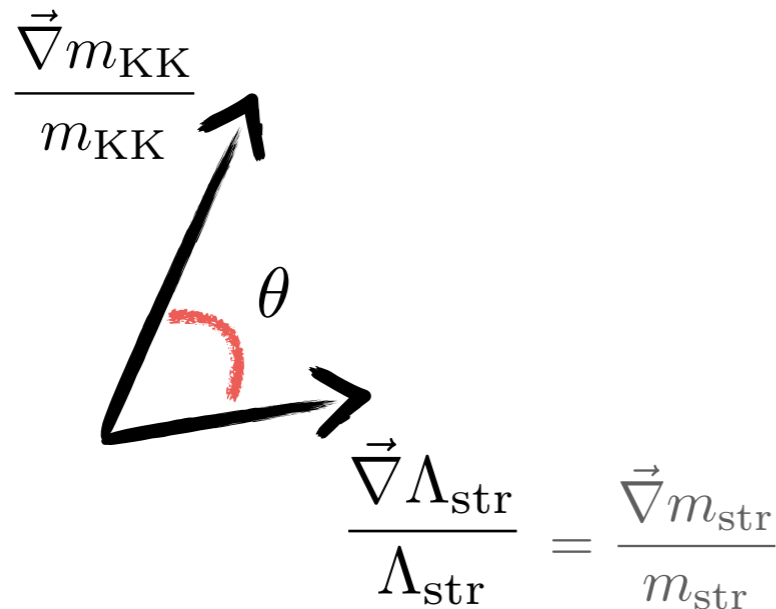
Taxonomy of infinite distance limits

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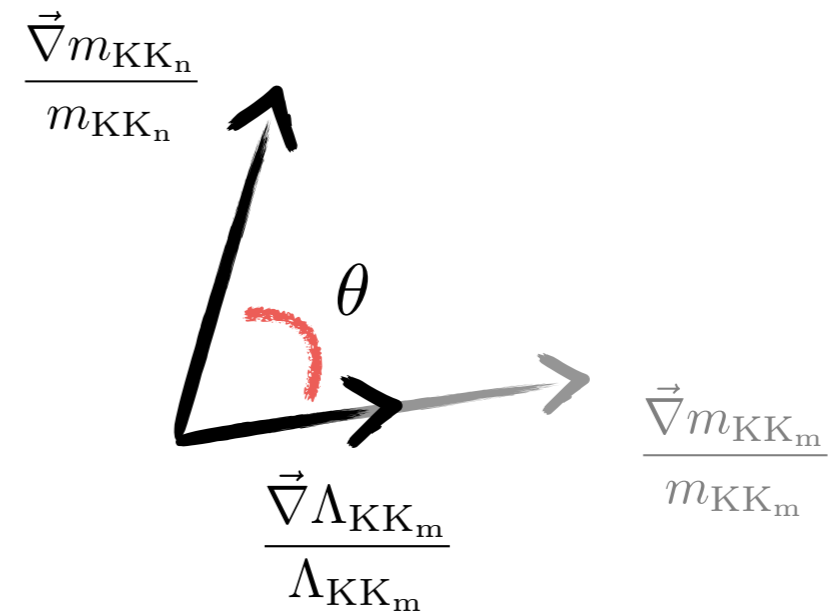
Angle between KK and string towers:



$$\cos \theta(\text{KK}_n, \text{str}) = \sqrt{\frac{n}{n + d - 2}}$$

↪ number of dimensions decompactifying

Angle between different KK towers:



$$\cos \theta(\text{KK}_n, \text{KK}_m) = \sqrt{\frac{nm}{(n + d - 2)(m + d - 2)}}$$

[Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, IV' ongoing]

(These angles can also be derived if assuming that we have only KK and string towers)

Taxonomy of infinite distance limits

Example: Consider a moduli space with a flat metric

Different asymptotic regimes are characterized by different leading towers of states, but the sum of all of the angles must be:

$$\sum \theta_i = 2\pi$$


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For a 2-dim moduli space in 9d:

$$\cos \theta = \left\{ \frac{1}{\sqrt{8}}, \frac{1}{8} \right\}$$


The diagram shows two arrows originating from the terms in the set. One arrow points from $\frac{1}{\sqrt{8}}$ to $\cos \theta_1$, and the other points from $\frac{1}{8}$ to $\cos \theta_2$.

