Quantum field theory, mathematics, and their recent interactions

Yuji Tachikawa (Kavli IPMU, U. Tokyo)

KAIST physics colloquium

Apr. 12, 2021

Life between mathematics and theoretical physics

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I am usually categorized as a string theorist.

Is what I do physics, or mathematics, or mathematical physics?

I only give a light-hearted analysis of this issue.

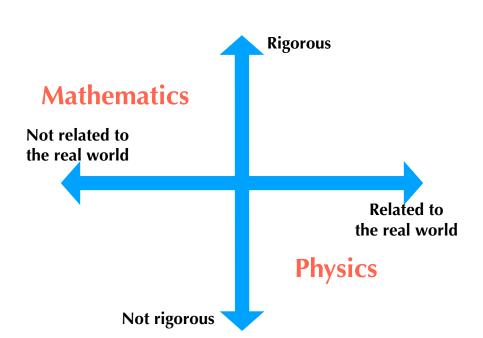
For more incisive discussions,

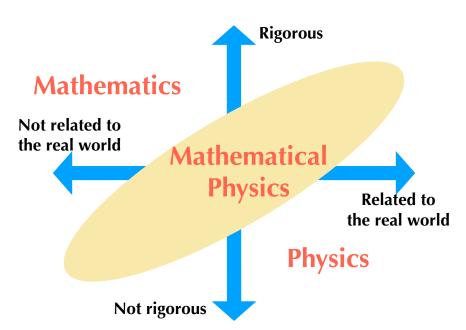
see essays such as Jaffe-Quinn: Theoretical mathematics

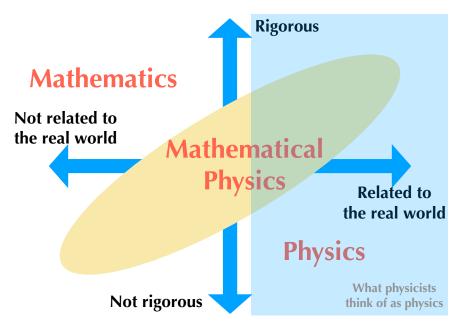
https://arxiv.org/abs/math/9307227

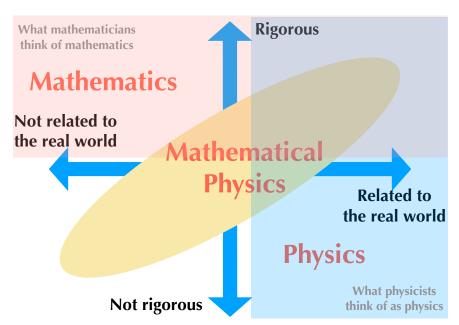
and Thurston: On proof and progress in mathematics

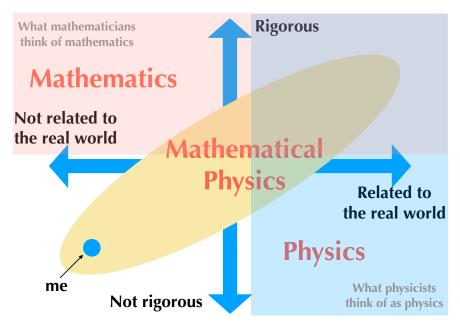
https://arxiv.org/abs/math/9404236











Is there any use in something which is not rigorous

and **not related to the real world** at the same time?

I want to say yes...

Well, I got invited to give a colloquium here by doing that!

Experimental Physics

Work hard

←

Show new experimental results

Experimental Physics

Think about the reason

Experimental Physics

Theoretical Experimental Physics Physics

Predict new phenomena

Experimental Physics

Work hard

Rigorous Mathematics

Work hard to prove something

•

Rigorous Mathematics

Present new theorems

Rigorous Mathematics

Think of their physical significance

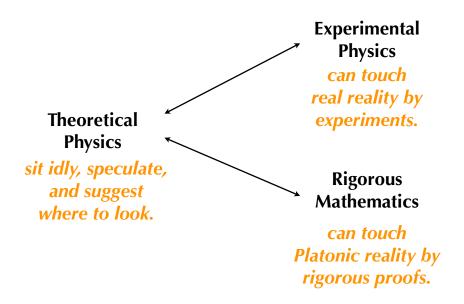
Theoretical _____ Physics

Rigorous Mathematics

Predict new theorems

Rigorous Mathematics

Work hard to prove something



Does it actually work?