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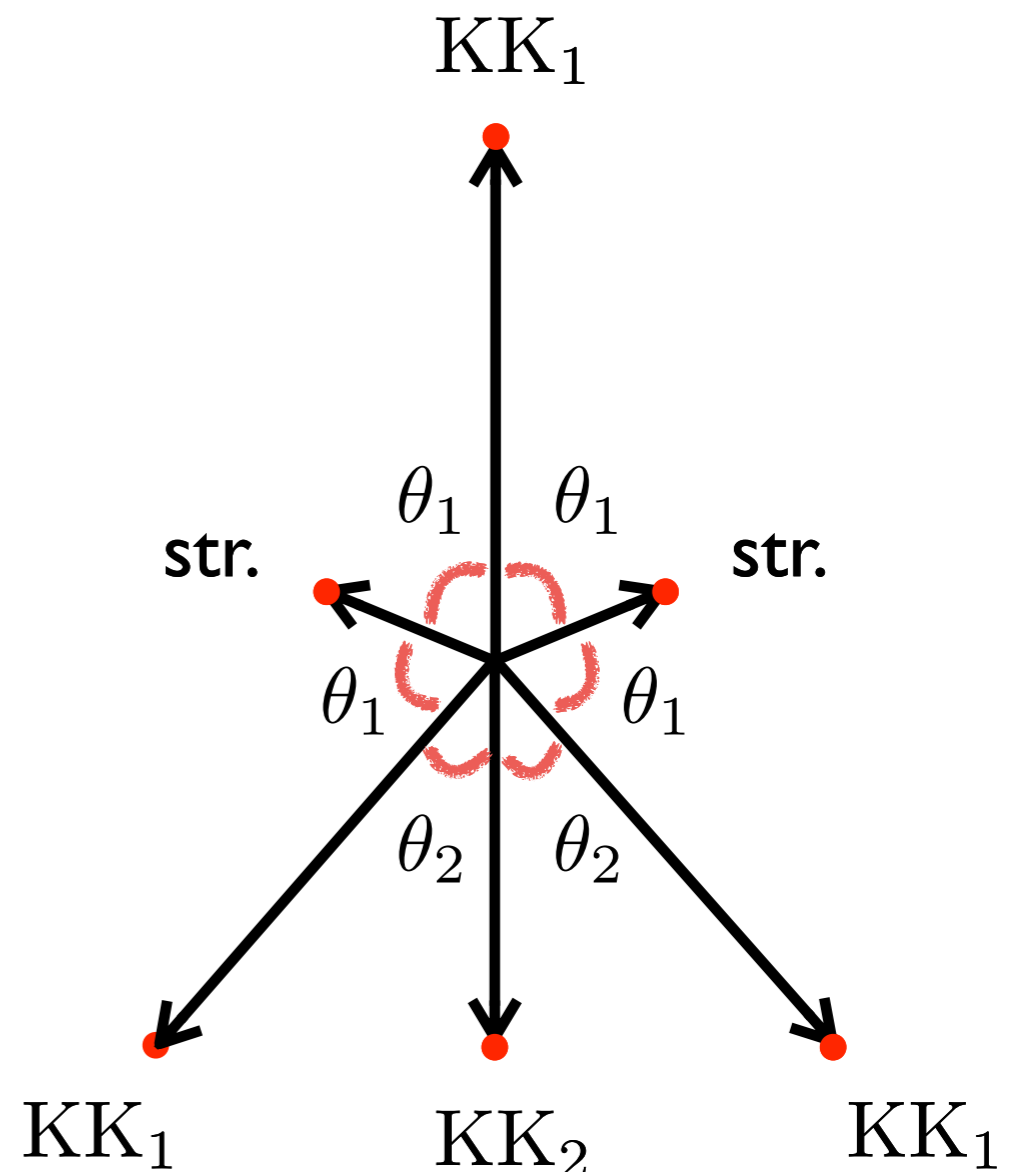
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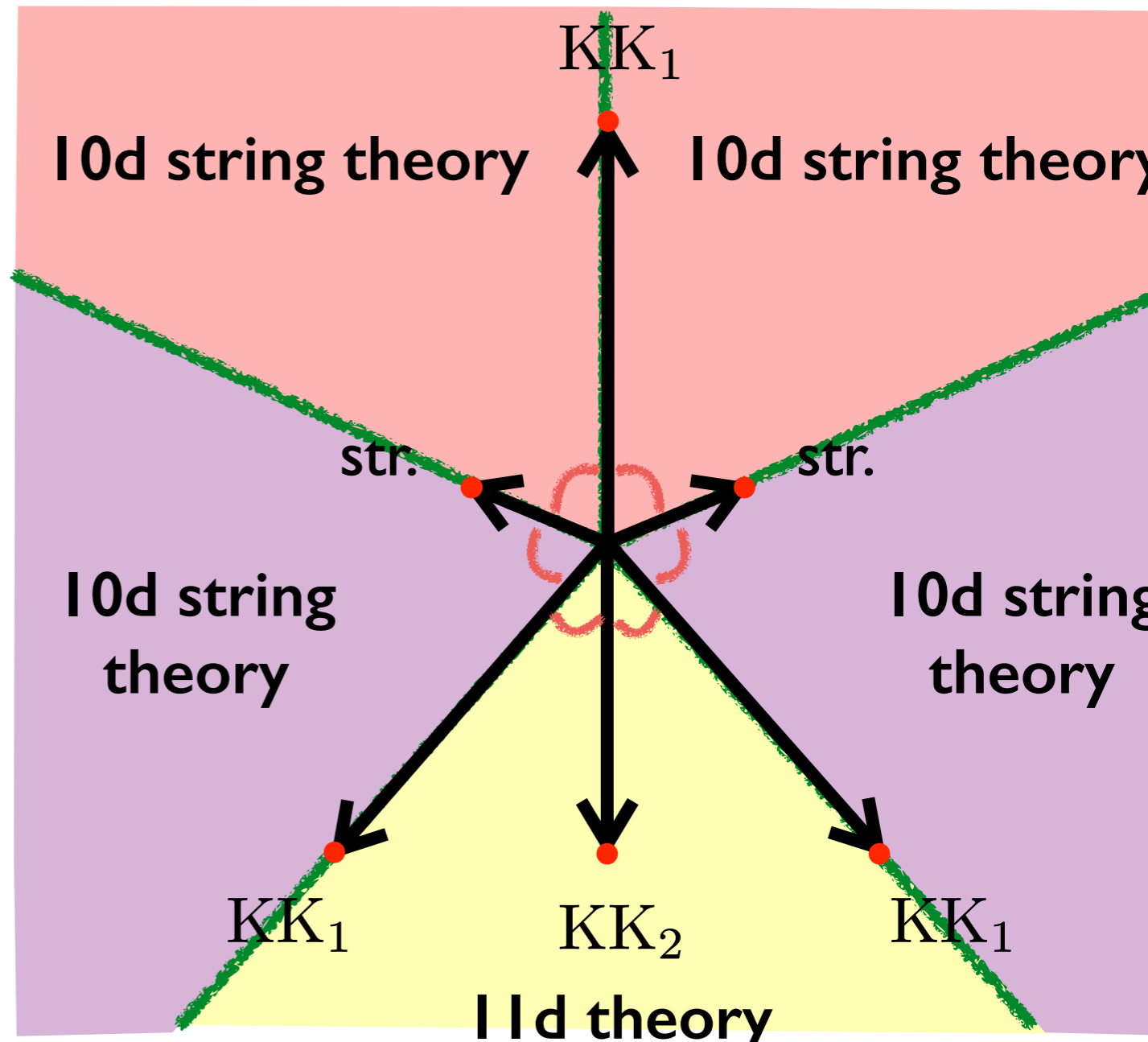
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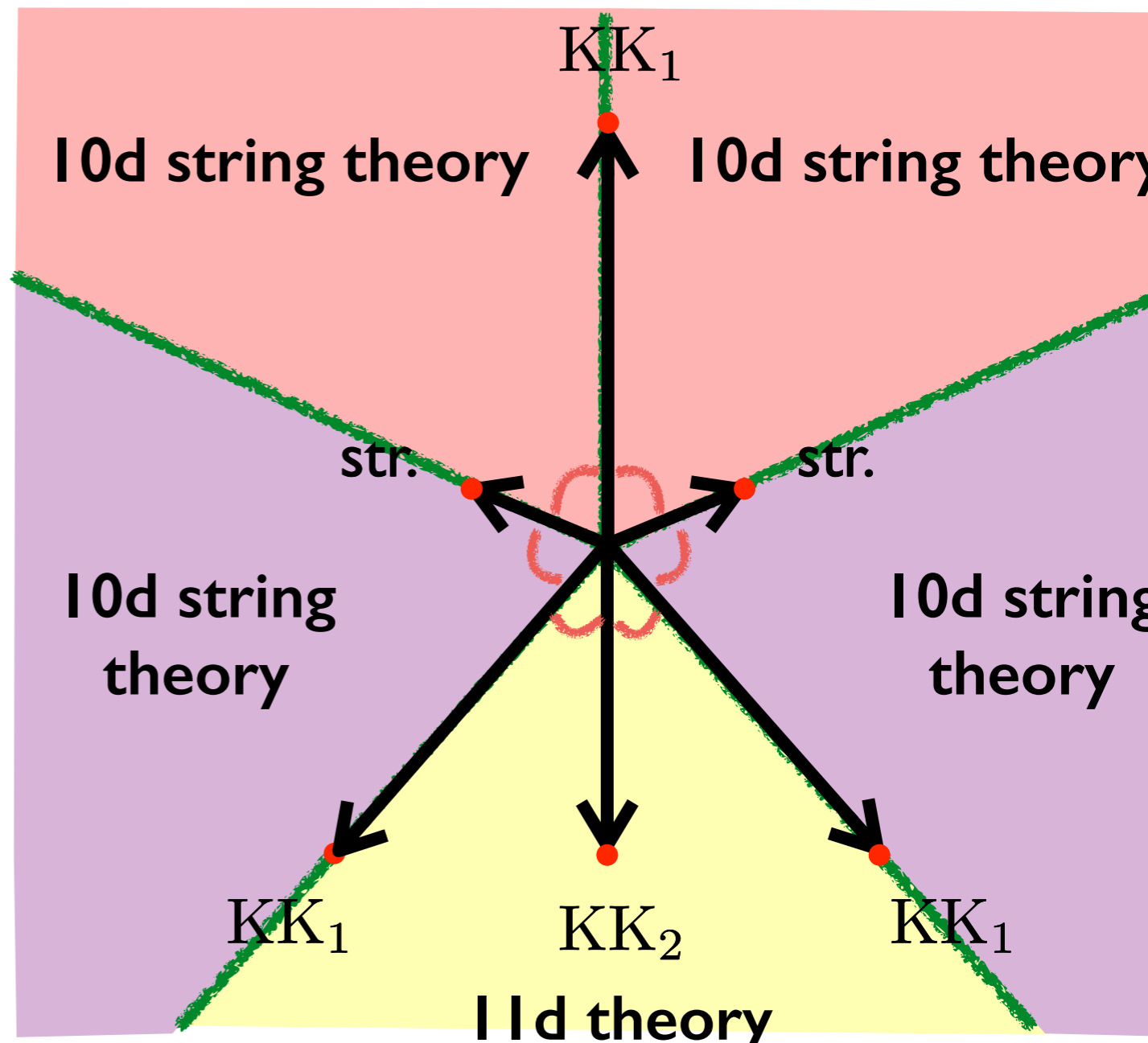
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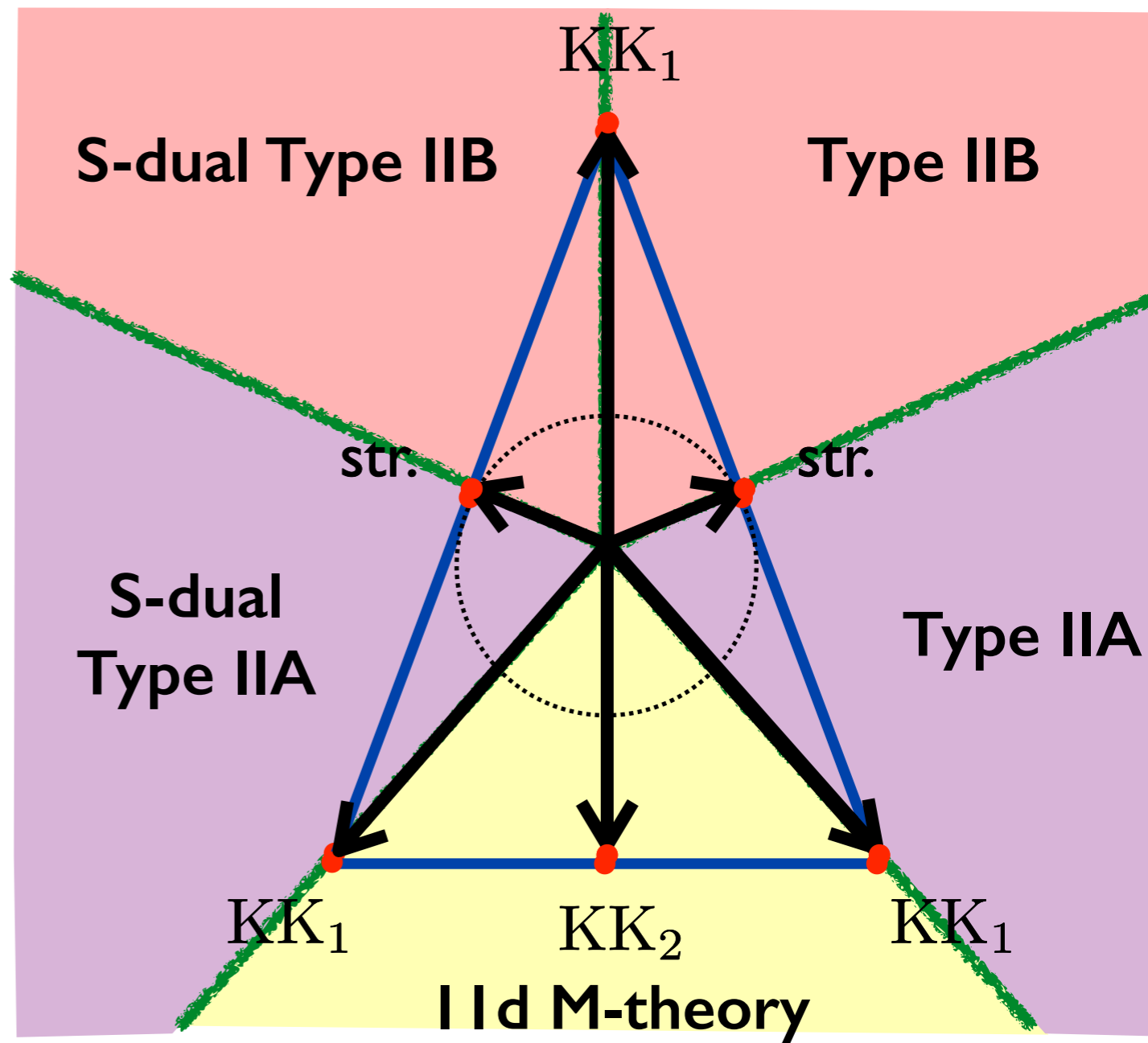
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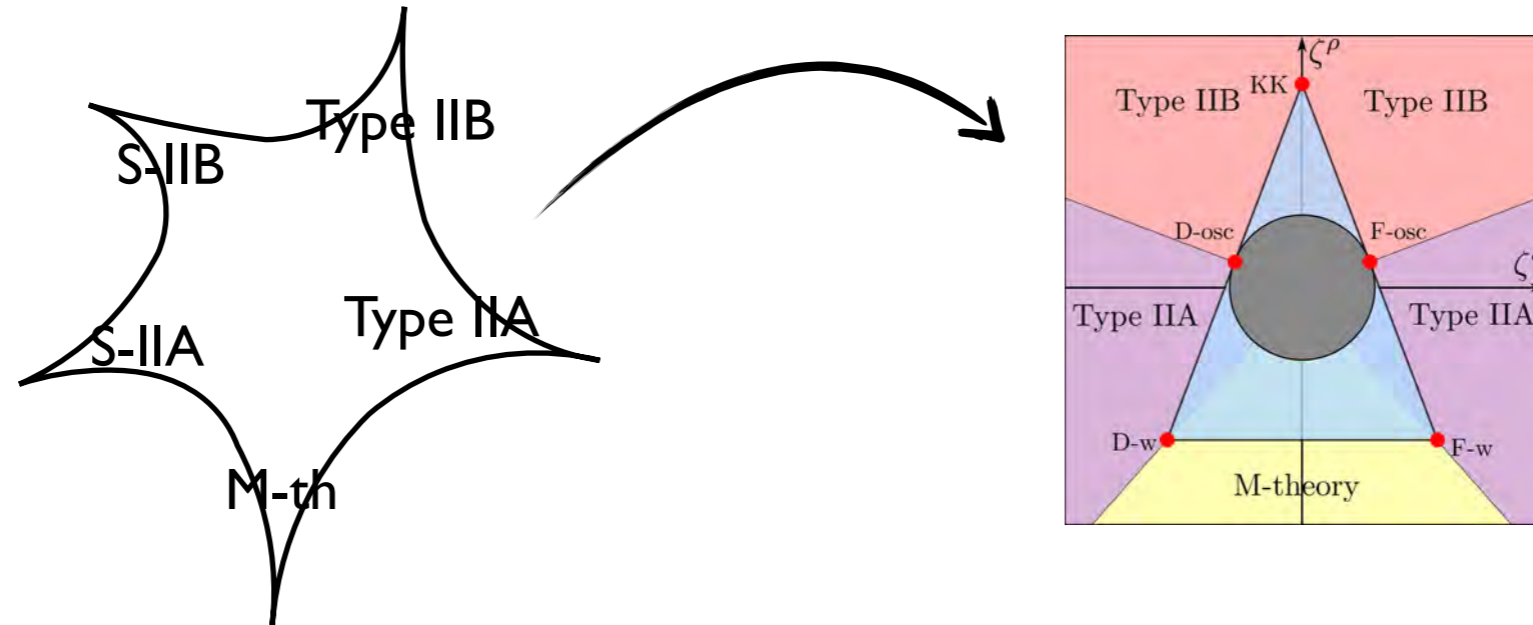
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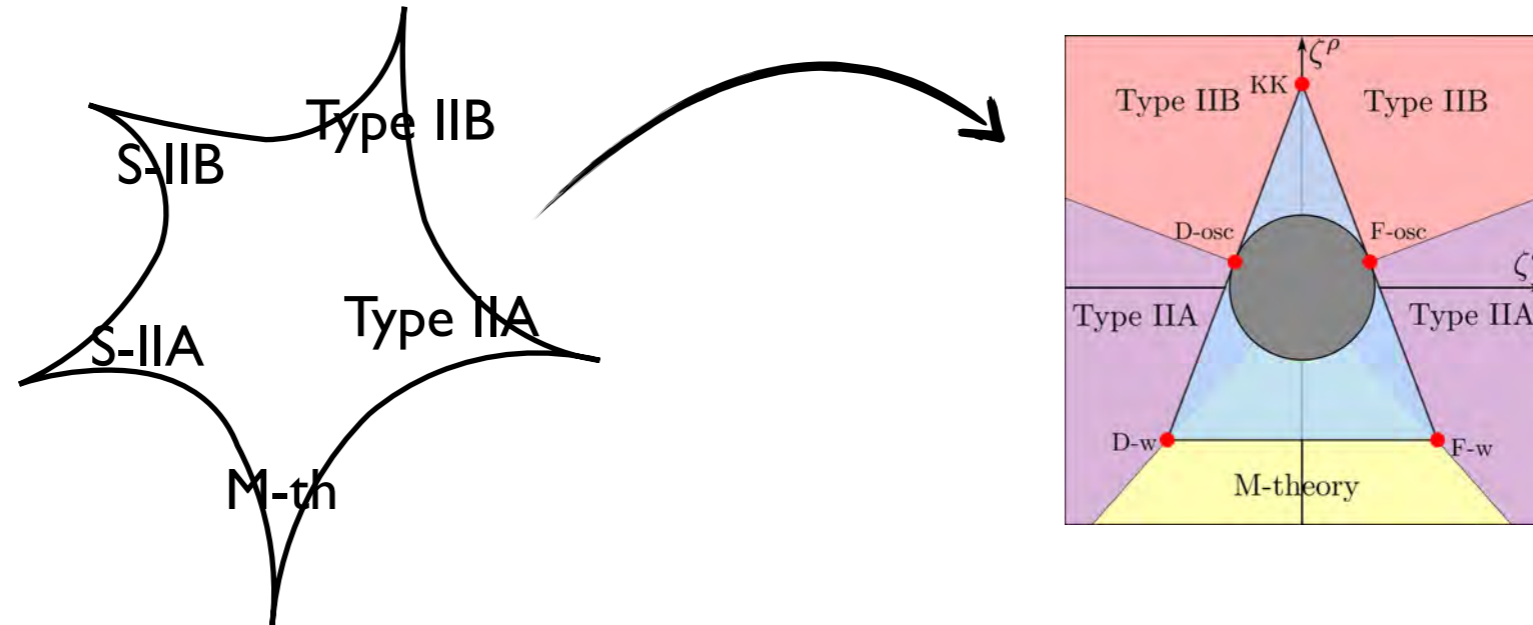
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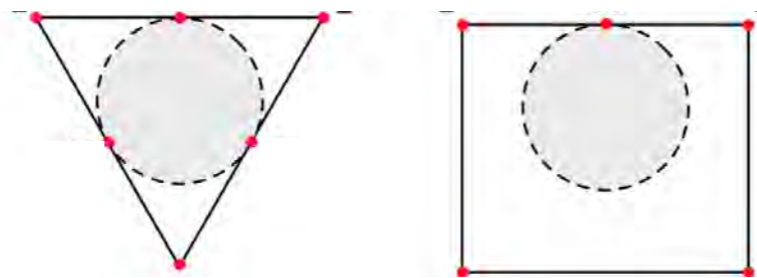


Can we classify all possibilities from bottom-up?

See Muldrow
Etheredge's poster

In $d=8$:

⋮



[Etheredge, Heidenreich, McNamara,
Rudelius, Ruiz, IV' ongoing]

[Etheredge, Heidenreich, McNamara, Rudelius, Ruiz, IV'23]

For lower SUSY setups: necessary to understand rules of jumping of states
(the convex hull can change as we move in moduli space)

Outline

- 1) Review of Distance Conjecture and drop-off of QG cut-off

- 2) Universal pattern underlying all known string theory examples
 - ➔ To sharpen the conjecture and provide a quantitative bound on the mass of the tower

- 3) Beyond moduli spaces: Implications for accelerated expansion and AdS scale separation
 - ➔ Non-perturbative test of DGKT AdS scale separated proposal

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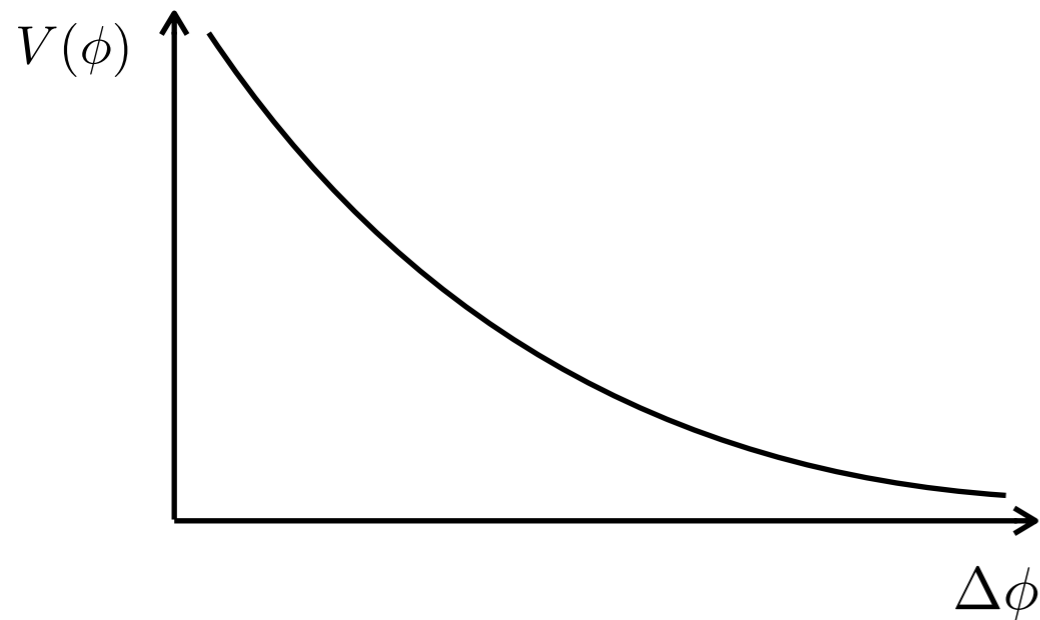
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Beyond Moduli Spaces

What happens in the presence of a potential?

If the infinite distance limit is not obstructed, there are two typical scenarios:

Positive runaway

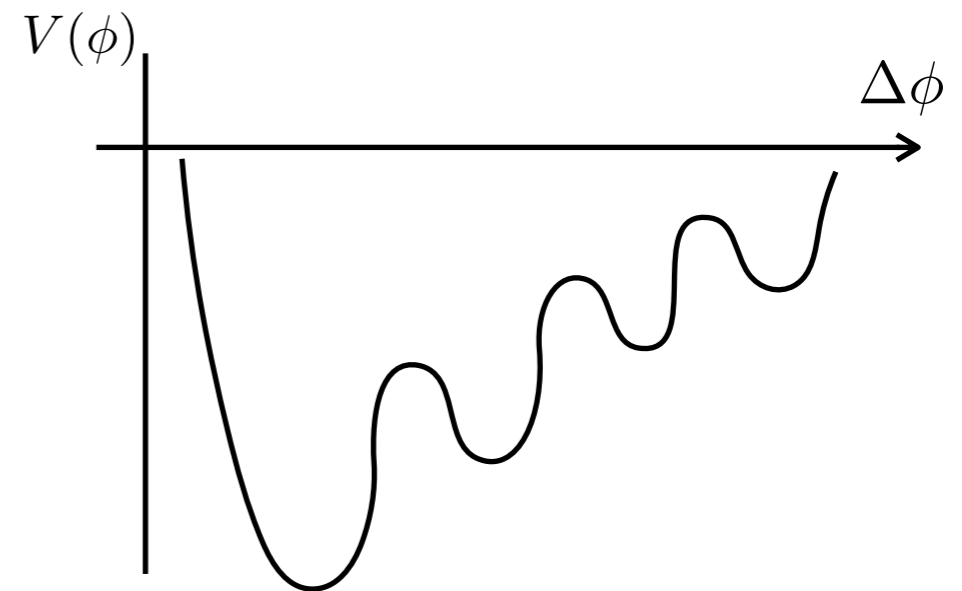


$$V_0 \sim m_{\text{tower}}^\gamma$$

[Ooguri et al'18]

Accelerated expansion?

Family of AdS vacua



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[Luest,Palti,Vafa'19]

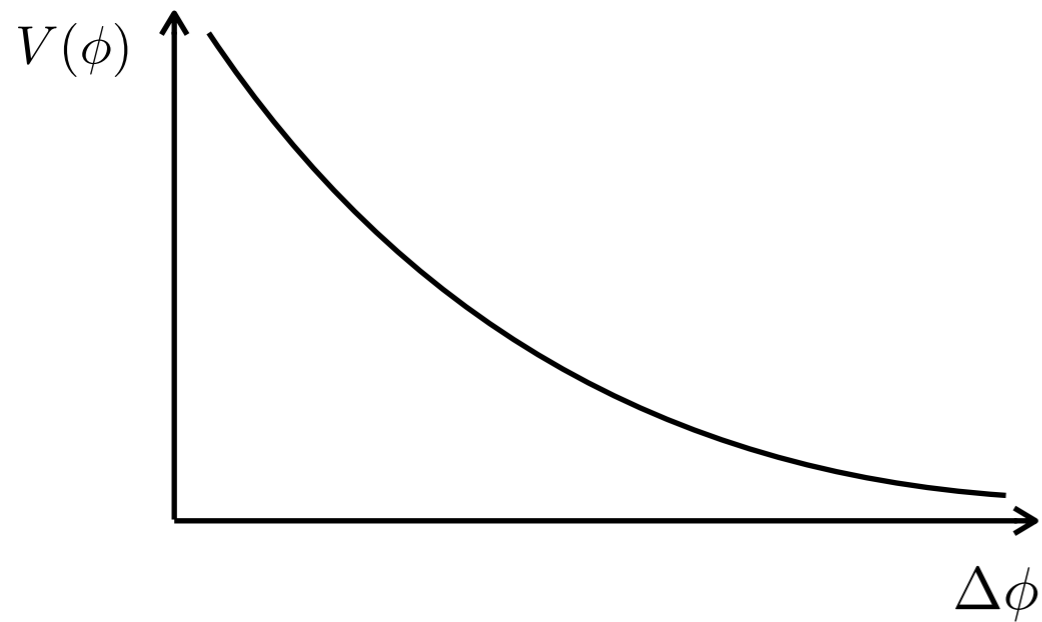
AdS scale separation?

(also important for Dark Dimension scenario, [Cumrun's talk](#))

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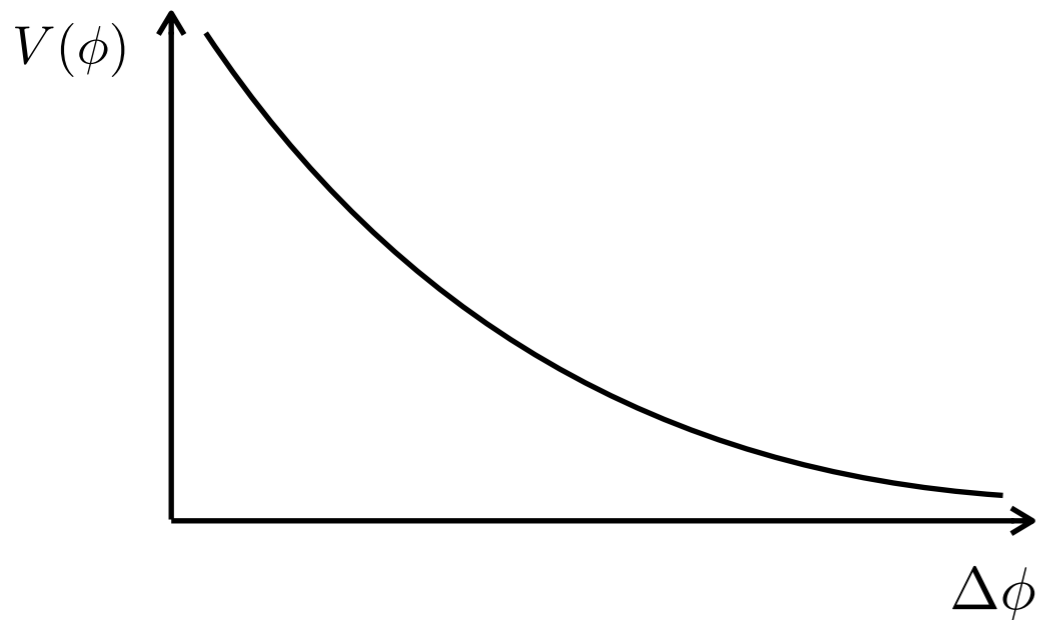
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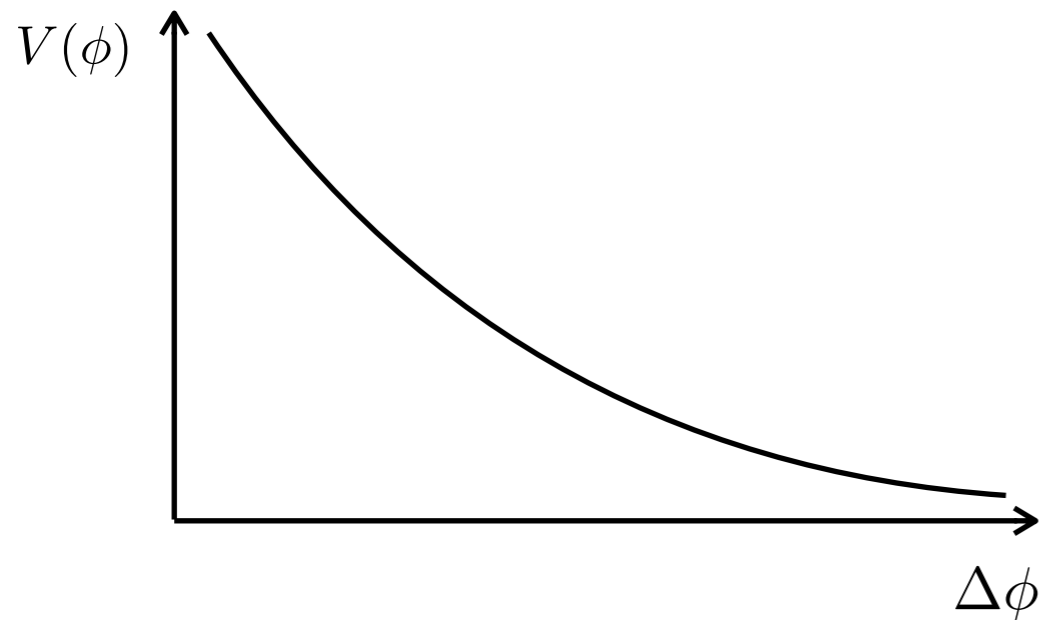
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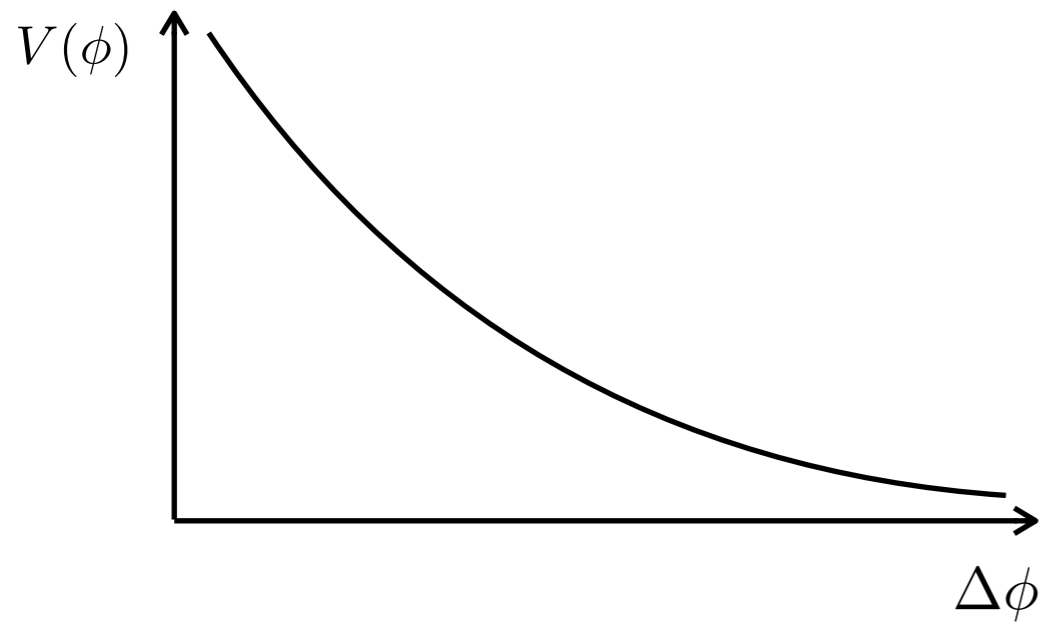
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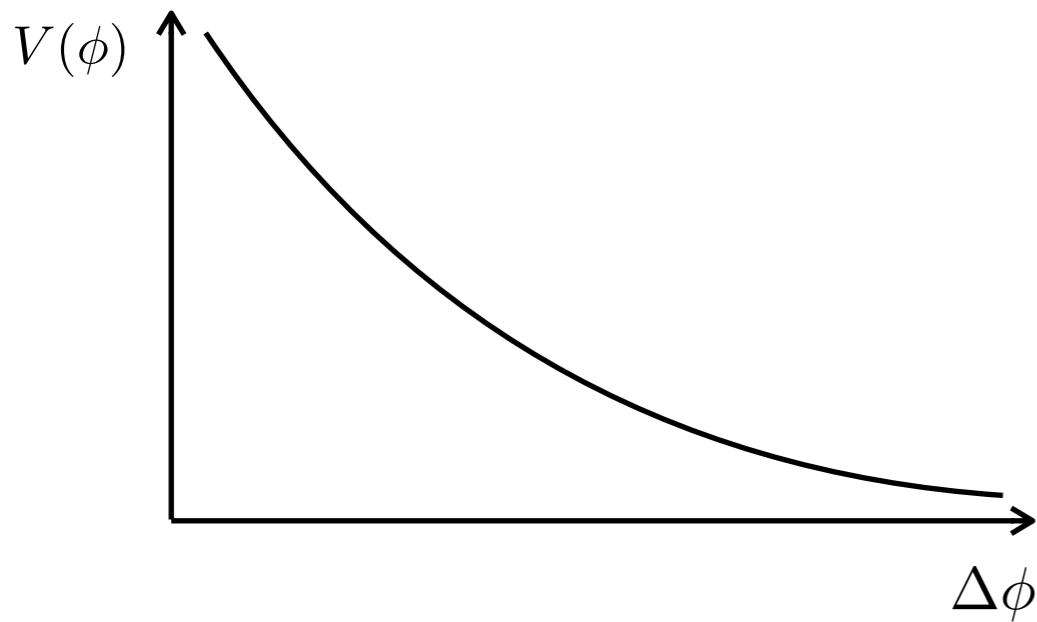
If the tower contains higher spin fields,

Higuchi bound implies $\gamma \geq 2$ [Montero,Vafa,IV'22]
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which combined with $\left| \frac{\vec{\nabla} m}{m} \right| \geq \frac{1}{\sqrt{d-2}}$

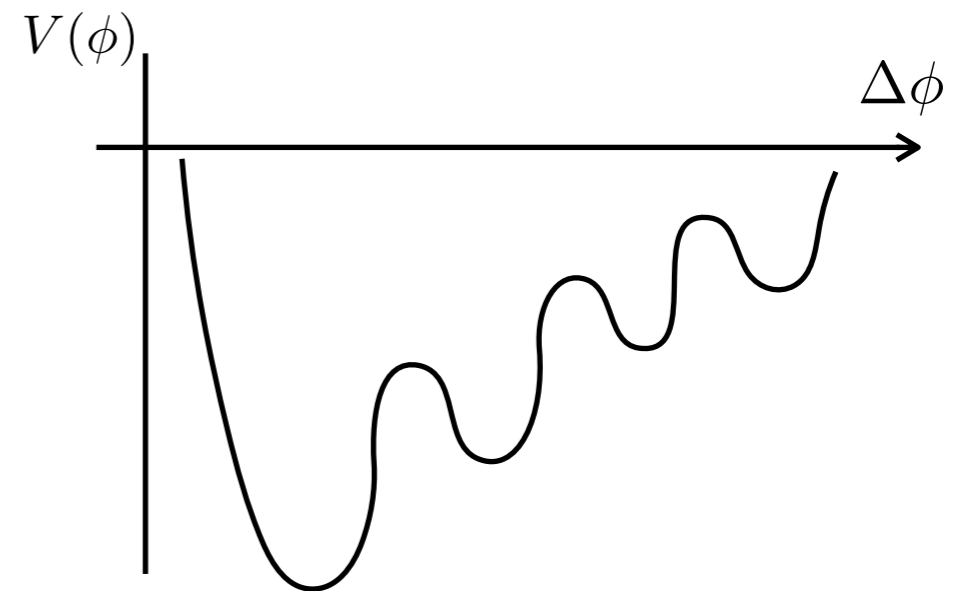
implies $\left| \frac{\vec{\nabla} V_0}{V_0} \right| \geq \frac{2}{\sqrt{d-2}}$ (no asymptotic accelerated expansion)