In some cases, yes!

Once upon a time (meaning that it was in the 1980s), mathematicians were **counting the numbers of spheres** in a particular **six-dimensional space**, called the **quintic Calabi-Yau**.

size	number	person	
1	2875	J. Harris	(1979)
2	609250	S. Katz	(1986)

It was a laborious process.

I don't know why they were interested in this question!

Around the same time, superstring theory was born.

It says the world is 9 + 1 dimensional.

To match with the fact that our world looks $\mathbf{3} + \mathbf{1}$ dimensional, we need to curl up the unwanted $\mathbf{9} - \mathbf{3} = \mathbf{6}$ dimensions into a very, very small space.

It was found that the same **Calabi-Yau space** is a nice choice for this purpose.

So physicists started studying them too.

Slices of the six-dimensional quintic Calabi-Yau you often see in books and TVs.

https://www.wolframcloud.com/obj/yuji.tachikawa/Published/calabi_yau_based_on_jakes_code3.nb

String theorists P. Candelas, X. de la Ossa, P. Green and L. Parkes found in early 1990 a vastly quicker but non-rigorous method to compute the number of spheres in the same quintic studied by mathematicians.

Recall:

```
size number person n_1 = 2875 J. Harris (1979) n_2 = 609250 S. Katz (1986) n_3 = ???
```

The method of four string theorists involved differential equations and expanding the solutions in a Taylor series, as physicists would naturally do. They predicted

$$n_3 = 317, 206, 375.$$

It fell to mathematicians **G. Ellingsrud** and **S. A. Strømme** to test it. In **June 1990**, they got

$$n_3 = 2,682,549,425.$$

There was a joint math-physics workshop in May 1991 to resolve the issue, so that each side can learn the other side.

There was some progress but the issue remained...

Finally in **July 1991**, Ellingsrud and Strømme found a bug in their computation, and reproduced the prediction by physicists.

$$n_3^{
m phys} = 317, 206, 375. \ n_3^{
m math} = 317, 206, 375.$$

This was when the mathematical field called the **mirror symmetry** was born.

(Details taken from P. Galison, *Mirror symmetry*, in "Growing Explanations," M. Norton Wise ed., Duke University Press, 2004.

https://doi.org/10.1515/9780822390084-002

I thank D. R. Morrison for information.)

There are many other examples of such interactions between mathematics and theoretical physics.

I was lucky to have been involved in one, called the **Mathieu Moonshine**. arXiv.org > hep-th > arXiv:1004.0956

Search...

Help | Advance

High Energy Physics - Theory

[Submitted on 6 Apr 2010 (v1), last revised 25 Jun 2010 (this version, v2)]

Notes on the K3 Surface and the Mathieu group M_24

Tohru Eguchi, Hirosi Ooguri, Yuji Tachikawa

We point out that the elliptic genus of the K3 surface has a natural decomposition in terms of dimensions of irreducible representations of the largest Mathieu group M_24. The reason is yet a mystery.

Comments: 10 pages. v2: published version

Subjects: High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Group Theory (math.GR);
Ouantum Algebra (math.OA)

Journal reference: Exper.Math.20:91-96,2011

DOI: 10.1080/10586458.2011.544585

Cite as: arXiv:1004.0956 [hep-th]

(or arXiv:1004.0956v2 [hep-th] for this version)

Search o

arXiv.org > math > arXiv:1211.5531

Mathematics > Representation Theory

Much ado about Mathieu

Terry Gannon

(Submitted on 23 Nov 2012 (v1), last revised 15 Mar 2013 (this version, v2))

Eguchi, Ooguri and Tachikawa have observed that the elliptic genus of type II string theory on K3 surfaces appears to possess a Moonshine for the largest Mathieu group. Subsequent work by several people established a candidate for the elliptic genus twisted by each element of M24. In this paper we prove that the resulting sequence of class functions are true characters of M24, proving the Equchi-Ooguri-Tachikawa conjecture. We prove the evenness property of the multiplicities, as conjectured by several authors. We also identify the role group cohomology plays in both K3-Mathieu Moonshine and Monstrous Moonshine; in particular this gives a cohomological interpretation for the non-Fricke elements in Norton's Generalised Monstrous Moonshine conjecture. We investigate the intriguing proposal of Gaberdiel-Hohenegger-Volpato that K3-Mathieu Moonshine lifts to the Conway group Co1.

arXiv.org > math > arXiv:2006.02922

Search...