[Bedroya,Vafa'19] [Rudelius'22]

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Family of AdS vacua



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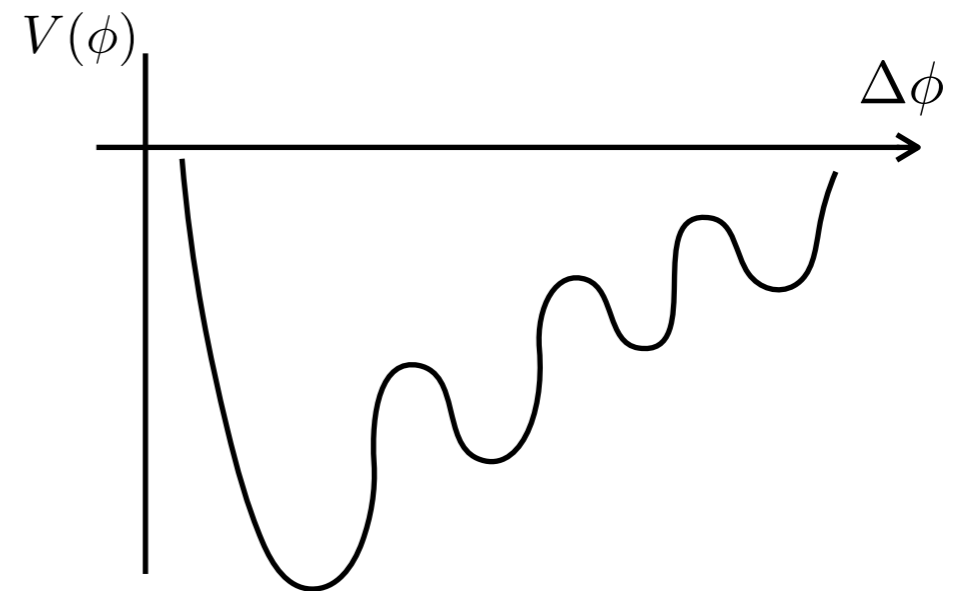
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(no scale separation between AdS length and internal dimensions)

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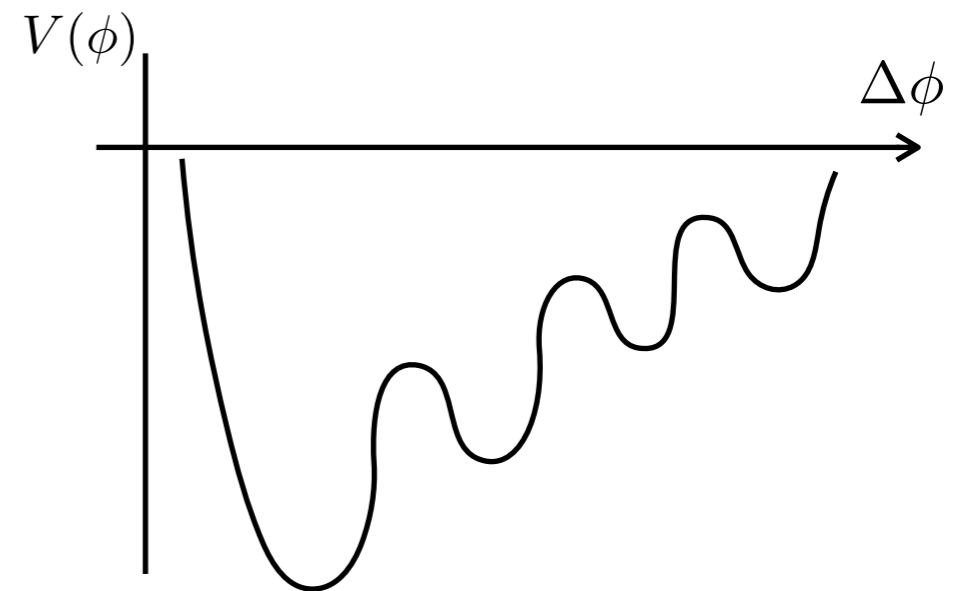
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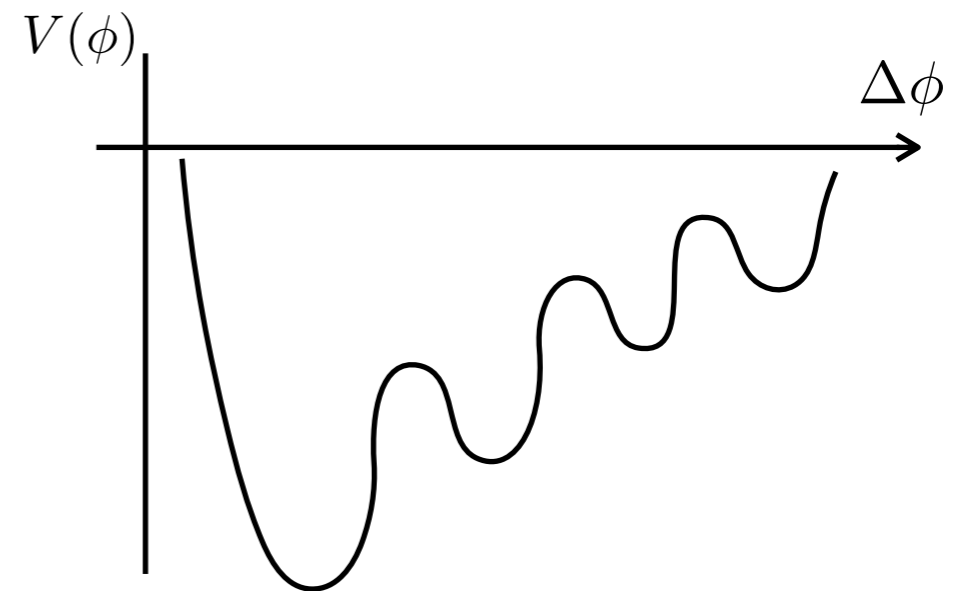
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New results regarding consistency of the vacuum
[Montero, Valenzuela 'ongoing]

Family of AdS vacua



$$V_0 \sim m_{\text{tower}}^\gamma$$

AdS scale separation if

$$\gamma \geq 2$$

DGKT vacuum

[De Wolfe, Giddings, Kachru, Taylor '05]

4d $N=1$ AdS vacuum arising from compactifying massive Type IIA on a CY3 with O6-planes and fluxes for

$$F_0, F_4, H_3 \quad \longrightarrow \quad \text{AdS}_4 \times \text{CY}_3$$

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By solving the 4d eoms, one finds a family of AdS vacua with

$$\begin{aligned} V_0 &\sim N^{9/2} \\ m_{\text{KK}}^{-2} &\sim L_{\text{KK}}^2 \sim N^{7/2} \end{aligned} \quad \rightarrow \quad \left(\frac{\ell_{\text{AdS}}}{L_{\text{KK}}} \right)^2 \sim N$$

So this solution is **scale-separated** in the large N limit.

DGKT vacuum

The consistency of the solution is not clear because we only solved **4d equations of motion** (zero mode of 10d eoms on CY3)

Lot of recent progress, everything seems fine so far, but no conclusive answer.

[Andriot, Apers, Casas, Castellano, Collins, Cribiori, Dall'Agata, De Luca, Emelin, Farakos, Graña, Herraez, Hoter, Ibañez, Junghans, Lust (x2), Marchesano, Marconnet, Montella, Moritsu, Ning, Palti, Plauschinn, Prieto, Quirant, Revello, Shiu, Shukla, Tomasiello, Tonioni, Toulukas, Tringas, Tsimpis, Vafa, Van Hemelryck, Van Riet, Walcher, Wiesner, Wrasse, Xu, Yau, Zatti, ...]

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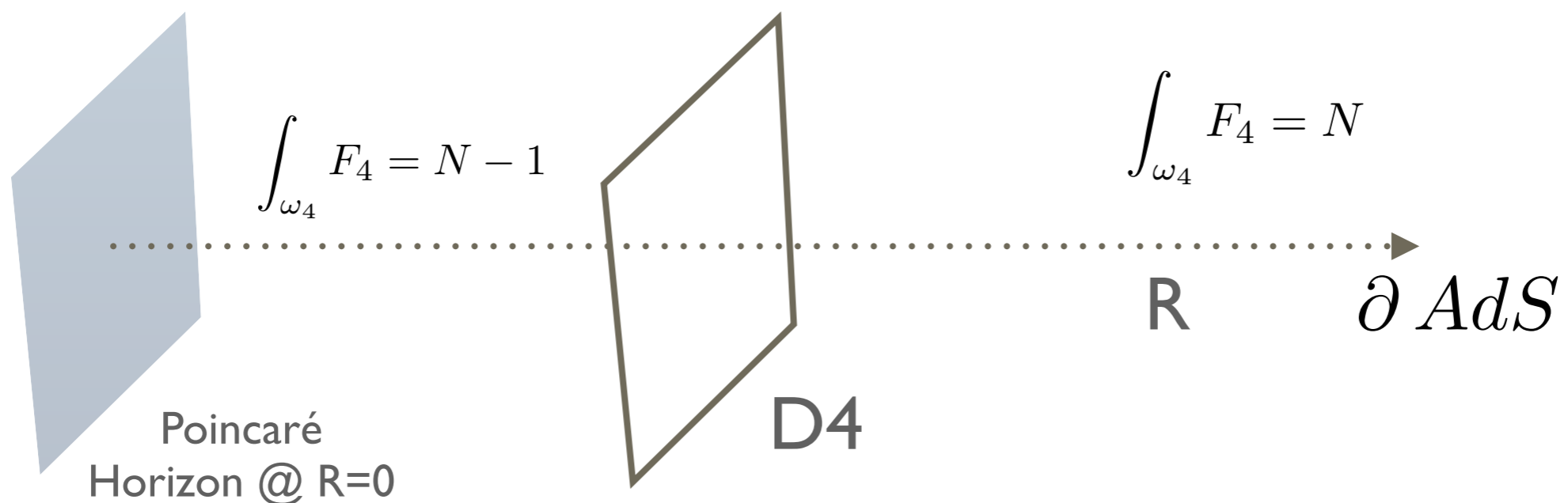
We will assume everything is OK, and study the fate of branes on DGKT vacuum, **to perform a non-perturbative consistency check**

[Montero, Valenzuela 'ongoing]

Test of DGKT vacuum

Consider a D4-brane wrapping a holomorphic 2-cycle dual to the large N flux

$$\int_{\omega_4} F_4 = N$$



In principle, it is **BPS**, so the position of the brane is a modulus

(see also [Aharony,Antebi,Berkooz '08])

Test of DGKT vacuum

However, at low energies the worldvolume theory is $3d \mathcal{N} = 1$

This is so little SUSY that there are no protected quantities

The superpotential can receive corrections (unless there is a parity symmetry)

[Gaiotto-Komargodski-Wu '18]

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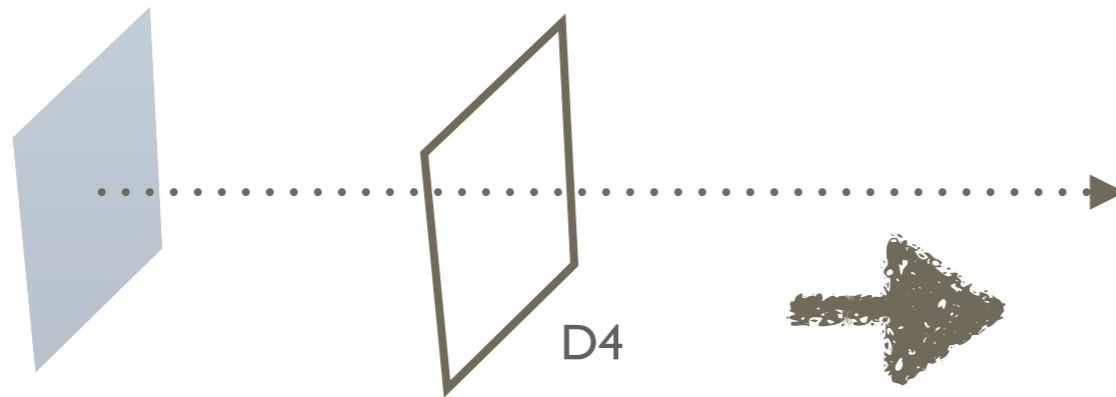
(SUSY is broken spontaneously on the brane)

➔ the position of D4 is not a modulus
(it feels an attractive force)

Test of DGKT vacuum

The result is in tension with the Weak Gravity Conjecture (WGC)

According to the WGC, a non-BPS brane should be repulsive



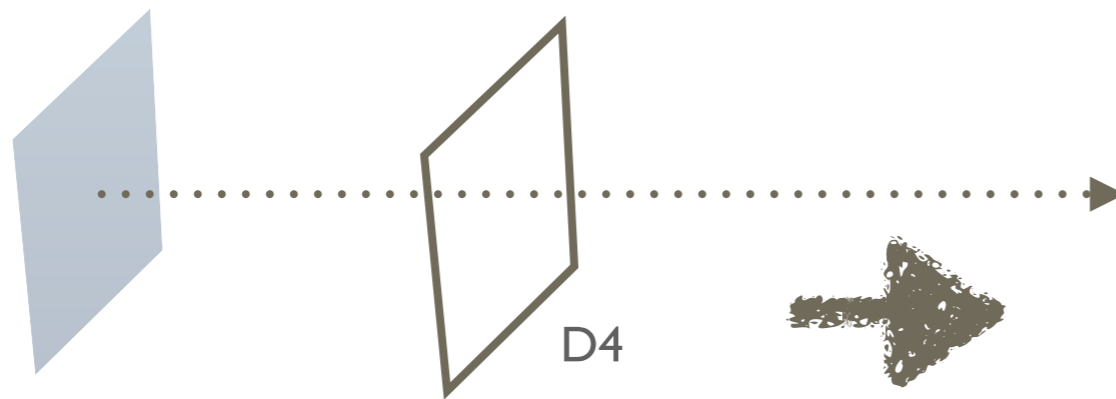
which would imply an instability of the DGKT vacuum

➡ No SUSY AdS scale-separated vacuum

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Ongoing work!