Help | Advance

Mathematics > Algebraic Topology

Submitted on 4 Jun 2020

Topological Mathieu Moonshine

Theo Johnson-Freyd

We explore the Atiyah-Hirzebruch spectral sequence for the $tmf^*[\frac{1}{2}]$ -cohomology of the classifying space BM_{24} of the largest Mathieu group M_{24} , twisted by a class $\omega \in H^4(BM_{24}; \mathbb{Z}[\frac{1}{2}]) \cong \mathbb{Z}_3$. Our exploration includes detailed computations of the \mathbb{Z}_3 cohomology of M_{24}^{-} and of the first few differentials in the AHSS. We are specifically interested in the value of $tmf_{\omega}(BM_{24})[\frac{1}{2}]$ in cohomological degree -27. Our main computational result is that $tmf_{\omega}^{-27}(BM_{24})[\frac{1}{2}] = 0$ when $\omega \neq 0$. For comparison, the restriction map $tmf_{\omega}^{-3}(BM_{24})[\frac{1}{2}] \to tmf^{-3}(pt)[\frac{1}{2}] \cong Z_3$ is nonzero for one of the two nonzero values of ω . Our motivation comes from Mathieu Moonshine, Assuming a well-studied conjectural relationship between TMF and supersymmetric quantum field theory, there is a canonicallydefined Co_1 -twisted-equivariant lifting $[\bar{V}^{fq}]$ of the class $\{24\Delta\} \in TMF^{-24}(pt)$, where Co_1 denotes Conway's largest sporadic group. We conjecture that the product $\lceil \bar{V}^{f^{\eta}} \rceil \nu$, where $\nu \in TMF^{-3}(pt)$ is the image of the generator of $tmf^{-3}(pt) \cong Z_{24}$, does not vanish Co_1 equivariantly, but that its restriction to M_{24} -twisted-equivariant TMF does vanish. This conjecture answers some of the questions in Mathieu Moonshine; it implies the existence of a minimally supersymmetric quantum field theory with M_{24} symmetry, whose twisted-and-twined partition functions have the same mock modularity as in Mathieu Moonshine. Our AHSS calculation establishes this conjecture "perturbatively" at odd primes. An appendix included mostly for entertainment purposes discusses " ℓ -complexes" and their relation to SU(2) Verlinde rings. The case $\ell=3$ is used in our AHSS calculations.

I would like to give some detail of this **Mathieu Moonshine**, but it would be a long way to go. I need to tell you

- what is quantum field theory,
- what is string theory,
- what are the Mathieu groups, and
- what is the moonshine.

Let me try.

What is **Quantum Field** Theory = QFT?

- Describes quantum properties of fields, where
- fields are anything which extend along space and time, such as
- electromagnetic fields (=light), crystal vibration, electron fields ...

The prototypical example is the **Quantum Electrodynamics** (QED):

- describes quantized electromagnetic fields interacting with charged particles
- was established around 1950s
- with many developments since then

Theory and experiment match extremely well in QED.

The prime example is the **anomalous magnetic moment of electron**:

```
a_e^{	ext{theory}} = 0.01\,159\,652\,181... a_e^{	ext{experiment}} = 0.01\,159\,652\,181...
```

I can say that:

- QFT is well researched,
- QFT predicts quantities with high precision, and
- QFT agrees very well with experiment.

But it is **mathematically incomplete**, in that no satisfactory formalization is known.

Quantum mechanics and general relativity are OK for mathematicians, but QFT is not.

It is analogous to the situation in the past:

Ancient Egyptians could build pyramids, although they had not formalized geometry.

Physicists can compute things although they have not formalized QFT.

The flip side of the coin is that QFT might produce new results in mathematics.

Let us move on to **string theory**.

It is a **quantum** theory of **strings moving relativistically**. It turns out that:

- it is consistent only in 9+1 dimensions
- it automatically contains quantum gravity

Reconciling gravity and quantum mechanics is one of the long-standing problems in physics. There are many competing approaches.

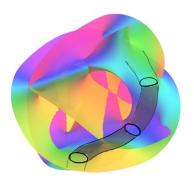
There are many who are passionately for string theory, and also many who are passionately against string theory.