We are exploring the implications of this clash between DGKT vacuum and the WGC for membranes

[Montero, Valenzuela 'ongoing]

Stay tuned

Conclusions

❖ We are entering an era of precision in the Swampland program

Fix all “order one” factors

$$m \sim m_0 \exp(-\alpha \Delta\phi) \quad \Lambda \sim M_p \exp(-\lambda \Delta\phi)$$
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❖ We found a universal pattern between the leading tower of states and the quantum gravity cut-off underlying all string theory examples

- It provides a quantitative bound for the exponential rate of the tower

$$\frac{\vec{\nabla} m}{m} \cdot \frac{\vec{\nabla} \Lambda}{\Lambda} = \frac{1}{d-2} \quad \longrightarrow \quad \alpha \geq \frac{1}{\sqrt{d-2}}$$

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- Implications for asymptotic accelerated expansion and AdS scale separation

- ❖ We perform a non-perturbative check of DGKT and the result is in tension with the Weak Gravity Conjecture

Thank you!

back-up slides

Scalar WGC

Can we reformulate the Distance Conjecture as a sharp local condition that the spectra of the theory must satisfy at a given point of the moduli space?

(in analogy to the Weak Gravity Conjecture $\frac{Q}{m} \geq \gamma_{\text{BH}}$) *

* The Distance Conjecture (a drop-off of the cut-off) resembles the “magnetic version” of the WGC, is there an analogous “electric version” for the Distance Conjecture?

[Grimm,Palti,IV'18]

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Notice! Moduli (massless scalars) induce Yukawa scalar forces:

$$\mathcal{L} \supset M^2(\phi)\chi^2 \simeq 2M\partial_\phi M \phi\chi^2 + \dots$$

modulus \rightarrow *scalar charge*

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scalar charge-to-mass ratios: $\vec{\zeta} \equiv -\frac{\vec{\nabla} m}{m}$

If applied to towers of states decaying exponentially: $\alpha = \vec{\zeta} \cdot \hat{\tau}$

exponential rate of the tower

unit tangent geodesic vector

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Convex Hull Distance Conjecture

The Distance Conjecture can be formulated as a convex hull SWGC applied to towers: [Calderon-Infante,Uranga,IV'20]

Distance conjecture with exponential rate $\alpha \geq \alpha_0$



convex hull of $\frac{\vec{\nabla} m}{m}$ of light towers remains outside ball of radius α_0

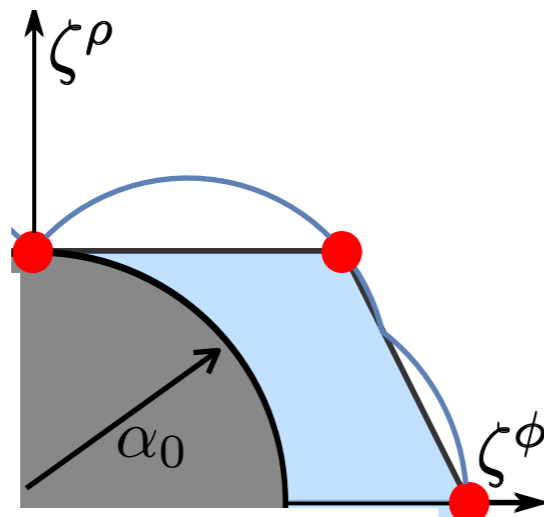
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It becomes sharp and powerful if combined with

$$\alpha \geq \frac{1}{\sqrt{d-2}} \equiv \alpha_0 \quad [\text{Etheredge et al'22}]$$

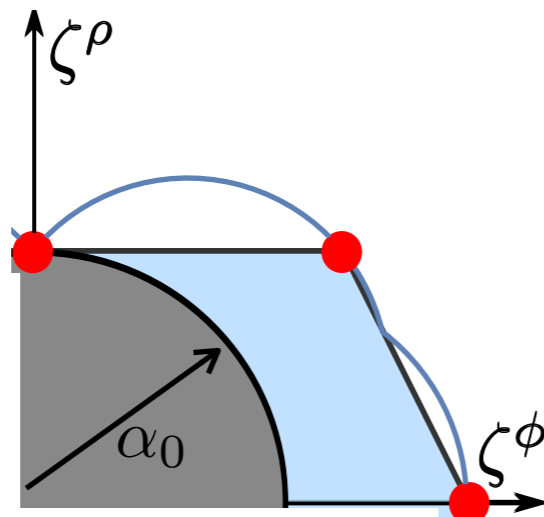
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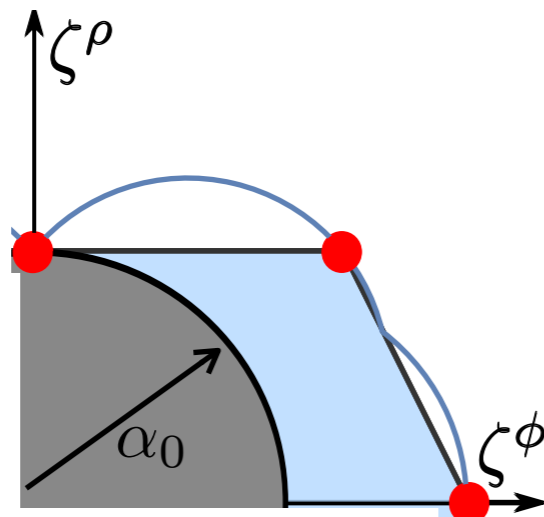
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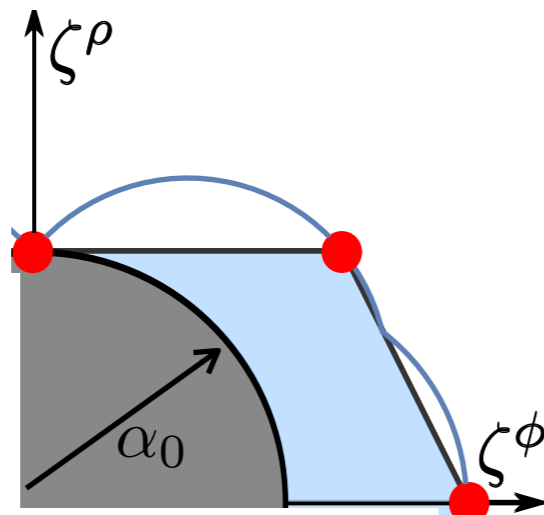
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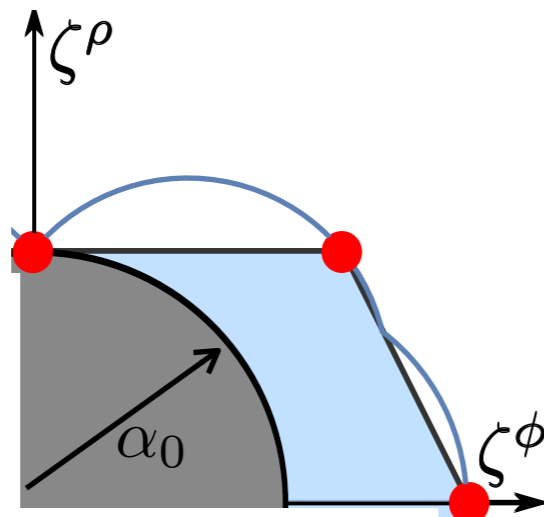
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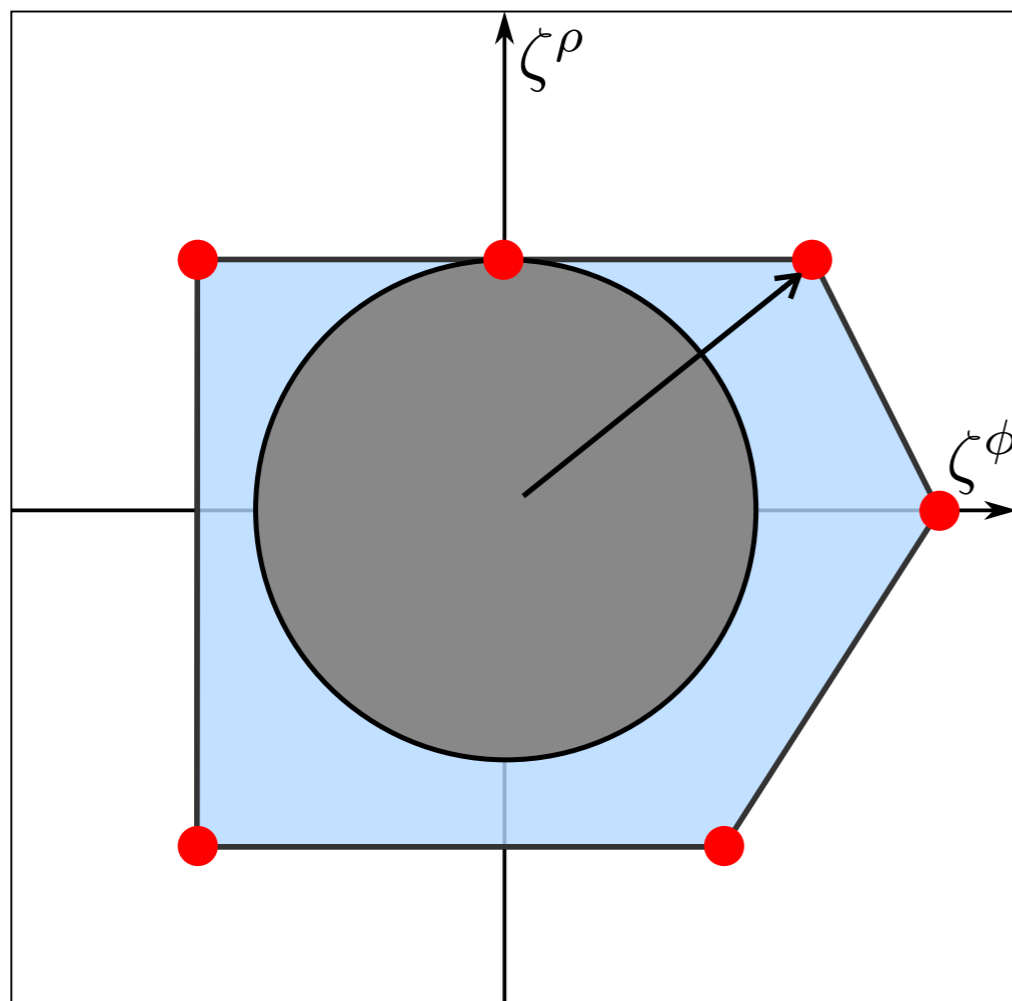
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...does it?

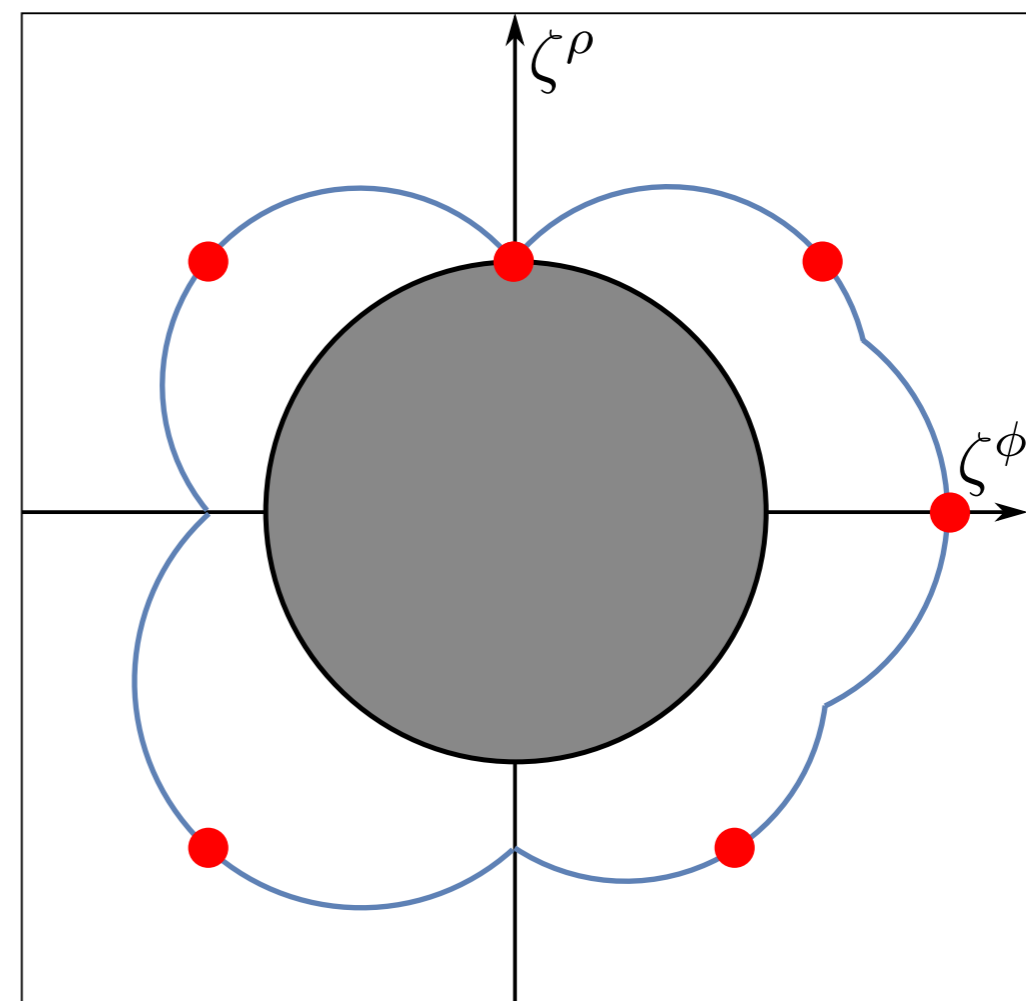
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Let us check these ideas in 9 dimensional theories with 16 supercharges
(e.g. circle compactification of heterotic string theory)