Does it describe the real world? I do not know.

For me what matters is whether it is Platonically consistent.

It seems it is. And many mathematical predictions have come out of it.

Suppose you want to study strings moving in a Calabi-Yau...



Please excuse my bad drawing.

It is done in terms of 1+1 dimensional QFT on the worldsheet.

Mathieu Moonshine concerns strings moving in K3 space.

It is a **closed four-dimensional** space satisfying

$$\frac{1}{2}\epsilon_{\mu\nu}{}^{\alpha\beta}R_{\alpha\beta\rho\sigma} = R_{\mu\nu\rho\sigma}$$

which is not completely flat.

You can think of it as a space which is **half flat**, in a precise technical sense.

In particular, it solves the vacuum Einstein equation with **zero cosmological constant**.

K3 is named by André Weil, honoring three mathematicians **Kähler**, **Kummer**, and **Kodaira**, and also after the beautiful mountain **K2**:



https://en.wikipedia.org/wiki/K2

The origin of Mathieu Moonshine goes back to the work by Eguchi and Ooguri in 1989.





Ooguri was preparing his PhD thesis under Eguchi, studying strings moving in K3.

A central result in his PhD thesis is this:

so are the numbers $N_{h,1} - 2N_{h,0}$.

$$F(\tau) = 90q + 462q^2 + 1540q^3 + 4554q^4 + 11592q^5 + 27830q^6 + 61686q^7 + 131100q^8 + \cdots$$
 (23)

This is basically the partition function of a string moving in K3.

What does it mean?

https://ooguri.caltech.edu/documents/8002/phd thesis.pdf

To understand it, we now need to turn to the role of **symmetry** in **quantum mechanics**.

Everybody will learn / learned that the angular momentum in quantum mechanics takes the value

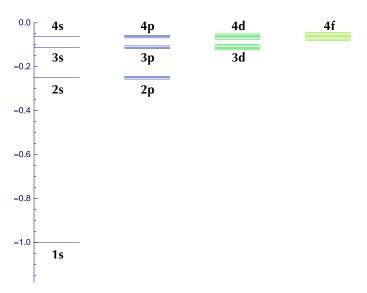
$$j_z = \underbrace{-j, -j+1, \ldots, j-1, j}_{2j+1 \text{ choices}}.$$

So

For integer j, there are also traditional names

$oldsymbol{j}$	0	1	2	3	• • •
name	s	\boldsymbol{p}	d	\overline{f}	• • •
degeneracy	1	3	5	7	• • •

The spectrum of a hydrogen atom, to the zeroth approximation, looks like this:



The **degeneracy** is related to the **symmetry**.

The angular momentum operators $L_{x,y,z}$ are infinitesimal generators of the three dimensional rotation group so(3), the symmetry of the hydrogen atom.

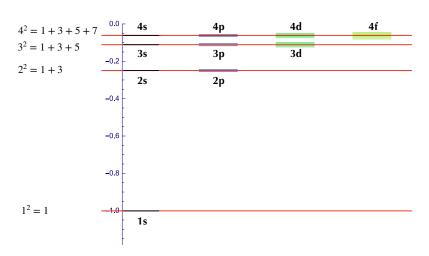
representation=	$oldsymbol{j}$	0	1	2	3	• • •
	name	s	\boldsymbol{p}	d	f	• • •
dimension=	degeneracy	1	3	5	7	• • •

The total angular momentum j specifies how the symmetry acts on the quantum states.

Equivalently, it specifies the representation of the symmetry group.

The degeneracy is also known as the **dimension**.

The spectrum of a hydrogen atom, to the zeroth approximation, shows accidental degeneracies in addition to the rotational symmetry:



It is known to reflect a hidden 4-dimensional rotational symmetry so(4) [Pauli, Fock, ...]

The rotation group is a continuous group.

There are also finite groups, e.g. the symmetry ${\it A}_{\it 5}$ of



which contains 60 elements.

For the symmetry A_5 to act on quantum mechanical systems, it needs to be **represented** by matrices.

$$A_5
ightarrow g(g)=egin{pmatrix} lackbox{ac}ab}ackbox{lackbox{lackbox{lackbox{lackbox{lackbox{lack$$

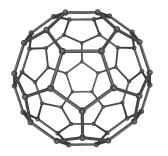
The size is known as the **dimension**, which is three in this example.