SWGCG plot



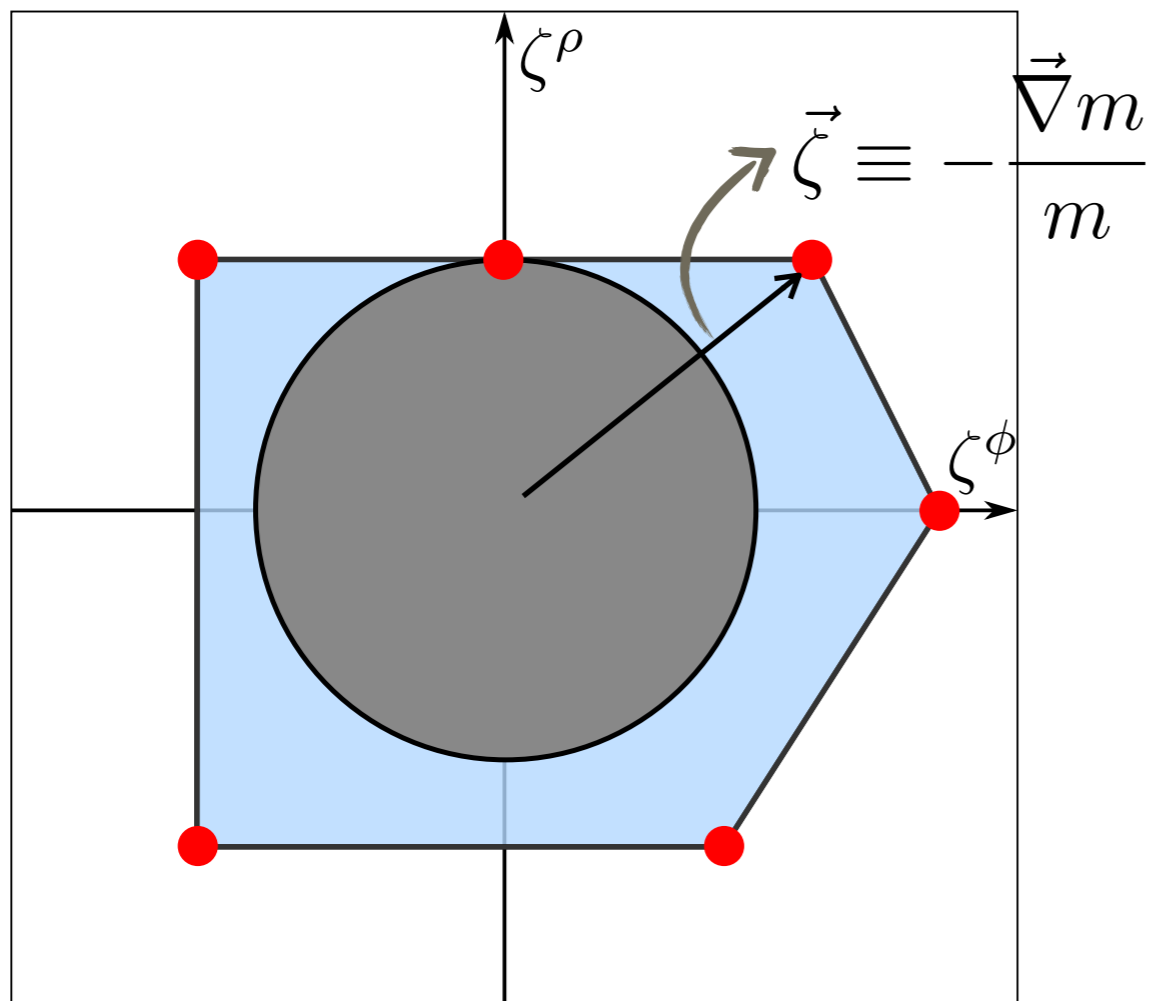
Max-alpha plot



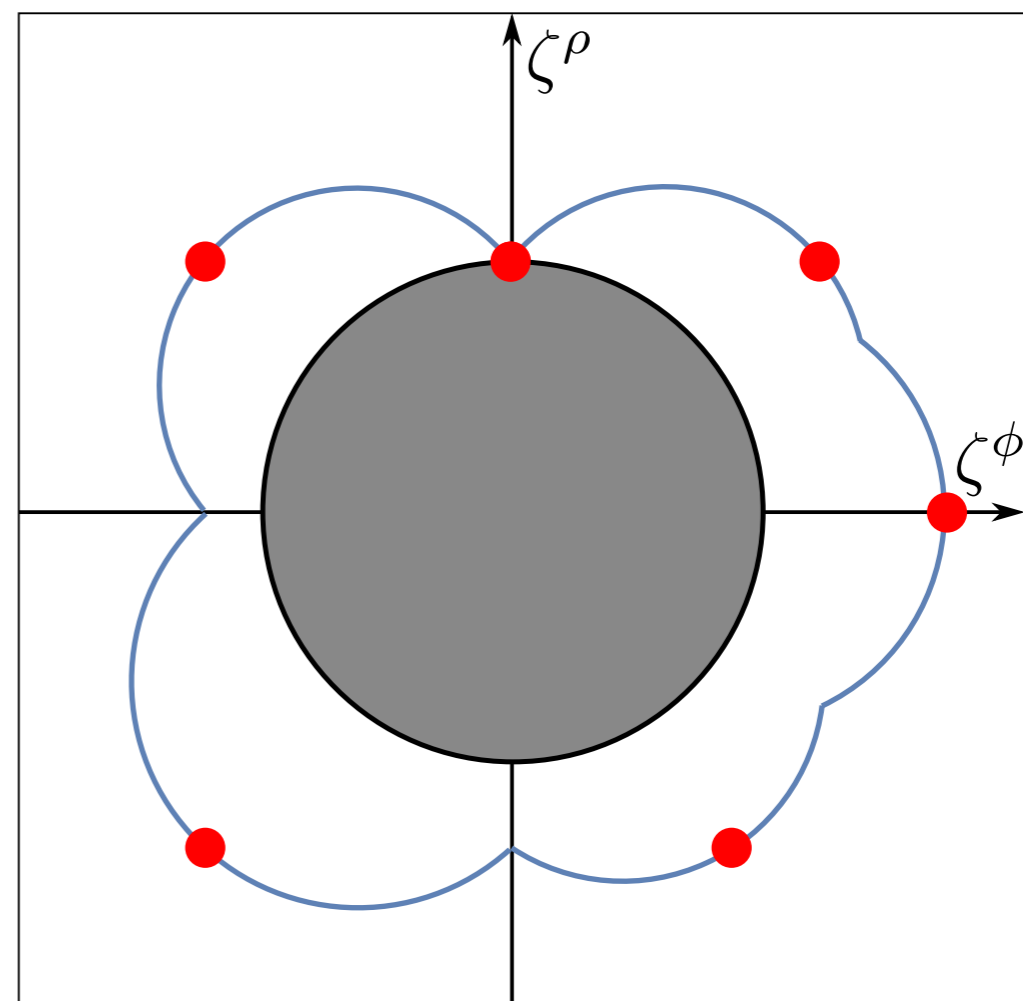
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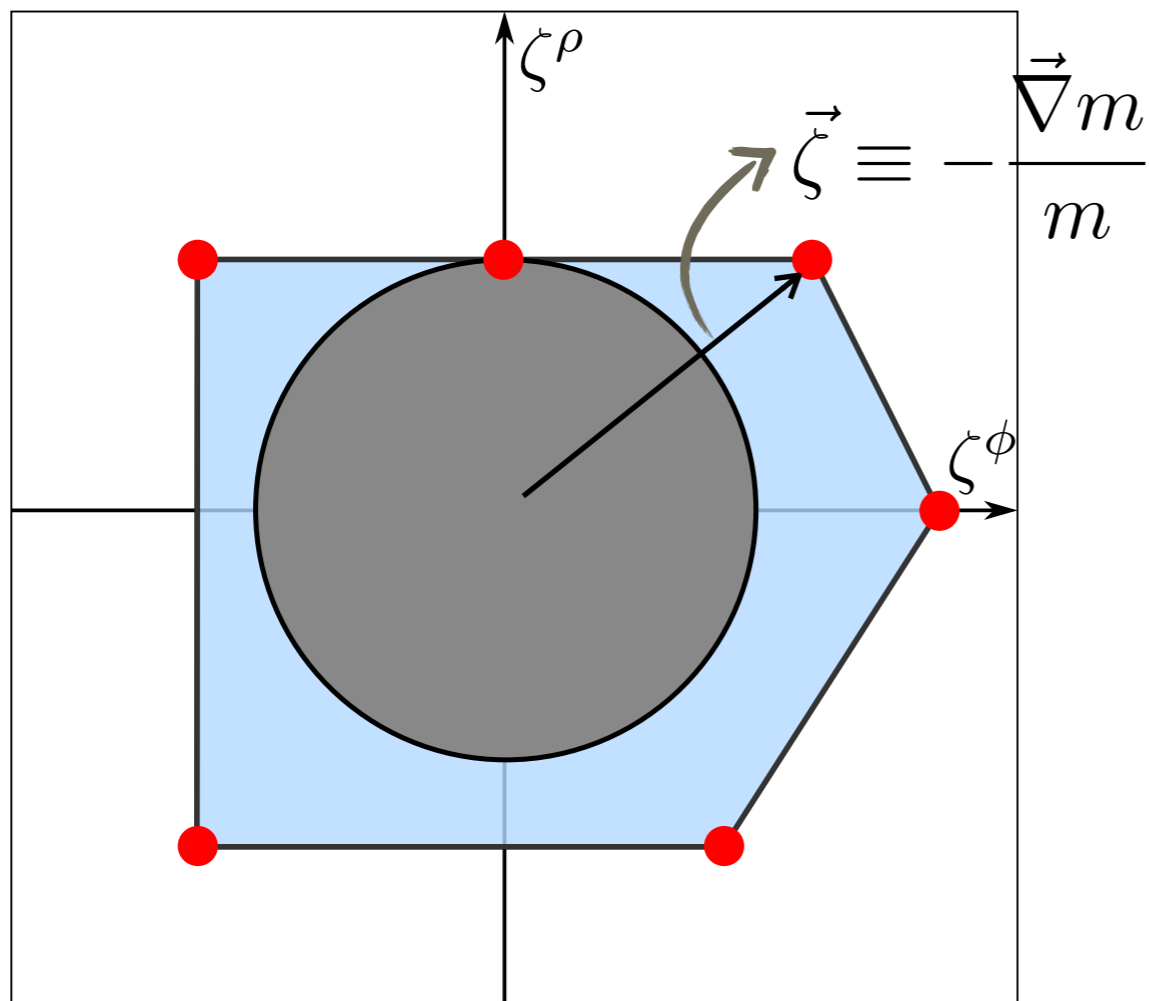
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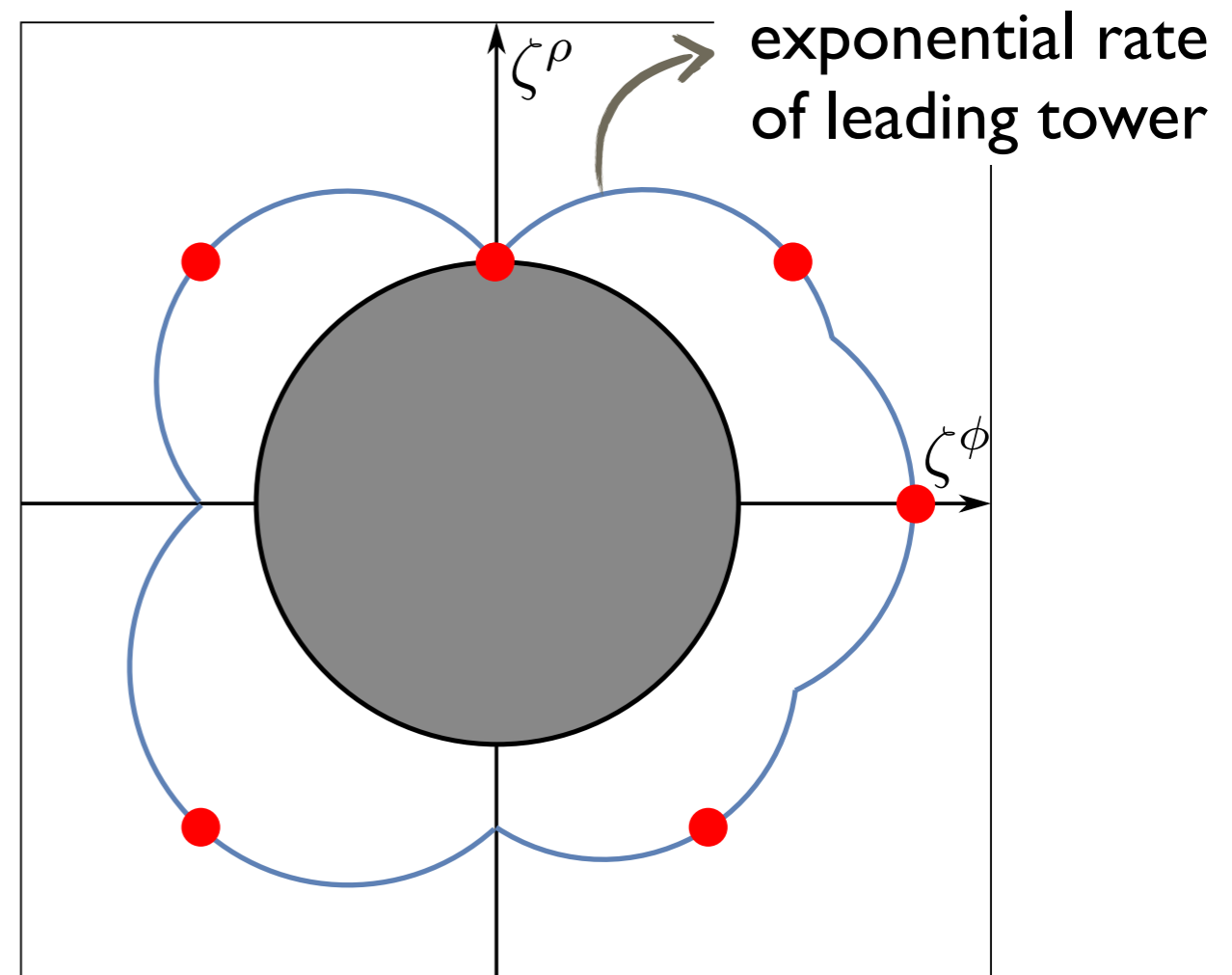
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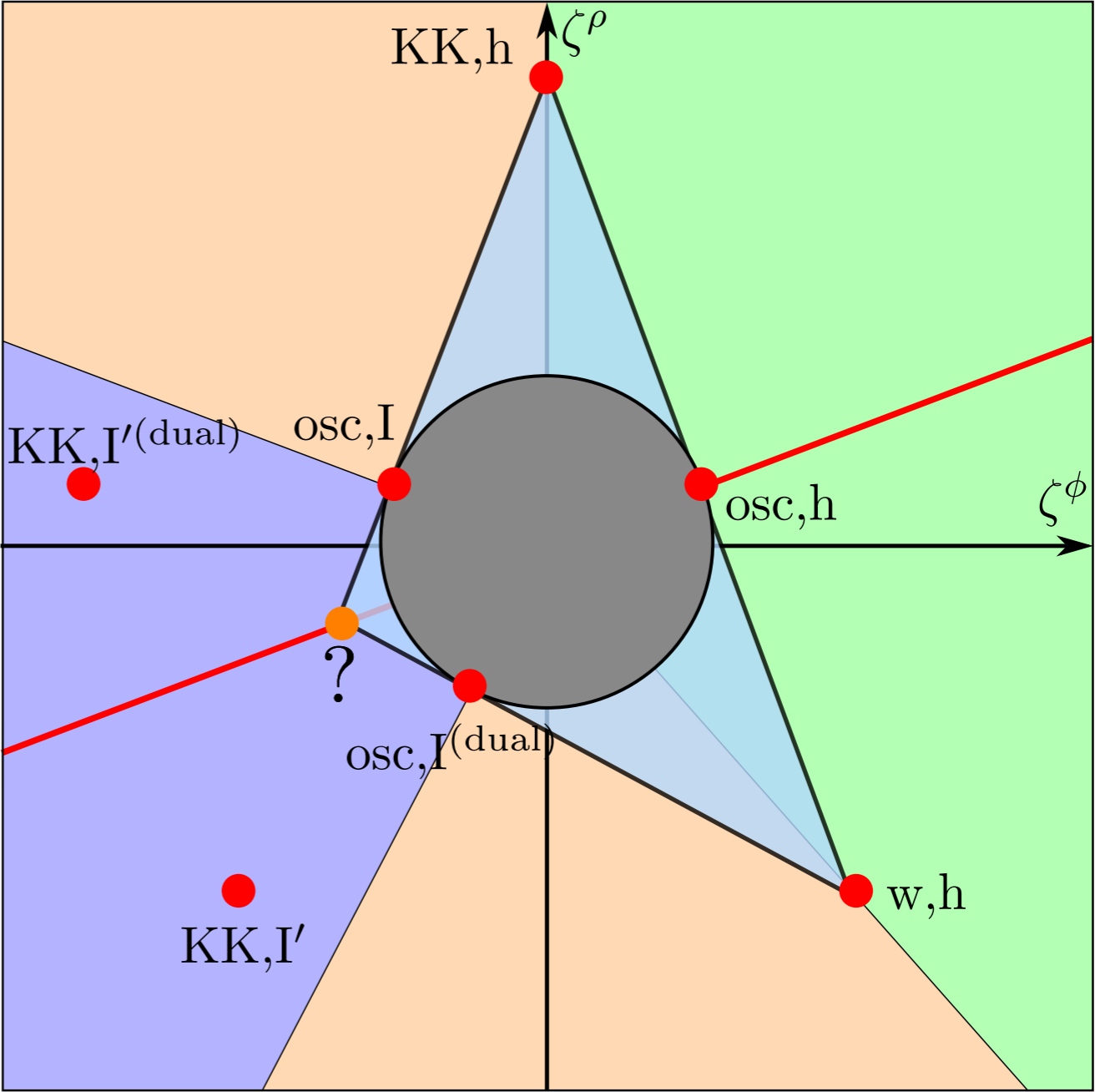
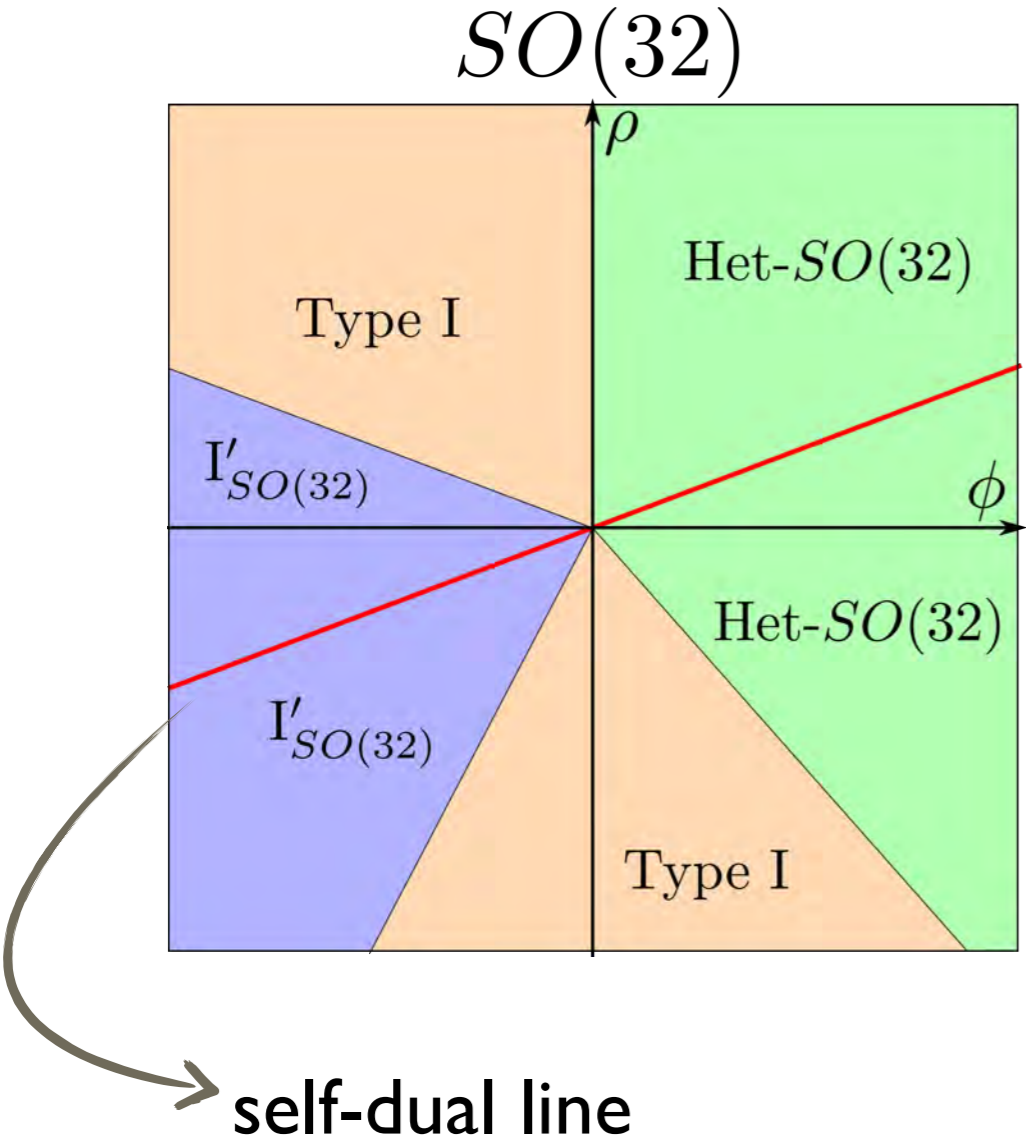