Irreducible representations and their dimensions of A_5 are known:

name	\boldsymbol{A}	T_1	T_2	\boldsymbol{G}	\boldsymbol{H}
dimension	1	3	3	4	5

(There are two different irreducible representations with the same dimension 3).

We physicists are perfectly happy studying various concrete systems with various concrete symmetries. This is the schematic structure of C_{60} , the fullerene, from Wikipedia:



https://en.wikipedia.org/wiki/Fullerene

The properties of ${m A}_5$ is very useful (and essential) when studying its electronic properties, etc.

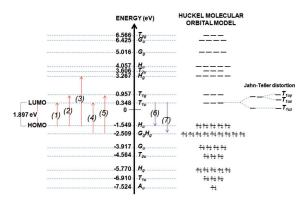


Figure 7: Schematic of the electronic structure of C_{∞} as calculated by the Hückel model [80, 81] and its possible electronic excitation transitions experimentally observed in this study. The electronic excitation and emission transitions were numbered in the bracket as (1) to (5) and (6) to (7), respectively.

from T. E. Saraswati et al., *The Study of the Optical Properties of C60 Fullerene in Different Organic Solvents*, Open Chem. 17 (2019) 119–1212

https://doi.org/10.1515/chem-2019-0117

Mathematicians think differently:

Let's classify all possible symmetries, say all finite groups.

Any **finite group** is made out of **finite simple groups**, just as any **integer** is a product of **prime numbers**.

So they say: let us classify finite simple groups first.

- Cyclic group of prime order \mathbb{Z}_p , $p = 2, 3, 5, \ldots$
- Alternating groups A_5 , A_6 , ...
- Finite groups of Lie type, obtained by considering continuous groups over finite fields,
- and finally, the 26 **sporadic groups**.

The proof is said to be the **longest in the history** of mathematics. Originally announced to be complete in the late 1970s to early 1980s with papers and preprints said to **total 5000 pages**.

https://doi.org/10.1090/S0273-0979-1979-14551-8

A streamlined rewrite of the entire proof in a single series of volumes is going on for decades. It already has about **3500 pages**, but is yet not complete.

https://www.ams.org/publications/authors/books/postpub/surv-40

https://www.ams.org/journals/notices/201806/rnoti-p646.pdf

- Cyclic group of prime order \mathbb{Z}_p , $p=2,3,5,\ldots$
- Alternating groups A_5 , A_6 , ...
- Finite groups of Lie type
- and finally, the 26 sporadic groups.

The last two series of finite groups of Lie type were found by Rimhak Ree (이임학, 李林學) in 1960/1961.

```
https://doi.org/10.1090/S0002-9904-1960-10523-X
https://doi.org/10.1090/S0002-9904-1961-10527-2
https://dx.doi.org/10.5169/seals-685366
```

They are called Ree groups.

Lie groups and Ree groups are difficult to distinguish for Koreans and Japanese alike, since we don't have distinctions between /r/ and /l/...



(Prof. Rimhak Ree, 1922-2005)

http://news.khan.co.kr/kh_news/khan_art_view.html?artid=201510302151325

https://horizon.kias.re.kr/13561/

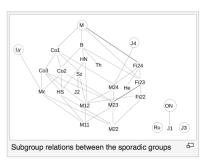
- Cyclic group of prime order \mathbb{Z}_p , $p=2,3,5,\ldots$
- Alternating groups A_5 , A_6 , ...
- Finite groups of Lie type
- and finally, the 26 sporadic groups.

Names of the sporadic groups [edit]

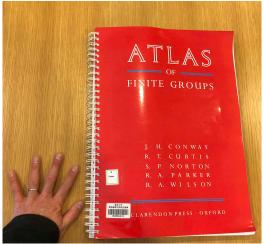
Five of the sporadic groups were discovered by Mathieu in the 1860s and the other 21 were found between 1965 and 1975. Several of these groups were predicted to exist before they were constructed. Most of the groups are named after the mathematician(s) who first predicted their existence. The full list is:

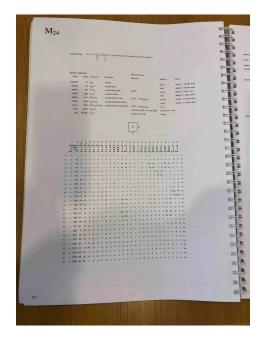
Mathieu groups M_{11} , M_{12} , M_{22} , M_{23} , M_{24}

- . Janko groups J₁, J₂ or HJ, J₃ or HJM, J₄
- . Conway groups Co_1 or F_{2-} , Co_2 , Co_3
- . Fischer groups Fi22, Fi23, Fi24' or F3+
- . Higman-Sims group HS
- . McLaughlin group McL
- . Held group He or F_{7+} or F_{7}
- . Rudvalis group Ru
- . Suzuki sporadic group Suz or F₃_
- . O'Nan group O'N
- . Harada-Norton group HN or F_{5+} or F_{5+}
- . Lyons group Ly
- . Thompson group Th or $F_{3|3}$ or F_3
- . Baby Monster group B or F_{2+} or F_2
- . Fischer-Griess Monster group If or F₁



When it comes to the data of finite groups, nothing can beat **the ATLAS**:





Although not as comprehensive as the Atlas, this Japanese math encyclopedia is also quite useful:



which has tables of representations of some finite groups.

Mathieu groups: M_{11} , M_{12} , M_{22} , M_{23} and M_{24}

Found by Mathieu in 1861 and 1873.

The largest of them, M_{24} , has 244823040 elements.

It is the **symmetry** of the extended binary **Golay code**, introduced in 1949, only one year after Shannon introduced the information theory.

Notes on Digital Coding*

The consideration of message coding as a seasof or approaching the theoretical capacisation of the communication channel, while reducing the probability of errors, has suggested the interesting number theoretical problem of devising lossless hinary (or other) coding schemes serving to insure the reception of a orrect, but reduced, message when an uper limit to the number of transmission errors is notatilated.

An example of lossless binary coding is treated by Shannon1 who considers the case of blocks of seven symbols, one or none of which can be in error. The solution of this rase can be extended to blocks of 2n-1-binary symbols, and, more generally, when coding shemes based on the prime number o are employed, to blocks of $\phi^n - 1/\phi - 1$ symbols which are transmitted, and received with complete equivocation of one or no symbol, each block comprising n redundant symbols designed to remove the equivocation. When encoding the message, the n redundant symbols xm are determined in terms of the message symbols Yo from the congruent relations

$$E_n \equiv X_m + \sum_{k=1}^{k = (p^n - 1)/p - 1) - n} a_{mk} \, Y_k = 0 \, \, (\text{mod } p).$$

In the decoding process, the E's are recalculated with the received symbols, and their ensemble forms a number on the base p which determines univocally the mistransmitted symbol and its correction.

In passing from n to n+1, the matrix with n rows and $p^n-1/p-1$ columns formed with the coefficients of the X's and Y's in the expression above is repeated p times horizontally, while an (n+1) st row added, consisting of $p^n-1/p-1$ zeroes, followed by as many one's etc. up to p-1; an added column of n zeroes with a one for the lowest term completes the new matrix for n+1.

If we except the trivial case of blocks of 2S+1 binary symbols, of which any group comprising up to S symbols can be received in error which equal probability, it does not appear that a search for lossless coding schemes, in which the number of errors is limited but larger than one, can be systematized so as to vield a family of solutions. A necessary but not sufficient condition for the existence of such a lossless coding scheme in the binary system is the existence of three or more first numbers of a line of Pascal's triangle which add up to an exact power of 2. A limited search has revealed two such cases: namely, that of the first three numbers of the 90th line, which add up to 212 and that of the first four numbers of the 23rd line, which add up to 211. The first case does not correspond to a lossless coding scheme, for, were such a scheme to exist, we could designate by r the number of Em ensembles corresponding to one error and having an odd number of 1's and by 90-r the remaining (even) ensembles. The odd ensembles corresponding to two transmission errors could be formed by re-entering term by term all the conbinations of one even and one odd ensemble corresponding each to one error, and would number r(90-r). We should have $r+r(90-r)=2^n$, which is impossible for integral values of r.

Me other side, the second case can be coded on as to yield 12 sure symbols, and the $a_{\rm sk}$ matrix of this case is given in Table I. A second matrix is also given, which is that of the only other lossless coding scheme encountered (in addition to the general class mentioned above) in which blocks of eleven ternary symbols are transmitted with no more than 2 errors, and out of which six sure symbols can be obtained.