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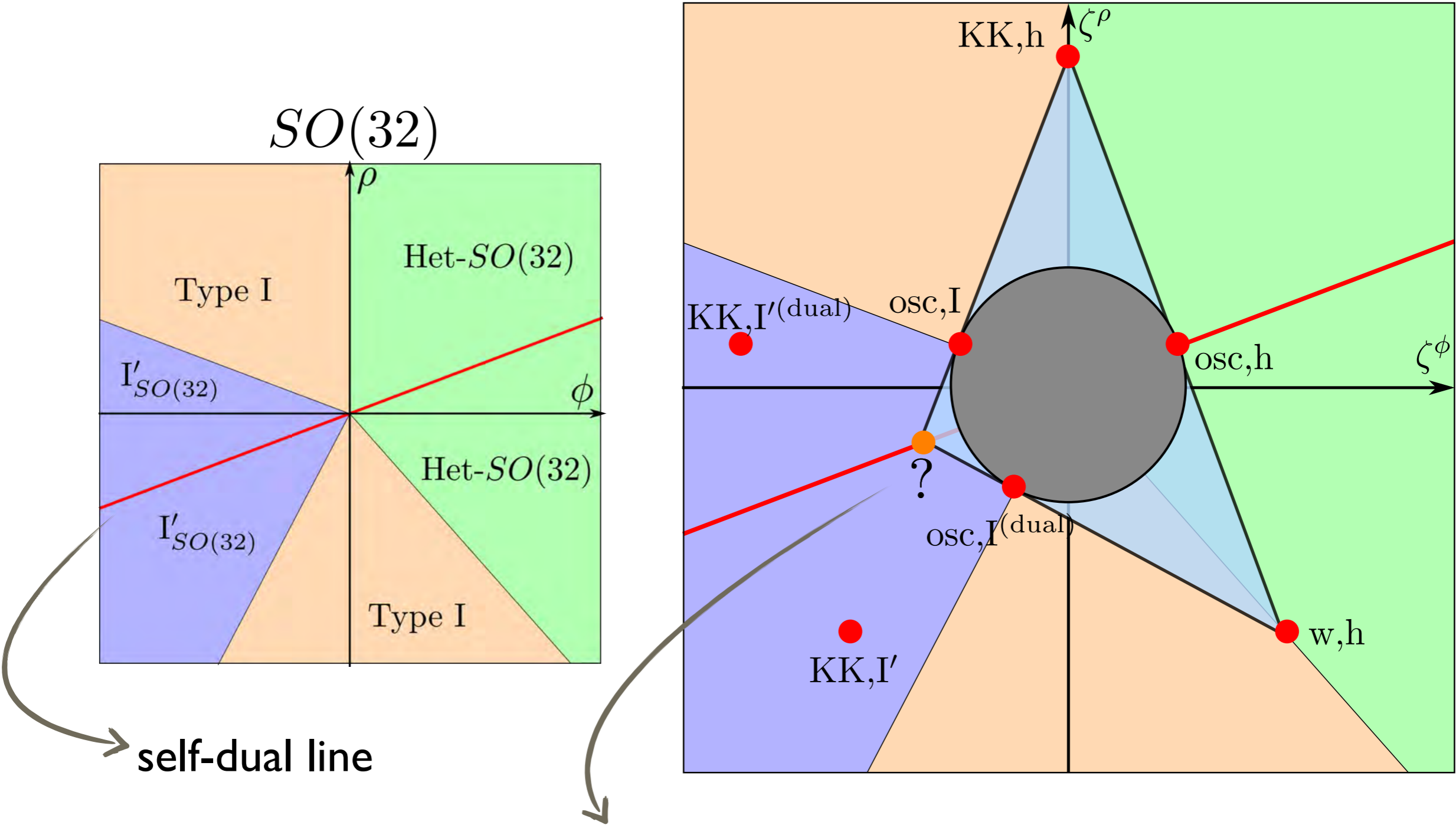
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SO(32) slice