It must be mentioned that the use of the ternary coding scheme just mentioned will always result in a power loss, whereas the coding scheme for 23 binary symbols and a maximum of three transmission errors yields a power saving of 1½ db for vanishing probabilities of errors. The saving realized with the properties of the properties of

MARCEL J. E. GOLAY Signal Corps Engineering Laboratories Fort Monmouth, N. J

TABLE I																			
	Y_1	Y_1	Y_1	Y_4	$Y_{\mathfrak{s}}$	$Y_{\mathfrak{s}}$	Y_{τ}	Y_{1}	Y,	Y 10	$Y_{\mathfrak{m}}$	Y 11		Y_{i}	Y_2	γ,	Y_{\bullet}	$Y_{\mathfrak{s}}$	Y .
X_1	1	0	0	1	1	1	0	0	0	1	1	1	X_1	1	1	1	2	2	0
X:	1	0	1	0	1	1	0	1	1	0	0	1	X_{\bullet}	1	1	2	1	0	2
X,	1	0	1	1	0	1	1	0	1	0	1	0	Xi	1	2	1	0	1	2
X.	1	0	1	1	1	0	1	1	0	1	0	0	X.	- 1	2	0	1	2	1
X.	1	1	0	0	1	1	ī	0	1	1	Ö.	0	X .	1	0	2	2	1	1
X .	1	1	0	1	O.	1	1	1	0	0	0	1							
X ₇	1	1	0	1	1	0	0	î	1	0	1	Ô							
X.	î	î	í	0	ő	1	0	1	ñ	1	1	0							
X,	1	1	1	0	1	0	1	Ô.	0	0	1	1							
Y	1	î	1	1	0	0	o.	n	1	1	ô	î							
XII	ô	î	î	î	ĭ	1	1	ĭ	î	î	ï	1							

Benrinted from Proc. IRF vol. 37 p. 657 June 1949.

^{*}Received by the Institute, February 23, 1949.

1 C. E. Shannon, "A mathematical theory of communication," Bell Sys. Tech. Jour., vol. 27, p. 418;

What is a **code**? Computers use strings of bits, such as

 $010110111\cdots$

You can't directly communicate them over long distance, because transmission errors might flip bits.

You need to add redundancies so that small number of errors per bit can be corrected.

One encoding is the Golay code, which encodes original 12 bits into 24 bits.

It is a particularly symmetric code: the symmetry is M_{24} , as noticed by Leech in 1967.

https://doi.org/10.4153/CJM-1967-017-0

It was actually used in the real world. One famous example is NASA's Voyager mission. Some of the scientific data from Jupiter was sent back using the Golay code.

This is the actual data of Jupiter from March 1979 which I took from the NASA website.

https://voyager.jpl.nasa.gov/mission/science/jupiter/

(If you read the documents from those days carefully, you find that the photographic image was *not* encoded by the Golay code, which was considered too wasteful.)

Before talking about the Mathieu Moonshine,

I need to talk about the original Monstrous Moonshine.

Modular **J** function

$$J(q) = \frac{1}{q} + \frac{196884q}{q} + 21493760q^2 + 864299970q^3 + \cdots$$

is known from 19th century.

McKay noticed the following in 1978:

The new finite simple group, the Monster, which is the largest of the sporadics, was being constructed at that time and has order $\sim 8 \cdot 10^{53}$.

The smallest nontrivial representation has dimension

196883

Modular **J** function

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196883

Connects two distant branches of mathematics

J function: classical complex analysis

Monster group : **finite group**

Sounded too crazy back then, and called the **Monstrous Moonshine**.

(The word moonshine means foolish thought.)

Mostly solved around the early 1990s [Frenkel-Lepowsky-Meurman], [Borcherds]

and many developments since then.

The proof used ideas from two-dimensional quantum field theories.

I can finally come back to the Mathieu Moonshine.

In his PhD thesis, Ooguri computed the partition function of a **string moving in K3**:

so are the numbers $N_{h,1} - 2N_{h,0}$.

$$F(\tau) = 90q + 462q^2 + 1540q^3 + 4554q^4 + 11592q^5 + 27830q^6 + 61686q^7 + 131100q^8 + \cdots$$
 (23)

This means that

- the first excited state has degeneracy 90,
- the second excited state has degeneracy 462,
- the third excited state has degeneracy **1540**, ...