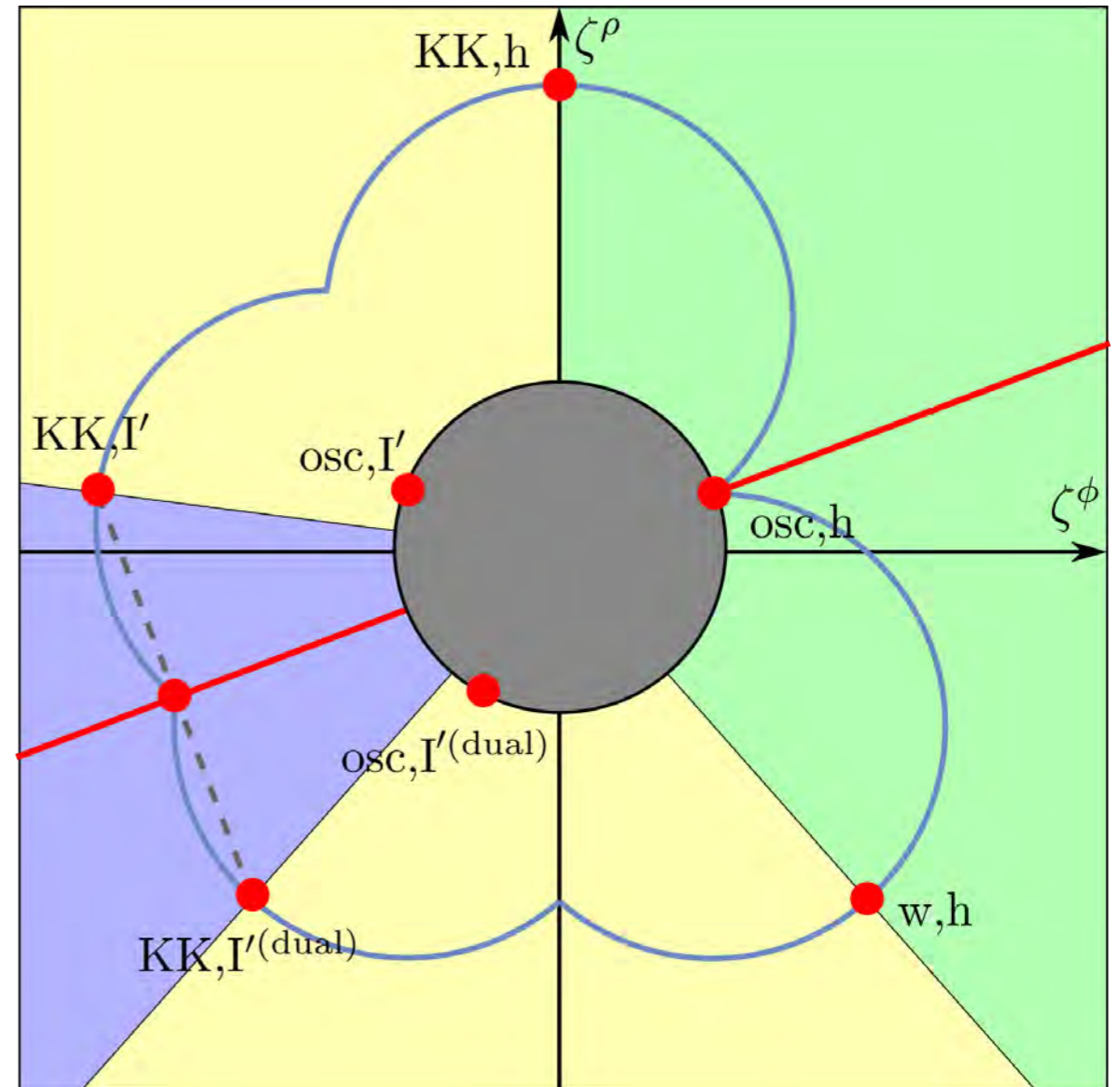
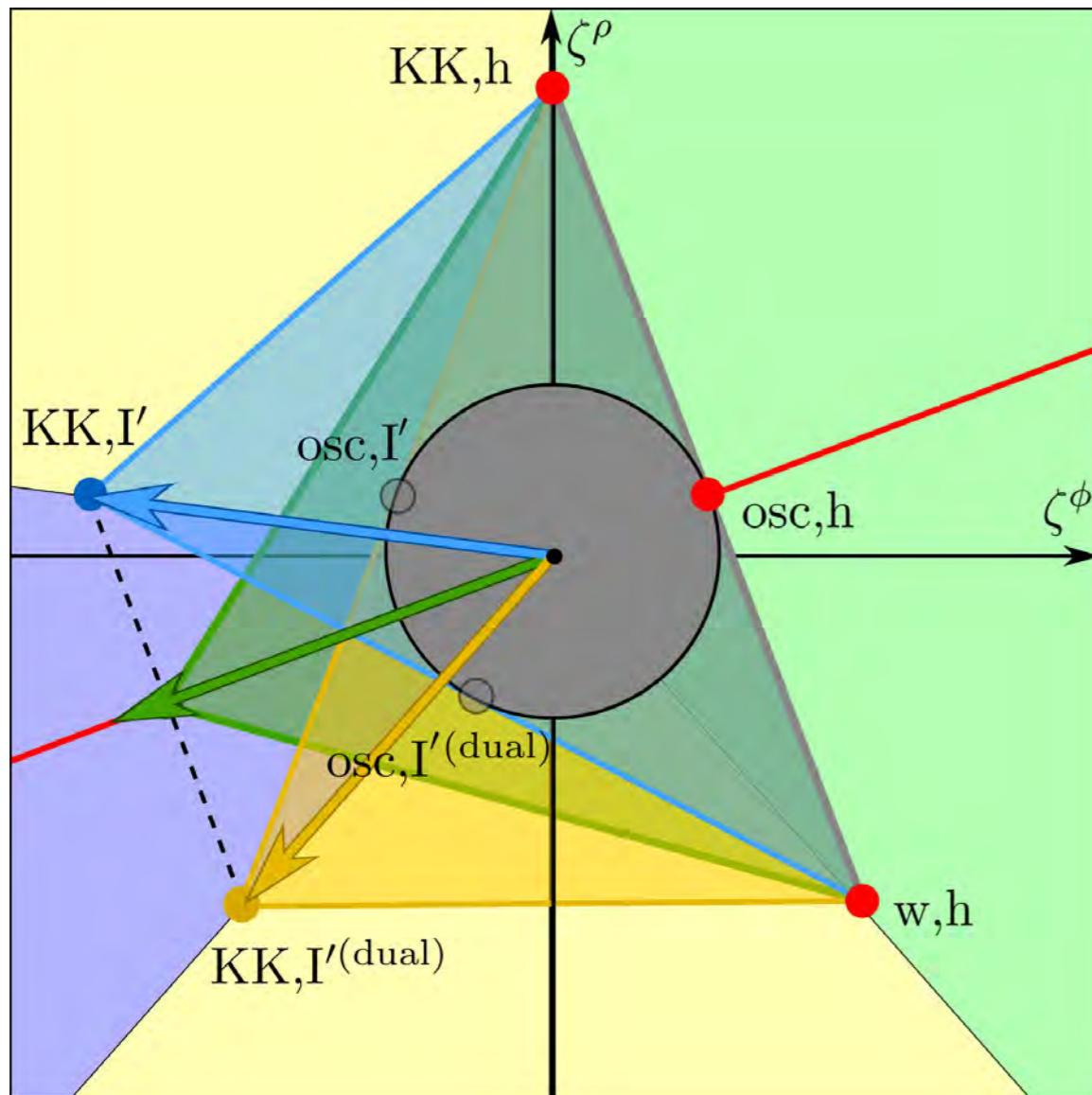
ζ

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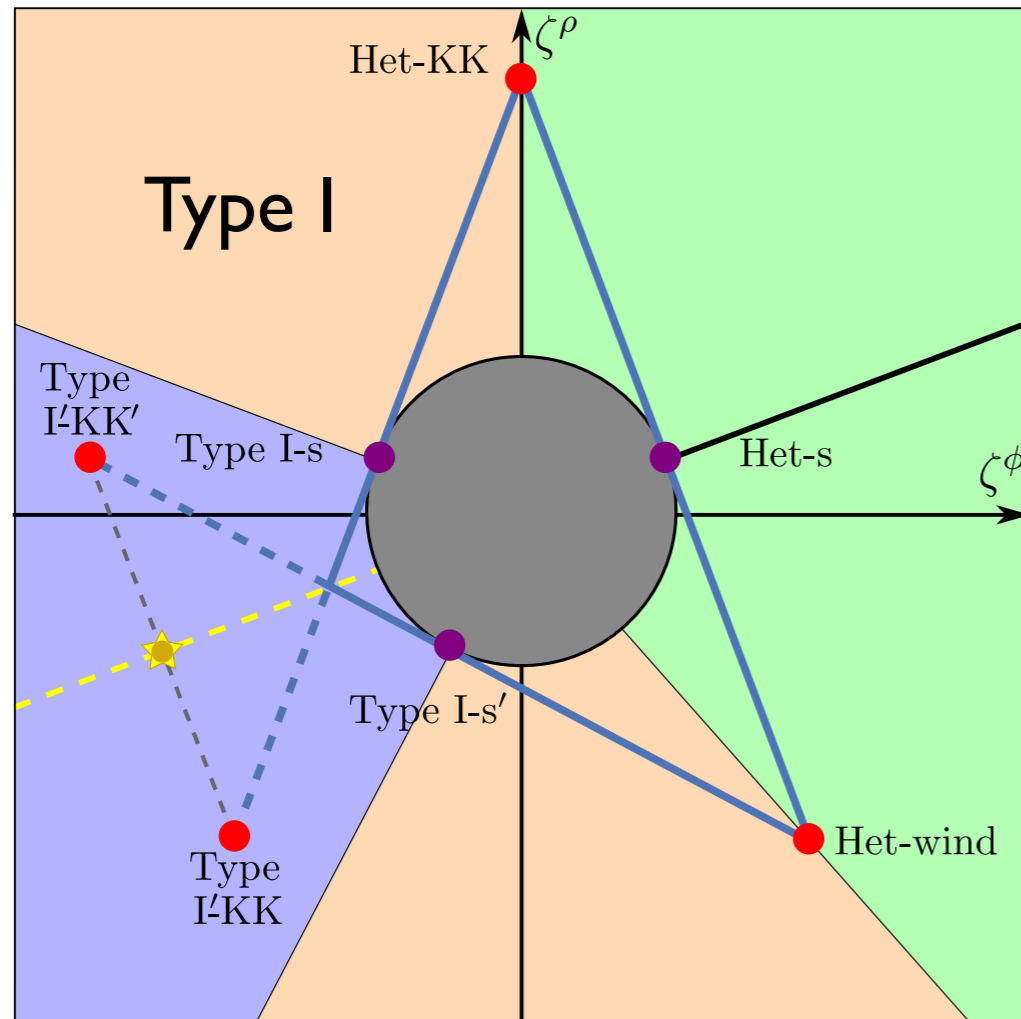
Nothing here, the $\vec{\zeta}$ -vectors of the towers change as we move in the moduli space!

E8xE8 slice

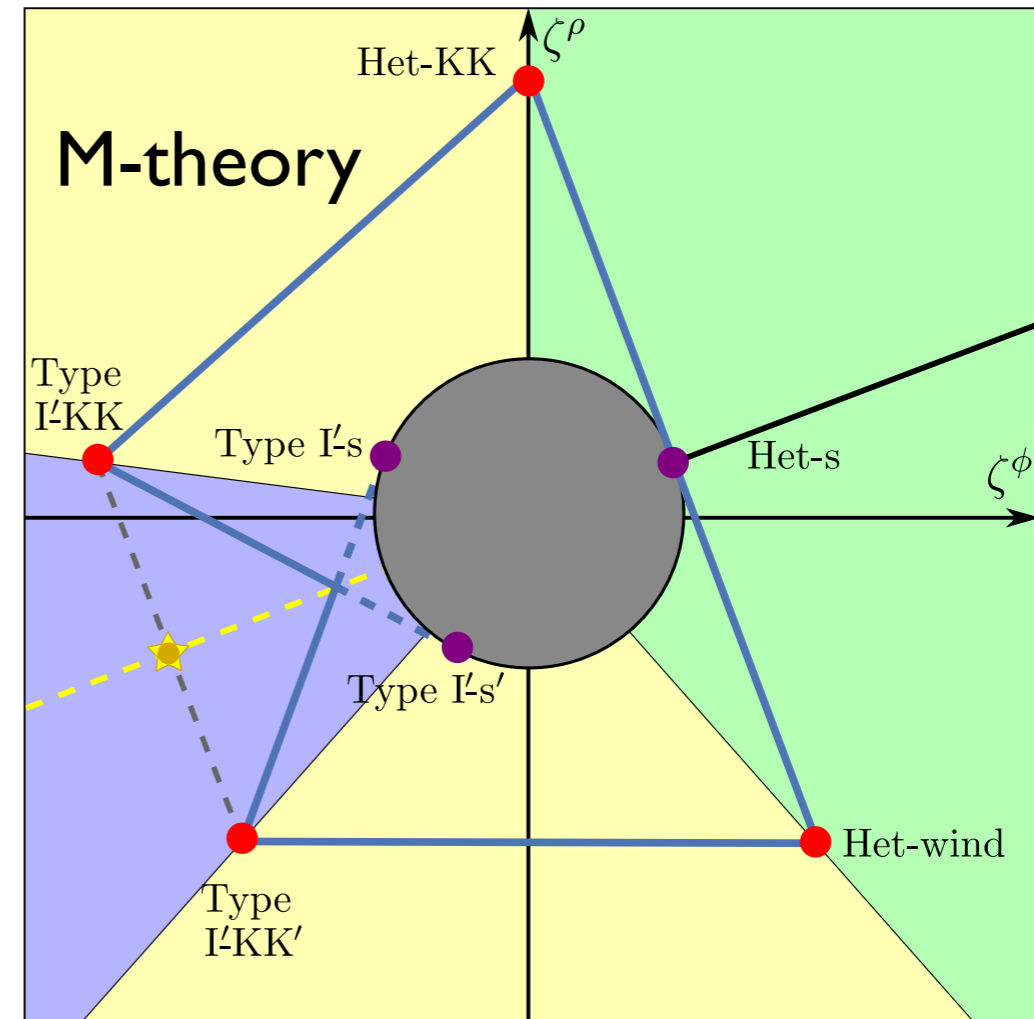


Jumping occurs in opposite direction than the $SO(32)$ case

SO(32) slice



E8xE8 slice



This contains the essential information to derive the “quantum gravity resolution” / UV completion

(i.e. whether we reach an string theory, a higher dimensional theory, etc.)

(see also [Bedroya,Raman,Tarazi'23])

Each asymptotic region is associated to a concrete species scale

WGC and SDC from Entropy Bounds

Take Einstein-Maxwell-Dilaton theory:

$$S = \int d^4x \sqrt{-g} \left[R + 2|d\phi|^2 + \frac{1}{2g(\phi)^2} |F|^2 \right] \quad \text{s.t.} \quad g(\phi) \rightarrow 0 \quad \text{as} \quad \phi \rightarrow \infty$$

There are electrically charged BH solutions with classical zero area (small BHs)

If $g(-\infty) \rightarrow 0$ then $A(-\infty) \rightarrow 0$: **Small BH**

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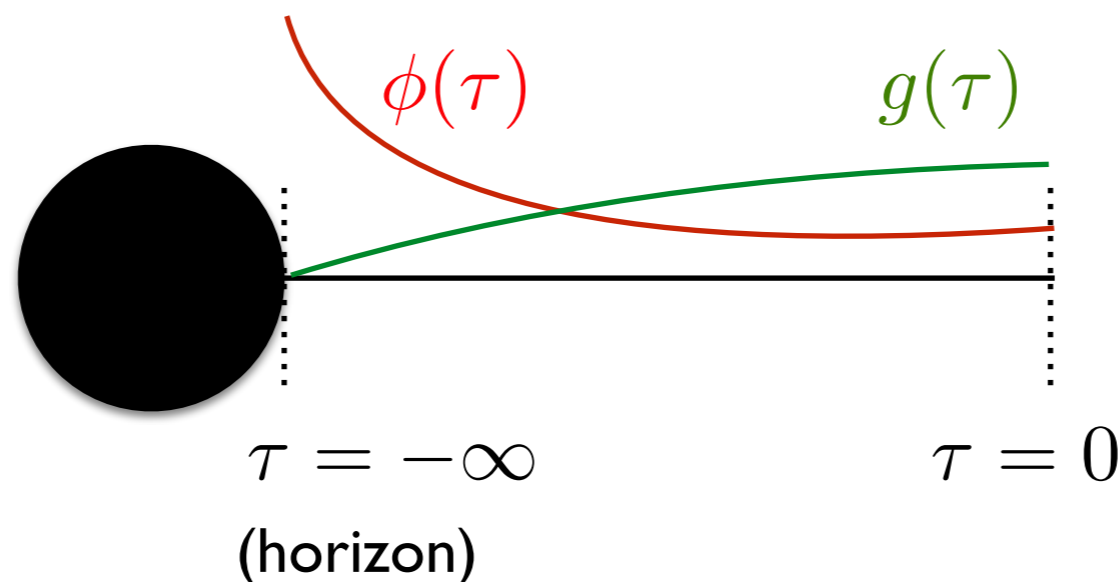
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BH induces a running of the scalar field and gauge coupling as approaching the horizon leading to:



large field range!
small gauge coupling!

WGC and SDC from Entropy Bounds

Small BHs lead to a violation of the Bekenstein bound, unless the EFT cutoff decreases as dictated by the SDC / WGC

Entropy Bound:

A region of size L cannot have more entropy than a Schwarzschild black hole of the same area $A = L^2$

$$N_{\text{species}} = Q_{\text{max}} \lesssim L^2 = A$$

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Using extremality condition and that EFT breaks down at $|d\phi|^2 \sim \Lambda^2$



$$\Lambda \lesssim g \quad \text{in Planck units}$$

due to an infinite tower of states