Around the same time, there was also the following paper:

http://eudml.org/doc/143625

Invent. math. 94, 183-221 (1988)



Finite groups of automorphisms of K3 surfaces and the Mathieu group

Dedicated to Professor Masayoshi Nagata on his 60th Birthday

Shigeru Mukai

Department of Mathematics, Nagoya University, Furō-chō Chikusa-ku, Nagoya 464 Japan

which says that the possible symmetries of K3 are certain small subgroups of the Mathieu group M_{24} .

Eguchi, his advisor, thought:

so are the numbers $N_{h,1} - 2N_{h,0}$.

$$F(\tau) = 90q + 462q^2 + 1540q^3 + 4554q^4 + 11592q^5 + 27830q^6 + 61686q^7 + 131100q^8 + \cdots$$
(23)

These coefficients should be related to Mathieu group.

And nothing happened for twenty years ...

In the meantime, I became a student of Eguchi obtained PhD in 2006, became a postdoc ...

T. Eguchi





Aspen, Colorado, Aug. 6th, 2009.



All three were in the workshop. We revisited the question.

I said:

Why don't we look up the table in the Iwanami math encyclopedia?

so are the numbers $N_{h,1} - 2N_{h,0}$.

$$F(\tau) = 90q + 462q^2 + 1540q^3 + 4554q^4 + 11592q^5 + 27830q^6 + 61686q^7 + 131100q^8 + \cdots$$
 (23)

Iwanami Math Encyclopedia, 4th ed.:

数 表 6 $M_{24} \qquad (1)^{24} \qquad g \qquad 1 \ 23 \ 7 \cdot 36 \ 23 \cdot 11 \ 23 \cdot 77 \ 55 \cdot 64 \ \overline{45} \ \overline{22 \cdot 45} \ \overline{23 \cdot 45} \ 23 \cdot 45 \ \overline{11 \cdot 21} \ \overline{770}$ $(1)^{24} \qquad g \qquad 23 \cdot 21 \ 23 \cdot 55 \ 23 \cdot 88 \ 23 \cdot 99 \ 23 \cdot 144 \ 23 \cdot 11 \cdot 21 \ 23 \cdot 7 \cdot 36 \ 77 \cdot 72 \ 11 \cdot 35 \cdot 27$

so are the numbers $N_{h,1} - 2N_{h,0}$.

$$F(\tau) = 90q + 462q^2 + 1540q^3 + 4554q^4 + 11592q^5 + 27830q^6 + 61686q^7 + 131100q^8 + \cdots$$
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so are the numbers $N_{h,1} - 2N_{h,0}$.

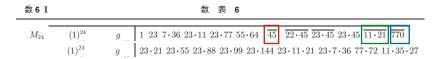
$$F(\tau) = 90q + 462q^{2} + 1540q^{3} + 4554q^{4} + 11592q^{5} + 27830q^{6} + 61686q^{7} + 131100q^{8} + \cdots$$
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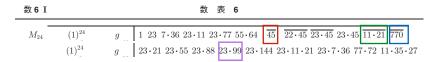
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Iwanami Math Encyclopedia, 4th ed.:



There is a correspondence!

We wrote a paper saying that there is a correspondence, and nothing more:

Experimental Mathematics, 20(1):91–96, 2011 Copyright © Taylor & Francis Group, LLC ISSN: 1058-6458 print DOI: 10.1080/10586458.2011.544585



Notes on the K3 Surface and the Mathieu Group M_{24}

Tohru Eguchi, Hirosi Ooguri, and Yuji Tachikawa

CONTENTS

- 1. Introduction and Conclusions
- 2. Appendix: Data on M24
- 3. Appendix: M₂₄ and the classical geometry of K3 Acknowledgments

References

We point out that the elliptic genus of the K3 surface has a natural decomposition in terms of dimensions of irreducible representations of the largest Mathieu group $M_{\rm 24}$. The reason remains a mystery.

https://doi.org/10.1080/10586458.2011.544585

https://arxiv.org/abs/1004.0956

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https://arxiv.org/abs/1004.0956

My contribution was literally only the suggestion that we should look up the table.

Yes it was **essential**. But it was also totally **trivial**.

Eguchi and Ooguri could have looked up the same table in 1989.

This became one of the most cited papers of mine, and both theoretical physicists and mathematicians still work on it.

